

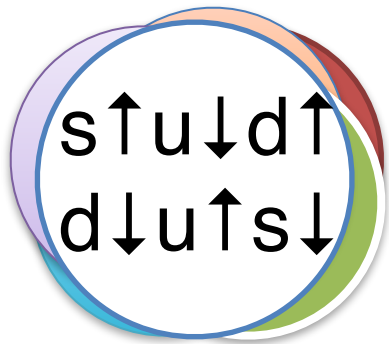
Stable Sexaquark as Dark Matter



Glennys R. Farrar
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Unique among multi-quark states:

Fermi statistics is compatible with a totally symmetric spatial wave function



S

6-quark, $Q=0$, $B=2$

Spin-0, scalar

Flavor singlet

$m \sim 1.7-2$ GeV

(Most-Attractive Channel)³ :

totally Antisymmetric in:

color

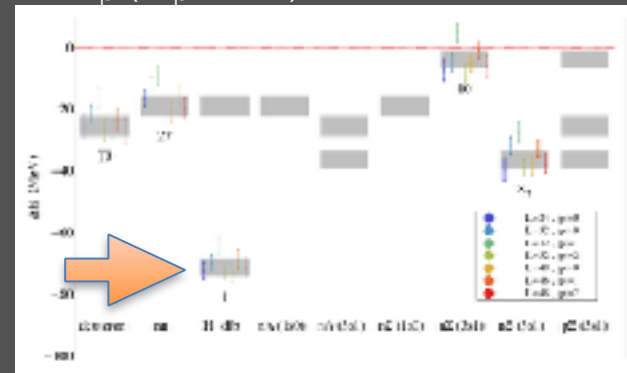
flavor

spin

totally Symmetric in space

Mass of S

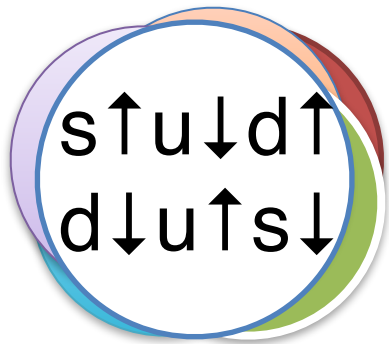
- Options for stability ($\tau > \tau_{\text{Univ}}$)
 - $M_S < 2 m_p + 2 m_e = 1877.6 \text{ MeV} \rightarrow$ **absolutely stable**
 - $M_S < m_p + m_e + m_\Lambda = 2054.5 \text{ MeV} \rightarrow$ decay rate $\sim G_F^2 \times$ wave function overlap \rightarrow lifetime up to 10^{27} yr.
- Is $M_S < 1.7\text{-}2 \text{ GeV}$ reasonable?
 - Hyperfine attraction $\Rightarrow M_S < 2 m_\Lambda = 2230 \text{ MeV}$ (Jaffe, 1977) **(Most-Attractive Channel)³**
 - Constituent quark model unreliable/inapplicable
 - $m_\pi = 135, 140 \text{ MeV}; m_{\eta'} = 958 \text{ MeV}$ (not 600 MeV)
 - Chiral symmetry breaking not necessary for mass of S (unlike baryons)
 - $\pi K K$: same quark content and total mass 1131 MeV
 - Need only 16% (10%) binding energy for $M_S < 2 m_p$ ($m_p + m_\Lambda$)
- Lattice predicts binding (Beane+13)**
- Experiments exclude decaying S**
 \Rightarrow **it must be STABLE ! ;-**



Stable Sexaquark

Same quark content as H-dibaryon (Jaffe 1977), but different physics: not a loosely bound di- Λ !

Deeply bound: $m_S \sim 2 m_p$ and stable (or $\tau \gg \tau_{\text{Univ}}$)



S

6-quark, $Q=0$, $B=2$

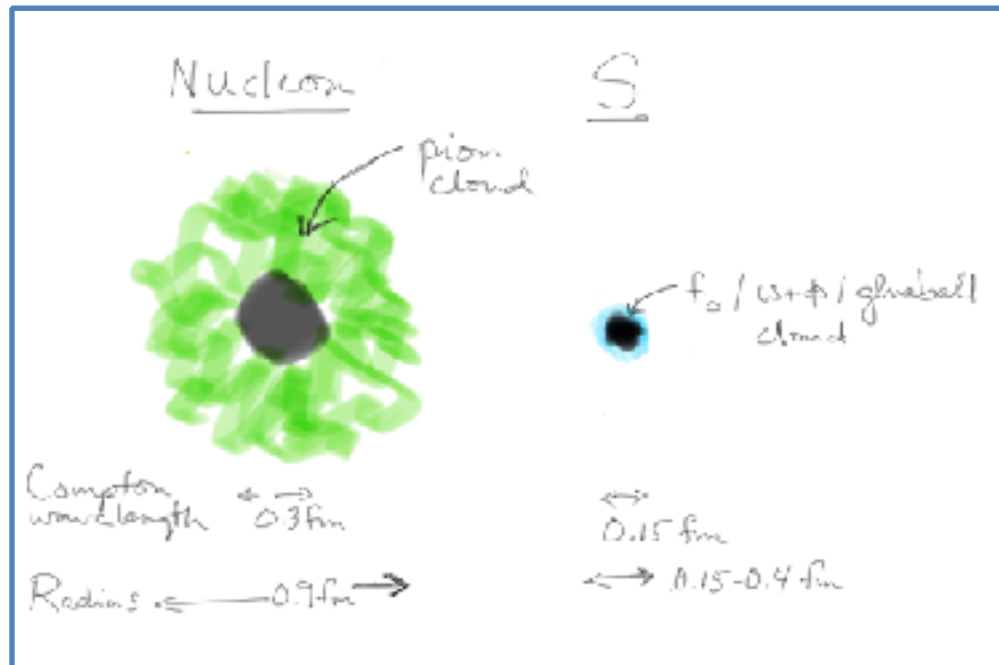
Spin-0, scalar

Flavor singlet

$m \sim 1.7-2 \text{ GeV}$

Crucial features:

S does not couple to pions and is much smaller than usual hadrons (proton, pion,...)



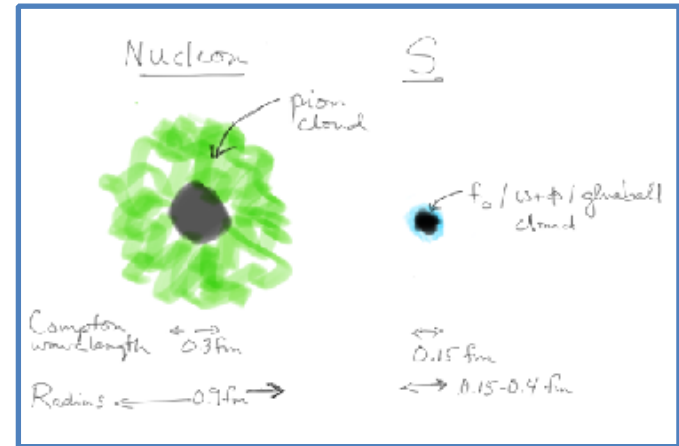
Structure of the S: $u_{\uparrow} u_{\downarrow} d_{\uparrow} d_{\downarrow} s_{\uparrow} s_{\downarrow}$

- Flavor SU(3) singlet \Rightarrow

- no coupling to SU(3) octets π, ρ
- **No pion cloud**

- No pion cloud \Rightarrow

- $r_S \sim$ Compton wavelength \sim **0.15 fm** (compare $r_p = 0.9$ fm)
- **S does not bind to nuclei (no exotic isotopes)**
 "Non-binding of Flavor-Singlet Hadrons to Nuclei", GRF and G. Zaharijas, Phys. Lett. B.559: 223-225, 2003.



- Couples to f_0 , glueball, SU(3) singlet $\omega-\phi$

- $r_S \sim$ **0.15 - 0.4 fm** depending on strength of f_0 coupling
- $\sigma_{SN} \sim 5$ to 20 mb (at $v/c \sim 1$; may grow as $v/c \rightarrow 0$, as σ_{NN} does)

"Nucleon and Nuclear Transitions of the H dibaryon", GRF and G. Zaharijas, Phys. Rev. D70:014008, 2004.

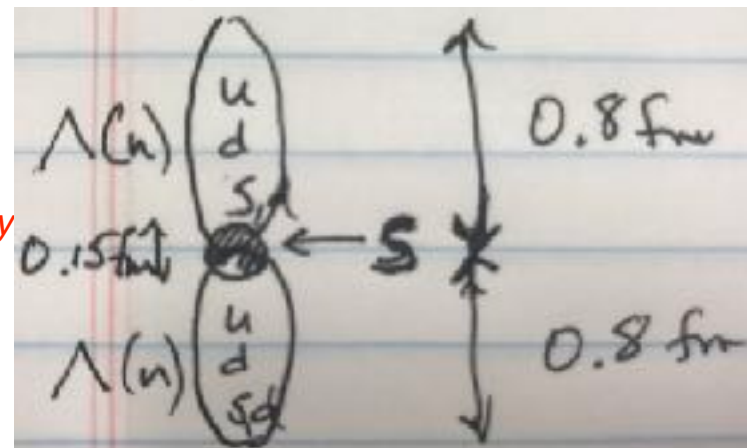
\Rightarrow **Small wave function overlap with 2 nucleons**

- not created in hypernuclear expts
- Amp (NN \rightarrow S) very suppressed (small wavefunction overlap)

\Rightarrow **Nuclei and neutron stars are stable;**

- **S lifetime can be > age of Universe even if not absolutely stable (doubly-weak decay & wfn overlap)**

- Scalar (spin-0, even parity)



"A STABLE H DIBARYON: DARK MATTER CANDIDATE WITHIN QCD?" Int. J. Theor. Phys. 49:1211-1218, 2010. Also in "Minneapolis 2010: Continuous advances in QCD" 582-590.

S has not been discovered because it is elusive

- **Many negative searches, but all are inapplicable.** They either:
 - looked for H-dibaryon through decays (but S is stable)
 - restricted to mass > 2 GeV (but $m_S < 2$ GeV)
 - required fast production in S=-2 hypernuclei (but small overlap with baryons)

- **Wavefunction overlap with baryons is very small.** Extremely rare fluctuation required for $S \leftrightarrow \Lambda\Lambda$; $S \leftrightarrow NN$ is G_F^2 smaller =>
 - nuclei are stable ($\tau > 10^{29}$ yr)
 - hard to produce in fixed target experiments

- **S is similar to** (much more copious) **neutrons**
- **Promising accelerator detection strategies**

- **Apparent lack of baryon number and strangeness conservation:**

- $\Delta B = \pm 2$ with $\Delta S = \mp 2$

- **Reconstruct missing mass, e.g.:**

- $\Upsilon \rightarrow \Lambda \Lambda \bar{S}$ (+ pions) $M_S^2 = (p_\Upsilon - p_{\Lambda 1} - p_{\Lambda 2} - \sum p_{\pi_i})^2$

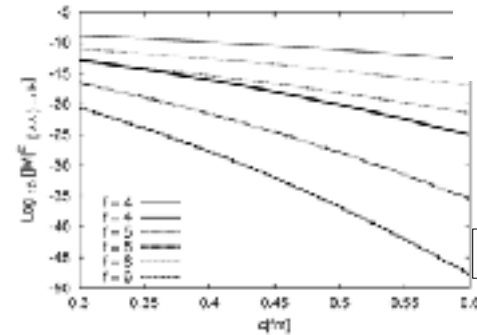
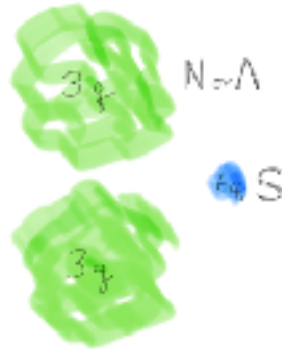
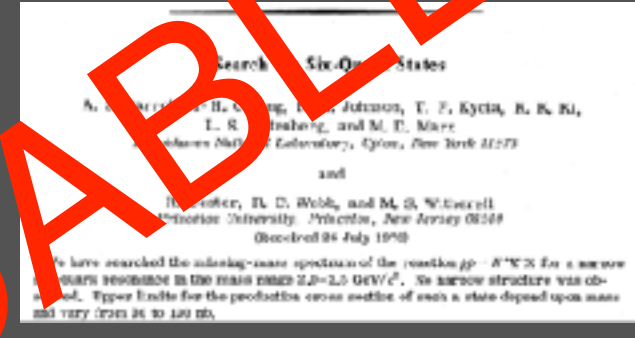


FIG. 1. Log of $[\text{Br}]_{\Upsilon \rightarrow \Lambda\Lambda}$ versus hadron mass ratio in different channels, for ratio $p/m_N = 0.2$ and for values of the hyperon mass $m_\Lambda/m_N = 0.8, 0.9, 1.0, 1.1, 1.2, 1.3, 1.4, 1.5, 1.6, 1.7, 1.8, 1.9, 2.0$ (from left) and $m_\Lambda/m_N = 0.8, 0.9, 1.0, 1.1, 1.2, 1.3, 1.4, 1.5, 1.6, 1.7, 1.8, 1.9, 2.0$ (from right).

GRF+G.Zaharijas 2004

Experimental Searches

- Require $M > 2$ GeV:
 - Gufstafson+ FNAL1976 : Beam-dump + tof *Limit on production of neutral stable strongly interacting particle with mass > 2 GeV.*
 - Carroll+ BNL 1978: No narrow missing mass peak above 2 GeV in $pp \rightarrow K K X$
- Require H-dibaryon decay:
 - Badier+ NA3 1986
 - Bernstein+ FNAL 1988: Limit on production of neutral with $10^{-8} < \tau < 2 \times 10^{-6}$ s
 - Belz+ BNL 1996: $H \rightarrow \Lambda \Lambda$ or $\Sigma \Sigma$ [c.f., issue raised by L. Littenberg]
 - Kim+ Belle 2013: no narrow resonance in $\Upsilon \rightarrow \Lambda \Lambda K$
- Limits on production in doubly-strange hypernuclei:
 - Ah+ BNL 2001
 - Takanashi+ KEK 2001



Discovery Strategy

- Apparent lack of B and S conservation:

- $\Delta B = \pm 2 \quad + \quad \Delta S = \mp 2$

- Reconstruct missing mass, e.g.:

- $\Upsilon \rightarrow \Lambda \Lambda \bar{S}$ (+ pions) $M_S^2 = (p_\Upsilon - p_{\Lambda 1} - p_{\Lambda 2} - \sum p_{\pi_i})^2$

- $K^- p \rightarrow \bar{\Lambda} S$ (+ pions) $M_S^2 = (p_{K^-} + p_p - p_{\bar{\Lambda}} - \sum p_{\pi_i})^2$

- $\bar{S} + p, n \rightarrow \bar{\Lambda} + K^{+,0}$ $M_S^2 = (p_K + p_{\bar{\Lambda}} - p_p)^2$

- or $\bar{E}^{+,0} + \pi$



- Υ is *localized* source of ggg

\Rightarrow production of S is (relatively) enhanced

- Many $\times 10^8$ events collected (CLEO, Babar, Belle)

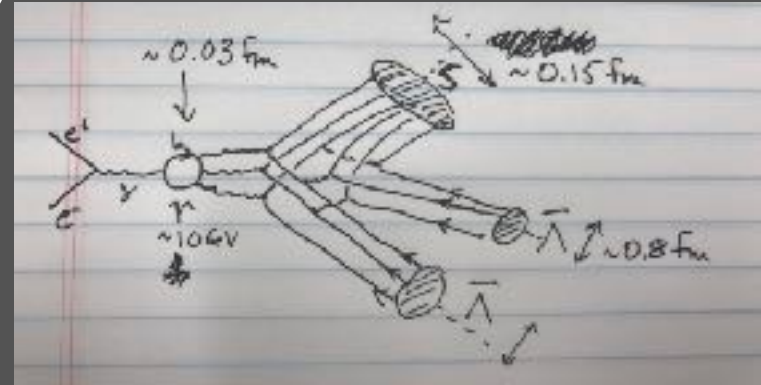
- detectors pretty hermetic, have good mass resolution, $O(10 \text{ MeV})$
- Λ decays quickly to $p\pi^-$ so easy to ID. $c\tau = 8 \text{ cm}$

- Can MEASURE m_S via missing mass

- *Very clean*

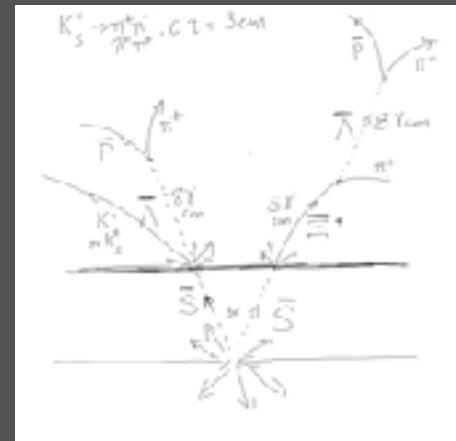
- Main bkg is $K_S K_S K_L K_L$ (+ pions)
 - K_S 's mis-ID'd as Λ 's and K_L 's escaping before decay : *negligible for Belle*
 - rare and can model accurately
 - $K_S K_S K_L K_L$ (+ pions) *is measurable*, from $K^+ K^+ K^- K^-$ (+ pions)
- "Conspiracy" of missed particles producing $\Delta B = \pm 2$, $\Delta S = \mp 2$ very hard

Background does not have narrow peak in missing mass!



Other discovery options besides Upsilon decay

- Production in accelerator expts
 - fixed target: too low production rate due to small wfn overlap
 - LHC strategies
 - statistical examination of correlation between $\Delta B = \pm 2, \Delta S = \mp 2$
 - 2nd exponential in scattering length distribution of n-like interactions, due to S
 - \bar{S} annihilation in tracker, followed by $\bar{E}^{+,0} \rightarrow \bar{\Lambda} \pi^{+,0}$ ($c\tau = 5\gamma$ cm) $\bar{\Lambda} \rightarrow \bar{p} \pi^+$ (63%, $c\tau = 8\gamma$ cm)



- Production at SuperK or SNOlab (nuclear decay) $p n \rightarrow S e^+ \dots$

Conditions on QCD Dark Matter

- ✓ $\tau_{\text{DM}} > \tau_{\text{Univ}}$, cold, neutral
- ✓ primordial nucleosynthesis
- ✓ Particle must not be already excluded
 - accelerator searches
 - exotic isotopes
 - **DM searches**
 - indirect impacts (heating planets, helioseismology,...)
 - stability of nuclei
 - equation of state of neutron stars (and their stability)
- ✓ Correct relic density (for natural m_{DM} - $\sigma_{\text{f.o.}}$)

Dark Matter Relic Abundance

- Reheating: production of photons & fragmentation of $\langle \phi_S \rangle$ into particles

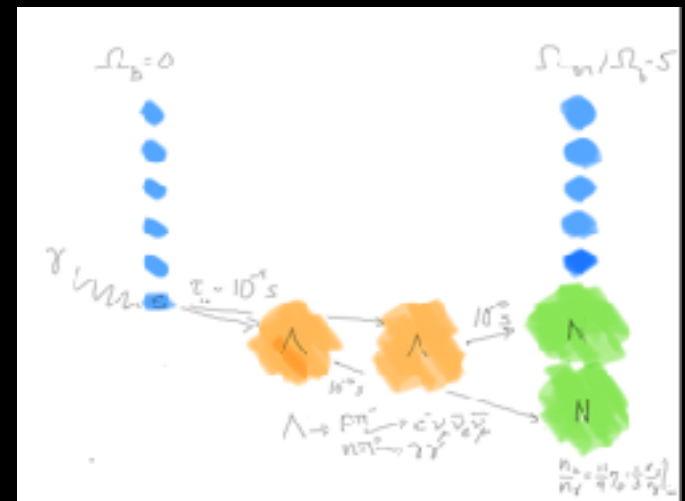
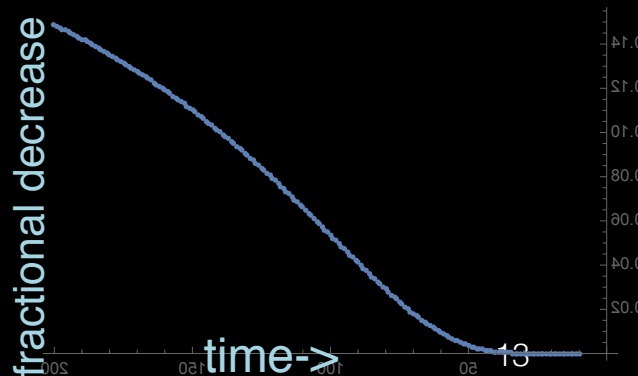
- Breakup: $\gamma + S \rightarrow \Lambda \Lambda$

$$\frac{dn_S}{n_S dT} = -\Gamma(\gamma + S \rightarrow \Lambda\Lambda) \frac{dt}{dT}$$

$$\Gamma(\gamma + S \rightarrow \Lambda\Lambda)(T) = \langle n_\gamma(T) \sigma_{\gamma+S \rightarrow \Lambda\Lambda}(T) v(T) \rangle$$

- ~16% of S's are broken up before freeze out \Rightarrow

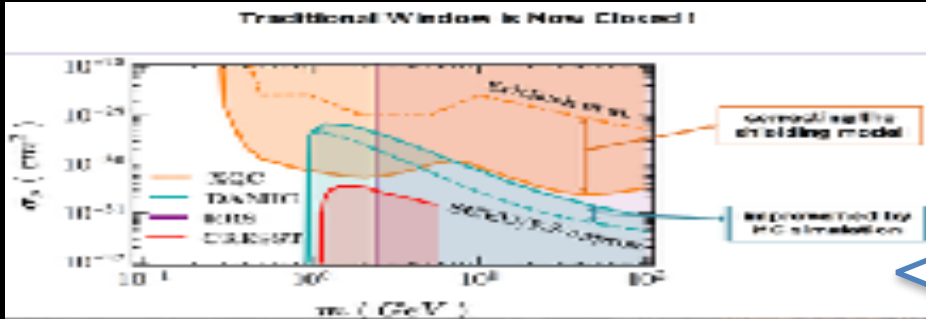
$$\Omega_{DM} / \Omega_b = 5. \text{ Natural!}$$



$10^{-26} - 10^{-25} \text{ cm}^2$



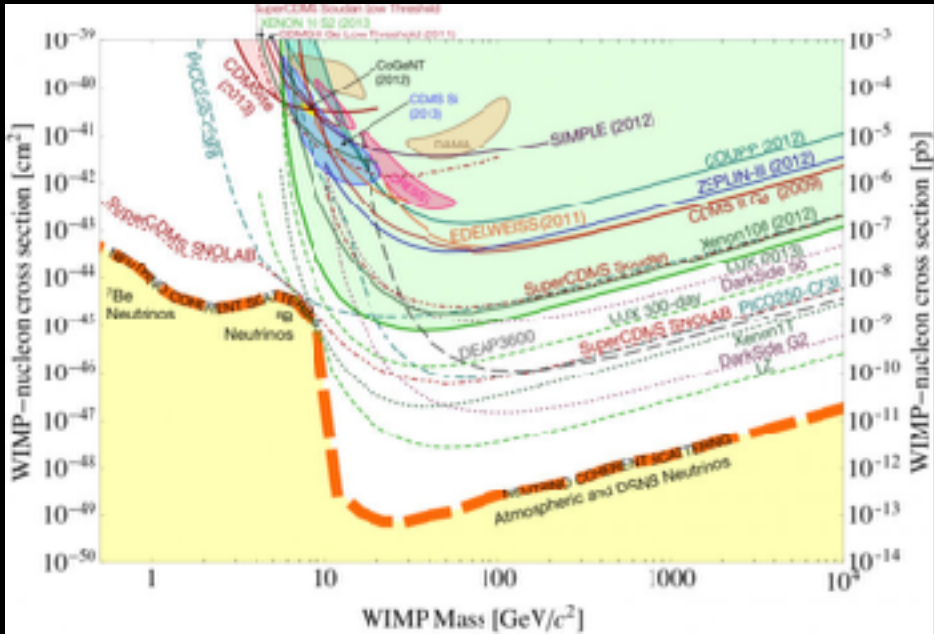
Stable S as Dark Matter



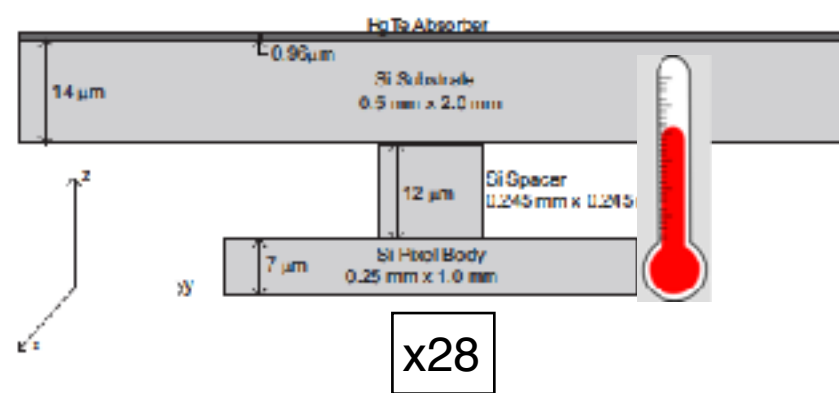
Closing the window on $\sim \text{GeV}$ Dark Matter with moderate ($\sim \mu\text{b}$) interaction with nucleons

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 Center for Cosmology and Particle Physics, Department of Physics, New York University,
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Abstract. We improve limits on the spin-independent scattering cross section of Dark Matter on nucleons, for DM in the 300 MeV – 100 GeV mass range, based on the DAMIC and XQC experiments. Our results close the window which previously existed in this mass range, for a DM-nucleon cross section of order $\sim \mu\text{b}$, assuming the standard velocity distribution.



X-Ray Quantum Calorimeter (XQC)



- XQC on sounding rocket, 200 km above earth
- Best limit for high x-secn (McCammon+02, Wandelt+02, GF+Zaharijas05, Erickcek+07, Mahdawi & GF 17)
 - sensitive to X-rays with $E \geq 30$ eV
 - 100 sec flight, ~ 100 events
 - nuclear recoil \Rightarrow X-rays, which thermalize (*assumption*)

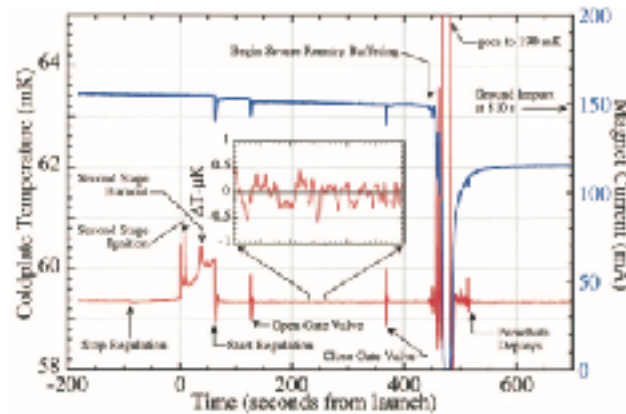
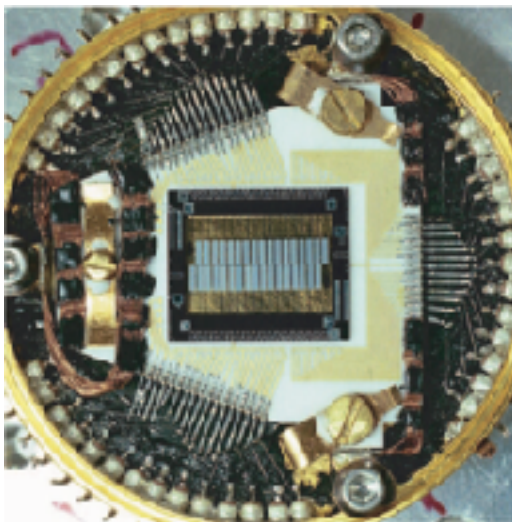


Fig. 7.—Highly performance of the temperature control system, showing the colloplate temperature and magnet current. Temperature fluctuations during data taking are about 2.5 mK rms. The gate-valve action is located on the current (blue) and caused the most serious thermal disturbance up to 800 μs. Accelerations during gateway (repeated) data taking are about ~ 4 Hz, but only higher than the colloplate frame rate can be removed. Temperature repetition is recovered once stabilizing steps, allowing calibration data to be obtained.

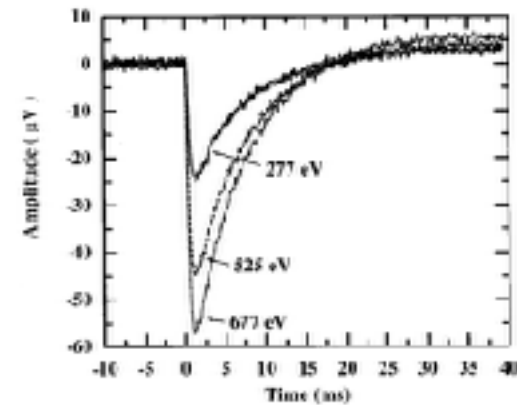


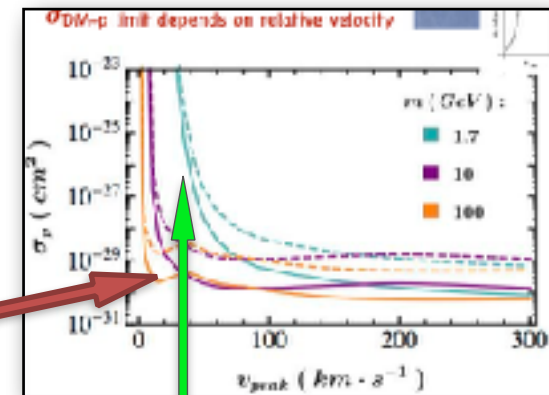
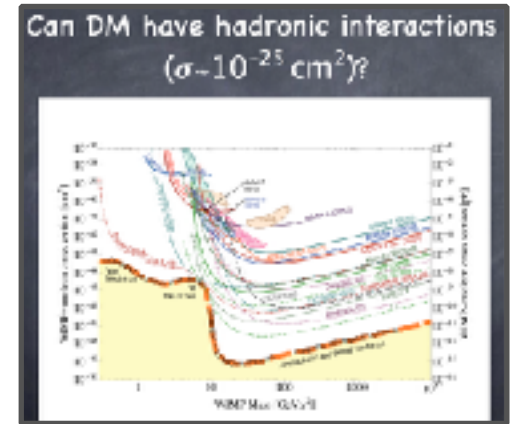
Fig. 8.—Unfiltered X-ray pulses from the gate-valve calibration source.

Evading Direct Detection II.

XQC applicability

Evading Direct Detection II. Co-Rotation

- Light, *slow-moving* DM is not visible in current detectors
 - $KE_{DM} = 500 \text{ eV} (m_{DM} / 2 m_p) v^2 / (220 \text{ km/s})^2$
 - $\langle E_{\text{deposit}} \rangle / KE_{DM} = 0.12$ for Si target
 $= 0.02$ for Hg target
 $= 0.44$ for H or He target
 - Energy-loss length in Earth crust: $\lambda = 2 \text{ cm} / \sigma_{10mb}$
 - ~~CRESST, Xenon1T, LUX, DAMIC~~
- **XQC above atmosphere is best; $E_{\text{thresh}} = 30 \text{ eV}$**
 (McCammon+02, Wandelt+02, GF+Zaharijas05, Erickcek+07, Mahdawi+GRF 2017)
 - Mahdawi+GRF 2017: $\sigma_{DM} < 10^{-29} \text{ cm}^2$ *for standard velocity dispersion* — *SDM has $\sigma_{DM} \sim 10\text{-}100 \text{ mb}$*
 - $v_{\text{min},XQC} = 107 \text{ km/s} (2 m_p / m_{DM})^{1/2}$



Co-rotation reduces SDM signal

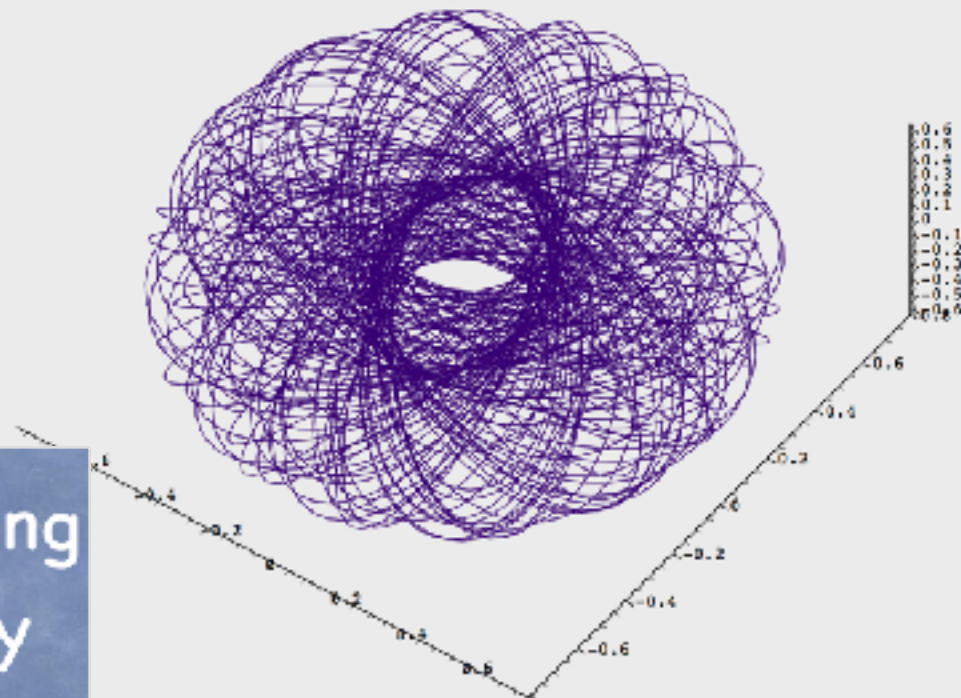
Viability of Co-Rotation Scenario

Schwarzschild Modelling

1. Choose a potential (NFW)
2. Follow an orbit in this potential

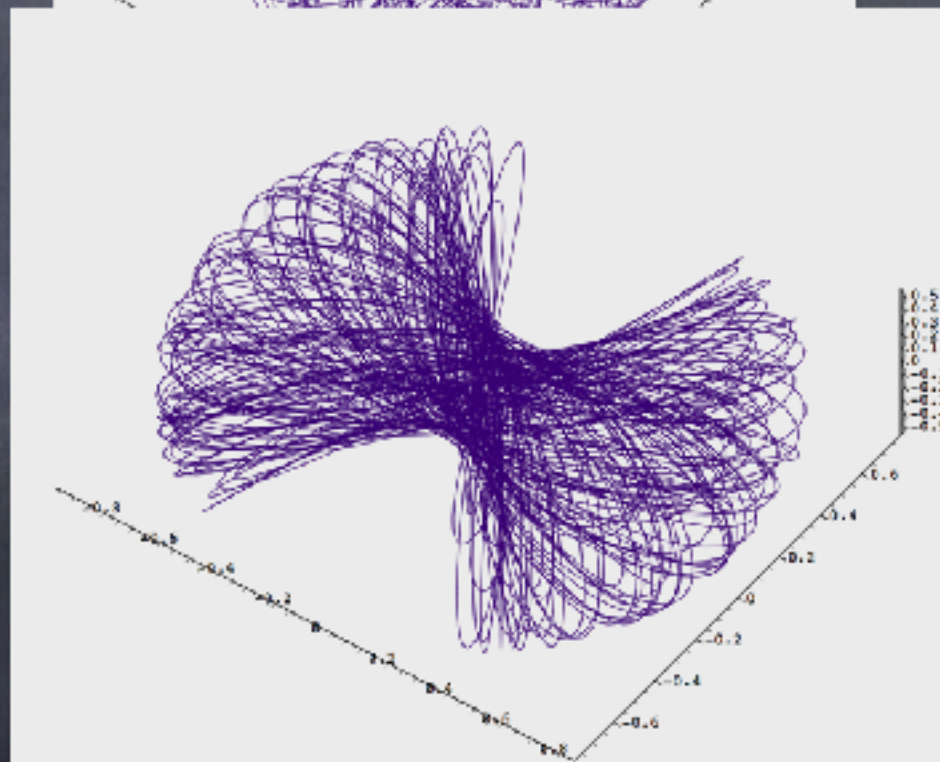
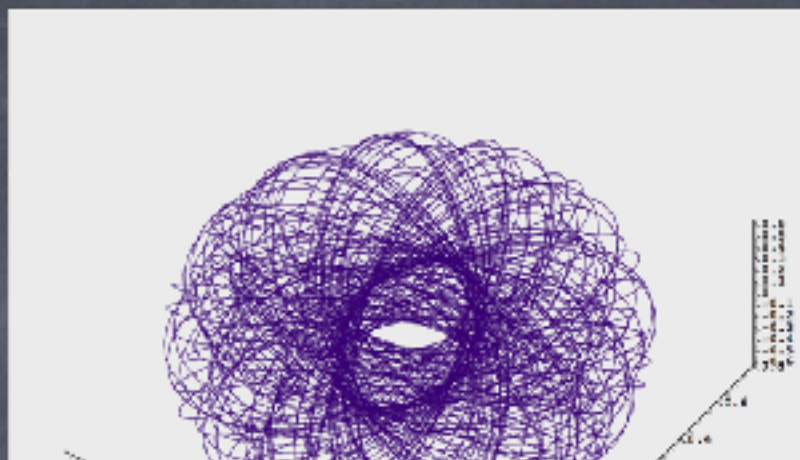
Simulations of DM interacting with gas in the Milky Way

(Jay) Digvijay Wadekar (NYU)
in collaboration with G. R. Farrar



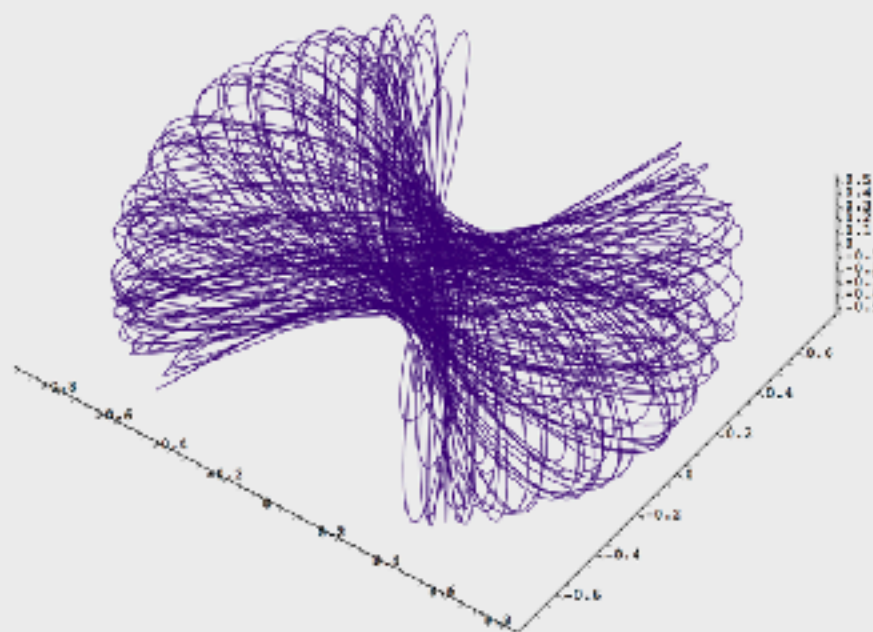
Schwarzchild

1. Choose a potential (NFW)
2. Follow an orbit in this potential
3. Generate a library of orbits



Schwarzschild

1. Choose a potential (NFW)
2. Follow an orbit in this potential
3. Generate a library of orbits
4. Weighted sum of orbits gives density corresponding to NFW

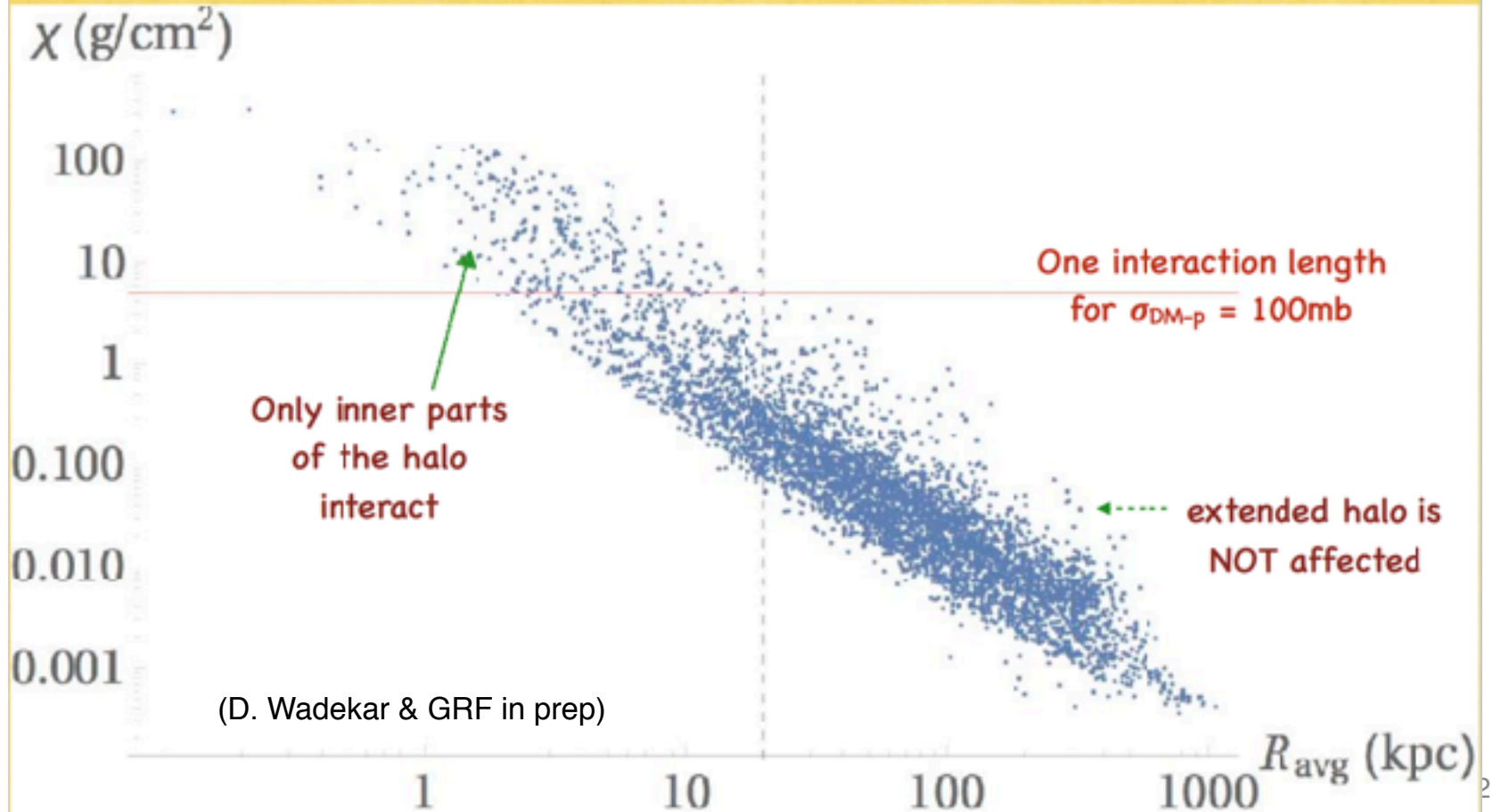


Dark Matter-Gas Scattering

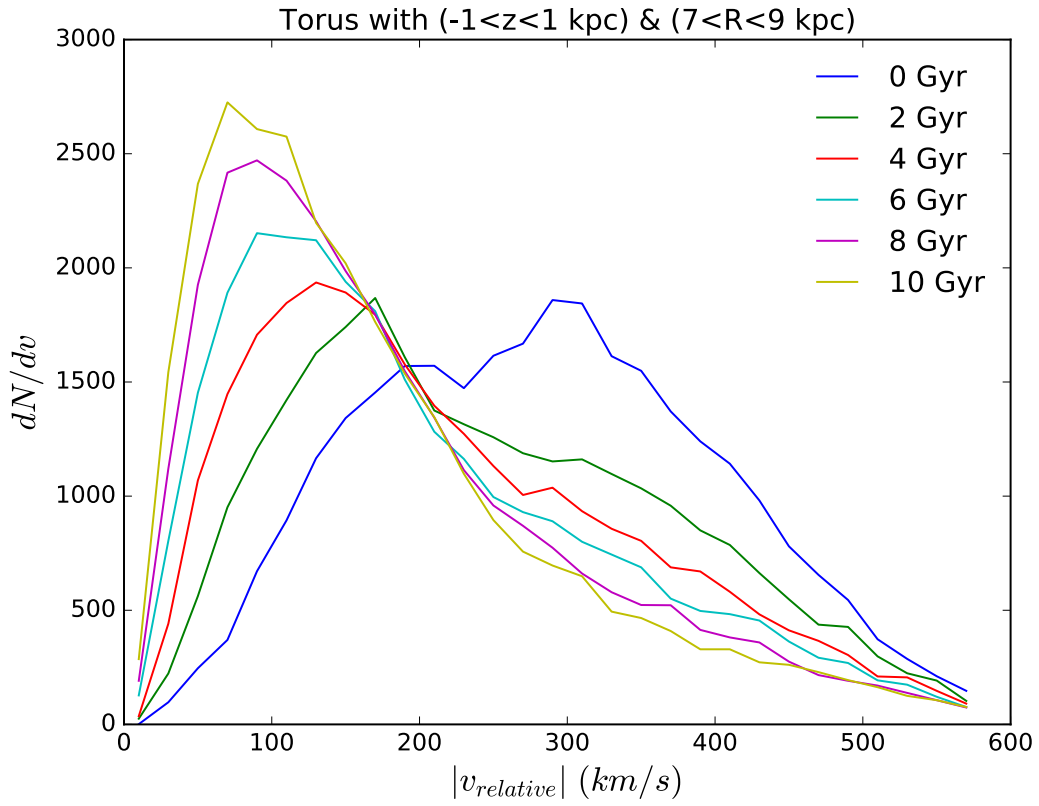
$v_f \sim (2/3)^N v_i \Rightarrow$ only a few interactions are enough for co-rotation scenario to work

Integrated Column Density
for NFW orbits 10 Gyr

$$\chi = \int_0^l dl' \rho_{\text{gas}}(r)$$



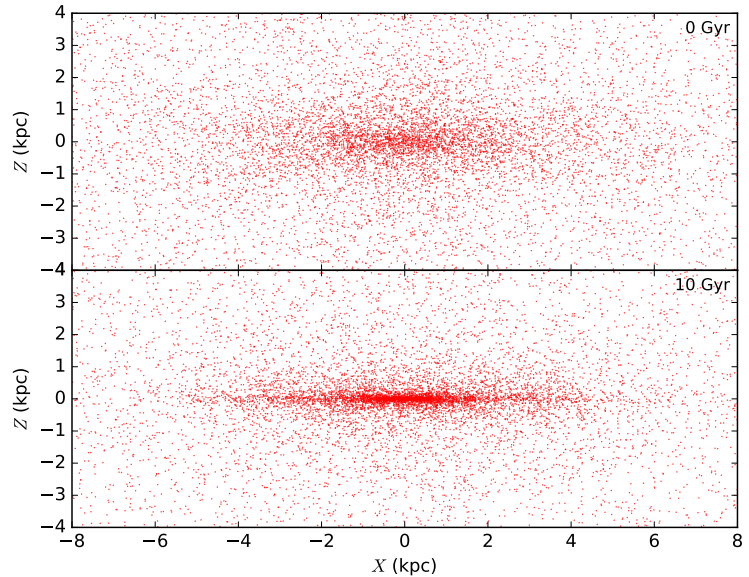
- Select a DM particle and follow its orbit
- Monte-carlo choose its scattering point on gas
- Isotropically scatter \rightarrow kinematics of new orbit



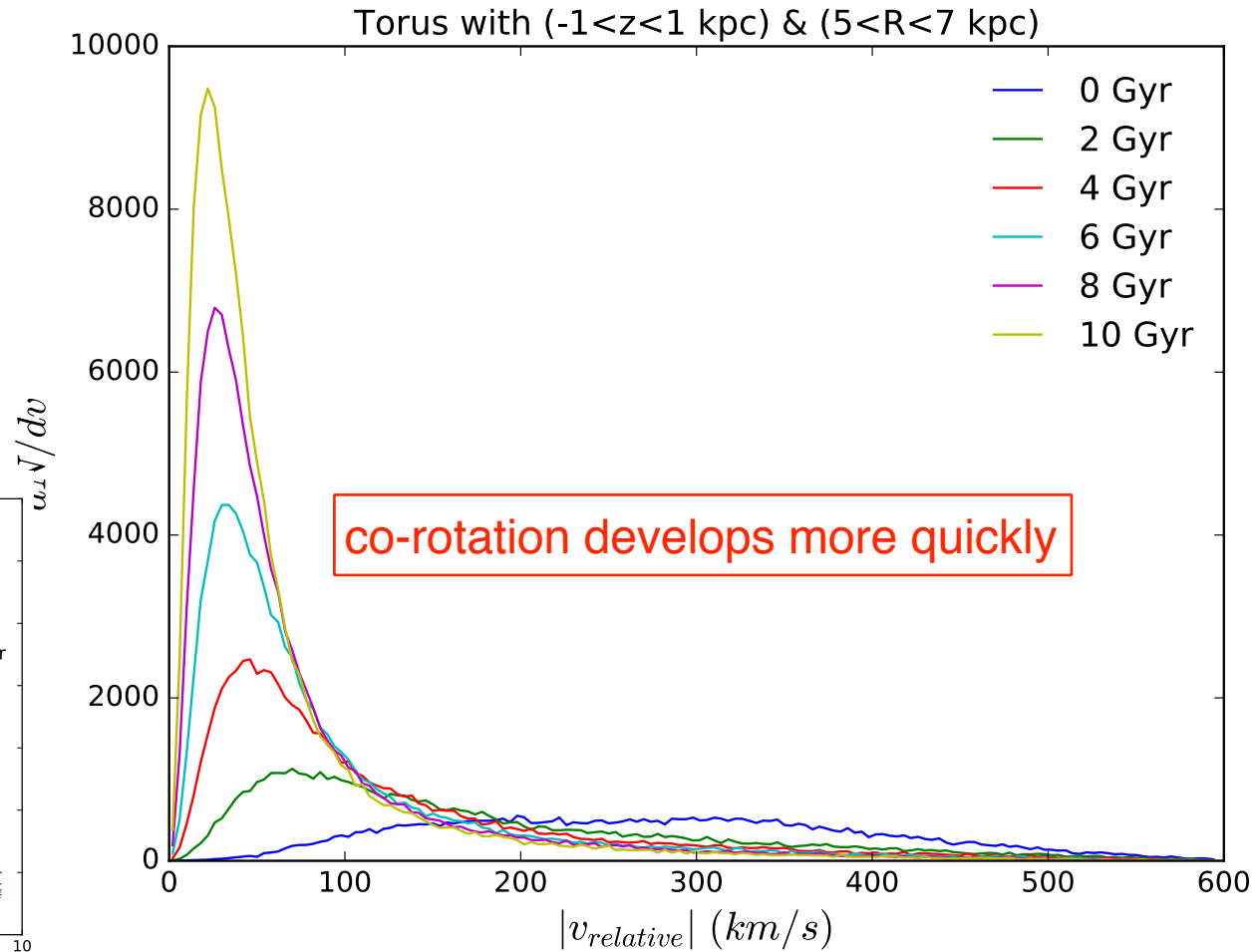
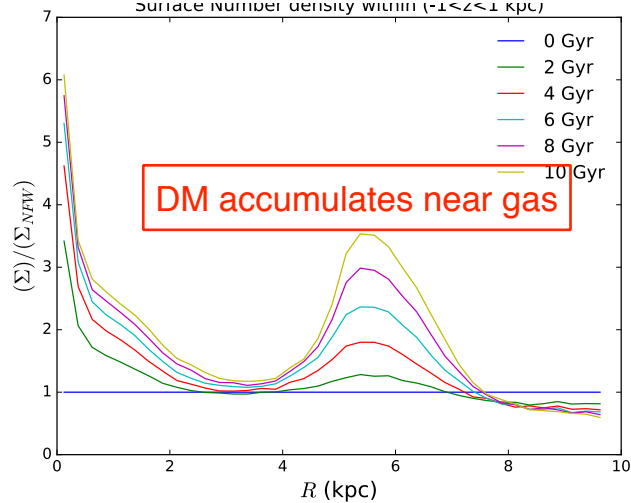
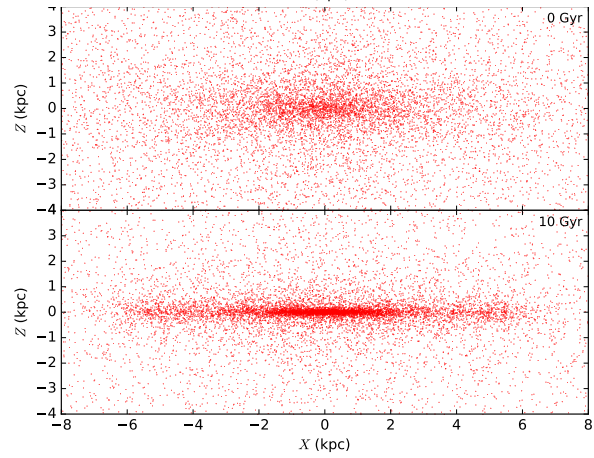
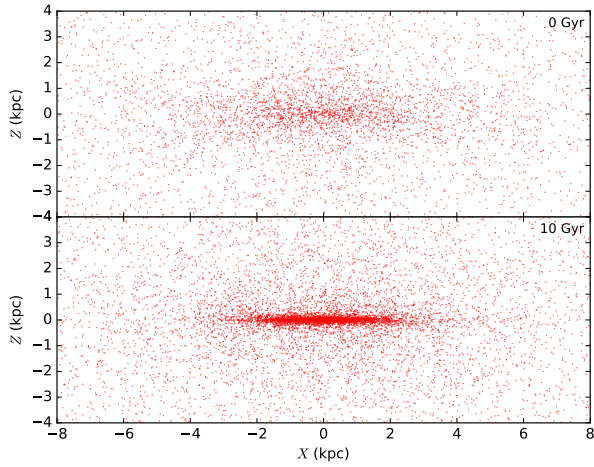
Co-rotation!

Highly idealized analysis:

- static potential & gas
- actually grows by accretion
- c.f. Bruch-Read+
- proof of concept, shows cross-section needed is reasonable.



in realistic galaxies, gas is inhomogeneous \Rightarrow
test with extra torus



Rotation Curves \Leftrightarrow DM reflects gas

Swaters+12: hundreds of galaxies of all types

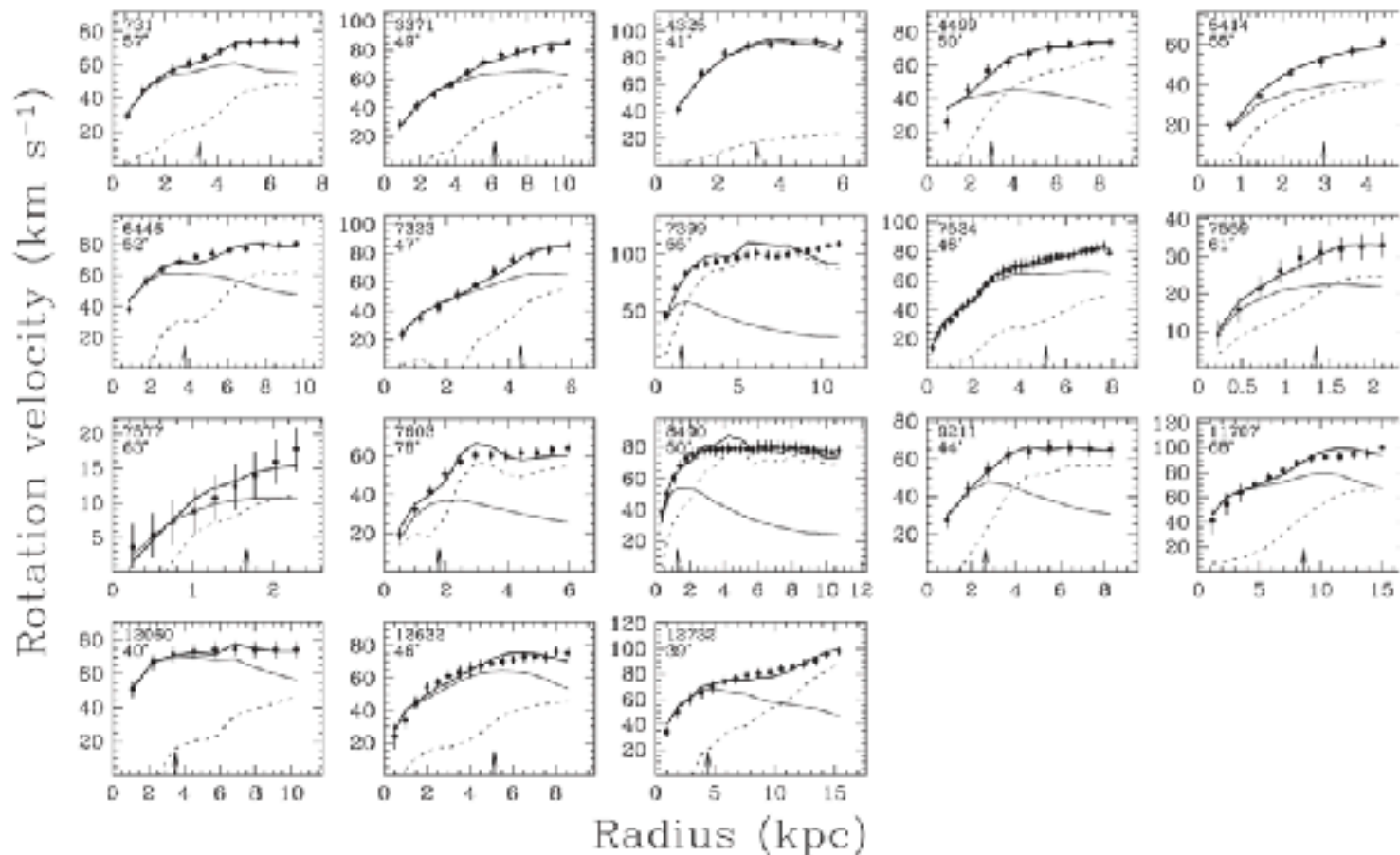
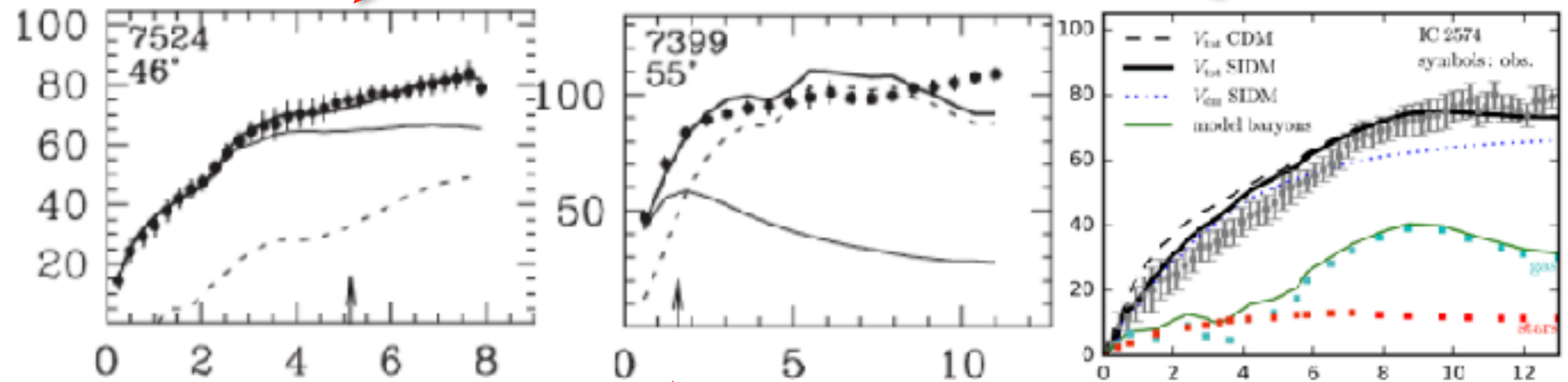


Figure 1. Mass models based on scaling the stellar disc and the H I component for the late-type dwarf galaxies in our sample. The filled circles represent the derived rotation curves. The thin full lines represent the contribution of the stellar discs to the rotation curves and the dotted lines that of the gas. The thick solid lines represent the best-fitting model based on scaling the contributions of the stars and the gas. The arrows at the bottom of each panel indicate a radius of two optical disc scale lengths. In the top left corner of each panel, the UGC number and the inclination are given.

Rotation Curves \Leftrightarrow DM reflects gas

Sharp increases in overall RC, amplifying structure in gas



But not always... Exceptions \Leftrightarrow galaxy experienced recent merger, as predicted by SDM!

“Exceptions prove the rule” SDM provides natural explanation...

Swaters+12

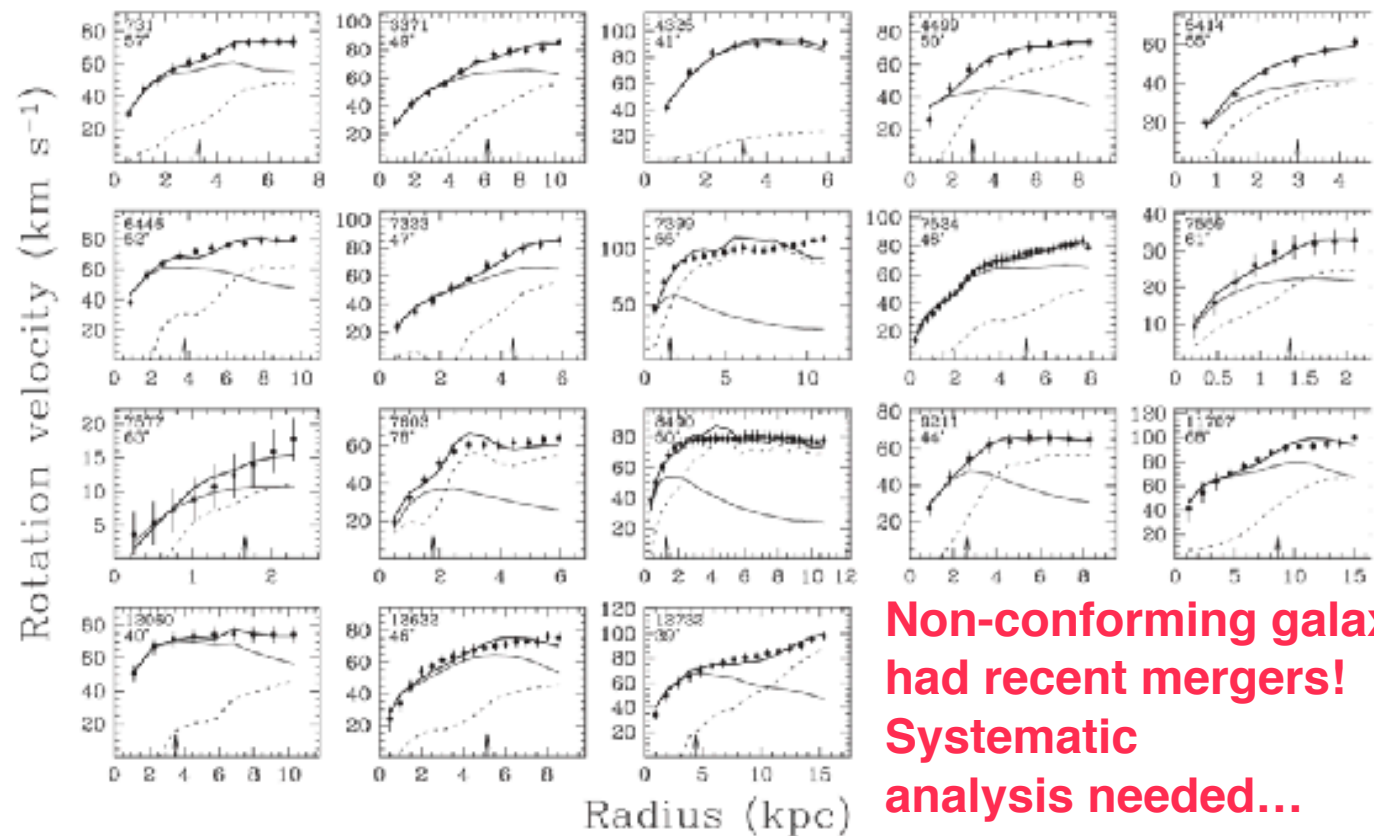


Figure 1. Mass models based on scaling the stellar disc and the H I component for the late-type dwarf galaxies in our sample. The filled circles represent the derived rotation curves. The thin full lines represent the contribution of the stellar discs to the rotation curves and the dotted lines that of the gas. The thick solid lines represent the best-fitting model based on scaling the contributions of the stars and the gas. The arrows at the bottom of each panel indicate a radius of two optical disc scale lengths. In the top left corner of each panel, the UGC number and the inclination are given.

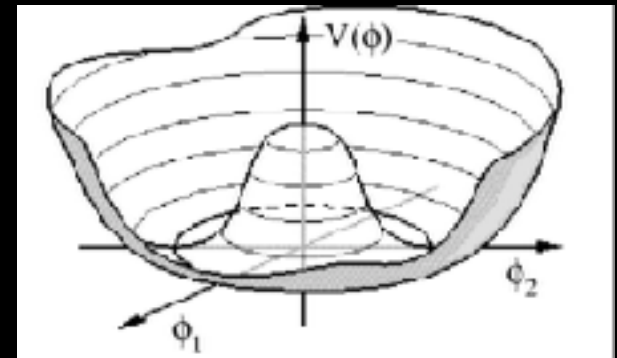
Cosmology & structure formation

- **DM-baryon interaction: momentum transfer => *slight drag on DM during structure formation***
 - Dvorkin, Blum, Kamionkowski (2014):
 - **Ly-alpha forest: $\sigma < \sim 10$ mb if v-indept**
 - Buen-Abad, Marques-Tavares, Schmaltz (2015):
 - **momentum transfer helps reconcile H_0 & σ_8**
 - Problem or opportunity? To be determined...
- **S-S self interactions + S-baryon interactions: *similar benefits as Self Interacting DM***
 - core-cusp, “too-big-to-fail”, subhalo problems, etc.

Baryon Asymmetry of the Universe

- S described by complex scalar field, like the Higgs
 - may have VEV in early Universe
 - phase rotates at no cost in energy \Rightarrow
 - **Non-zero baryon number density is generic**

$$n_B \sim \langle \phi_S^\dagger \overleftrightarrow{\partial}_t \phi_S \rangle$$



Baryon number is non-zero. After reheating it looks like it was just an initial condition.

Strategies to detect DM

if DM is comprised of S's

- With $v_{rel} = 30 \text{ km/s}$, $KE_{DM} \sim 10 \text{ eV}$

- $\langle E_{deposit} \rangle / KE_{DM} = 0.12$ for Si target
 - $= 0.02$ for Hg target
 - $= 0.44$ for H or He target

- Energy-loss length of S in Earth crust:

$$\lambda = 2 \text{ cm} / \sigma_{10mb}$$

- Present detectors shielded or too high threshold (new CRESST expt has $E_{th} = 19 \text{ eV}$, but 30 cm shielding \Rightarrow not sensitive)

- Heating rate liquid He: $\sim \text{nW/mol} \sim \text{CR muon energy deposit rate.}$

- can't *shield* muons & other CRs: veto? but what about neutrons?

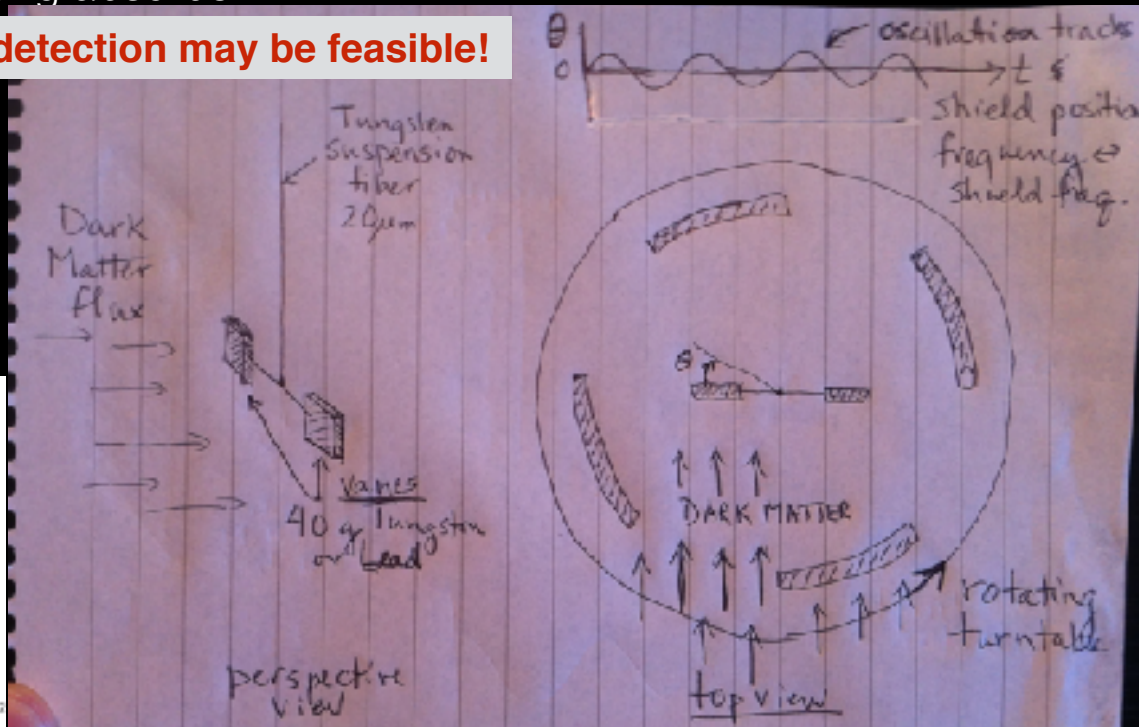
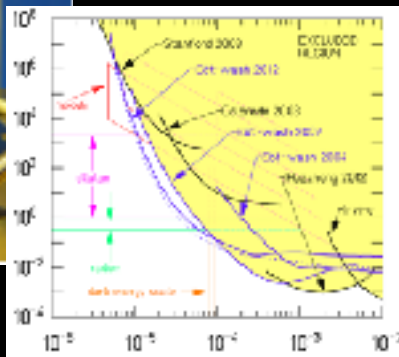
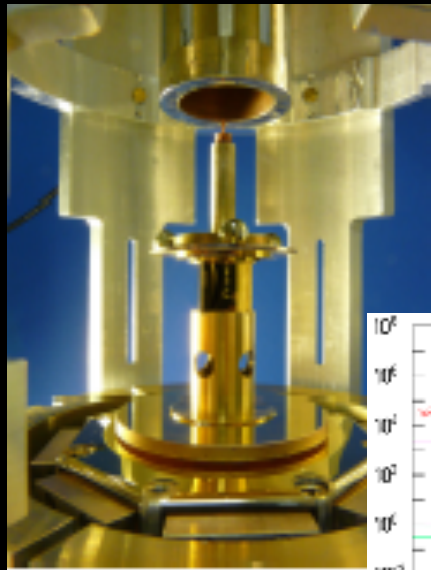


S dark matter detection with a torsion balance on the ISS

W. Terrano & GRF, in preparation

- Individual S collision deposits too little E to detect, but an S flux exerts a tiny pressure.
- Torsion balance: (Eotwash) 1 yr torque sensitivity $\sim 2 \cdot 10^{-11}$ dyne-cm (erg)
- Modulate DM pressure by rotating absorber

directional S-detection may be feasible!



Key points to take home

- ***There may a tightly bound 6-quark state $S = uuddss$***
 - **Unique, symmetric structure \Rightarrow other hadrons don't provide guidance**
 - mass is not driven by chiral symmetry breaking (unlike baryons)
 - constituent quark model probably completely misleading
 - ***If $M_S < 2 m_p + 2 m_e$, S is absolutely stable***
- ***If S is stable, its an excellent Dark Matter candidate***
 - Relic abundance is natural.
 - Usual WIMP detection strategy isn't applicable.
 - Baryon asymmetric Universe is expected
 - May reconcile tension in H_0 & σ_8 and explain astrophysics puzzles ("quenching", core-cusp, DM rotation curves...)
- ***S may be waiting to be discovered in existing Υ -decays or LHC experiments... mass can be accurately measured in Υ -decay.***
- ***SDM will be challenging to detect, but not impossible. Astrophysical and cosmological effects may allow it to be constrained, excluded or confirmed.***

Backup Slides

Sexa vs Hexa...

- https://en.wikipedia.org/wiki/Numeral_prefix
- If sexaquark is DM it, should be renamed R for Rubin!

Number	Latin prefixes				Greek prefixes**			Sanskrit ^[1]
	Cardinal	Multiple	Distributive	Ordinal	Cardinal	Multiple Proportional Quantitative	Ordinal	
6	sexa- ^[19]	–	sen- ^[20]	sext- ^[21]	hex- ^[22]	hexakis- hexaplo- hexad- e.g. hexahedron	hect- ^[23] hectaic-	shat-

^{a b} Sometimes Greek *hexa-* is used in Latin compounds, such as *hexadecimal*, due to **taboo avoidance** with the English word **sex**.

Potential impediments to SDM scenario

- ~~Direct detection limits from X-ray Quantum Calorimeter — co-rotation!~~
- ~~Existence of Neutron Stars, Eqn of State NS~~

PHYSICAL REVIEW D **85**, 023519 (2012)

Constraints on scalar asymmetric dark matter from black hole formation in neutron stars

Samuel D. McDermott, Hai-Bo Yu, and Kathryn M. Zurek

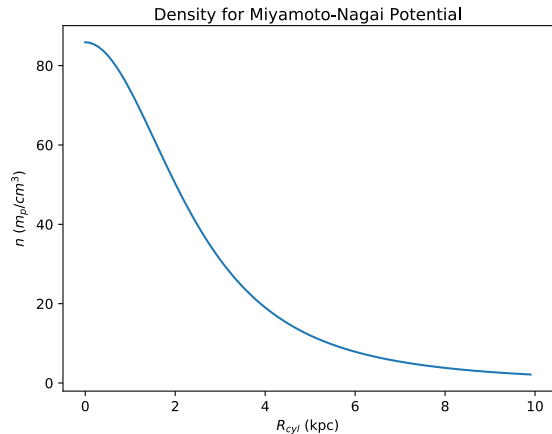
Michigan Center for Theoretical Physics, Department of Physics, University of Michigan, Ann Arbor, Michigan 48109, USA

(Received 21 April 2011; published 13 January 2012)

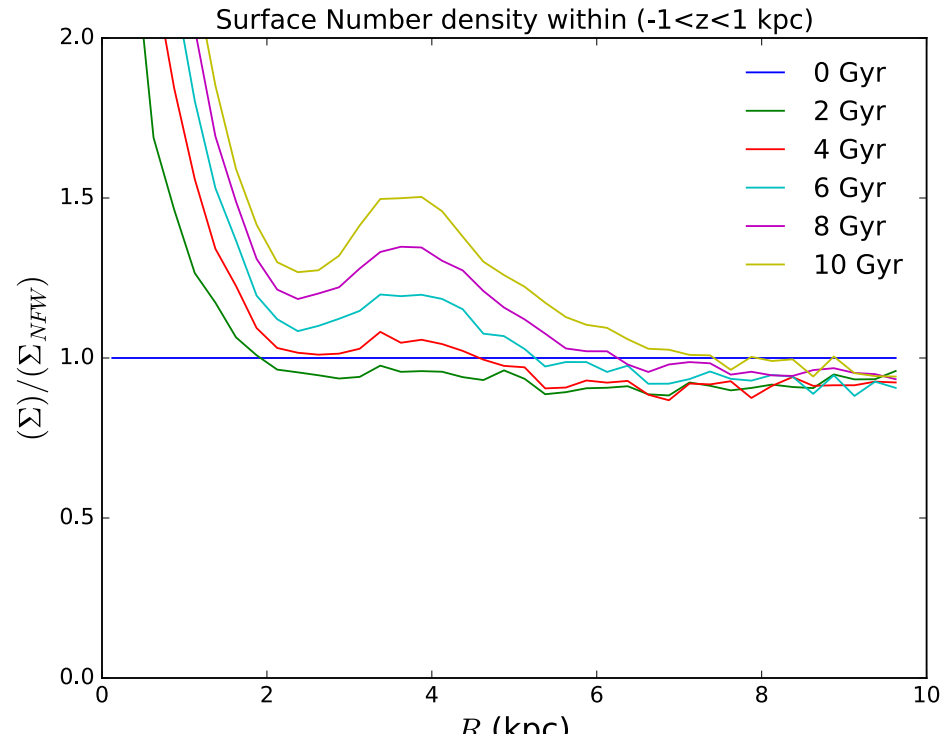
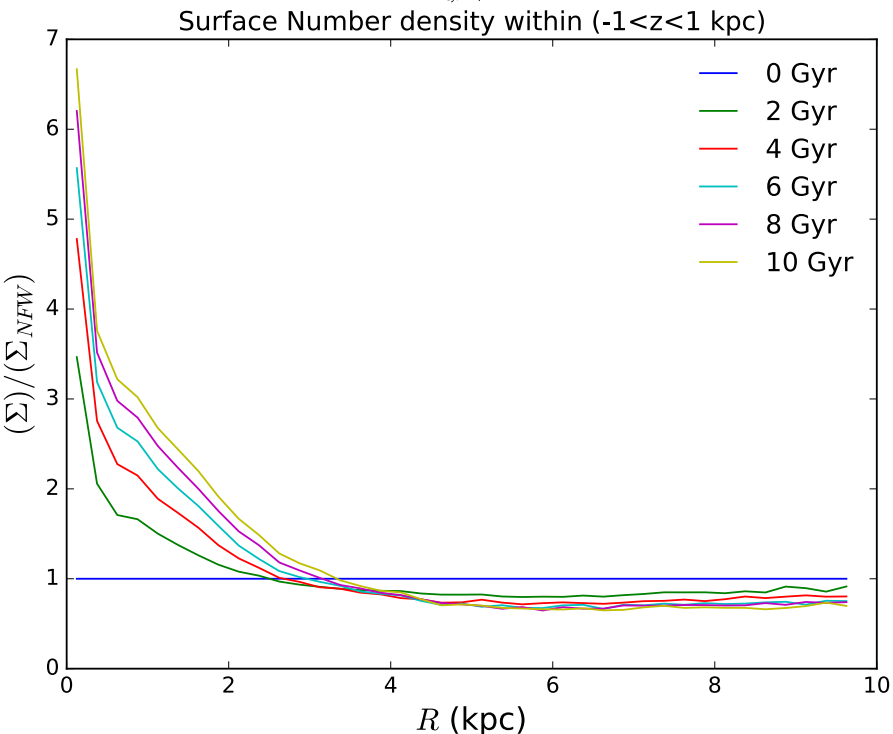
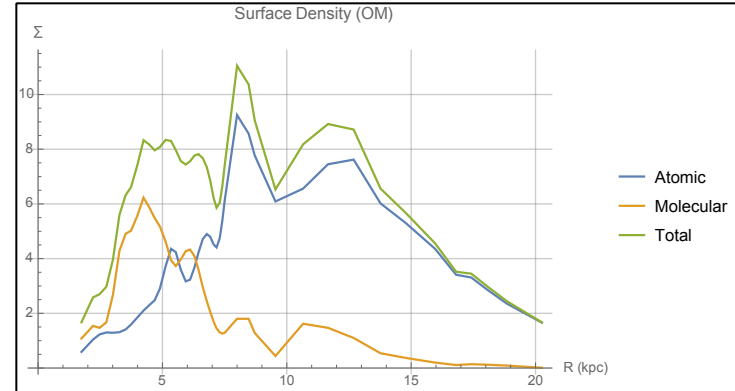
We consider possibly observable effects of asymmetric dark matter (ADM) in neutron stars. Since dark matter does not self-annihilate in the ADM scenario, dark matter accumulates in neutron stars, eventually reaching the Chandrasekhar limit and forming a black hole. We focus on the case of scalar ADM, where the constraints from Bose-Einstein condensation and subsequent black hole formation are most severe due to the absence of Fermi degeneracy pressure. We also note that in some portions of this constrained parameter space, nontrivial effects from Hawking radiation can modify our limits. We find that for scalar ADM with mass between 5 MeV and 13 GeV, the constraint from nearby neutron stars on the scattering cross section with neutrons ranges from $\sigma_n \approx 10^{-45} \text{ cm}^2$ to 10^{-47} cm^2 .

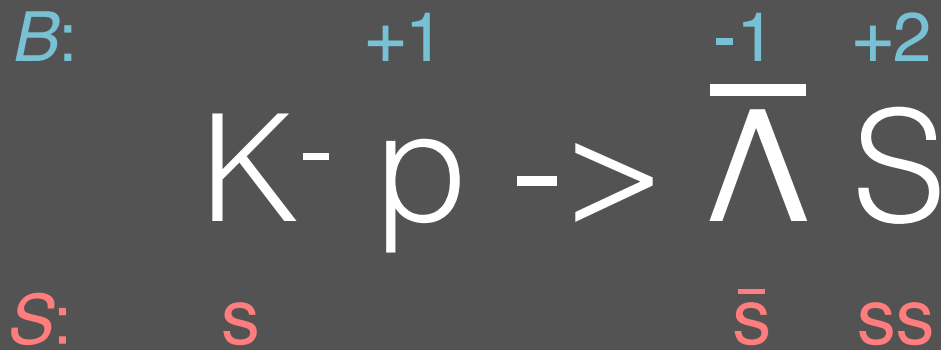
Sensitivity to gas density structure

Galpy Baryons



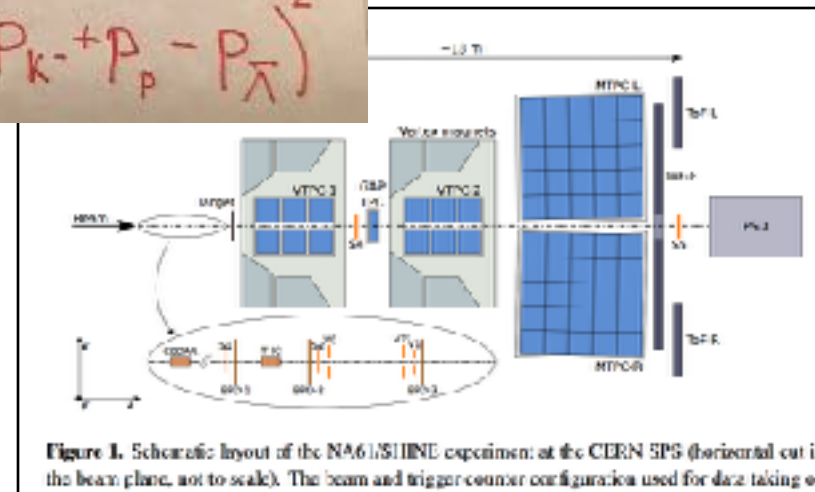
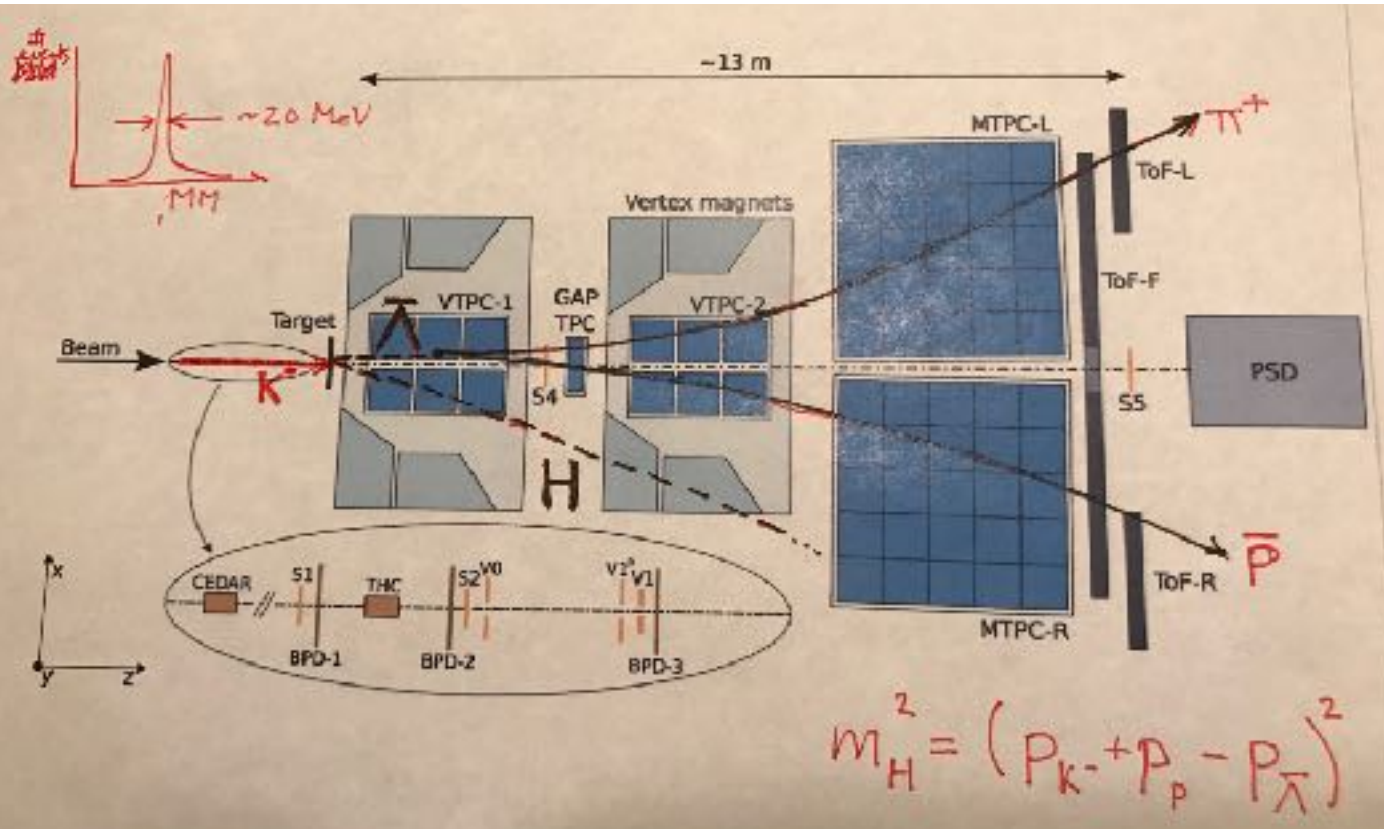
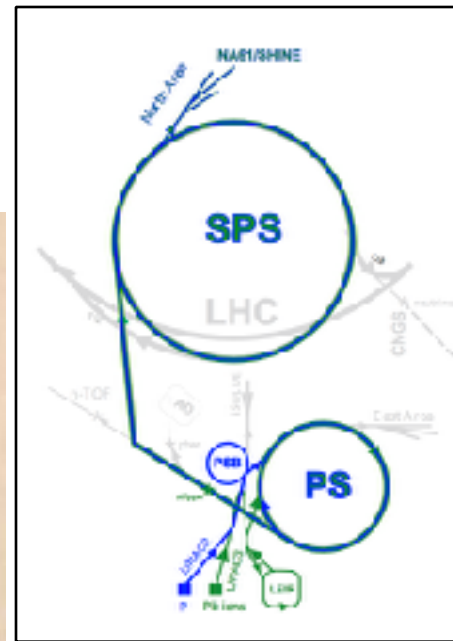
Olling-Merrifield





- $\overline{\Lambda}$ is a gold-plated signature : $\overline{\Lambda} \rightarrow \pi^+ \overline{p}$
 - Easy to ID & reconstruct 4-momentum
 - $c\tau = 8 \text{ cm}$ all $\overline{\Lambda}$ are ID'd
- S: undetected, but 4 momentum determined
 - $p_S = p_K + p_p - p_{\overline{\Lambda}}$
 - NA61: est. $\sim 20 \text{ MeV}$ accuracy on “missing-mass” of S
 - For $p_{\text{beam}} < 5.35 \text{ GeV}/c$, no conventional source of $\overline{\Lambda}$'s
- NA61: 9 GeV/c K- beam, need trigger to

NA61/SHINE



NA61

- Trigger rate ~ 100 Hz $\Rightarrow 10^7$ events per day
 - GEANT: $\sim 0.5\%$ K^0 n + neutrals \Rightarrow must refine trigger
- Schedule mostly fixed till shutdown in 2018; restarts 2020.
- *? short K-p run at 9 GeV/c before shutdown, to evaluate rejection efficiency and background?*
- Maybe 9 GeV/c beam is ok! \Rightarrow longer run in 2020...

Background to $K^- p \rightarrow \bar{\Lambda} R_H$

- $K^- p \rightarrow K^0 n + \text{neutrals}$

↪ $\pi^+ \pi^-$

DANGER: mis-ID π^- as \bar{p} & interpret $n + \text{neutrals}$ as **H**.

- NA61: *good rejection of K^0 faking $\bar{\Lambda}$*
 - ToF, dEdX, kinematic cuts to reject in dangerous regions
 - GEANT sims running to quantify...

“Armenteros Podolanski” plot

