

# The shape of New Physics in $B$ -meson decays

**J. Martin Camalich**



**VUB** Workshop on flavour anomalies

March 29th 2018

# Outline

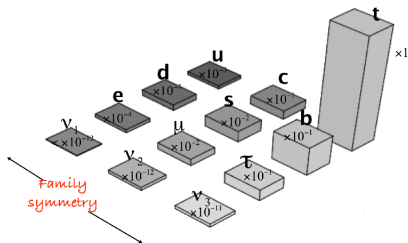
- 1 Flavor in Particle Physics
- 2 LUV in  $b \rightarrow c\tau\nu$  transitions
- 3 LUV  $b \rightarrow s\ell\ell$  transitions
- 4 Conclusions

# The flavor puzzle of the SM

- The matter content of the SM comes in **3 almost identical generations**
  - Identical:** Same quantum numbers under  $SU(3)_c \times SU(2)_L \times U(1)_Y$
  - Almost:** Symmetry broken by Yukawa couplings!

$$-\mathcal{L}_Y = \bar{q}_L Y_d d_R H + \bar{q}_L Y_u u_R \tilde{H} + \bar{\ell}_L Y_e e_R H + h.c.$$

- Masses**  $M_u = L_u Y_u R_u^\dagger$ ,  $M_d = L_d Y_d R_d^\dagger$  and **mixing**,  $V_{CKM} = L_u L_d^\dagger$  are **hierarchical!**



**CKM**

$$|V| = \begin{matrix} u \\ c \\ t \end{matrix} \begin{bmatrix} \text{large} & \text{small} & \cdot \\ \text{small} & \text{large} & \text{very small} \\ \cdot & \text{very small} & \text{large} \end{bmatrix} \begin{matrix} d \\ s \\ b \end{matrix}$$

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  - ▶ **Almost:** Symmetry broken by Yukawa couplings!

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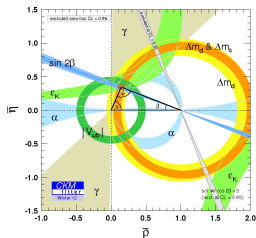
- **Rotations** are “not physical” in the SM ,only **masses** and  $V_{CKM}$

## Where does flavor come from?

- ▶ **Geometrical:** *à la* Warped x-dimensions  $\Lambda_{NP} \sim \Lambda_{\text{Natural}}$  Randall&Sundrum'98
- ▶ **Dynamical:** Horizontal symmetries  $\Lambda_{NP} \gg \Lambda_{\text{Natural}}$  Froggatt&Nielsen '78
- ▶ ...

# Flavor BSM phenomenology

- **Masses**  $\Rightarrow$  **9 parameters**
- **Complex** and **Unitary** matrix  $\Rightarrow$  **3 angles** and **1 phase**

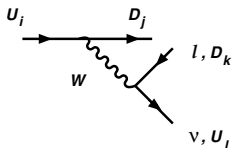


$$V_{CKM} \simeq \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

$$\lambda = 0.2253(7), \quad A = 0.808(22), \\ \bar{\rho} = 0.132(22), \quad \bar{\eta} = 0.341(13)$$

- Few parameters if we count the **thousands of processes** they describe! (**PDG**)

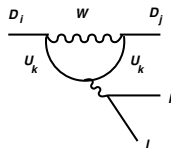
- **CC**  $U_i \rightarrow D_j$ : **Tree level**



- $\mathcal{M} \sim G_F V_{ij} U_{kl}^*$ ,  
 $V_{ij} U_{kl}^*$  can be  $\mathcal{O}(1)$

- In the SM, FCNCs are suppressed w.r.t. CC interactions: “**Rare**” decays!

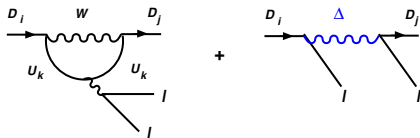
- **FCNC**  $D_i \rightarrow D_j$ : **Loop**



- $\mathcal{M} \sim G_F \sum_k V_{ki} V_{kj}^* \frac{m_k^2}{m_W^2} \frac{\alpha}{4\pi}$ ,  
**GIM** and **loop** suppression

# Flavor BSM phenomenology

- Suppressed processes are very sensitive to **virtual exchange** of new particles!
- **Example: FCNC  $b \rightarrow sll$ :**

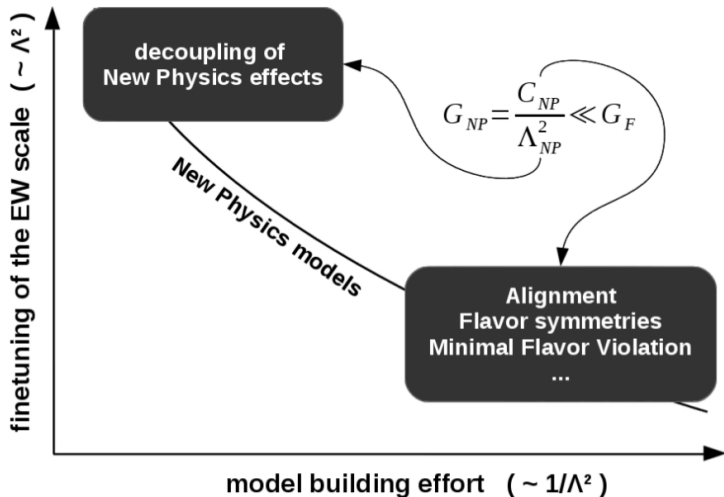


$$\mathcal{M} \sim G_F V_{tb} V_{ts}^* \frac{\alpha}{4\pi} (C^{\text{SM}} + \overbrace{\frac{4\pi}{\alpha} \frac{1}{V_{tb} V_{ts}^*} \frac{v^2}{M^2} g^2}^{\sim 200^2}) \times \langle \bar{s}b \otimes \bar{l}l \rangle$$

Rare  $b$  decays sensitive to  $M \sim 100$  TeV !!

► Neutral meson mixing up to few  $10^4$  TeV!!

# Approaches to the New Physics Flavor Puzzle



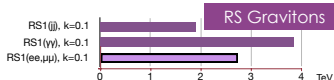
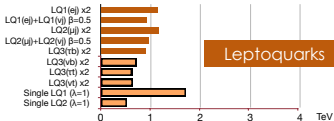


● No **New Physics** at colliders ...

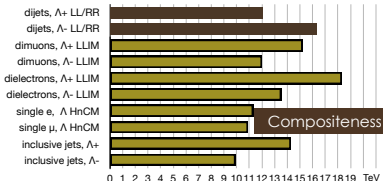
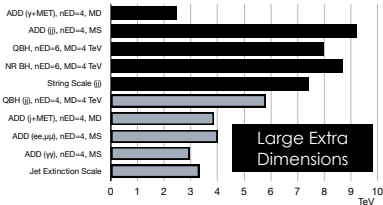
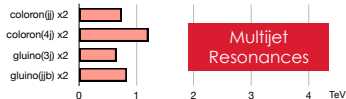
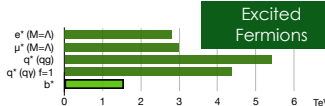
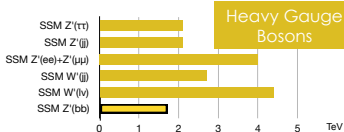
<https://twiki.cern.ch/twiki/bin/view/CMSPublic/>

<https://twiki.cern.ch/twiki/bin/view/AtlasPublic/>

13 TeV 8 TeV

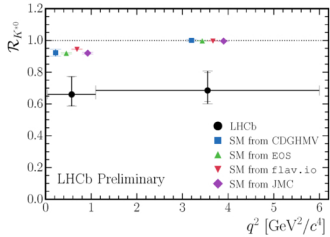
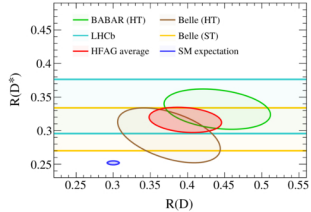
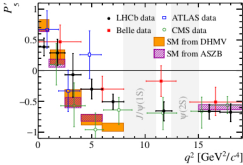
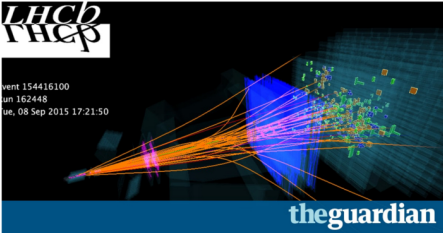


**CMS Preliminary**



CMS Exotica Physics Group Summary – ICHEP, 2016

# B-decay anomalies



**“Evidence” for lepton-universality violation in *b*-quark decays!**

# Lepton-flavor symmetries is a hallmark of the SM

- **Universality of gauge interactions**: Same reps. under  $SU(2)_L \times U(1)_Y$
- **Yukawa interactions** are not lepton-flavor symmetric

**Charged-lepton mass basis**  $\implies U(1)_\tau \times U(1)_\mu \times U(1)_e$  survives!\*

\* Up to tiny effects in charged-lepton processes produced by neutrino masses

## Lepton Flavor Symmetries of the SM

SM is **Charged-lepton flavor universal** up to ...

- 1 **Higgs mediated** (Negligible)
- 2 **Kinematic effects** due to different masses (process dependent)

... and **Charged-lepton flavor symmetric**

- Many experimental tests:

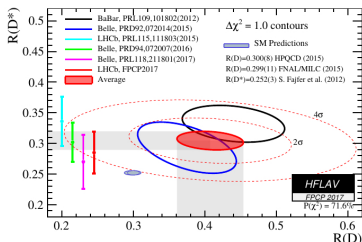
$$\mu \rightarrow e\gamma \quad \mathcal{O}(10^{-13}), \quad Z \rightarrow \ell\ell \quad \mathcal{O}(10^{-4})$$

...

# Lepton-universality violation in $b \rightarrow c\tau\nu$ decays

## Lepton-universality ratios

$$R_{D^{(*)}} = \frac{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \tau^- \bar{\nu})}{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \ell^- \bar{\nu})} \quad \text{where } \ell = e, \mu$$



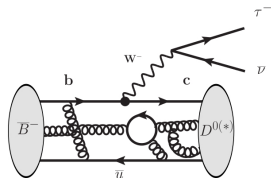
Belle: Hadronic tag, leptonic  $\tau$  Semileptonic tag, leptonic  $\tau$  Hadronic tag, hadronic  $\tau$   
 BaBar: Hadronic tag, leptonic  $\tau$  LHCb: leptonic  $\tau$  hadronic  $\tau$

- Excesses** reported by **3 different experiments** in **2 channels** at  $\sim 4\sigma$ 
  - 15% enhancement of the tau SM amplitude:

LUV in  $b \rightarrow c\tau\nu$

$$\frac{\Lambda}{g} = \frac{v}{\sqrt{|V_{cb}|} \times 0.15} \sim 3 \text{ TeV}$$

# Hadronic uncertainties (Form factors)



- **QCD is lepton universal!**

- ▶ **However:** Important kinematic effects

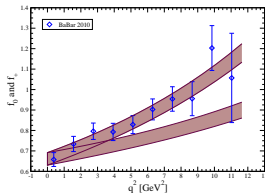
- **Fit Form Factors to experimental  $B \rightarrow D^{(*)}(\mu, e)\nu$  data**

Boyd, Grinstein & Lebed '96, Caprini, Lellouch & Neubert'98

- **Example:  $B \rightarrow D_T \nu$  with LQCD**

$$\langle D(k) | \bar{c} \gamma^\mu b | \bar{B}(\rho) \rangle = (\rho + k)^\mu f_+(q^2) + q^\mu \frac{m_B^2 - m_D^2}{q^2} (f_+(q^2) - f_0(q^2))$$

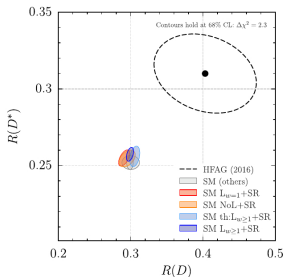
- ▶ Scalar  $f_0(q^2)$  enters rate  $\propto m_\ell^2$
- ▶ **CVC** implies  $f_0(0) = f_+(0)$



Na *et al.* PRD92(2015)no.5,054510 (see also Bailey *et al.* PRD92,034506)

## Hadronic uncertainties (Form factors)

- Upcoming **LQCD** calculation of the  $B \rightarrow D^*$  FFs at **non-zero recoil!**
- **Current prediction relies on HQET** relations including  $\Lambda_{\text{QCD}}/m_{c,b}$  corrections
  - ▶ Contribution to the  $B \rightarrow D^* \tau \nu$  rate of (pseudo)scalar FF is small  $\sim 10\%$ !



Bernlocher *et al.* arXiv: 1703.05330

$$R_{D^*} = 0.257 \pm 0.003$$

Bernlocher *et al.* arXiv: 1703.05330

$$R_{D^*} = 0.260 \pm 0.008$$

Bigi *et al.* arXiv: 1707.09509

**Hadronic uncertainties cannot explain the  $R_{D^{(*)}}$  anomalies**

# EFT of new-physics in $b \rightarrow c\tau\nu$

- Low-energy effective Lagrangian (no RH  $\nu$ )

$$\mathcal{L}_{\text{eff}}^{\ell} = -\frac{G_F V_{cb}}{\sqrt{2}} [(1 + \epsilon_L^{\ell}) \bar{\ell} \gamma_{\mu} (1 - \gamma_5) \nu_{\ell} \cdot \bar{c} \gamma^{\mu} (1 - \gamma_5) b + \epsilon_R^{\ell} \bar{\ell} \gamma_{\mu} (1 - \gamma_5) \nu_{\ell} \bar{c} \gamma^{\mu} (1 + \gamma_5) b \\ + \bar{\ell} (1 - \gamma_5) \nu_{\ell} \cdot \bar{c} [\epsilon_S^{\ell} - \epsilon_P^{\ell} \gamma_5] b + \epsilon_T^{\ell} \bar{\ell} \sigma_{\mu\nu} (1 - \gamma_5) \nu_{\ell} \cdot \bar{c} \sigma^{\mu\nu} (1 - \gamma_5) b] + \text{h.c.},$$

**Wilson coefficients:**  $\epsilon_{\Gamma}$  decouple as  $\sim v^2 / \Lambda_{\text{NP}}^2$

- Matching to high-energy Lagrangian – SMEFT

- ▶ Symmetry relations for  $\epsilon_{\Gamma}$

- ★ In charged-currents  $\epsilon_R^{\ell}$ :

$$\mathcal{O}_{Hud} = \frac{i}{\Lambda_{\text{NP}}^2} \left( \tilde{H}^{\dagger} D_{\mu} H \right) \left( \bar{u}_R \gamma^{\mu} d_R \right)$$



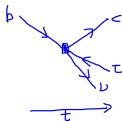
- RHC is lepton universal:  $\epsilon_R^{\ell} \equiv \epsilon_R + \mathcal{O}\left(\frac{v^4}{\Lambda_{\text{NP}}^4}\right) \Rightarrow$  **Cannot explain LUR  $R_{D^{(*)}}$ !**

Down to 4 operators to explain  $R_{D^{*}}$ :  $\epsilon_L, \epsilon_S, \epsilon_P, \epsilon_T$



## The constraint of the $B_C$ -lifetime

- $B \rightarrow D^* \tau \nu$  receives a contribution from  $\epsilon_P$



$$\epsilon_P \langle D^*(k, \epsilon) | \bar{c} \gamma_5 b | \bar{B}(p) \rangle = -\frac{2\epsilon_P m_{D^*}}{m_b + m_c} A_0(q^2) \epsilon^* \cdot q$$

- $B_C \rightarrow \tau \nu$  **also** receives a **helicity-enhanced** contribution from  $\epsilon_P$ !



$$\frac{\text{Br}(B_C^- \rightarrow \tau \bar{\nu}_\tau)}{\text{Br}(B_C^- \rightarrow \tau \bar{\nu}_\tau)^{\text{SM}}} = \left| 1 + \epsilon_L + \frac{m_{B_C}^2}{m_\tau (m_b + m_c)} \epsilon_P \right|^2$$

- Use the lifetime of  $B_C$

- ▶ Very high experimental precision (1.5%):

$$\tau_{B_C} = 0.507(8) \text{ ps}$$

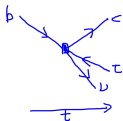
- ▶ **QCD**: “Most of the  $B_C$  lifetime comes from  $\bar{c} \rightarrow \bar{s}$  ( $\sim 65\%$ ) and  $b \rightarrow c$  ( $\sim 30\%$ )”

Bigi PLB371 (1996) 105, Beneke *et al.* PRD53(1996)4991,...

$$\tau_{B_C}^{\text{OPE}} = 0.52_{-0.12}^{+0.18} \text{ ps}$$

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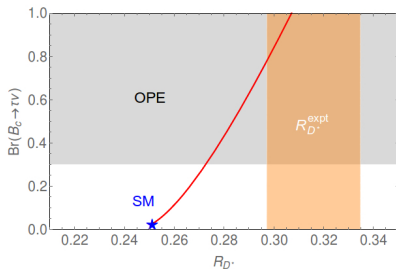


$$\epsilon_P \langle D^*(k, \epsilon) | \bar{c} \gamma_5 b | \bar{B}(p) \rangle = -\frac{2\epsilon_P m_{D^*}}{m_b + m_c} A_0(q^2) \epsilon^* \cdot q$$

- $B_C \rightarrow \tau \nu$  **also** receives a **helicity-enhanced** contribution from  $\epsilon_P$ !



$$\frac{\text{Br}(B_C^- \rightarrow \tau \bar{\nu}_\tau)}{\text{Br}(B_C^- \rightarrow \tau \bar{\nu}_\tau)_{\text{SM}}} = \left| 1 + \epsilon_L + \frac{m_{B_C}^2}{m_\tau(m_b + m_c)} \epsilon_P \right|^2$$



Alonso, Grinstein&JMC, arXiv: 1611.06676

$\tau_{B_C}$  makes **highly implausible ANY**  
 “scalar solution”  
 (e.g. 2HDM) to the  $R_{D^*}$  anomaly!

- Bound  $\text{BR}(B_C \rightarrow \tau \nu)$  from LEP!

Akeroyd&Chen, 1708.04072

# The left-handed operator

- Left-handed  $\epsilon_L = 0.13$ : *Universal* enhancement of the  $b \rightarrow c\tau\nu$  rates by 30%

**SMEFT operators:**  $Q_{\ell q}^{(1)} = \frac{1}{\Lambda^2} (\bar{Q}_L \gamma^\mu Q_L) (\bar{L}_L \gamma_\mu L_L), \quad Q_{\ell q}^{(3)} = \frac{1}{\Lambda^2} (\bar{Q}_L \gamma^\mu \vec{\tau} Q_L) \cdot (\bar{L}_L \gamma_\mu \vec{\tau} L_L)$

- Warning** ☠️ **Radiative LUV contributions in  $\tau$  and  $Z$  decays!**

Ferruglio *et al.* PRL118 (2017), 011801



- Problem with 3<sup>rd</sup> generation:** Non-trivial flavor str. Buttazzo *et al.* arXiv:1706.07808
  - Model dependence:** EFT only gives log parts (mixing)
- It can also solve  $b \rightarrow s\ell\ell$  anomaly! Bhattacharya *et al.* '14, Alonso, JMC & Grinstein. '15, ...
    - Lepton flavor structure:**
      - Large enhancements  $\tilde{C}_{\tau\tau} \gg \tilde{C}_{\mu\mu}$  ruled out by  $B \rightarrow K^{(*)}\nu\nu$  unless  $C_{\tau\tau}^{(1)} \simeq C_{\tau\tau}^{(3)}$
      - Vector Leptoquark**  $U_1^\mu (\bar{3}, 1, 2/3)$  produces this! Alonso, JMC & Grinstein. '15, Barbieri *et al.* '15, ...
  - More on **Dario's talk!**

## Tensor and scalar operators

- ▶ Mixing in  $H^3\psi^2$  operators that **modify Yukawas** [Jenkins et al., arXiv: 1310.4838](#)

- **Tensor  $\epsilon_T = 0.38$**

- ▶ **EW+QED corrections:** Large mixing tensor into scalars

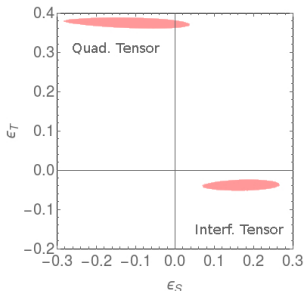
$$\begin{pmatrix} w_{ledq} \\ w_{lequ} \\ w_{lequ}^{(3)} \end{pmatrix}_{(\mu = m_Z)} = \begin{pmatrix} 1.19 & 0. & 0. \\ 0. & 1.20 & -0.185 \\ 0. & -0.00381 & 0.959 \end{pmatrix} \begin{pmatrix} w_{ledq} \\ w_{lequ} \\ w_{lequ}^{(3)} \end{pmatrix}_{(\mu = 1 \text{ TeV})}$$

[Gonzalez-Alonso, JMC & Mimouni arXiv: 1706.00410](#)

- ▶ **No explicit models** that give *only* tensor operators

- **Tensor & Scalar**

- ▶ Fit to current values of  $R_{D^{(*)}}$



- ▶ **New solution:**  $\epsilon_T$  interferes constructively in  $R_{D^*}$

- ★ **Best Fit:**  $\epsilon_S = 0.17$ ,  $\epsilon_T = -0.04$
- ★ **Scalar Leptoquark  $S_1$  ( $\bar{3}, 1, 1/3$ )** produces

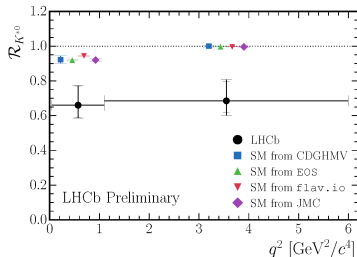
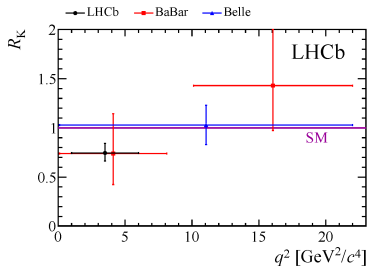
$$\epsilon_T(M) = -\frac{\epsilon_{S_L}(M)}{4}$$

- ★  $\epsilon_P \sim 0.2$  produces  $\text{BR}(B_c \rightarrow \tau\nu) \sim 6\%$

# Lepton-universality violation in $b \rightarrow sll$ decays

## Lepton Universality Ratios

$$R_{K^{(*)}} = \frac{\mathcal{B}(\bar{B} \rightarrow K^{(*)} \mu^+ \mu^-)}{\mathcal{B}(\bar{B} \rightarrow K^{(*)} e^+ e^-)}$$



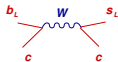
- Skewed  $\mu$ -to- $e$  ratios reported by **LHCb** in **2 channels** at  $\sim 4\sigma$ 
  - ▶ Anomalies in **muonic BRs** and **angular observables**: **Global analyses**  $\sim 5\sigma$
  - ▶ 25% deficit (enhancement) of the SM muon (electron) amplitude:

LUV in  $b \rightarrow sll$

$$\frac{\Lambda}{g} = \frac{v}{\sqrt{|V_{ts}||V_{tb}|} \times \frac{\alpha_{em}}{4\pi}} \sim 30 \text{ TeV}$$

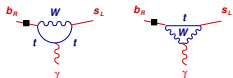
# Effective field theory approach to $b \rightarrow sll$ decays

- **CC** (Fermi theory):

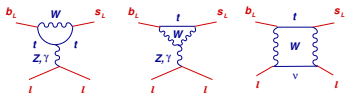

 $\Rightarrow$ 

$$G_F V_{cb} V_{cs}^* C_2 \bar{c}_L \gamma^\mu b_L \bar{s}_L \gamma_\mu c_L$$

- **FCNC**:

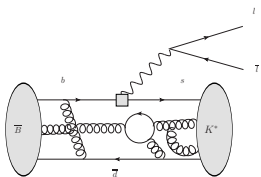

 $\Rightarrow$ 

$$\frac{e}{4\pi^2} G_F V_{tb} V_{ts}^* m_b C_7 \bar{s}_L \sigma_{\mu\nu} b_R F^{\mu\nu}$$


 $\Rightarrow$ 

$$G_F V_{tb} V_{ts}^* \frac{\alpha}{4\pi} C_{9(10)} \bar{s}_L \gamma^\mu b_L \bar{l} \gamma_\mu (\gamma_5) l$$

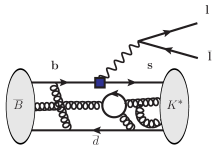
- ▶ **New-Physics** also in  $C_i$  or e.g.  $\mathcal{O}'_i$  obtained  $P_L \rightarrow P_R$  in  $\bar{s}_L b$



- ▶ Light fields active at long distances  
**Nonperturbative QCD!**

- ★ Factorization of scales  $m_b$  vs.  $\Lambda_{\text{QCD}}$   
HQEFT, QCDF, SCET,...

# The complex example: $B \rightarrow K^*(\rightarrow K\pi)ll$



$$\begin{aligned} \frac{d^{(4)}\Gamma}{dq^2 d(\cos\theta_l)d(\cos\theta_k)d\phi} &= \frac{9}{32\pi} (I_1^S \sin^2\theta_k + I_1^C \cos^2\theta_k) \\ &+ (I_2^S \sin^2\theta_k + I_2^C \cos^2\theta_k) \cos 2\theta_l + I_3 \sin^2\theta_k \sin^2\theta_l \cos 2\phi \\ &+ I_4 \sin 2\theta_k \sin 2\theta_l \cos\phi + I_5 \sin 2\theta_k \sin\theta_l \cos\phi + I_6 \sin^2\theta_k \cos\theta_l \\ &+ I_7 \sin 2\theta_k \sin\theta_l \sin\phi + I_8 \sin 2\theta_k \sin 2\theta_l \sin\phi + I_9 \sin^2\theta_k \sin^2\theta_l \sin 2\phi \end{aligned}$$

- Anomalies in the angular observables . . .

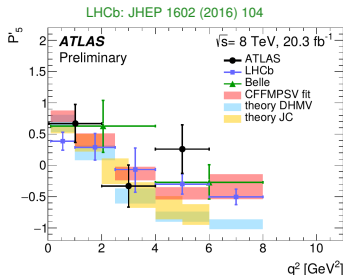
$$P'_5 = \frac{I_5}{2\sqrt{-I_{2S}I_{2C}}}$$

- ▶ Cancel leading theory uncertainties

New physics?

$$\delta C_9^\mu \simeq -1$$

Descotes-Genon *et al.* PRD88,074002



- ▶ Interpretation blurred by hadronic uncertainties



# Anatomy of the amplitude in a nutshell

Jäger and JMC, PRD93 (2016) no.1, 014028

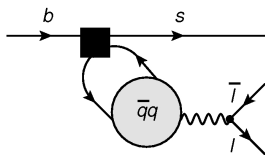
- Helicity amplitudes  $\lambda = \pm 1, 0$

$$H_V(\lambda) = -iN \left\{ \overbrace{C_9 \tilde{V}_{L\lambda} + \frac{m_B^2}{q^2} h_\lambda}^{C_9^{\text{eff}}} - \frac{\hat{m}_b m_B}{q^2} C_7 \tilde{T}_{L\lambda} \right\},$$

$$H_A(\lambda) = -iN C_{10} \tilde{V}_{L\lambda}$$

- Hadronic form factors:** 7 independent  $q^2$ -dependent nonperturbative functions

## “Charm” contribution



$$h_\lambda \propto \int d^4 y e^{iq \cdot y} \langle \bar{K}^* | T \{ J^{\text{em, had}, \mu}(y), \mathcal{O}_{1,2}(0) \} | \bar{B} \rangle$$

- Charm and  $\mathcal{O}_9$  are tied up by renormalization  
**Only  $C_9^{\text{eff}}$  is observable!**

# The lepton-universality ratios...

- **QCD interactions are lepton universal\***

- ▶ \* EM corrections are lepton-dependent but at  $\sim$  % level Bordone et al. EPJC76(2016),8,440

- ... In  $B \rightarrow K \ell \ell$

$$\frac{d\Gamma_K}{dq^2} = \mathcal{N}_K |\vec{k}|^3 f_+(q^2)^2 \left( |C_{10}^\ell + C_{10}^{\prime\ell}|^2 + |C_9^\ell + C_9^{\prime\ell} + 2 \frac{m_b}{m_B + m_K} C_7 \frac{f_T(q^2)}{f_+(q^2)} - 8\pi^2 h_K|^2 \right) + \mathcal{O}\left(\frac{m_\ell^4}{q^4}\right) + \dots$$

- ... in  $B \rightarrow K^* \ell \ell$

$$\frac{d\Gamma_0}{dq^2} = \mathcal{N}_{K^*0} |\vec{k}|^3 V_0(q^2)^2 \left( |C_{10}^\ell - C_{10}^{\prime\ell}|^2 + |C_9^\ell - C_9^{\prime\ell} + \frac{2m_b}{m_B} C_7 \frac{T_0(q^2)}{V_0(q^2)} - 8\pi^2 h_{K^*0}|^2 \right) + \mathcal{O}\left(\frac{m_\ell^2}{q^2}\right)$$

$$\frac{d\Gamma_\perp}{dq^2} = \mathcal{N}_{K^*\perp} |\vec{k}|^2 V_\perp(q^2)^2 \left( |C_{10}^\ell|^2 + |C_9^{\prime\ell}|^2 + |C_{10}^{\prime\ell}|^2 + |C_9^\ell + \frac{2m_b m_B}{q^2} C_7 \frac{T_\perp(q^2)}{V_\perp(q^2)} - 8\pi^2 h_{K^*\perp}|^2 \right) + \mathcal{O}\left(\frac{m_\ell^2}{q^2}\right)$$

Wilson coefficients in the SM

$$C_9^{\text{SM}}(m_b) \simeq -C_{10}^{\text{SM}} = +4.27 \quad C_7^{\text{SM}}(m_b) = -0.333$$

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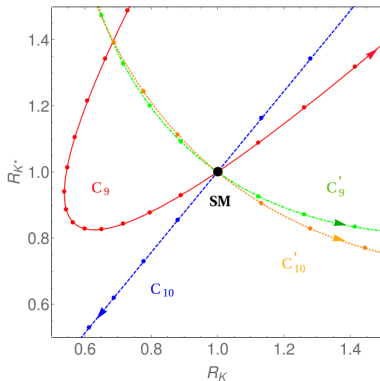
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## Wilson coefficients in the SM

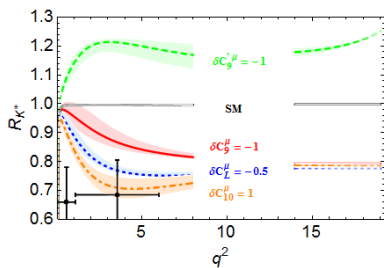
$$C_9^{\text{SM}}(m_b) \simeq -C_{10}^{\text{SM}} = +4.27 \qquad C_7^{\text{SM}}(m_b) = -0.333$$

- **New physics in muons**

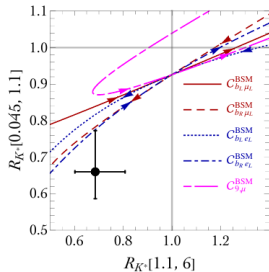


Geng, Grinstein, Jäger, Martin Camalich, Ren, Shi, arXiv: 1704.05446

- Nodes indicate steps of  $\Delta C^\mu = +0.5$ 
  - ▶ **Primed operators**  $C'_{9,10}$ : Monotonically decreasing dependence  $R_{K^*}(R_K)$ !
- **New physics in electrons**  $\sim$  mirror image of above (see D'Amico *et al.* 1704.05438)



Geng, Grinstein, Jäger, Martin Camalich, Ren, Shi, arXiv: 1704.05446



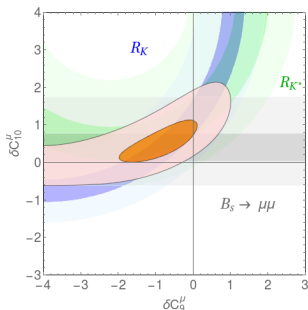
D'Amico *et al.* 1704.05438

**Very clean null-tests of the SM!**

- **Warning** 🚫: Value at ultralow- $q^2$  is difficult to accommodate with UV physics

# Fits with clean observables only

- Assume NP is  $\mu$ -specific



Coeff.	best fit	$\chi^2_{\min}$	$p$ -value	SM exclusion [ $\sigma$ ]	$1\sigma$ range	$3\sigma$ range
$\delta C_9^\mu$	-1.64	5.65	0.130	3.87	[-2.31, -1.12]	[<-4, -0.31]
$\delta C_{10}^\mu$	0.91	4.98	0.173	3.96	[0.66, 1.18]	[0.20, 1.85]
$\delta C_L^\mu$	-0.61	3.36	0.339	4.16	[-0.78, -0.46]	[-1.14, -0.16]
Coeff.	best fit	$\chi^2_{\min}$	$p$ -value	SM exclusion [ $\sigma$ ]	parameter ranges	
$(\delta C_9^\mu, \delta C_{10}^\mu)$	(-0.76, 0.54)	3.31	0.191	3.76	$C_9^\mu \in [-1.50, -0.16]$	$C_{10}^\mu \in [0.18, 0.92]$

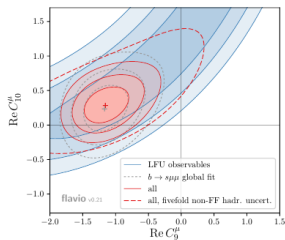
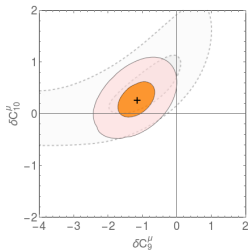
- Deviation of the SM:  $p$ -value of  $3.7 \times 10^{-4}$  ( $3.6\sigma$ )
- Best fit suggests a leptonic left-handed scenario  $\delta C_L^\mu$

- **Include 70-100 observables**

Coeff.	best fit	$\chi^2_{\min}$	$p$ -value	SM exclusion [ $\sigma$ ]	$1\sigma$ range	$3\sigma$ range
$\delta C_9^\mu$	-1.37	61.98 [64 dof]	0.548	4.37	[-1.70, -1.03]	[-2.41, -0.41]
$\delta C_{10}^\mu$	0.60	71.72 [64 dof]	0.237	3.06	[0.40, 0.82]	[-0.01, 1.28]
$\delta C_L^\mu$	-0.59	63.62 [64 dof]	0.490	4.18	[-0.74, -0.44]	[-1.05, -0.16]
Coeff.	best fit	$\chi^2_{\min}$	$p$ -value	SM exclusion [ $\sigma$ ]	parameter ranges	
$(\delta C_9^\mu, \delta C_{10}^\mu)$	(-1.15, 0.28)	60.33 [63 dof]	0.572	4.17	$C_9^\mu \in [-1.54, -0.81]$	$C_{10}^\mu \in [0.06, 0.50]$

- $C_9$  in global fits is subject to hadronic uncertainties!

- ▶ Results in the  $(\delta C_9^\mu, \delta C_{10}^\mu)$  plane



Altmannshofer *et al.* arXiv:1704.05435

- More on **Nazila's talk!**

# Top-down approach: Knocking on the flavor-puzzle's door?

## ● Gauged flavor symmetries Altmannshofer *et al.* arXiv:1403.1269, Alonso *et al.* arXiv:1704.08158,...

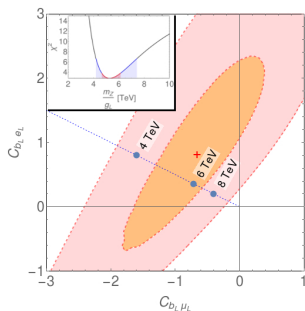
### ▶ Gauge $SU(3)_L \times SU(3)_R$ Cline&JMC arXiv:1706.08510

- ★ Dynamically generate Flavor and UV consistent (no Landau poles)
- ★ *Baroque* field content for gauge anomalies/SSB-structure

### ▶ Couplings to the leptons

$$V_l^L T^8 V_l^{L\dagger} \cong \frac{1}{2\sqrt{3}} \begin{pmatrix} 1 & \epsilon_1 & \epsilon_2 \\ \epsilon_1^* & -2 & \epsilon_3 \\ \epsilon_2^* & \epsilon_3^* & 1 \end{pmatrix}$$

**Coupling to both electrons and muons!**



- Couplings to electrons opens up much more stringent phenomenology!



## Efforts from the TH community

- “Instant” workshop at CERN last May

### Instant workshop on B meson anomalies

 17 May 2017, 09:00 → 19 May 2017, 16:30 Europe/Zurich

 4-3-006 - TH Conference Room (CERN)

 Jorge Martin Camalich (CERN) , Jure Zupan (University of Cincinnati) , Marco Nardecchia (CERN)

**Description** In light of recent anomalies in B physics there is an increased interest in the theory community on its implications. As a quick response we are organizing an “Instant workshop on B meson anomalies” at CERN from May 17-May 19 2017.

- Check recordings @ <https://indico.cern.ch/event/633880/>

**CERN-TH Institute** programmed for the next year

“From flavor anomalies to direct discovery of New Physics”

*Oct. 22nd to Nov. 2nd 2018 (tentative)*



“Extraordinary claims require Extraordinary evidence”  
– C. Sagan

## • Looking (indirectly) for the Extraordinary Evidence!

- 1 **More data on the measured channels!**  $\Rightarrow 5\sigma$  in a single measurement!
  - ★  $5\sigma$ -evidence should be soon available in single measurements of  $R_{D^{(*)}}$  and  $R_{K^{(*)}}$
- 2 **Measure NEW channels:** Different channels, different backgrounds
  - ★ Many new channels will be soon available at **LHCb**:  $R_{J/\psi}$ ,  $R_{\Lambda_c}$ , ...
- 3 **Measure NEW observables:** Consistency, different sensitivities to NP
  - ★ New observables will be soon available at **Belle II**: Polarizations, distributions, ...
- 4 **Look ELSEWHERE:** Low energy & **High  $p_T$  signatures**
  - ★ New flavor sectors at reach:  $b \rightarrow s\nu\nu$ ,  $b \rightarrow d\mu\mu$ ,  $b \rightarrow s\tau\tau$

**Indirect discovery or Ruling-out of the  $B$ -anomalies in 5-10 yrs!**

**IF true  $\implies$  THEN Discovery case** for next-generation collider (FCC-ee, FCC-hh, ...)

## Searches for $B_c \rightarrow \tau \nu$ at LEP

- BR( $B_c \rightarrow \tau \nu$ ) measured in a  $e^+ e^-$  collider at the Z pole Akeroyd&Chen, 1708.04072
  - Searches of  $B^- \rightarrow \tau^- \nu$  above  $B_c \bar{B}_c$  threshold really measure

Mangano&Slabospitsky, PLB410(1997)299

$$\underbrace{\text{BR}}_{\text{LEP}}^{\text{eff}} = \underbrace{\text{BR}(B \rightarrow \tau \nu)}_{\text{Belle \& BaBar}} + \underbrace{\frac{f_c}{f_u}}_{\text{TH. input}} \text{BR}(B_c \rightarrow \tau \nu)$$

- $B_c$  contribution suppressed by  $f_c/f_u \sim 10^{-3}-10^{-2}$  but enhanced by  $\frac{|V_{cb}|^2}{|V_{ub}|^2} \frac{f_{B_c}^2}{f_B^2} \sim 700$
- $f_c/f_u$ : Fraction of hadronization into  $B_c$  over  $B$ 
  - Traded by experimental data and **computable TH. input**

$$R_\ell = \frac{f_c}{f_u} \frac{\text{BR}(B_c \rightarrow J/\psi \mu \nu)}{B \rightarrow J/\psi K}$$

- $R_\ell$  measured by **CDF** and reconstructed from **LHCb** data

# Searches for $B_c \rightarrow \tau \nu$ at LEP

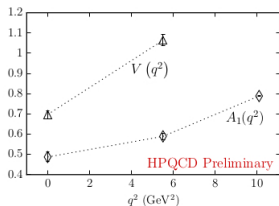
- Model calculations predict  $\text{BR}(B_c \rightarrow J/\psi \mu \nu) \in 1 - 7\%$

	pQCD	WSL [9]	EFG [7]	ISK [6]	HNV [5]	DV [4]
$V^{B_c \rightarrow J/\psi}$	0.42	0.74	0.49	0.83	0.61	0.91
$A_0^{B_c \rightarrow J/\psi}$	0.59	0.53	0.40	0.57	0.45	0.58
$A_1^{B_c \rightarrow J/\psi}$	0.46	0.50	0.50	0.56	0.49	0.63
$A_2^{B_c \rightarrow J/\psi}$	0.64	0.44	0.73	0.54	0.56	0.74

Wang, Fang&Xiao, arXiv: 1212.5903

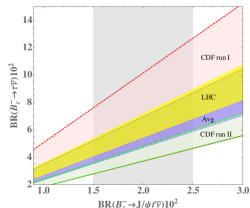
- Ongoing efforts in LQCD!

- Preliminary results to select models



HPQCD Collaboration, PoS LATTICE2016 (2016) 281

- Constrains  $\text{BR}(B_c \rightarrow \tau \nu) < 10\%$



Akeroyd&Chen, 1708.04072

## Adding new channels: $B_c \rightarrow J/\psi \tau \nu$

$$R_{J/\psi}^{\text{LHCb}} = 0.71 \pm 0.17 \pm 0.18$$

Greg's talk yesterday

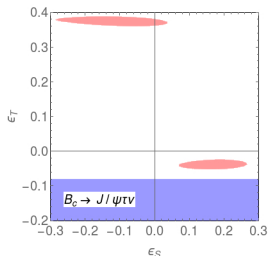
- Comparison with SM **NOW** is subtle because of **model dependence**

Mode	This paper	[8, 30]	[11]	[15]	[16]	[31]	[32]
$B_c^- \rightarrow J/\psi \ell \nu$	$6.7^{+2.1+1.0+0.9}_{-1.2-0.4-0.6}$	1.9	2.37	1.5	1.49	1.20	2.07
$B_c^- \rightarrow J/\psi \tau \nu$	$0.52^{+0.16+0.08+0.08}_{-0.09-0.03-0.05}$	0.48	0.65	0.4	0.37	0.34	0.49

$$R_{J/\psi}^{\text{SM}^*} \sim 0.24 - 0.29$$

Qiao&Zhu, 1208.5916

- Goes in the *right* direction of NP but effect is **large**



- For the LH solution one predicts

$$R_{J/\psi}^{\text{LH}^*} \sim 0.35 - 0.4$$

(see also Watanabe, arXiv: 1709.08644)

- Besides more data, **LQCD input urgently needed!**