

Astrophysical uncertainties on Dark Matter Capture by the Sun and the connection with Cosmological Simulations Supervisors: Emmanuel Nezri Vincent Bertin





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Cosmological simulations

Stars

Cosmological hydrodynamical Zoom-in Simulations



Stars Dark matter DM capture by the Sun $\frac{f(u)}{d}w\Omega(w)$ dC u_m du^{J} dVCapture rate in the Sun Astrophysics and Particle Physics come together (Press and Spergel 1985) (A. Gould 1987) (G. Jungmann, M. Kamionkowski, K. Griest 1996) (Ling 2010) (Choi et al 2014) (Garani et al 2017)

DM capture by the Sun







Press and Spergel 1985 (A. Gould 1987) (G. Jungmann, M. Kamionkowski, K. Griest 1996) (Ling 2010) (Choi et al 2014) (Garani et al 2017) Main source of uncertainties:

- Local Dark Matter velocity distribution
- Local Dark Matter density



Is there equilibrium between the capture and annihilation?



The Standard Halo Model

Uncertainties in Dark Matter distribution features

Standard Assumptions:

- Mass profile = NFW (DMO motivated)
- Phase space distribution = Maxwellian Distribution



...Often used as input for dark matter detection limits and theoretical predictions but not really agreeing with galactic dynamics and/or cosmological simulations...





Equilibrium Are ρ_{DM} and f(v) constant in time throughout the life of the solar system?

Time evolution in the solar system lifetime



Uncertainties in ρ_{DM}

From Simulations



Uncertainties in f(v)

Velocity distribution

• Fitting the velocity distribution: Fitting on the simulation data fails to reproduce the tail and/or the hat of the distribution

Generalized Gaussian

$$f(v) = \frac{1}{N(v_0, \alpha)} e^{-((v-\mu)^2/v_0^2)^c}$$

Generalized Maxwellian



Tsallis

$$f(\vec{v}) = \frac{1}{N(v_0, q)} \left(1 - (1 - q)\frac{\vec{v}^2}{v_0^2}\right)^{q/(1 - q)}$$

[arxiv:9511007],[arxiv:0909.2028]

Having several MW-like simulations would allow to do statistics to understand the behavior of $f(\boldsymbol{\nu})$



Velocity distribution

- Fitting the velocity distribution: Fitting on the simulation data fails to reproduce the $\frac{2}{5}^{0.002}$ tail and/or the hat of the distribution
- Using directly the Simulation data:

Each galaxy has different dynamics, therefore extrapolating the results of one to the other has no real meaning



Eddington inversion Dynamically self consistent approach Velocity distribution inferred from gravitational potential



Eddington inversion



Predictions from the Eddington method as studied by Lacroix et al. (Lacroix et al and from Binney - Tremaine) of f(v)

VS

 $\begin{aligned} \text{fully consistent objects build in a Zoom-in Cosmological Simulation.} \\ f(\vec{r}, \vec{v}) &= f(\mathcal{E}, L) \end{aligned} \\ f(\vec{r}, \vec{v}) &= f(\mathcal{E}, L) \end{aligned} \\ \begin{aligned} \mathcal{L} &= \Psi(r) - \frac{v^2}{2} \end{aligned} \\ \Psi &= \Phi(r) - \Phi(r_{max}) \end{aligned} \qquad f(\mathcal{E}) &= \frac{1}{\sqrt{8}\pi^2} \left(\frac{1}{\sqrt{\mathcal{E}}} \left[\frac{d\rho}{\Psi} \right]_{\Psi=0} + \int_0^{\mathcal{E}} \frac{d\Psi}{\sqrt{\mathcal{E} - \Psi}} \frac{d^3\rho}{d\Psi^3} \right) \end{aligned}$

Statistical moments of the velocity distributions

Direct and Indirect DM detection relevant quantities







Nunez-Castineyra, Bertin, Nezri in prep.

Eddington inversion



Eddington inversions on dynamical Models of the Milky way



Nunez-Castineyra, Bertin, Nezri in prep.

Summary

 ρ_{local} fluctuations in time fall within the uncertainty bands, f(v) fluctuations are not dramatic. (beware of recent merger histories.)

The effect of changing f(v) is of ~5-10 % while ρ_{local} uncertaities are higher thatn 30-40%. this could have a much bigger effect on the detection limits. (waiting for improvements from GAIA data)

• Extrapolating directly the DM features in MW-like Simulations (= first step)

 Comfronting Simulations and dynamical methods is much more consistent (= educated use of simulations.)







$$\Psi = \Phi(r) - \Phi(r_{max})$$



(Blue) The escape velocity computed from the potential.

(Red) the mean velocity of the fastest DM particles present in slices of θ and φ of a shell with radius r.

Shape of the potential



