

SOX – searching **sterile neutrinos** with Borexino



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*All evidence from SM in favor of just 3 light neutrinos.
Why should there be more?*

Fundamental (from theory)

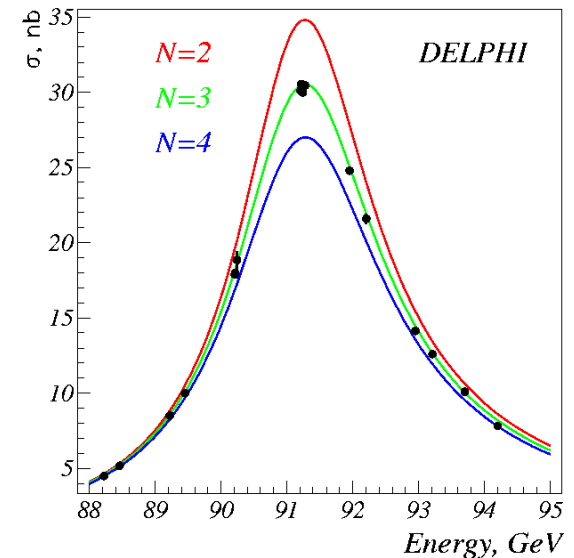
- arguably, the most simple extension of the SM
→ addition of inactive singlet state(s)
- excellent Dark Matter candidate
- required for See-Saw mechanism
→ light active neutrino masses
→ leptogenesis for M/AM asymmetry
→ νMSM ...

Agnostic (from experiments)

- short-baseline oscillation anomalies (eV)
- unexplained X-ray lines:
from keV-DM annihilation?

| | Fermions | | | Bosons | |
|---------|------------------------------|----------------------------|----------------------------|--------------------|----------------|
| Quarks | u up | c charm | t top | γ photon | Force carriers |
| | d down | s strange | b bottom | Z Z boson | |
| Leptons | ν_e electron neutrino | ν_μ muon neutrino | ν_τ tau neutrino | W W boson | |
| | e electron | μ muon | τ tau | g gluon | |
| | | | | Higgs boson | |

Source: AAAS



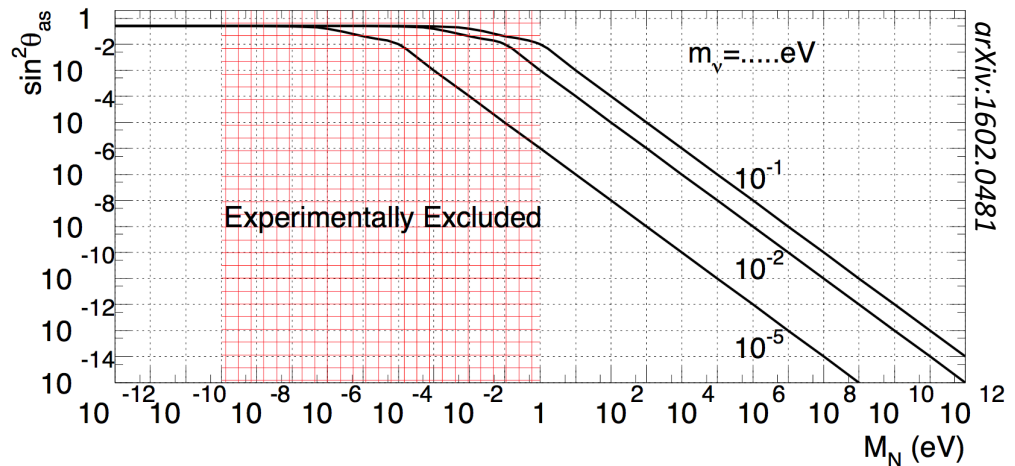
- no interactions with SM particles, but **mixing** with active neutrinos:

extended PMNS matrix

$$\begin{array}{c} \text{weak eigenstates} \end{array} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_{s1} \\ \vdots \end{pmatrix} = \begin{pmatrix} U_{11} & U_{12} & U_{13} & U_{14} & \dots \\ U_{21} & U_{22} & U_{23} & U_{24} & \dots \\ U_{31} & U_{32} & U_{33} & U_{34} & \dots \\ U_{41} & U_{42} & U_{43} & U_{44} & \dots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \\ \vdots \end{pmatrix} \begin{array}{c} \text{mass eigenstates} \end{array}$$

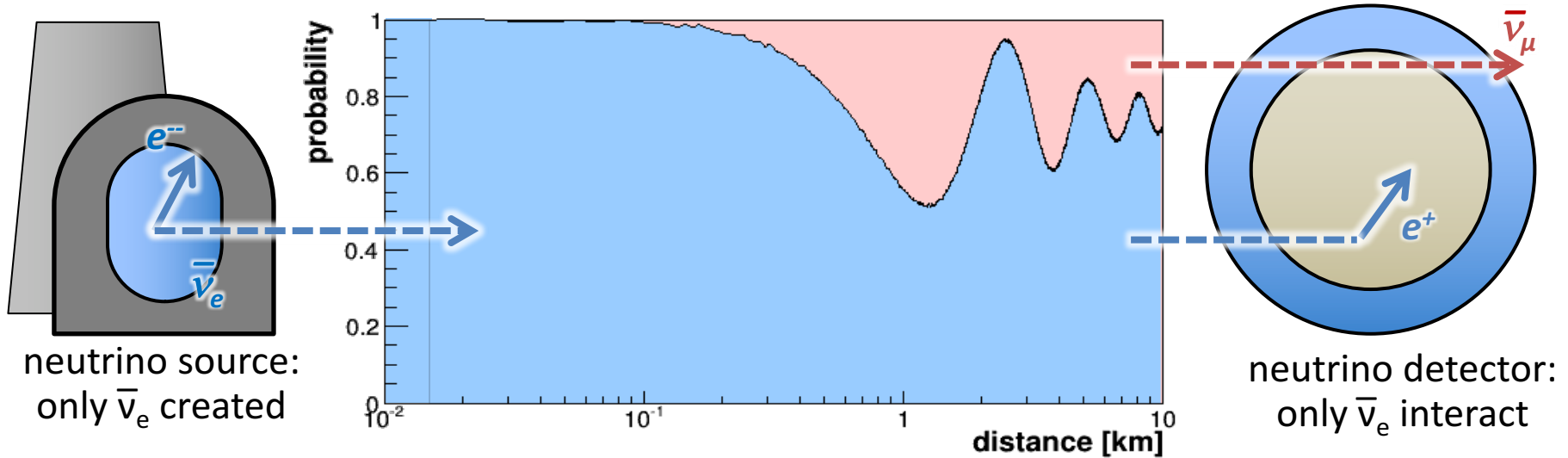
- in See-Saw: Natural scale for active-sterile mixing

$$\Theta \sim \frac{m_D}{M_R}$$



▪ Active states only:

$$\begin{pmatrix} \nu_a \end{pmatrix} = \begin{pmatrix} U_{ai} \end{pmatrix} \begin{pmatrix} \nu_i \end{pmatrix}$$



Survival probability: $P(\nu_e \rightarrow \nu_e) = 1 - \sin^2(2\theta_{12}) \sin^2\left(\frac{\Delta m_{21}^2 L}{4E}\right)$ *distance/energy*

- mixing angle \rightarrow amplitude
- mass² difference \rightarrow frequency
- neutrino energy \rightarrow oscillation length

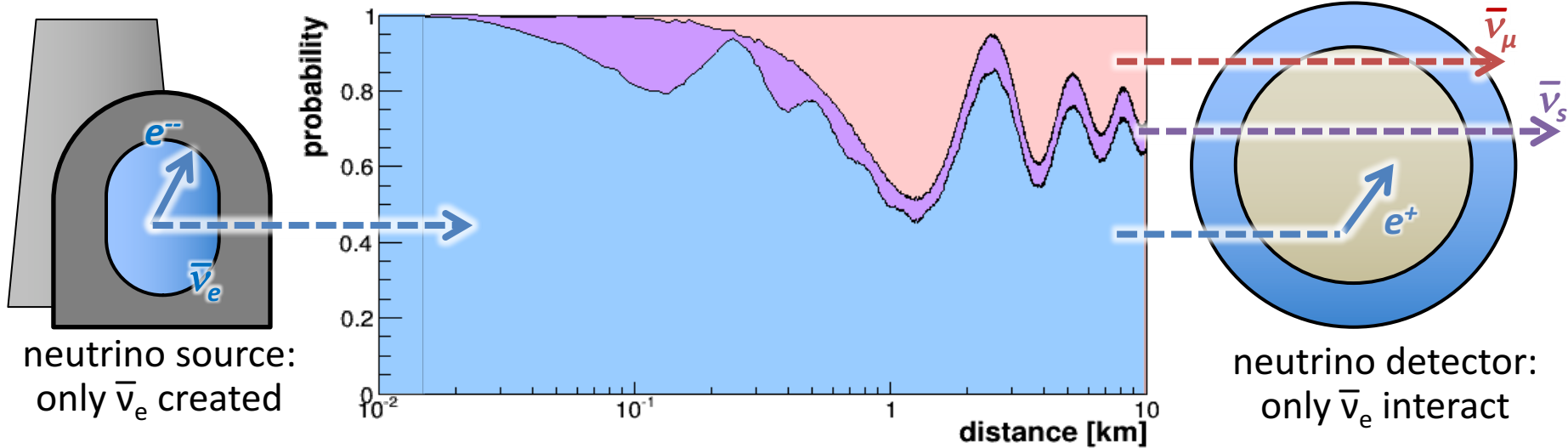
oscillation amplitude *oscillation frequency*

■ Active states only:

$$\begin{pmatrix} \nu_a \end{pmatrix} = \begin{pmatrix} U_{ai} \end{pmatrix} \begin{pmatrix} \nu_i \end{pmatrix}$$

■ Adding sterile states

$$\begin{pmatrix} \nu_a \\ \nu_s \end{pmatrix} = \begin{pmatrix} U_{ai} & \Theta_{aj} \\ \Theta_{si} & U_{sj} \end{pmatrix} \begin{pmatrix} \nu_i \\ \nu_j \end{pmatrix}$$



Survival probability: $P(\nu_e \rightarrow \nu_e) = 1 - \sin^2(2\theta_{12}) \sin^2\left(\frac{\Delta m_{21}^2 L}{4E}\right)$ *distance/energy*

- mixing angle → amplitude
- mass² difference → frequency
- neutrino energy → oscillation length

oscillation amplitude

oscillation frequency

New mass states & ordering schemes

- As a necessity, new neutrino flavor states imply **new neutrino mass states**, e.g. one further sterile state $\nu_s \rightarrow$ mass state ν_4
- Different mass ordering schemes might be realized:

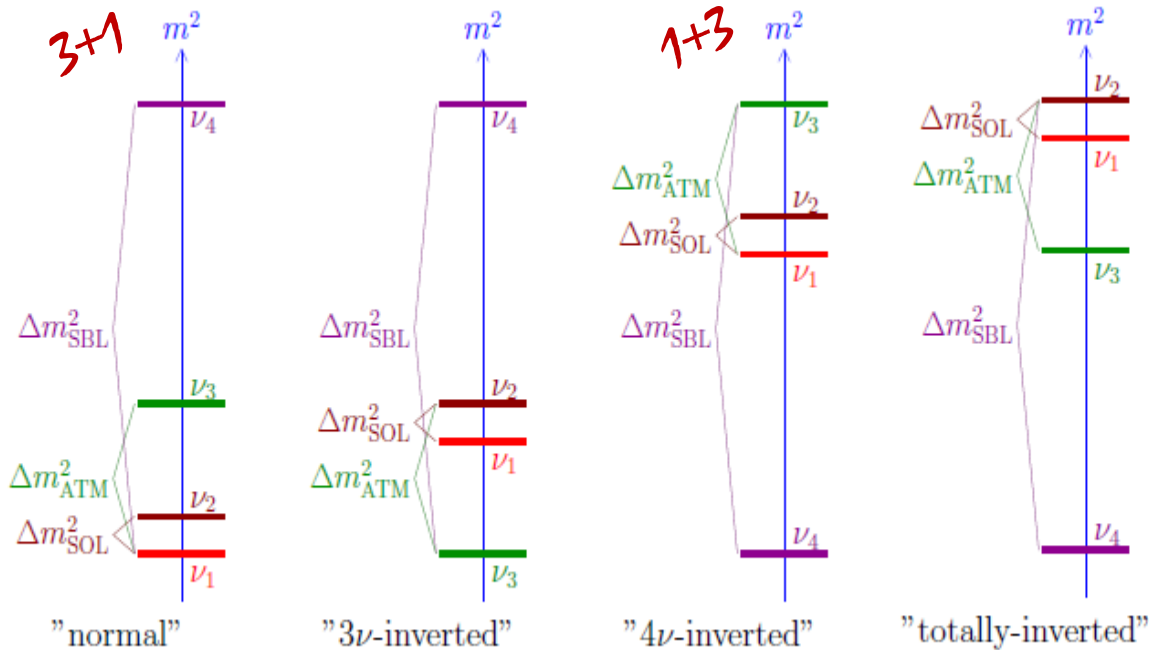
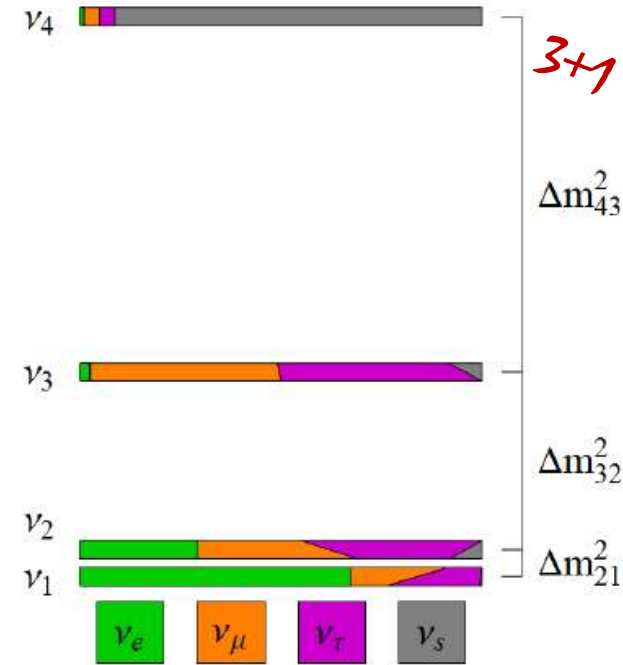


Figure 1: 3+1 four-neutrino schemes.



More complicated schemes possible:

- 3+2, 3+3 ...**
- 1+3+1 etc.**

- Active-sterile mixing matrix \rightarrow **new mixing amplitudes**

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_{s1} \\ \vdots \end{pmatrix} = \begin{pmatrix} U_{11} & U_{12} & U_{13} & U_{14} & \dots \\ U_{21} & U_{22} & U_{23} & U_{24} & \dots \\ U_{31} & U_{32} & U_{33} & U_{34} & \dots \\ U_{41} & U_{42} & U_{43} & U_{44} & \dots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \\ \vdots \end{pmatrix}$$

- new masses \rightarrow **new Δm^2 values:** $\Delta m_{41}^2, \dots > \Delta m_{21}^2, \Delta m_{31}^2, \Delta m_{32}^2$

\rightarrow occurrence of oscillation phenomena at **new (shorter) baselines**, e.g.

active \rightarrow sterile disappearance $P(\nu_e \rightarrow \nu_s) = \sin^2 2\theta_{ee} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right); \quad \sin^2 2\theta_{ee} = 4|U_{e4}|^2(1 - |U_{e4}|^2)$

active \rightarrow active appearance $P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta_{e\mu} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right); \quad \sin^2 2\theta_{e\mu} = 4|U_{e4}|^2|U_{\mu 4}|^2$

▶ (Long-standing) electron neutrino appearance anomalies

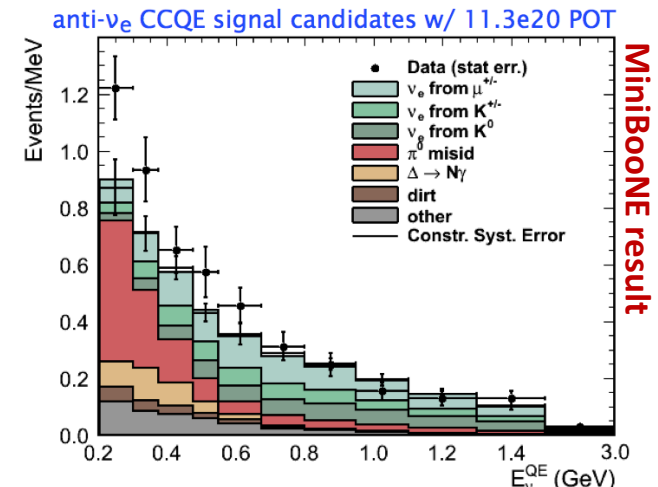
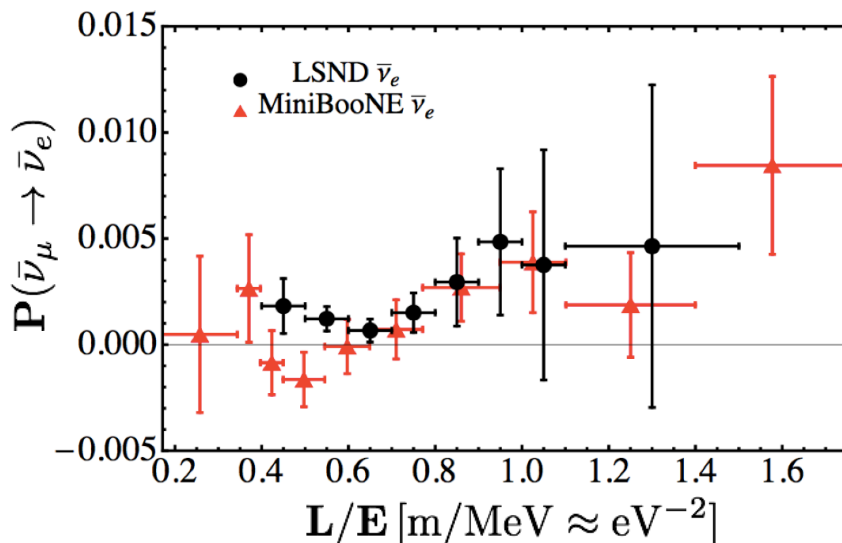
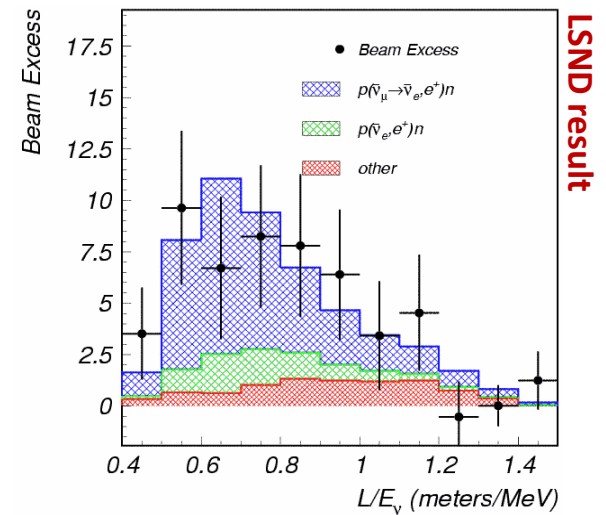
■ LSND result

ν_e appearance signal
in a low-energy $\bar{\nu}_\mu$ beam
from stopped pions

■ MiniBooNE result

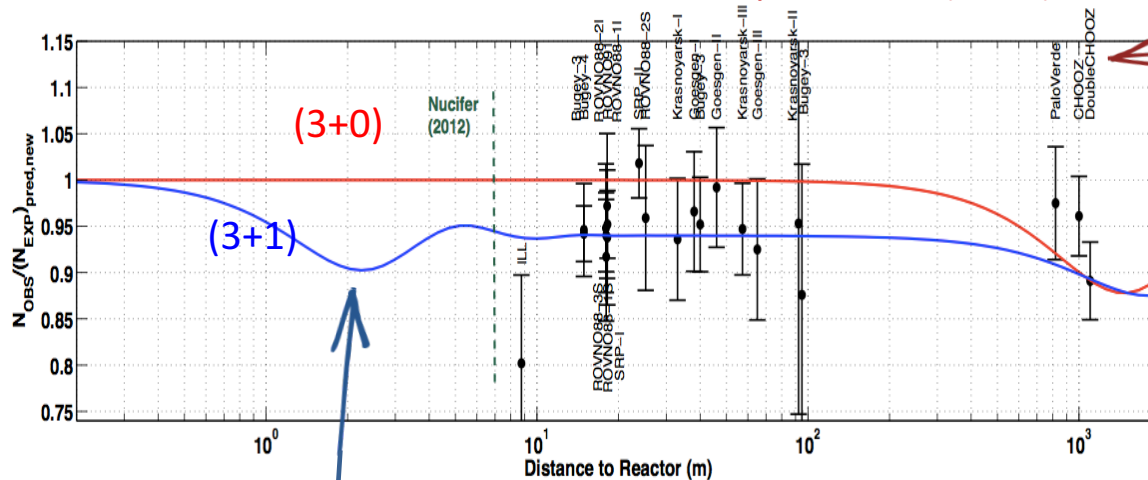
ν_e appearance signal
in a GeV $\nu_\mu/\bar{\nu}_\mu$ beam
at similar L/E ratio

→ interpretation as $\nu_\mu \rightarrow \nu_e$ appearance oscillations
via a **new Δm^2 on eV^2 scale**



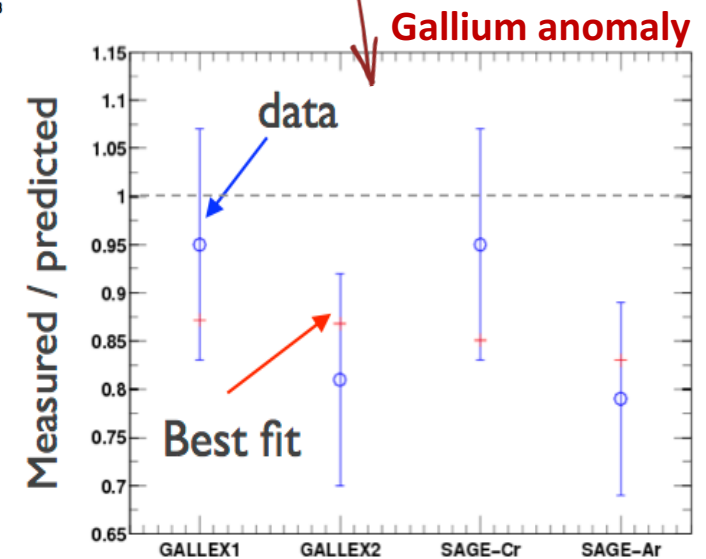
► (More recent) electron neutrino disappearance anomalies

Reactor short-baseline data vs. new rate prediction (2011)



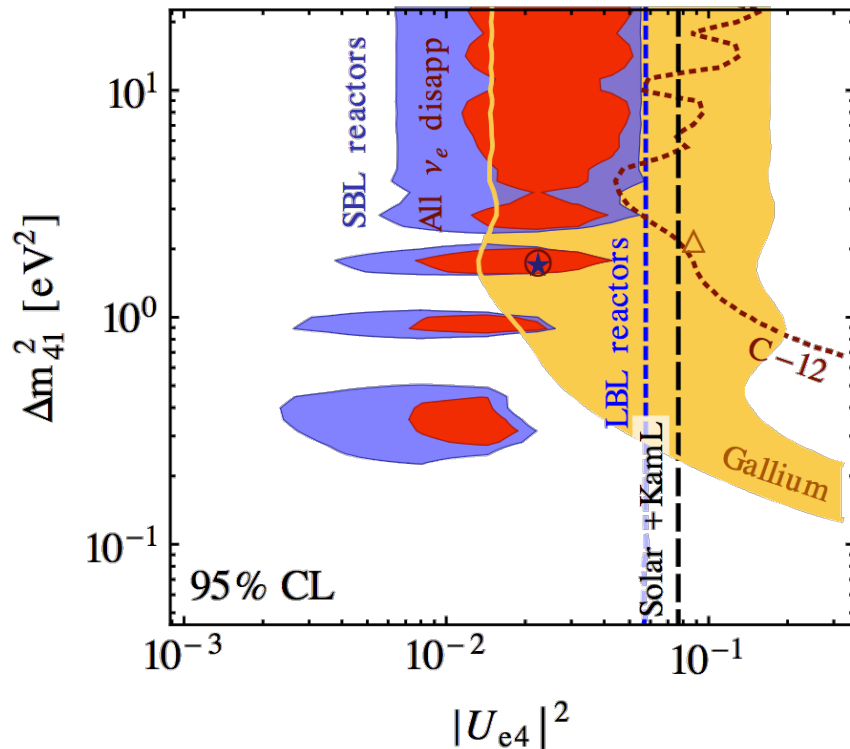
- reactor experiments (7.3±2.3)% deficit in $\bar{\nu}_e$ -rate at short distances (<100m)
- Gallium calibration data (14±5)% rate deficit close to a radioactive ν_e -source

- possible interpretation as **short-baseline disappearance $\nu_e \rightarrow \nu_s$**
- MeV energies, oscillation length $L_{osc} \leq 10m$
„sterile“ $\Delta m^2_{new} \geq 1eV^2$
- Required oscillation amplitude:
 $\sin^2 2\theta_{new} \geq 0.1$



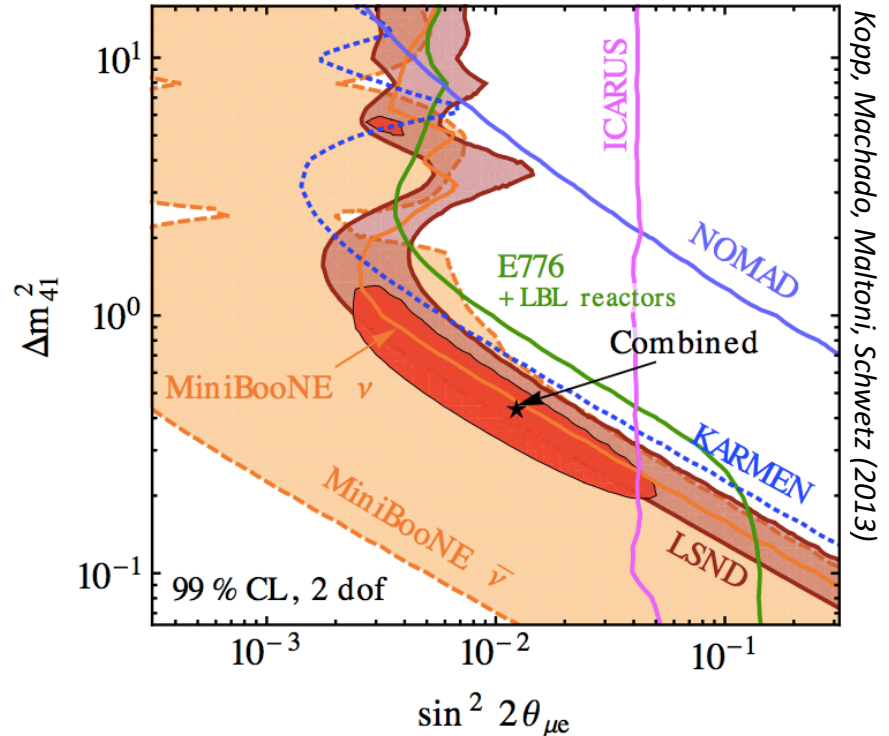
- All anomalies can be described by a **(3+1) scheme** adding a single eV-mass sterile neutrino

$\nu_e \rightarrow \nu_s$ disappearance



- reactor antineutrino anomaly
- gallium anomaly

$\nu_\mu \rightarrow \nu_e$ appearance



- LSND
- MiniBooNE

Kopp, Machado, Maltoni, Schwetz (2013)

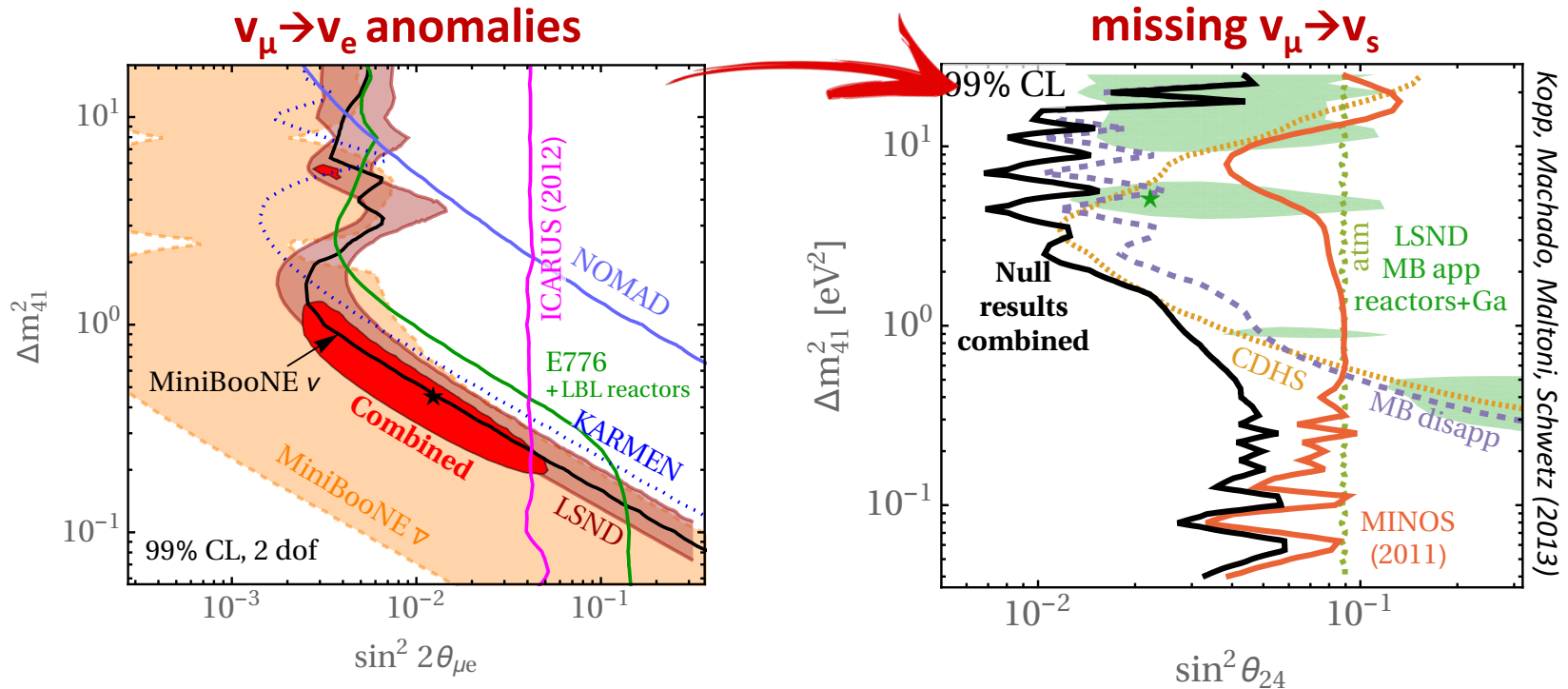
Results on $\nu_\mu \rightarrow \nu_s$ disappearance ^{1/2}

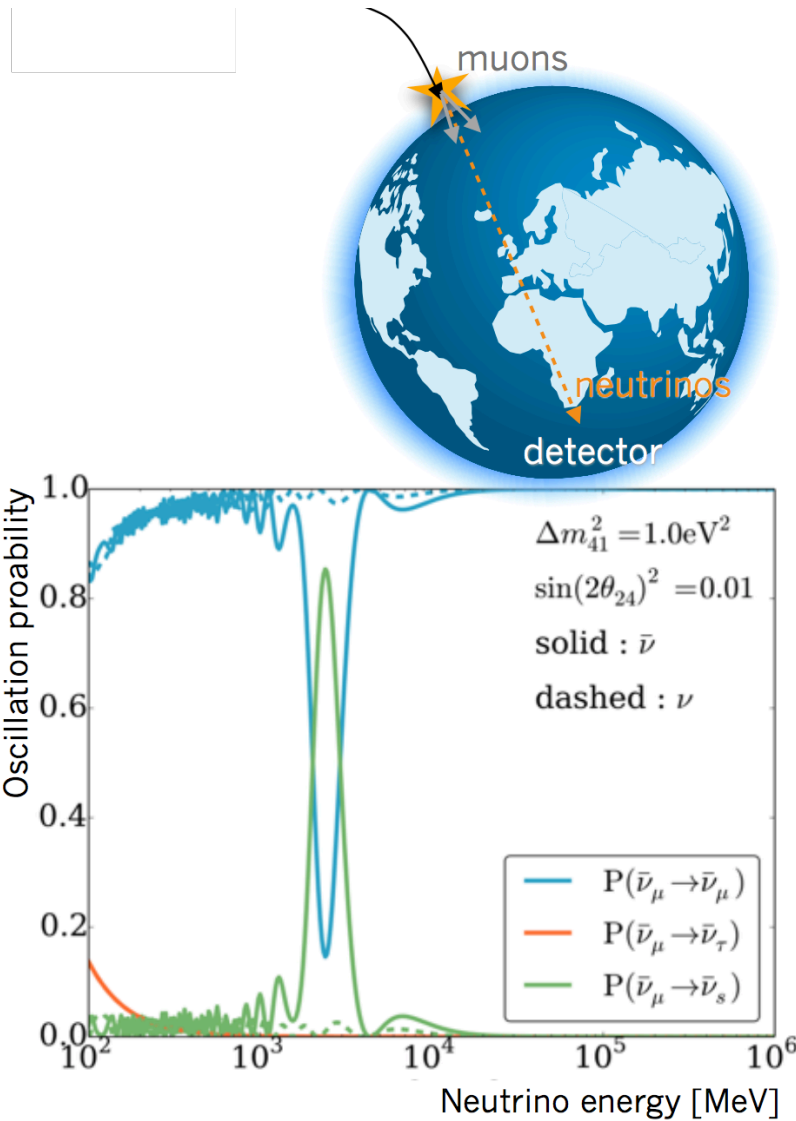
▪ **Note:**

Disappearance and **appearance amplitudes** are interlinked

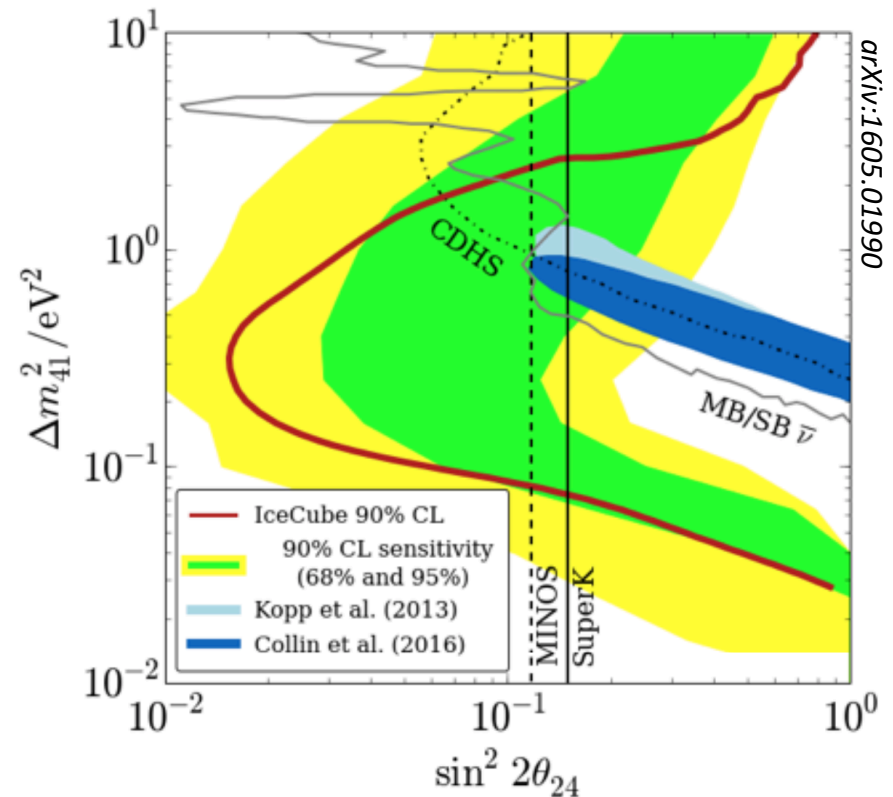
$$\begin{aligned} \sin^2 \theta_{ee} &= |U_{e4}|^2 \\ \sin^2 \theta_{\mu\mu} &= |U_{\mu4}|^2 \end{aligned} \longleftrightarrow \sin^2 \theta_{\mu e} = |U_{e4}| \cdot |U_{\mu4}|$$

▪ In 2011, already some tension between $\nu_\mu \rightarrow \nu_e$ and $\nu_\mu \rightarrow \nu_s$ results:





- Probe: **Atmospheric ν 's** crossing the Earth
- **matter potential** affects only active ν 's
- **No resonant conversion of $\nu_\mu \rightarrow \nu_s$ found at TeV energies, i.e. $\Delta m_{41}^2 \sim 1 \text{ eV}^2$**



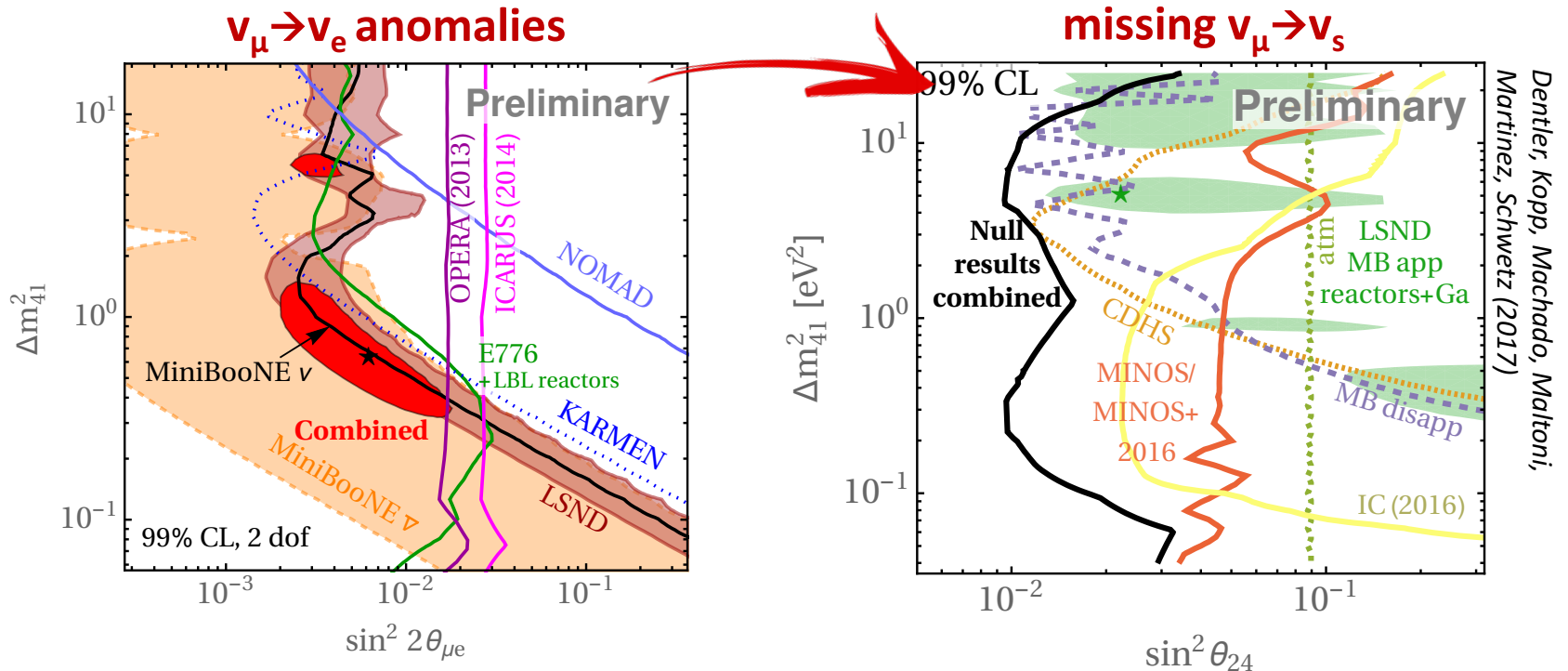
Results on $\nu_\mu \rightarrow \nu_s$ disappearance ^{2/2}

▪ **Note:**

Disappearance and **appearance amplitudes** are interlinked

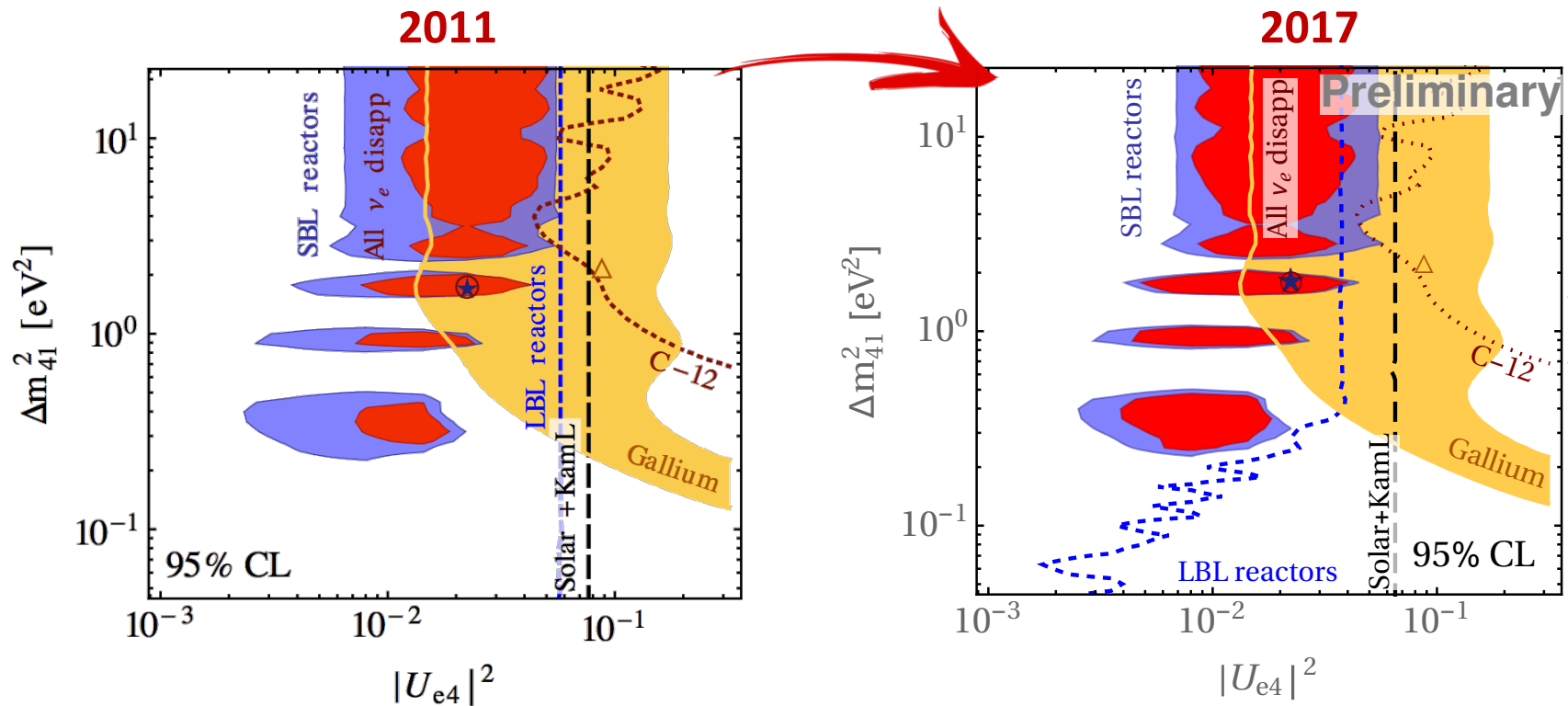
$$\begin{aligned} \sin^2 \theta_{ee} &= |U_{e4}|^2 \\ \sin^2 \theta_{\mu\mu} &= |U_{\mu4}|^2 \end{aligned} \longleftrightarrow \sin^2 \theta_{\mu e} = |U_{e4}| \cdot |U_{\mu4}|$$

▪ Now, new results by **MINOS+/IceCube** on $\nu_\mu \rightarrow \nu_s$ further increased the tension:



The open issue: $\nu_e \rightarrow \nu_s$ disappearance

- **Note:** No oscillation data directly contradicts the reactor/Ga anomalies!

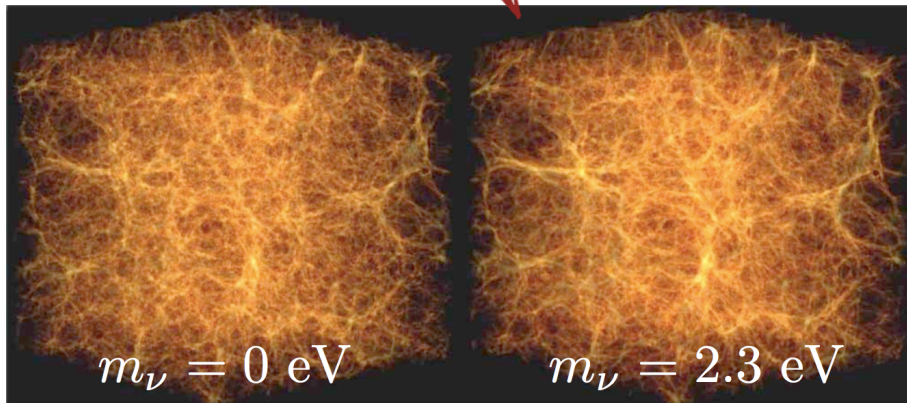
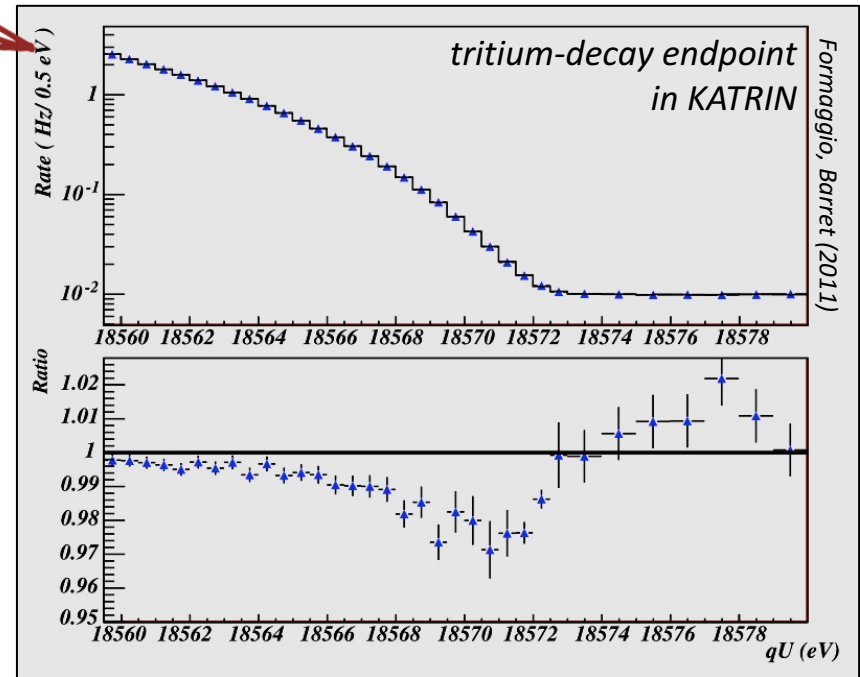
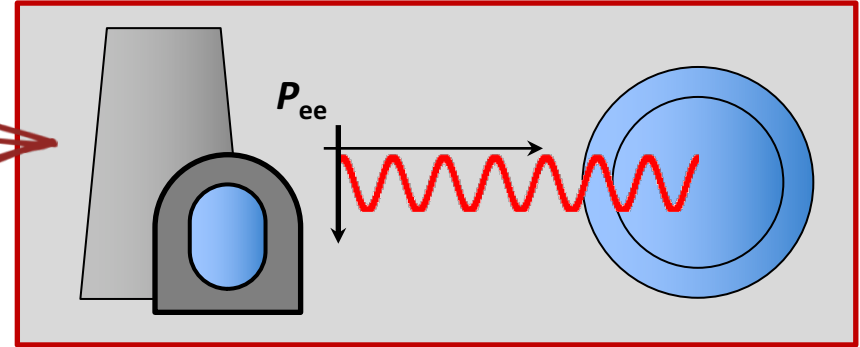


Dentler, Kopp, Machado, Maltoni,
Martinez, Schwetz (2017)

→ need for dedicated experiments

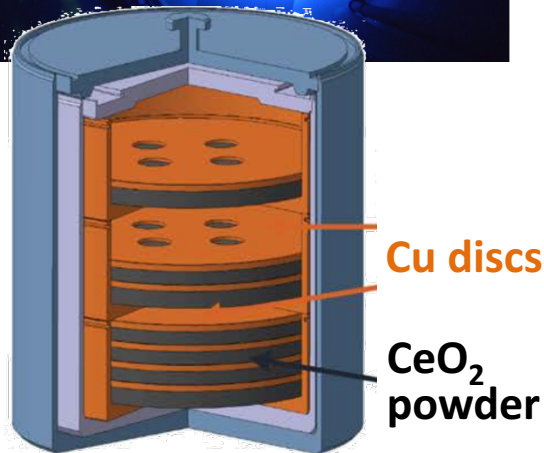
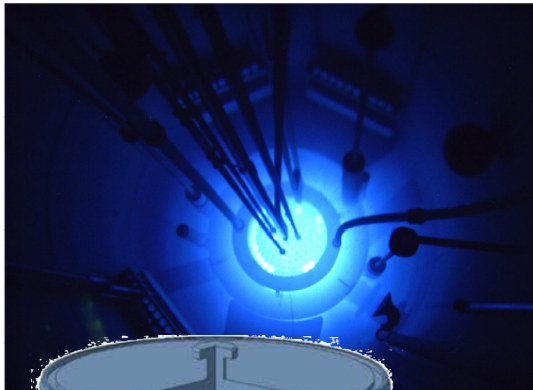
New experimental approaches

- very-short baseline experiments for observing $\nu_e \rightarrow \nu_s$ oscillation disappearance pattern
- β -decay ν mass experiments to find spectral deformation from eV-mass eigenstate ν_4
- cosmological limits on N_{eff} & Σm_ν (CMB, BBN, BAO ...)



Basic approach:

- search for $\nu_e \rightarrow \nu_s$ **disappearance** oscillations
- Intrinsicly **pure beam**: only ν_e or $\bar{\nu}_e$
- **oscillometry**: oscillation waves inside the detector
- energy range: 1-10 MeV \rightarrow **distance 1-10 m**
- **well-known cross-sections** at MeV energies



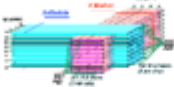



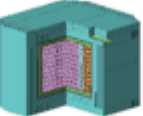
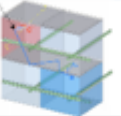

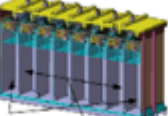
Reactor experiments

- intense, stable source of antineutrinos
- extended reactor core \rightarrow research reactor
- large intrinsic background levels

Radioactive source

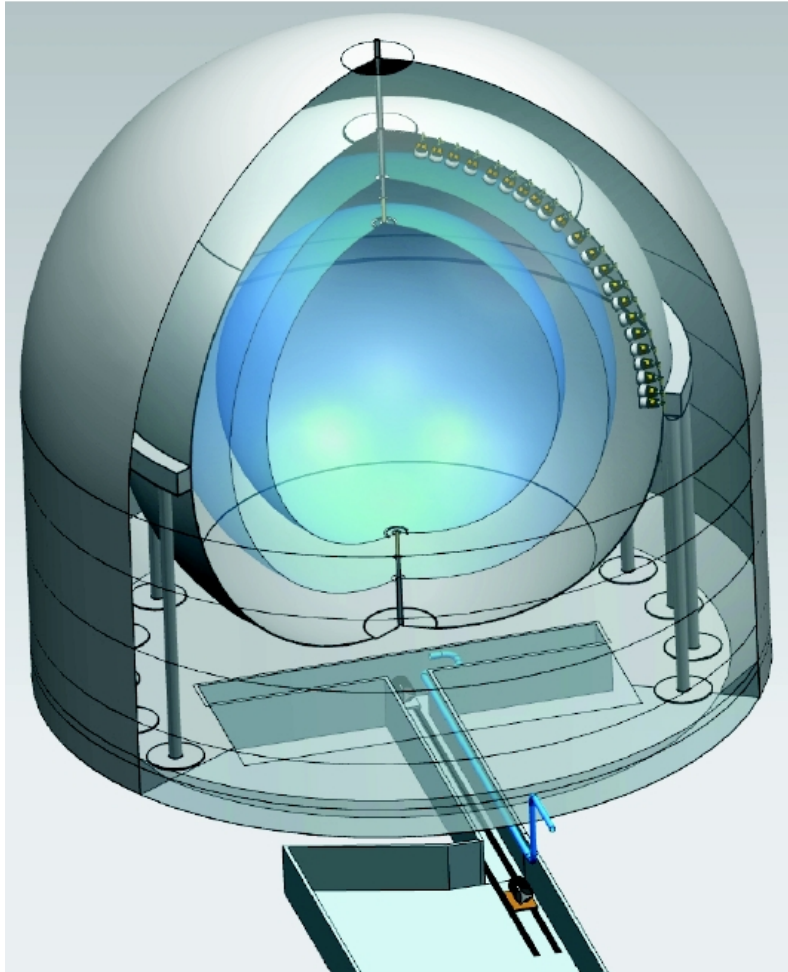
- low-background levels (nearly background-free)
- well-defined and well-localized source activity
- decaying source \rightarrow limited measuring time
- bureaucratic challenge

Short-baseline reactor experiments

| Experiment | Reactor Power/Fuel | Overburden (mwe) | Detection Material | Segmentation | Optical Readout | Particle ID Capability |
|---|------------------------------|------------------|---------------------------------------|---------------------------------|-------------------------|--------------------------------|
| DANSS (Russia)  | 3000 MW LEU fuel | ~50 | Inhomogeneous PS & Gd sheets | 2D, ~5mm | WLS fibers. | Topology only |
| NEOS (South Korea)  | 2800 MW LEU fuel | ~20 | Homogeneous Gd-doped LS | none | Direct double ended PMT | recoil PSD only |
| nuLat (USA)  | 40 MW ²³⁵ U fuel | few | Homogeneous ⁶ Li doped PS | Quasi-3D, 5cm, 3-axis Opt. Latt | Direct PMT | Topology, recoil & capture PSD |
| Neutrino4 (Russia)  | 100 MW ²³⁵ U fuel | ~10 | Homogeneous Gd-doped LS | 2D, ~10cm | Direct single ended PMT | Topology only |
| PROSPECT (USA)  | 85 MW ²³⁵ U fuel | few | Homogeneous ⁶ Li-doped LS | 2D, 15cm | Direct double ended PMT | Topology, recoil & capture PSD |
| SoLid (UK Fr Bel US)  | 72 MW ²³⁵ U fuel | ~10 | Inhomogeneous ⁶ LiZnS & PS | Quasi-3D, 5cm multiplex | WLS fibers | topology, capture PSD |
| Chandler (USA)  | 72 MW ²³⁵ U fuel | ~10 | Inhomogeneous ⁶ LiZnS & PS | Quasi-3D, 5cm, 2-axis Opt. Latt | Direct PMT/ WLS Scint. | topology, capture PSD |
| Stereo (France)  | 57 MW ²³⁵ U fuel | ~15 | Homogeneous Gd-doped LS | 1D, 25cm | Direct single ended PMT | recoil PSD |

from N. Bowden's talk at Nu16

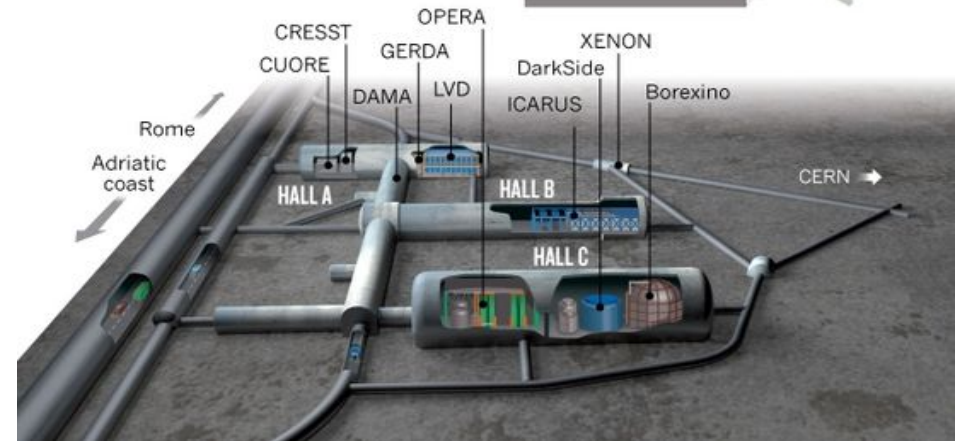
Schematic of Borexino



Start: May 2007

THE A, B AND C OF GRAN SASSO

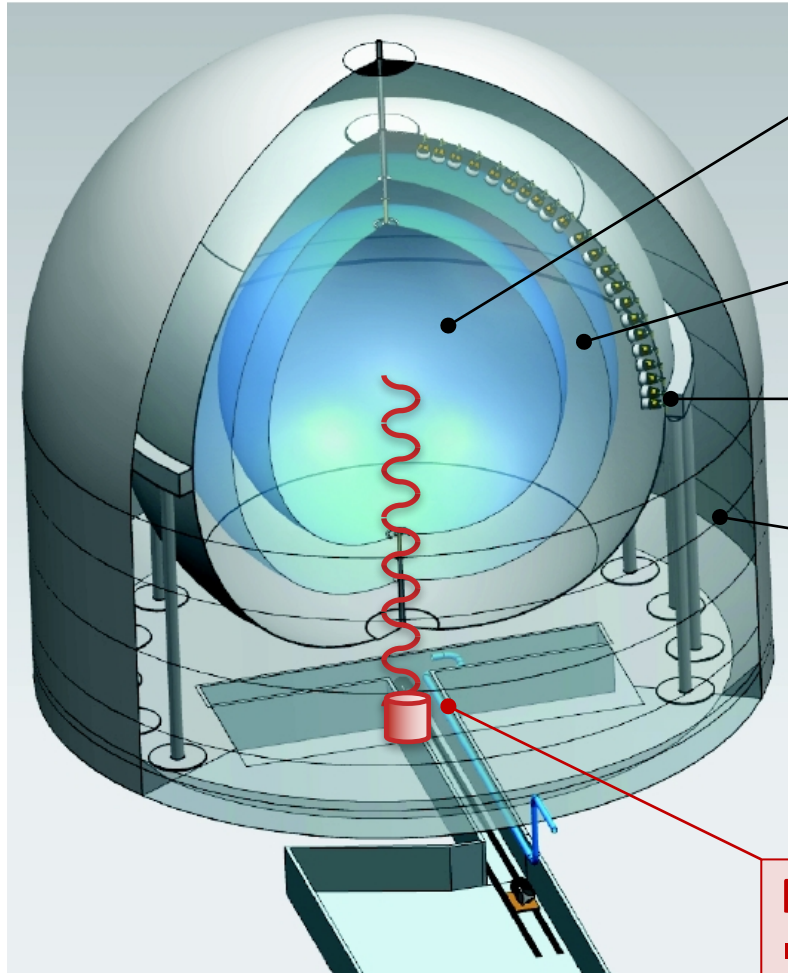
Experiments at the Gran Sasso National Laboratory are housed in and around three huge halls carved deep inside the mountain, where they are shielded from cosmic rays by 1,400 metres of rock.



Borexino @ Gran Sasso Laboratories

- low-energy solar neutrino experiment
- organic liquid-scintillator detector
- since 2007: ${}^7\text{Be}$, pep, pp, geo-neutrinos
- ultra-low background conditions:
 - rock shielding: 1.4 km
 - intrinsic radiopurity: 10^{-18} g/g U/Th

Schematic of Borexino



neutrino target

- diameter: 8.5 m
- pseudocumene+PPO: 300 tons

buffer volume

- shielding from external radioactivity

steel sphere with 2212 PMTs

- diameter: 13.7 m

water Cherenkov muon veto

Pit for neutrino source

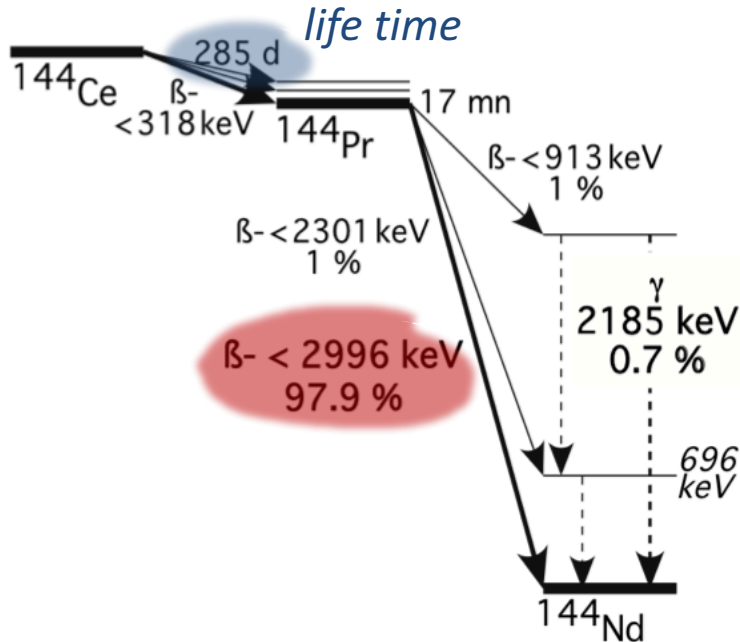
- 8.5m from center

Start: Spring 2018 – duration: 1.5 yrs

SOX Pit below Borexino



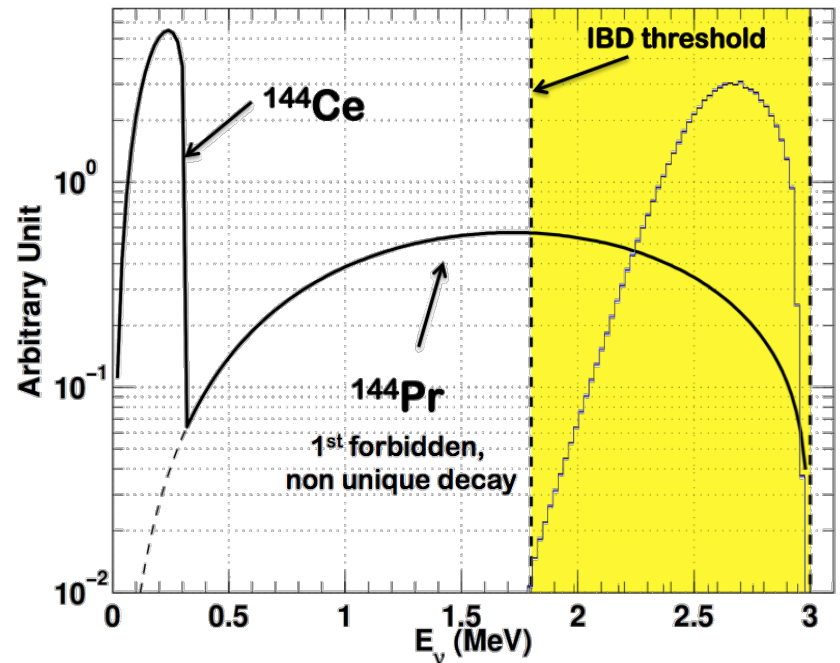
^{144}Ce - ^{144}Pr decay scheme



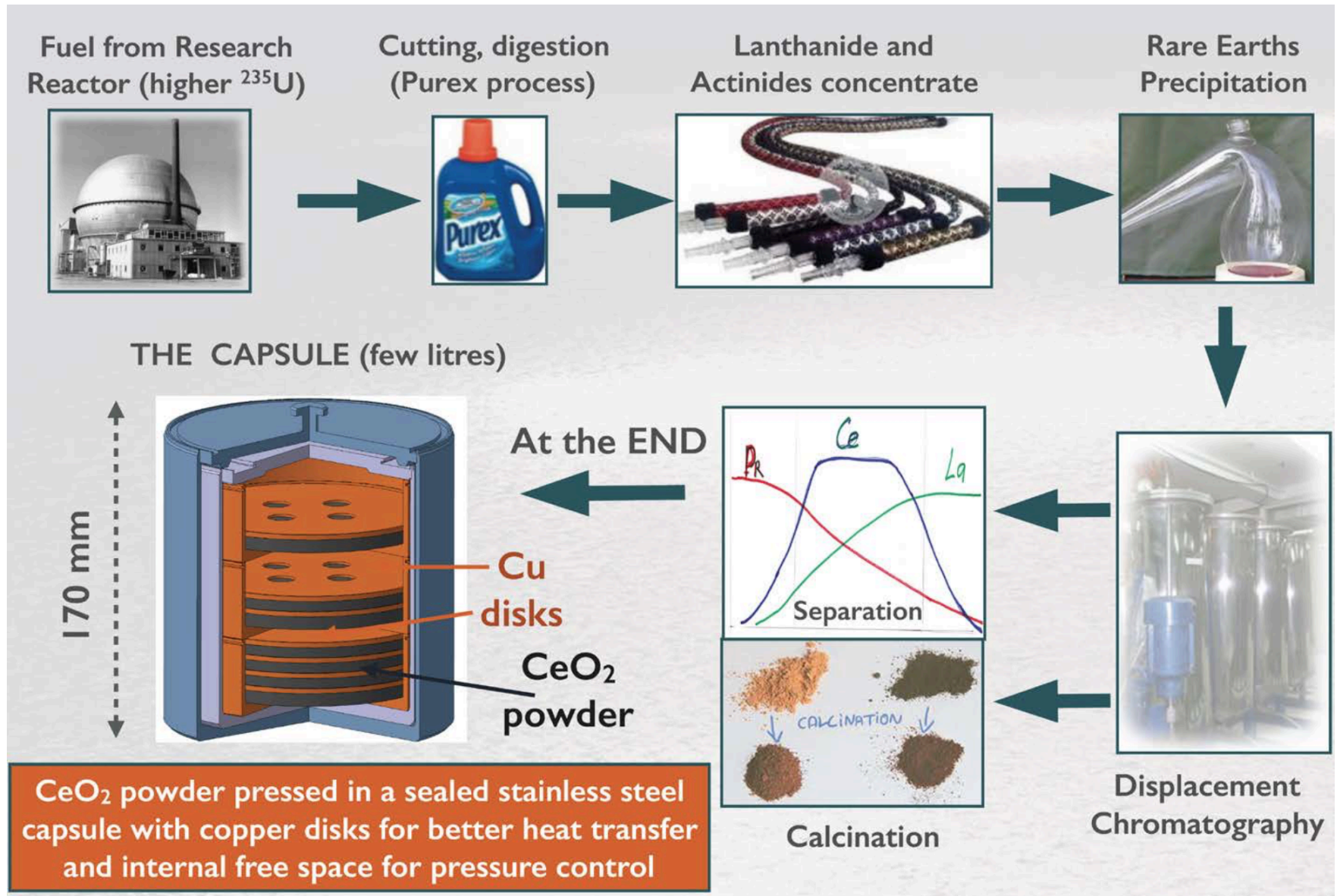
Initial activity: 100-150 kCi
 4-6 PBq
 heat power: 0.9 – 1.3 kW

Inverse Beta Decay (IBD) cross section:
 $\sigma_{\text{IBD}} \approx 9.5 \cdot 10^{-45} \text{ cm}^2 (E - 1.8 \text{ MeV})^2$

β -spectrum & cross-section



^{144}Ce source production

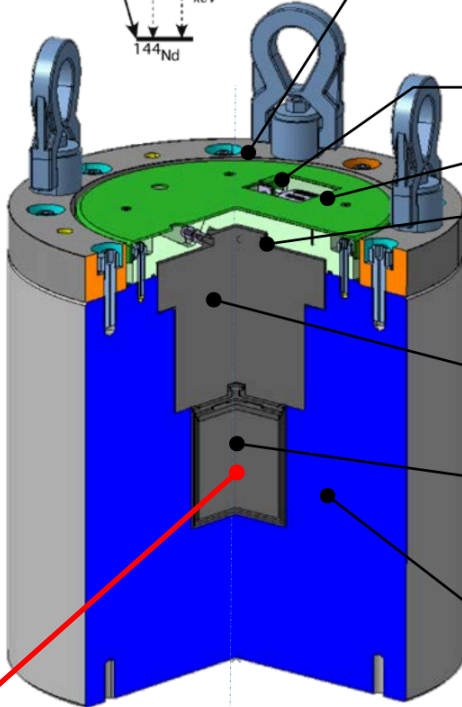


Source transport

3 weeks of transport



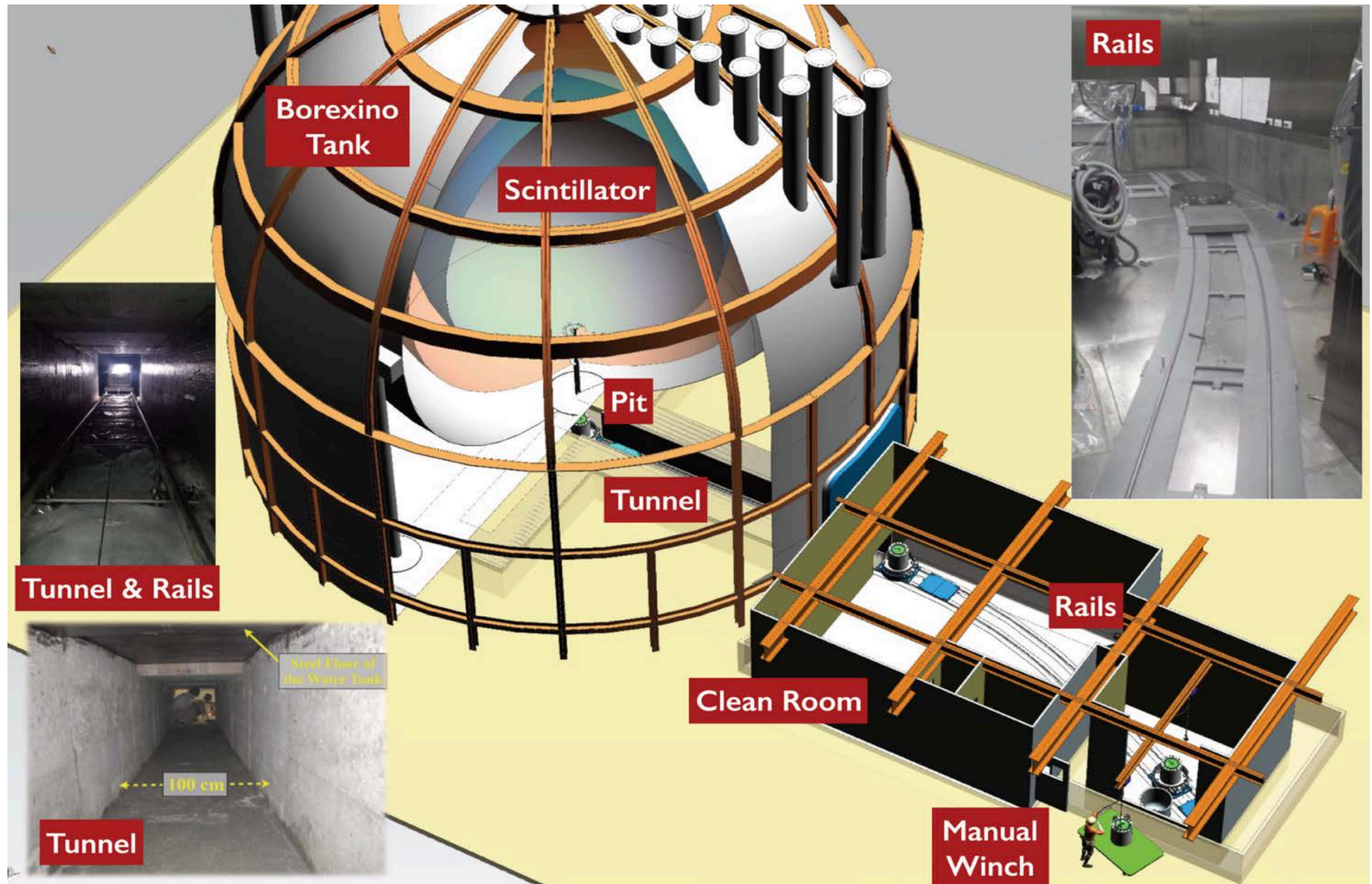
Source shielding: reduces 4PBq to 200Bq surface activity



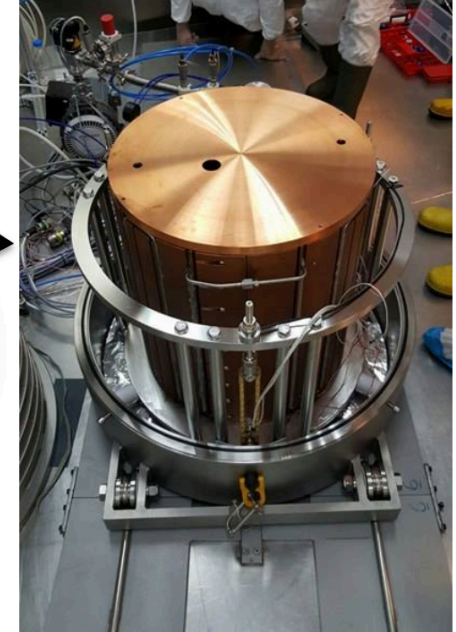
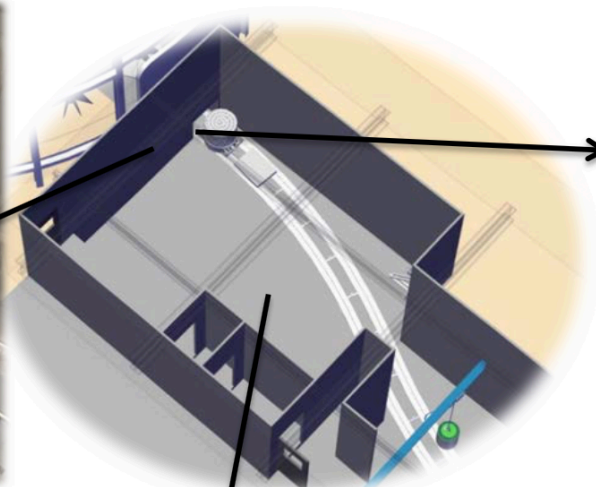
internal temperature: ~500°C



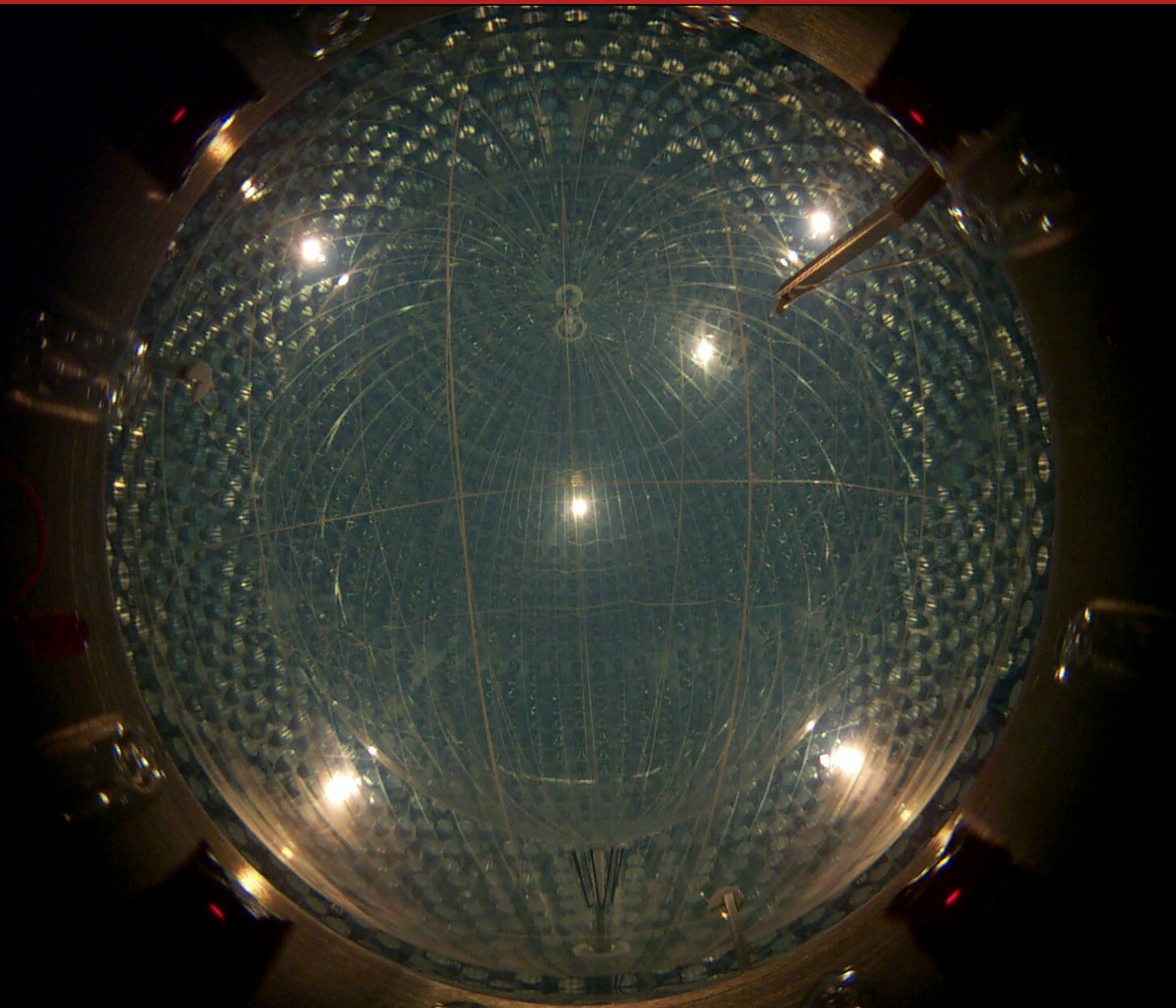
SOX experimental layout



SOX source insertion system



Antineutrino detection in Borexino



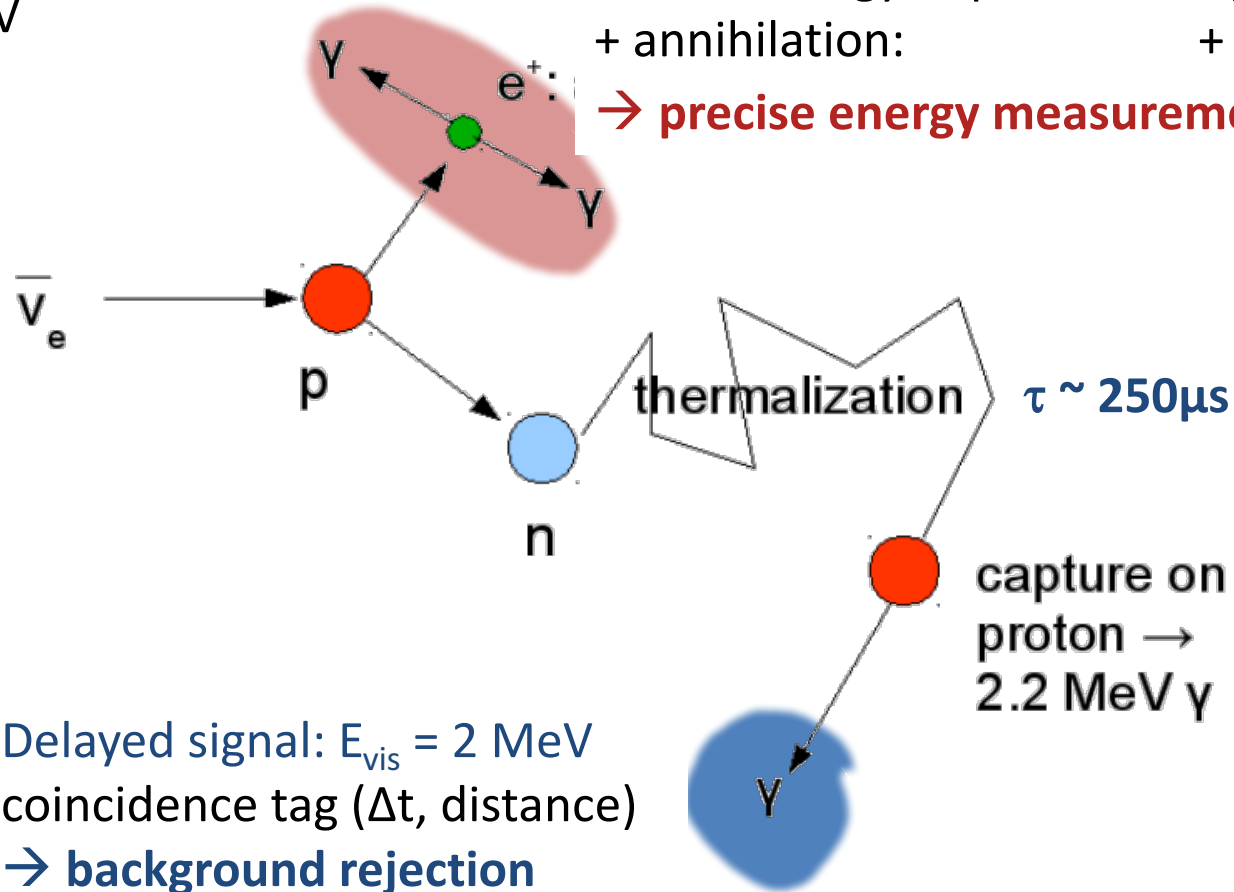
Threshold:

$$Q = m(n) + m(e^+) - m(p) = 1.8 \text{ MeV}$$

Prompt signal: $E_{\text{vis}} = 1 - 3 \text{ MeV}$

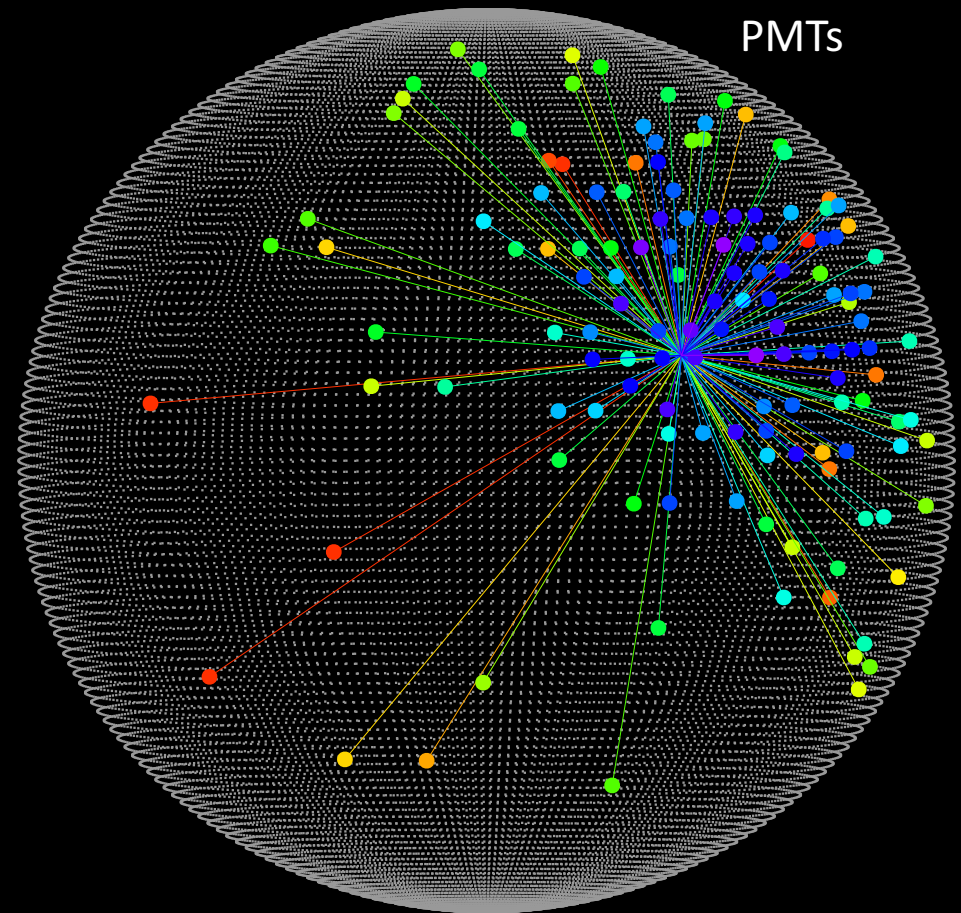
Kinetic energy of positron: $E(v) - Q$
+ annihilation: $+ 2m(e^\pm)$

→ precise energy measurement

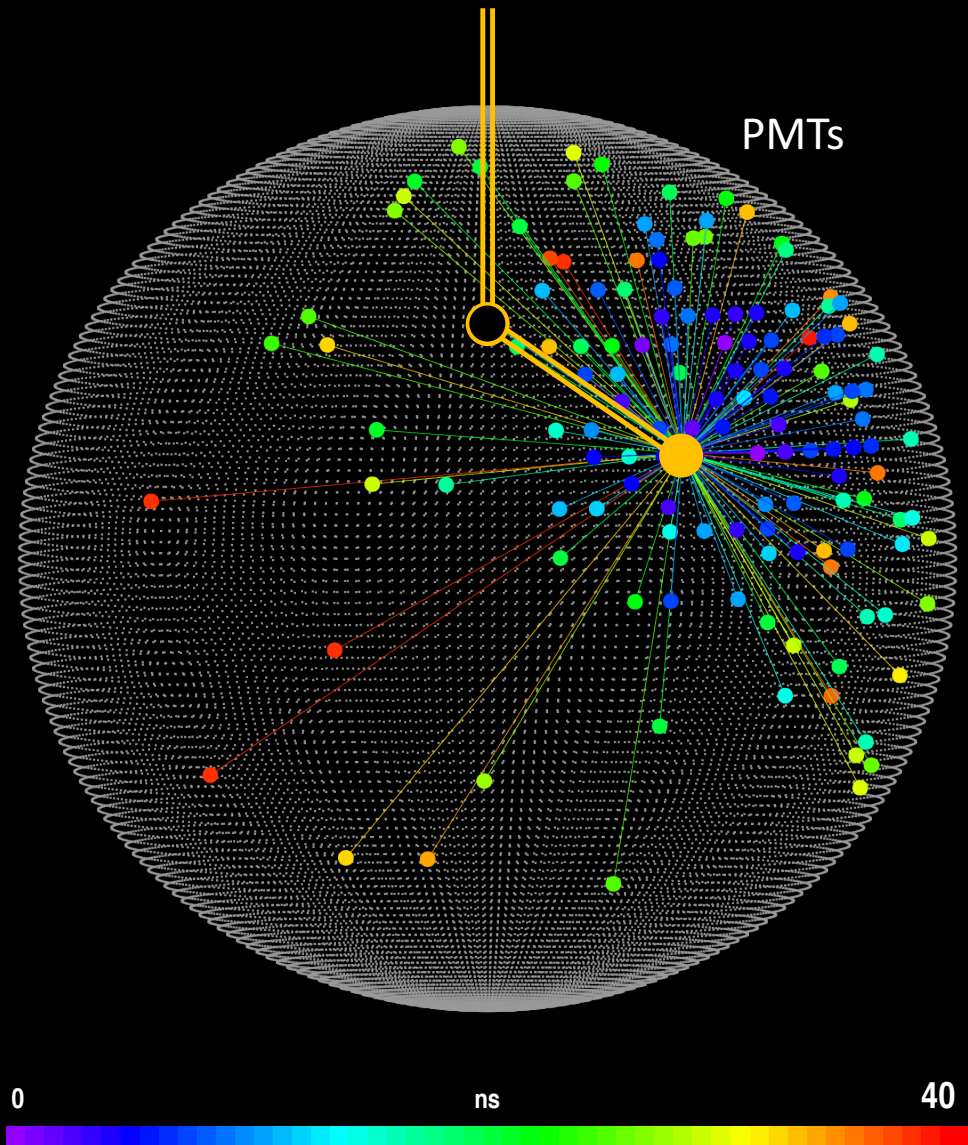


Delayed signal: $E_{\text{vis}} = 2 \text{ MeV}$
coincidence tag (Δt , distance)
→ background rejection

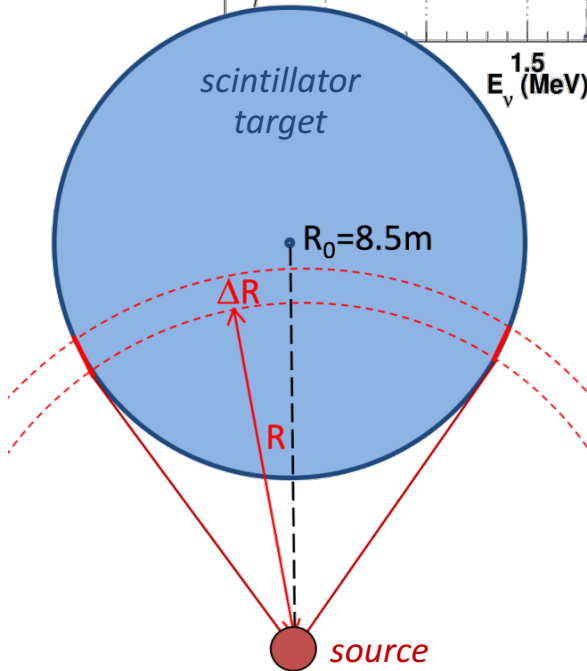
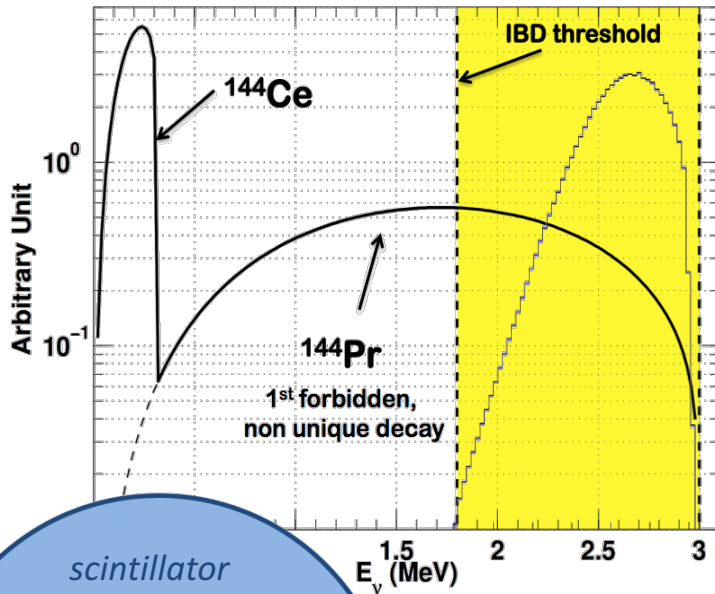
- **Scintillator light yield:**
~10k photons per MeV
→ 5% detected by PMTs
- **Energy resolution**
~500 p.e. per MeV
→ $\Delta E/E \sim 5\% @ 1 \text{ MeV}$
- **Energy threshold**
instrumental: ~50 keV
solar analysis: ~150 keV
- **Spatial reconstruction**
from photon time-of-flight
→ $\Delta x \sim 10 \text{ cm} @ 1 \text{ MeV}$



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- **Calibration campaign with sources inside IV planned for autumn.**

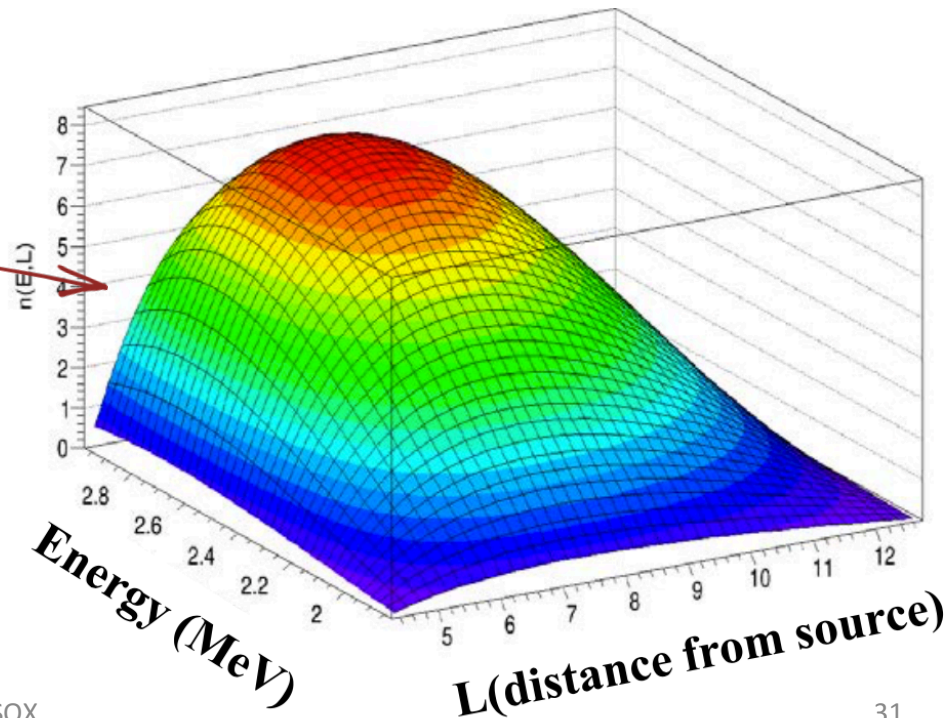


Expected antineutrino signal

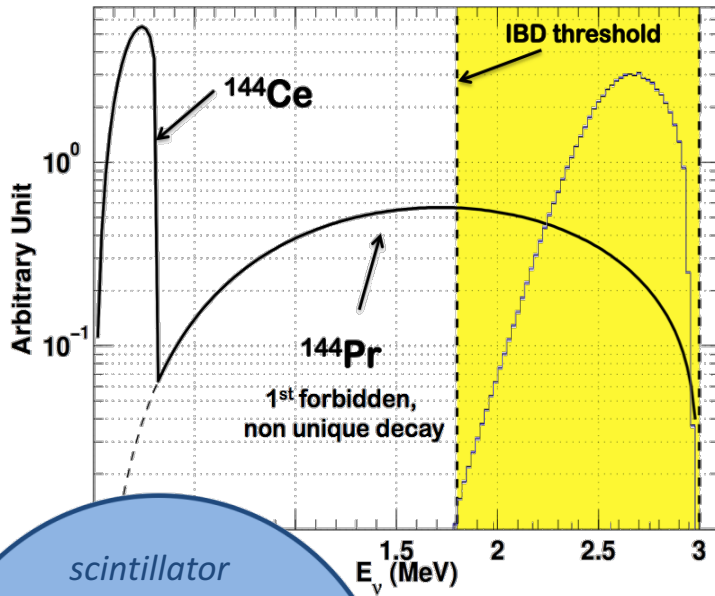


antineutrino spectrum
+
experimental geometry

L-E distribution
without oscillations



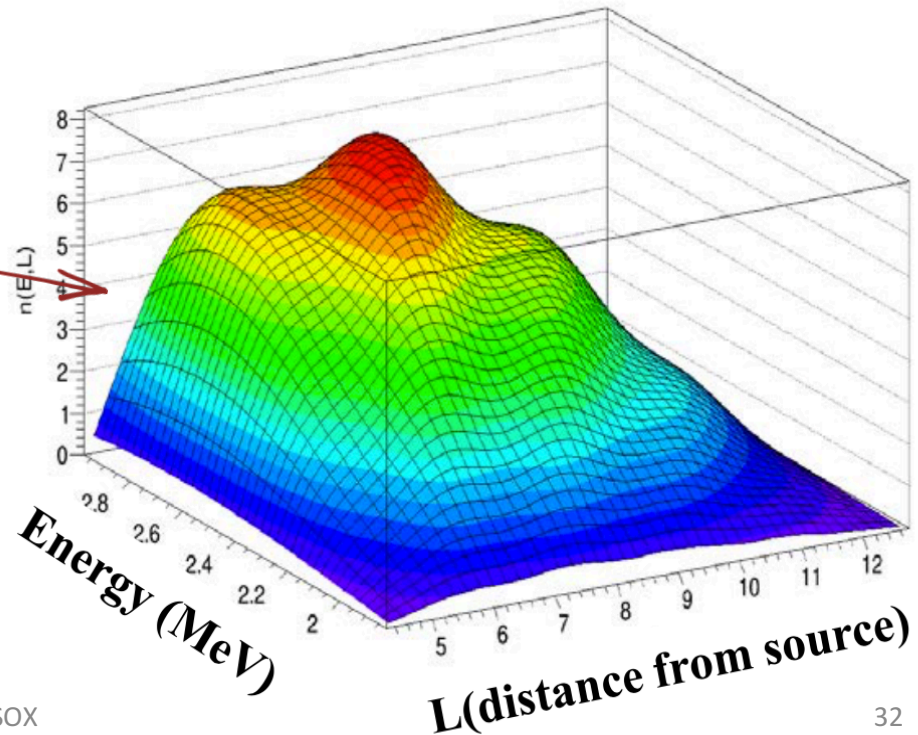
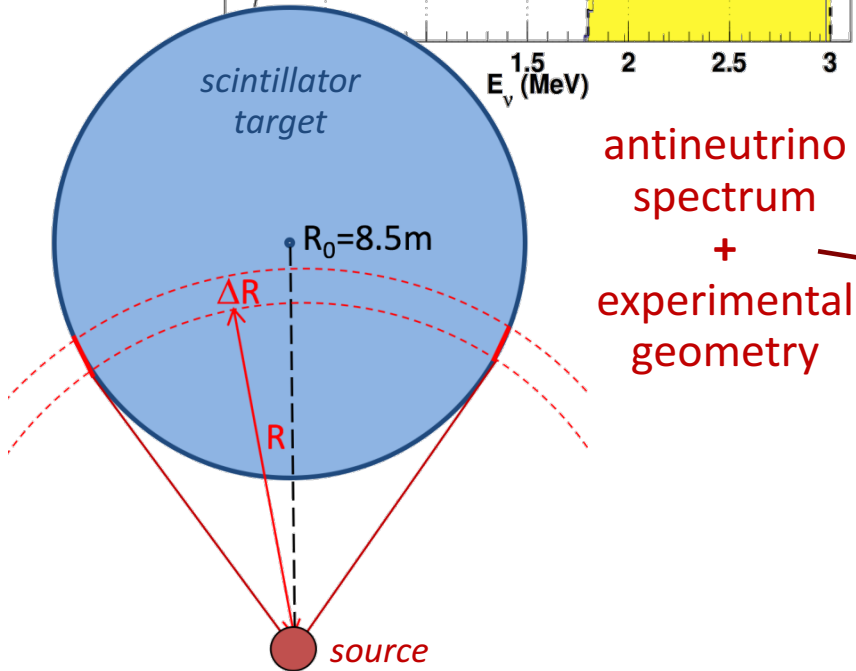
Expected antineutrino signal



- signature of disappearance oscillations**
- rate deficit (source activity)
 - oscillation waves (statistics: 10^4 IBDs)

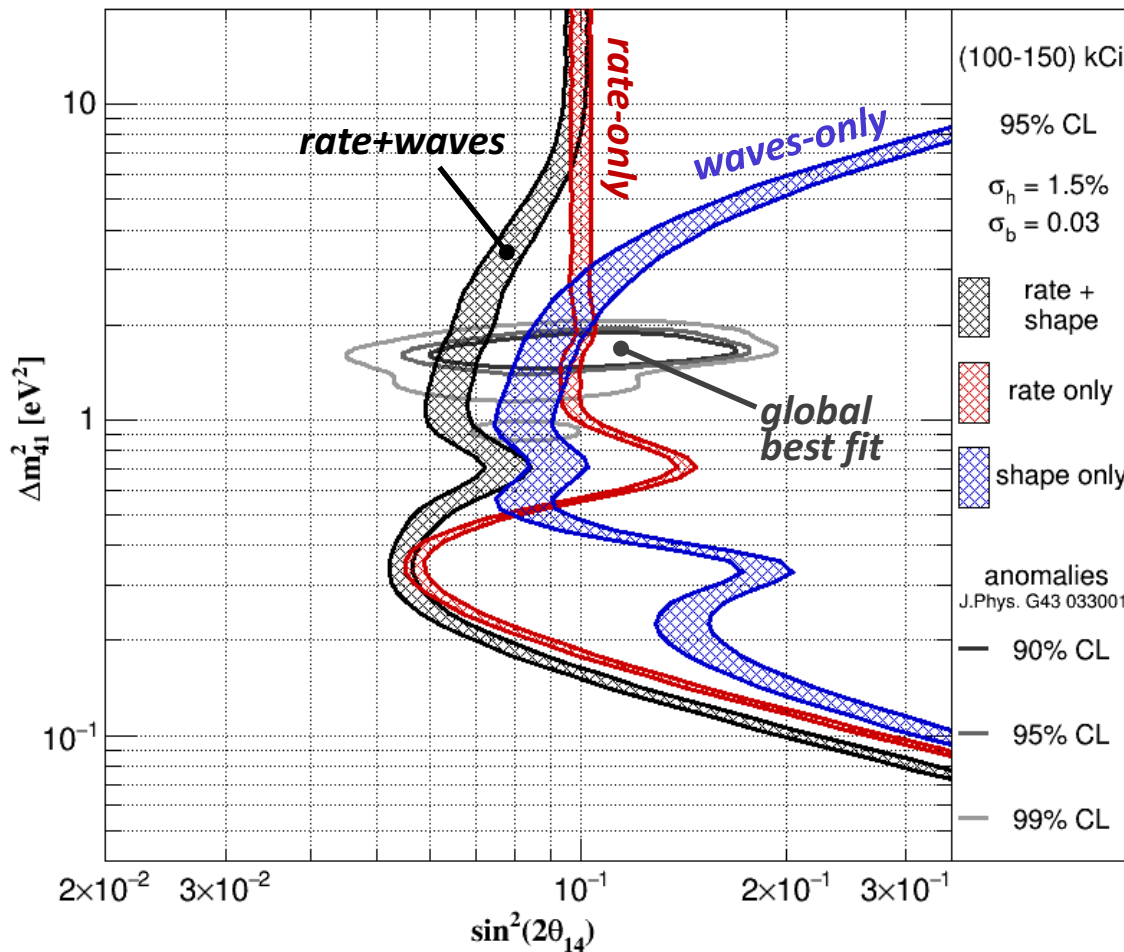
L-E distribution

with oscillations: $\sin^2 2\theta = 0.14$, $\Delta m^2 = 2.5 \text{ eV}^2$



Expected sensitivity vs. rate

$$P(\nu_e \rightarrow \nu_s) = 1 - \sin^2(2\theta_{14}) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$



Experimental parameters

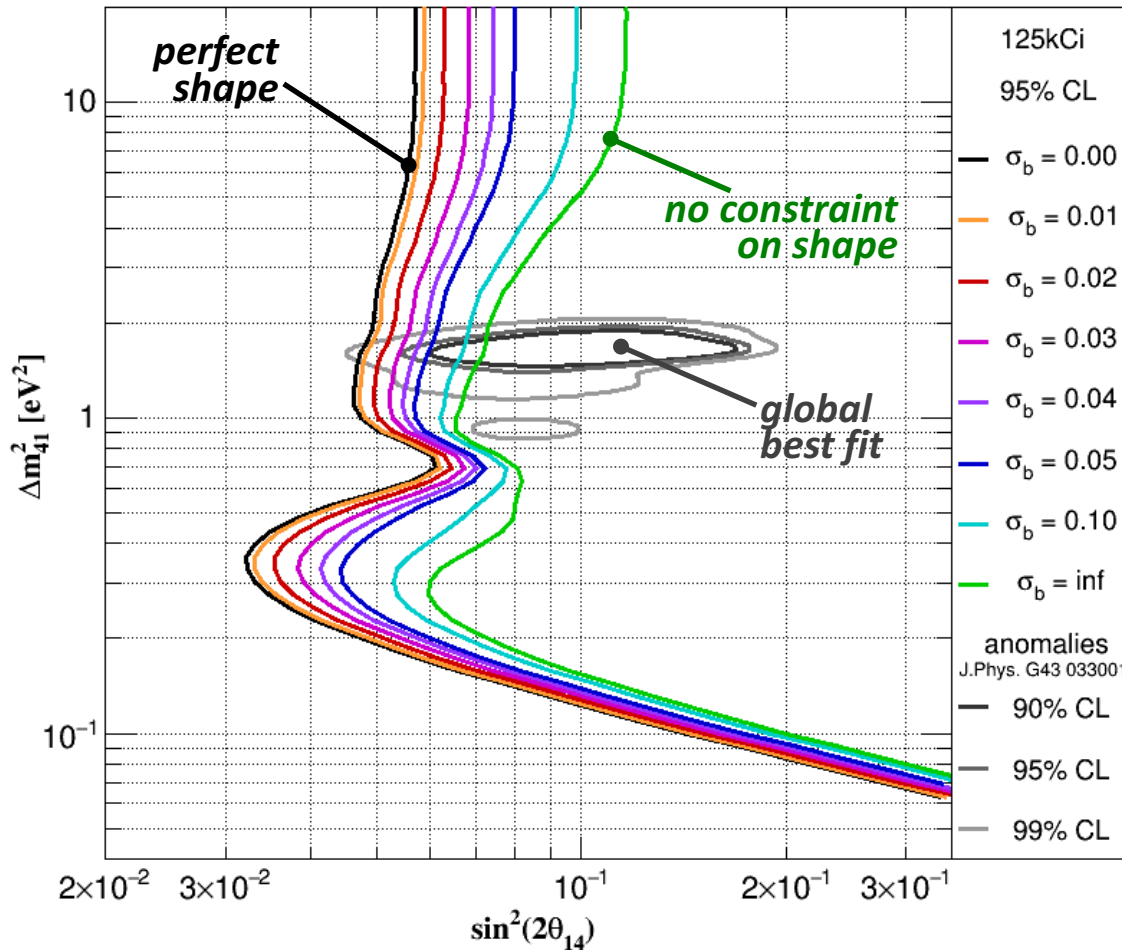
- **activity:** 100→150 kCi
- exposure: 1.5 yrs
- fiducial radius: 4 m
- uncertainties
 - on activity: 1%
 - on fiducial volume: 1%
 - on spectral shape b : 3%
- no background

→ maximum sensitivity for oscillation waves in region of the anomalies

Expected sensitivity vs. spectral shape

^{144}Pr β -spectrum with shape correction factor:

$$\frac{dN}{dE} = \sum_i \text{BR}_i \cdot S_{\beta i}(E) \cdot C(E) \text{ with } C(E) = 1 + b \cdot m_e/E$$



Experimental parameters

- activity: 125 kCi
- exposure: 1.5 yrs
- fiducial radius: 4 m
- uncertainties
 - on activity: 1%
 - on fiducial volume: 1%
 - on spectral shape b : 0% \rightarrow ∞
- no background

\rightarrow maximum sensitivity for oscillation waves in region of the anomalies

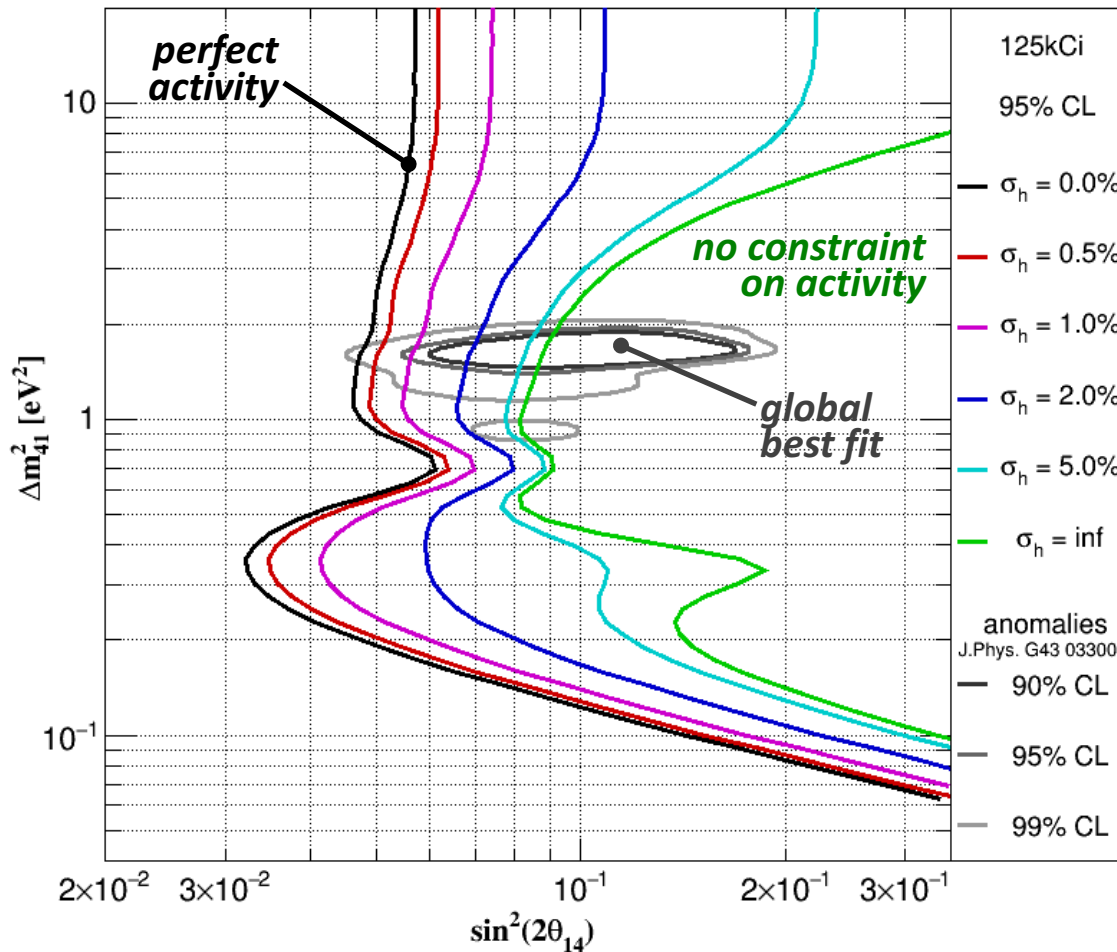
\rightarrow shape uncertainty matters but is under control

Expected sensitivity vs. spectral shape

Activity closely linked to heat power released by β -decays



expected event rate



Experimental parameters

- activity: 125 kCi
- exposure: 1.5 yrs
- fiducial radius: 4 m
- uncertainties
 - on activity: 0% → ∞
 - on fiducial volume: 1%
 - on spectral shape b : 3%
- no background

- maximum sensitivity for oscillation waves in region of the anomalies
- shape uncertainty matters but is under control
- error on activity matters! better 5% needed for gain over wave-only analysis

Source heat power measurement

Two calorimeters for independent measurements of **thermal power** (~1%)

- **Calorimeter inside SOX-Pit**
German groups/Genova
- **Calorimeter outside PIT/in Mayak**
CEA Saclay

*Mounting
mock-up source
in TUM/Genova
calorimeter*

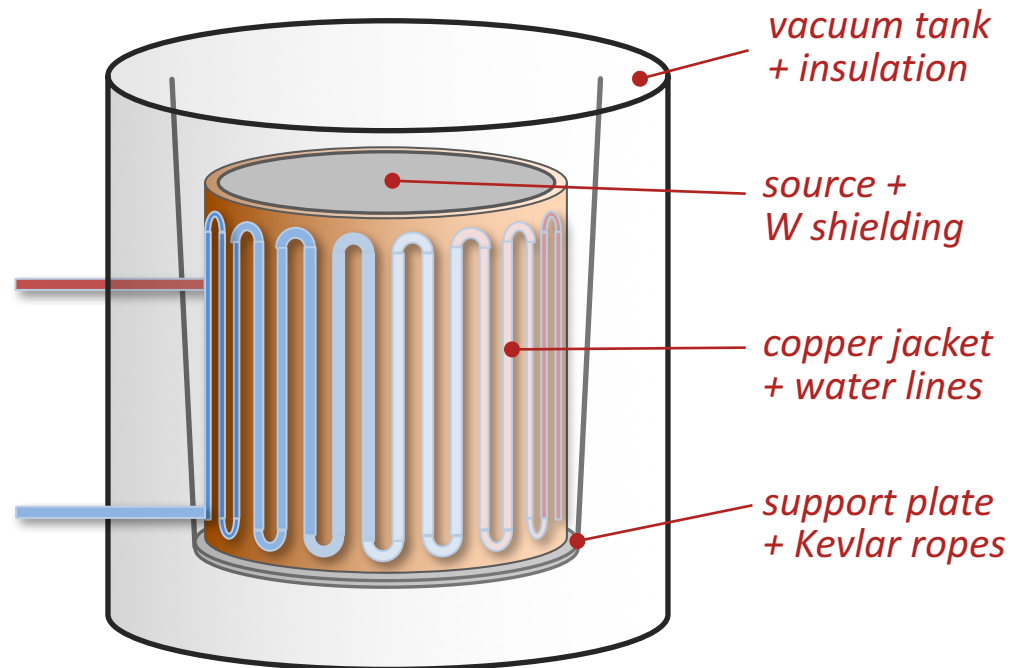


Measurement strategy

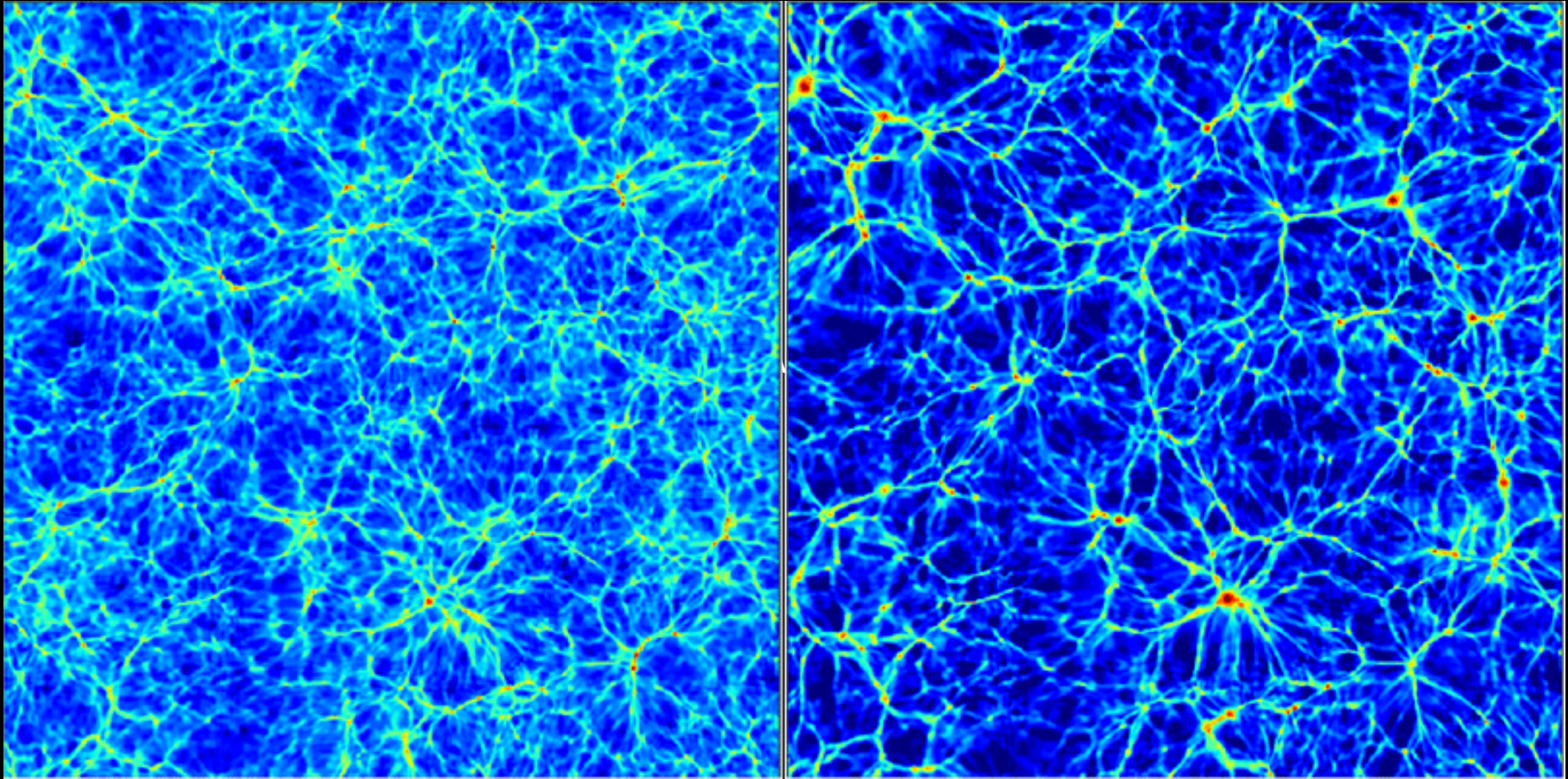
- insulate source from surroundings
- circulate water through loop around W shielding
- measure mass flow Φ and temperature increase ΔT

$$P = \Phi \cdot C_{\text{H}_2\text{O}} \cdot \Delta T + P_{\text{loss}}$$

- **sub-% accuracy reached**
in test measurements



Complementary information on sterile neutrinos



$m_\nu = 0$

$m_\nu = 1.9 \text{ eV}$

Measuring the electron neutrino mass

- Effect of mass is a shift of the endpoint/spectral deformation
- Effective mass is incoherent sum

$$m^2(\nu_e) := \sum |U_{ei}^2| m^2(\nu_i)$$

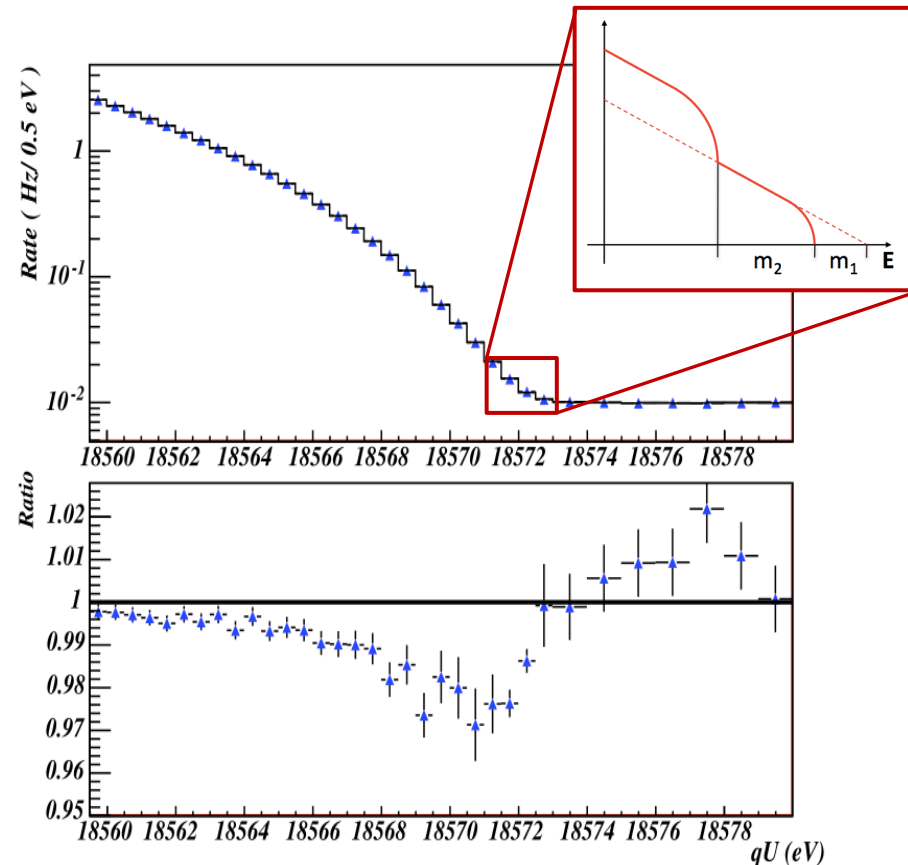
- 3 known mass eigenstates could in principle be resolved but mass differences very small

$$\Delta m_{31} < 50 \text{ meV}$$

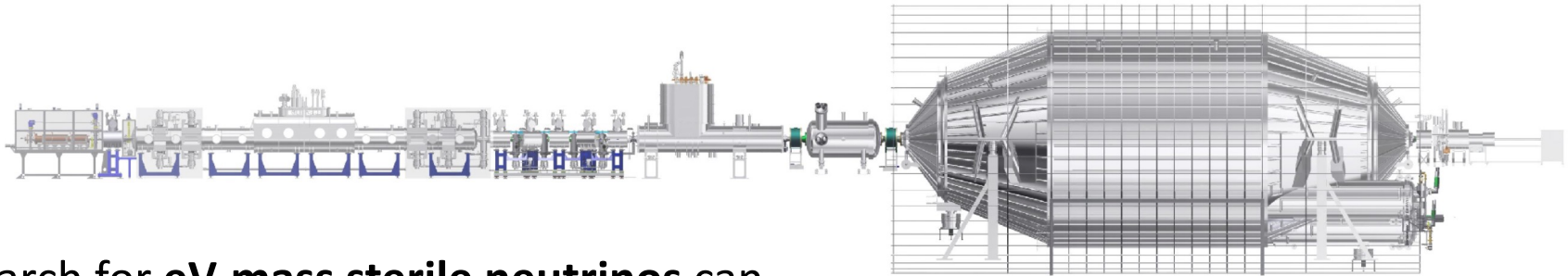
- sterile mass splitting** much larger
 $\Delta m_{41} \sim 1 \text{ eV}$ for light steriles

- Size of effect depends on ν_4 admixture to ν_e flavor state: $|U_{e4}|^2$

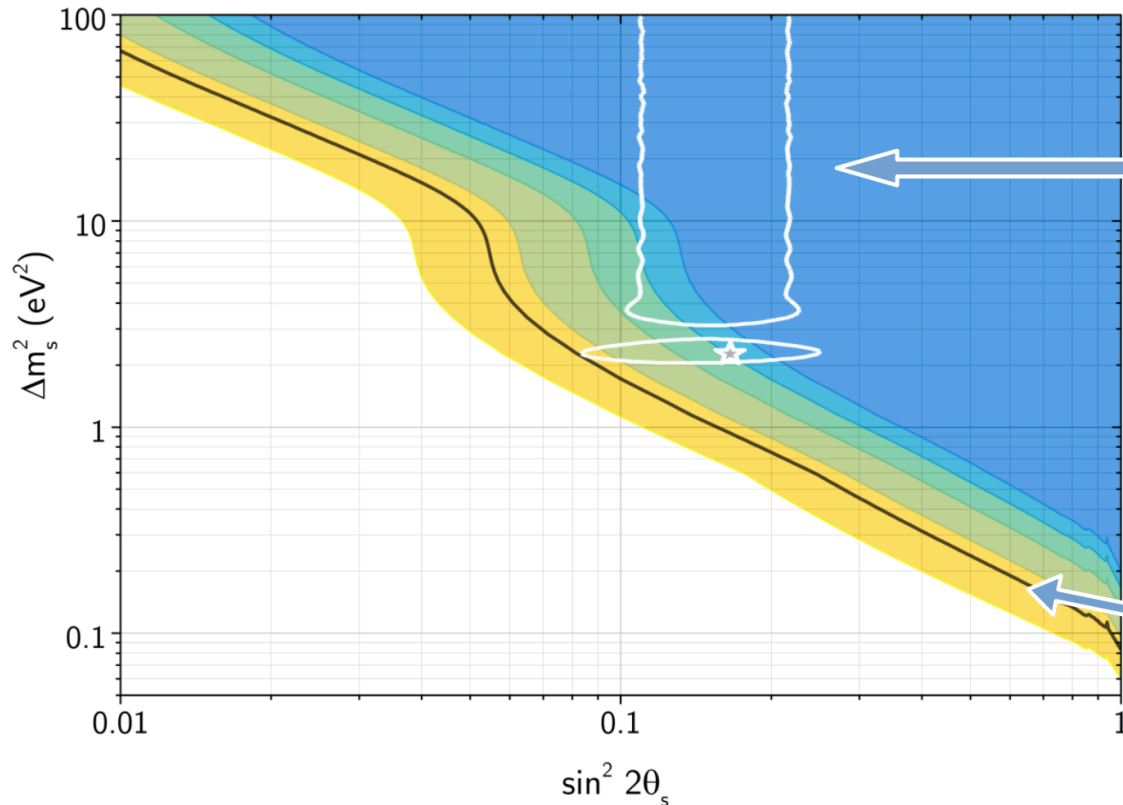
→ **observable in upcoming experiments?**



J. A. Formaggio, J. Barret, PLB 706 (2011) 68
Spectral deformation of tritium decay spectrum
(3-year measurement in KATRIN)
Sterile ν parameters: $\Delta m^2 = 2 \text{ eV}^2$, $|U_s|^2 = 0.067$



- search for **eV-mass sterile neutrinos** can be performed based on regular setup



reactor anomaly
combined fit 90% C.L.
K. N. Abazajian et al. 2012

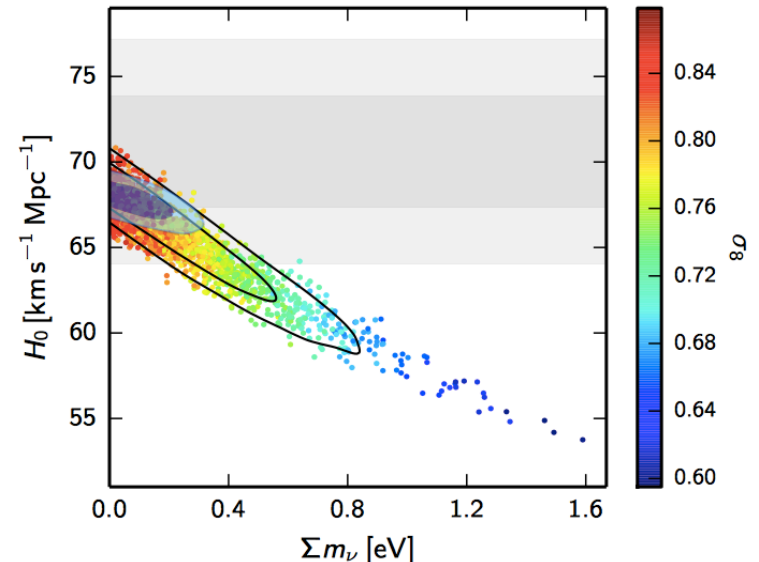
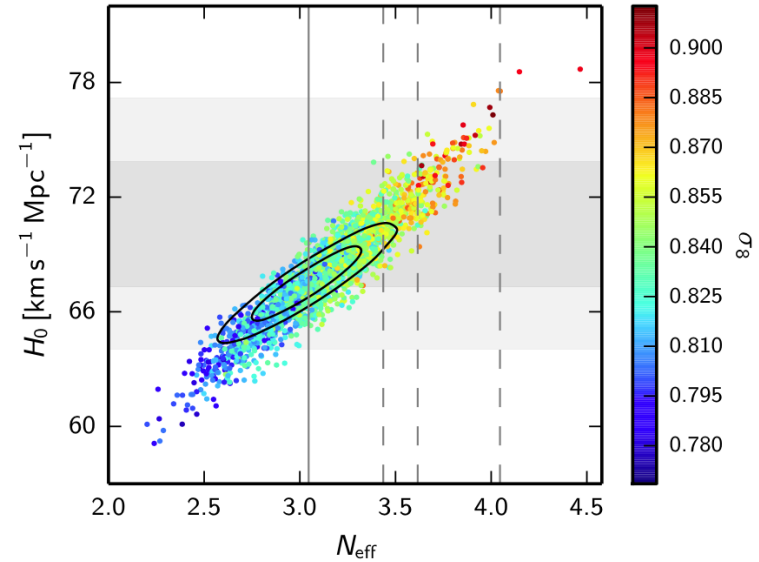
exclusion curve confidence levels

- KATRIN 5σ
- KATRIN 4σ
- KATRIN 3σ
- KATRIN 95%
- KATRIN 90%
- KATRIN 68%

talk by K. Valerius @ PhysStat-v 16

Constraints on light sterile neutrinos

- Cosmological observations able to place stringent bounds on the **number** N_{eff} and **mass sum** Σm_ν of light (i.e. **thermalizing**) sterile neutrinos
- Most important observables
 - Cosmic Microwave Background
 - Big Bang Nucleosynthesis
 - Large-scale structure
- Bounds from PLANCK (+BAO):
 - $N_{\text{eff}} = 2.99 \pm 0.20$
 - $\Sigma m_\nu < 0.49$ (**0.17**) eV (95% C.L.)
- These limits can be avoided by introducing additional physics, e.g. sterile neutrino self-interactions
Dasgupta, Kopp [arXiv:1310.6337]

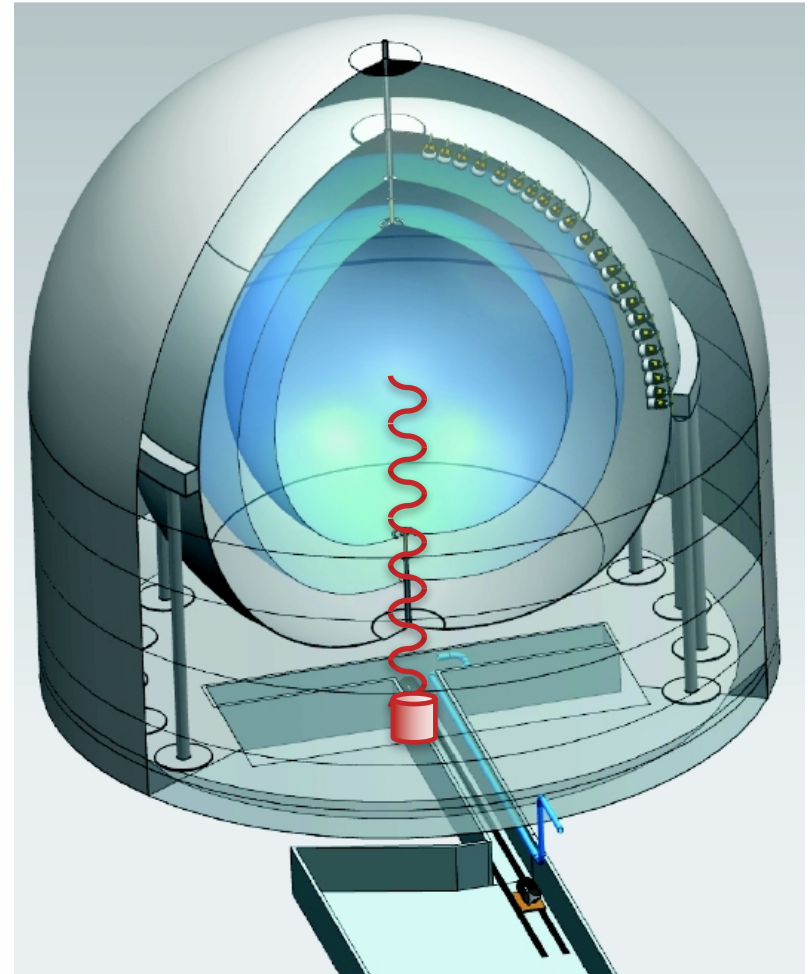


SOX is getting ready:

- Contract with Mayak for the ^{144}Ce source has been signed.
- Experimental site is ready (Borexino, clean room ...)
- Tungsten shield has arrived at LNGS.
- Source calorimeters are in commissioning phase.
- Summer: Complete test of procedures with mock-up source.
- Autumn: Calibration run with radioactive source inside the target.

Start of data taking in early 2018.

- Most of statistics acquired in $\frac{1}{2}$ year
→ stay tuned for first results



Thank you for listening!

Borexino and SOX Collaborations



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