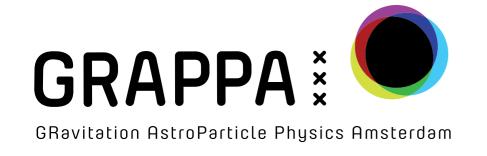
Dark Ghosts
Brussels
13 November 2018

Modeling dark matter substructure and annihilation boost for indirect searches

Shin'ichiro Ando

U. Amsterdam / U. Tokyo





Why subhalos?

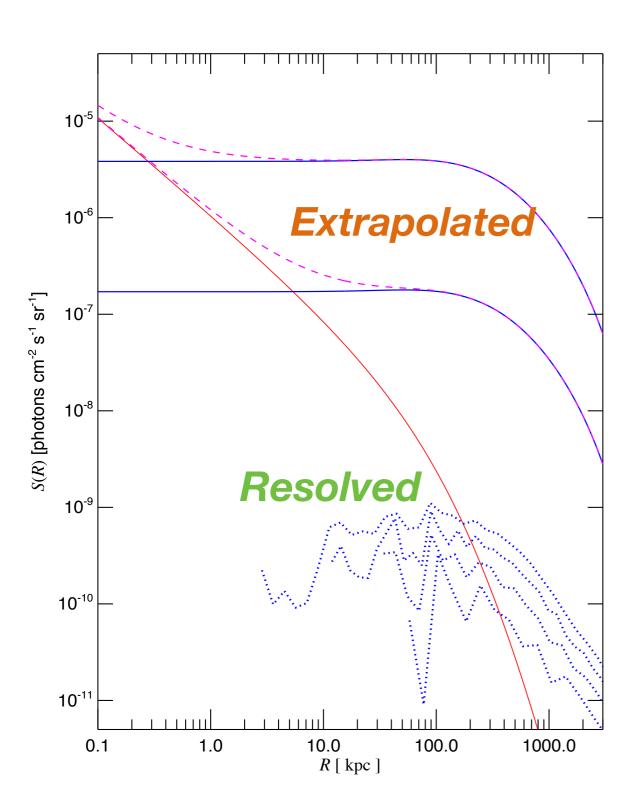
- Dark matter halos contain lots of subhalos (as CDM predicts), so all the extragalactic halos are subject to the substructure boost of dark matter annihilation
- Dwarfs are subhalos (but condition of forming galaxies in them is not yet well understood)
- WIMP accumulation in the Sun may be affected
- Hence subhalos are relevant for all the DM annihilation searches except for Galactic center region

Annihilation boost

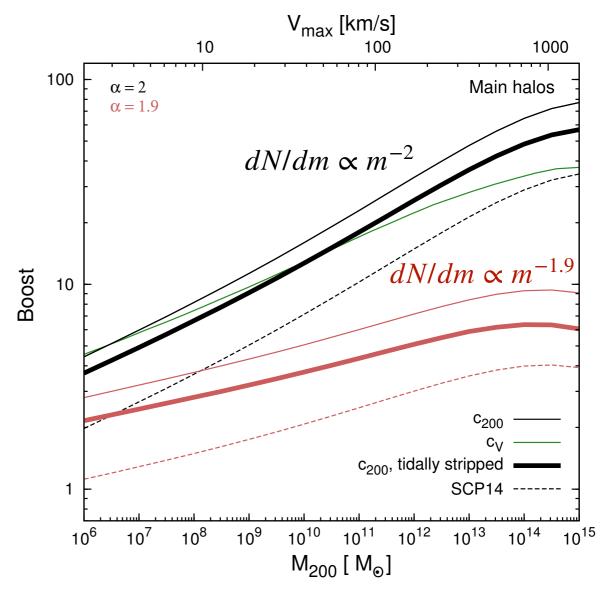
$$L(M) = [1 + B_{\rm sh}(M)]L_{\rm host}$$

$$B_{
m sh}(M) = rac{1}{L_{
m host}} \int rac{dN_{
m sh}}{dm} L_{
m sh}(m)$$

How uncertain is the boost?



Gao et al., Mon. Not. R. Astron. Soc. 419, 1721 (2012)



Moliné et al., Mon. Not. R. Astron. Soc. 466, 4974 (2017)

- Very uncertain, of which we don't even have good sense
- No way that it can be solved with numerical simulations

Analytic model of subhalo evolution

- Complementary to numerical simulations
- Light, flexible, and versatile
- Can cover large range for halo masses (micro-halos to clusters) and redshifts (z ~ 10 to 0)
- Physics-based extrapolation
- Reliable if it is tested compared with simulations at resolved scales

Analytic model: Recipe

Structures start to form

Initial condition:
Primordial power spectrum



Smaller halos merge and accrete to form larger ones

Extended Press-Schechter formalism



Subhalos experience mass loss

Modeling for tidal stripping and mass-loss rate

$$L_{\rm sh}^{\rm total}(M,z) = \int dm \frac{dN}{dm} L_{\rm sh}(m)$$

Conventional formula

$$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Accretion

$$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Accretion

$$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Number of subhalos accreted at z_{acc} with mass m_{acc}

Accretion

Evolution

$$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Number of subhalos accreted at z_{acc} with mass m_{acc}

Accretion

Evolution

$$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Number of subhalos accreted at z_{acc} with mass m_{acc}

Luminosity of

the subhalo at z

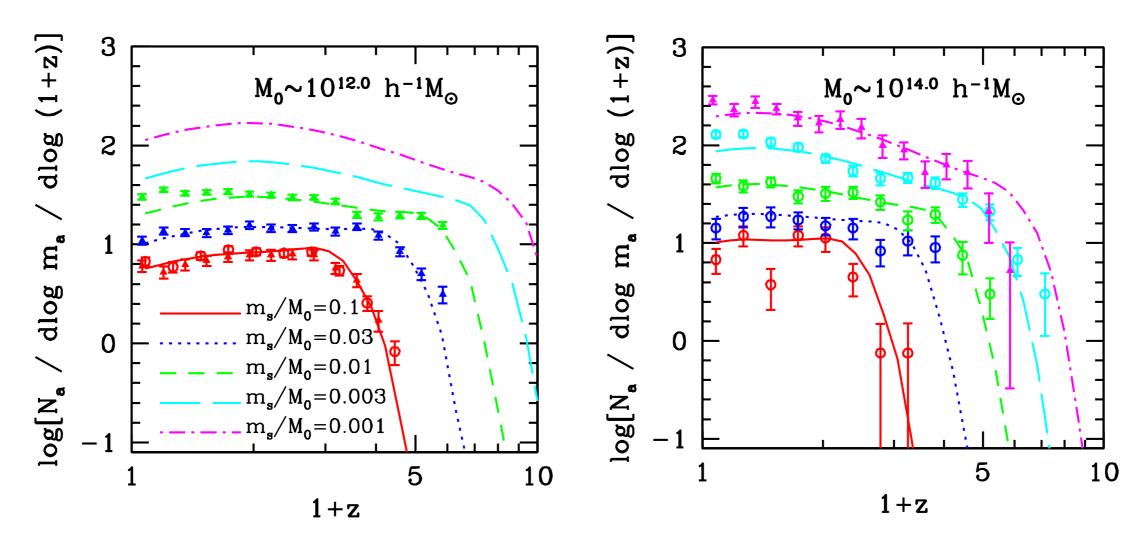
Halo formation and accretion history

 Based on spherical collapse model and extended Press-Schechter formalism (Yang et al. 2011)

$$\frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} \propto \frac{1}{\sqrt{2\pi}} \frac{\delta(z_{\rm acc}) - \delta_M}{(\sigma^2(m_{\rm acc}) - \sigma_M^2)^{3/2}} \exp\left[-\frac{(\delta(z_{\rm acc}) - \delta_M)^2}{2(\sigma^2(m_{\rm acc}) - \sigma_M^2)}\right]$$

 Primordial power spectrum + cutoff scale will change rms over-density σ(M)

Subhalo accretion rate



Yang et al., Astrophys. J. 741, 13, (2011)

Infall distribution of subhalos:

Extended Press-Schechter formalism

$$\frac{d^2N}{d\ln m_a d\ln(1+z_a)}$$

Accretion

Evolution

$$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Number of subhalos accreted at z_{acc} with mass m_{acc}

Luminosity of

the subhalo at z

$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$

Number of subhalos accreted at $z_{\rm acc}$ with mass $m_{\rm acc}$

Evolution

$$L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Luminosity of the subhalo at z

$$L_{\rm sh}(z\,|\,m_{\rm acc},z_{\rm acc}) \propto \rho_s^2(z\,|\,m_{\rm acc},z_{\rm acc}) r_s^3(z\,|\,m_{\rm acc},z_{\rm acc}) \left\{ \, 1 - \frac{1}{[1 + r_t(z\,|\,m_{\rm acc},z_{\rm acc})/r_s(z\,|\,m_{\rm acc},z_{\rm acc})]^3} \, \right\}$$

Evolution

$$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Number of subhalos accreted at $z_{\rm acc}$ with mass $m_{\rm acc}$

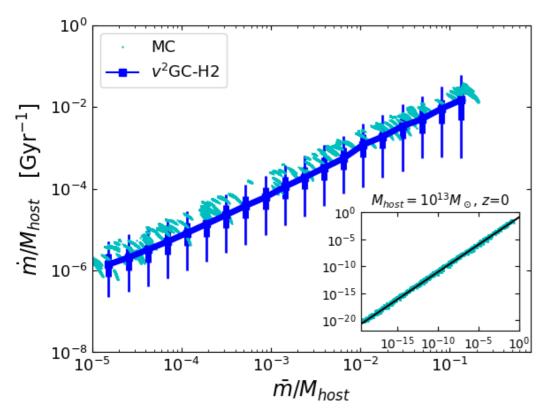
Luminosity of the subhalo at z

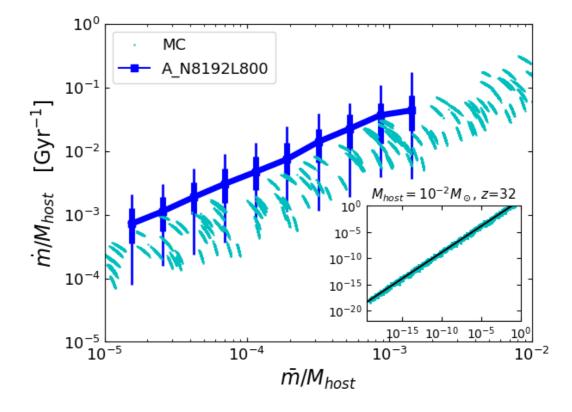
$$L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc}) \propto \rho_s^2(z \mid m_{\rm acc}, z_{\rm acc}) r_s^3(z \mid m_{\rm acc}, z_{\rm acc}) \left\{ 1 - \frac{1}{[1 + r_t(z \mid m_{\rm acc}, z_{\rm acc})/r_s(z \mid m_{\rm acc}, z_{\rm acc})]^3} \right\}$$

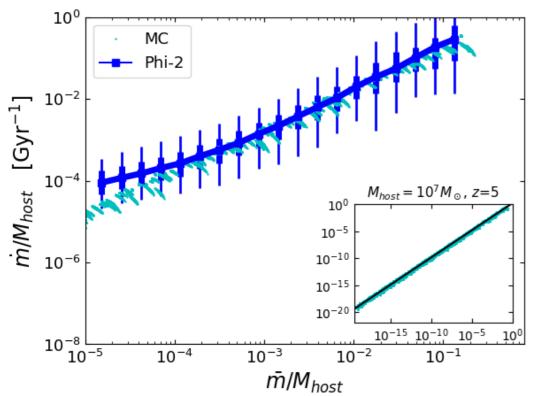
Parameters subhalo density profile after tidal mass loss

Subhalo mass loss

Hiroshima, Ando, Ishiyama, Phys. Rev. D 97, 123002 (2018)



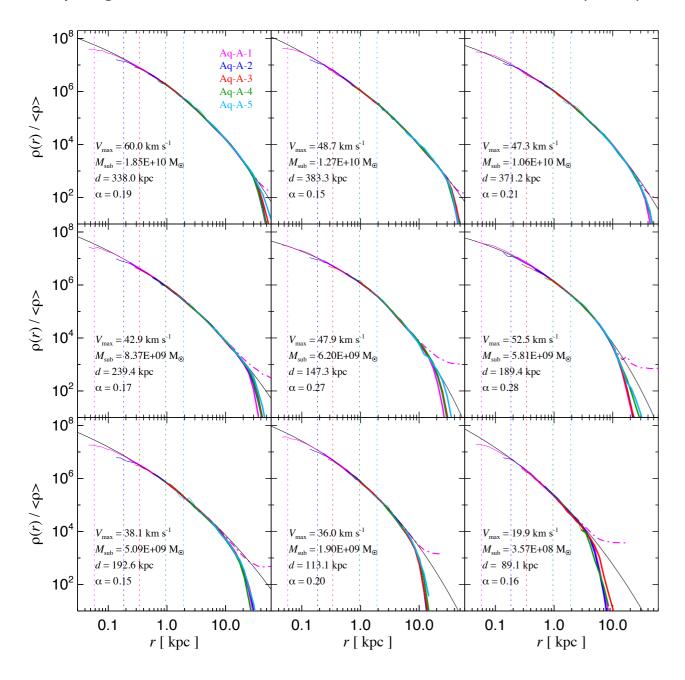




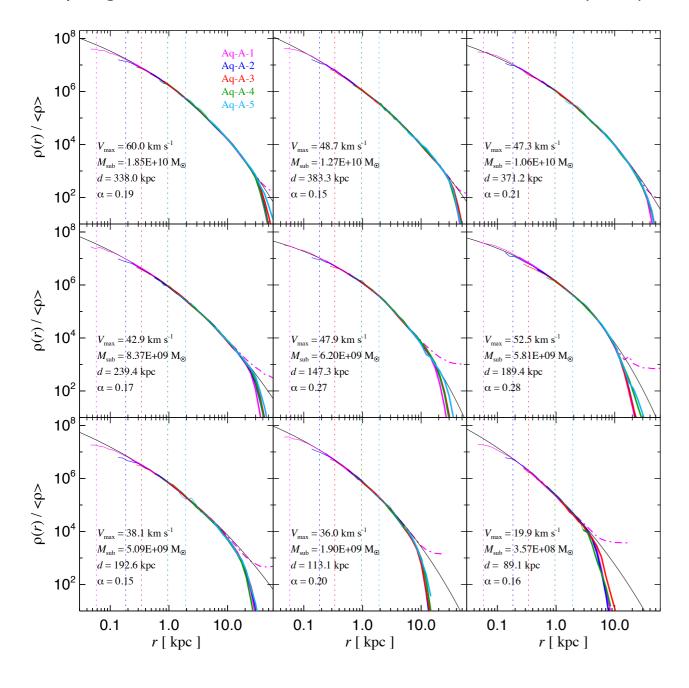
10° r

- Monte Carlo approach following Jiang & van den Bosch (2016)
 - Determine orbital energy and angular momentum
 - Assume the subhalo loses all the masses outside of its tidal radius instantaneously at its peri-center passage
- Mass-loss rate follows power law for wide range of m/M

Springel et al., Mon. Not. R. Astron. Soc. 391, 1685, (2008)

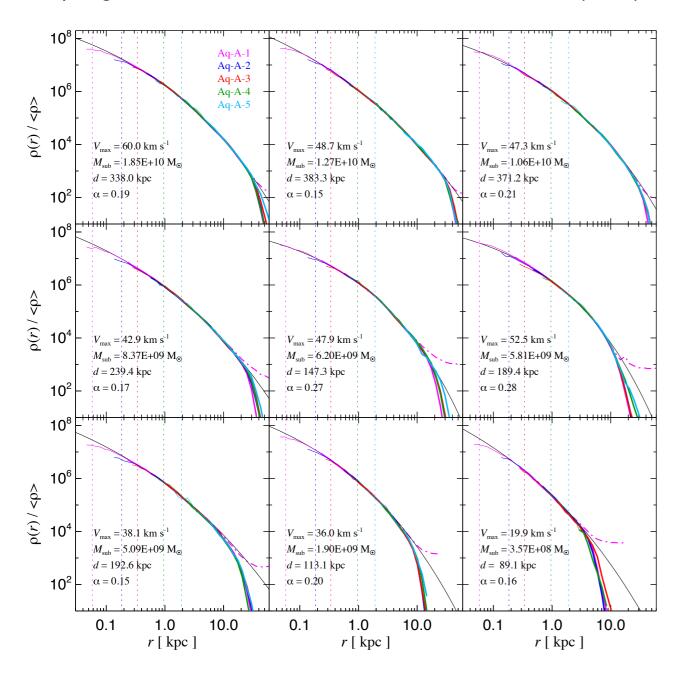


Springel et al., Mon. Not. R. Astron. Soc. 391, 1685, (2008)



Procedure

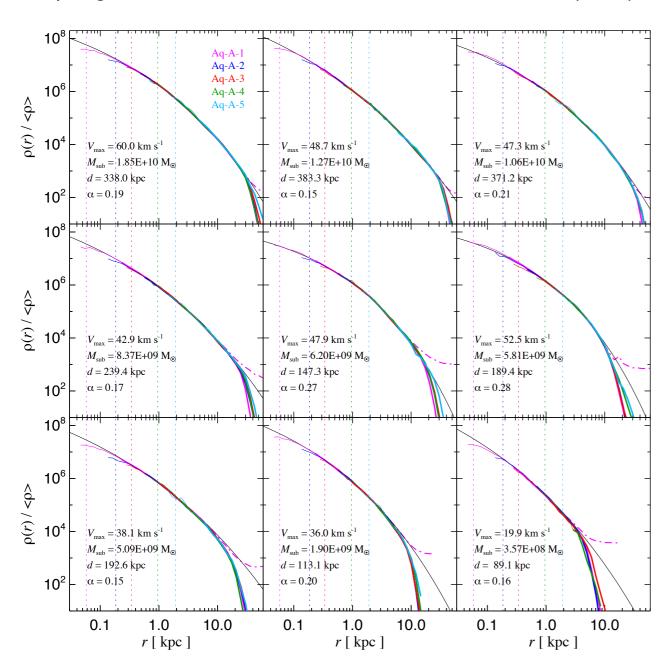
Springel et al., Mon. Not. R. Astron. Soc. **391**, 1685, (2008)



Procedure

1. Solve the differential equation from z_{acc} to z to get m

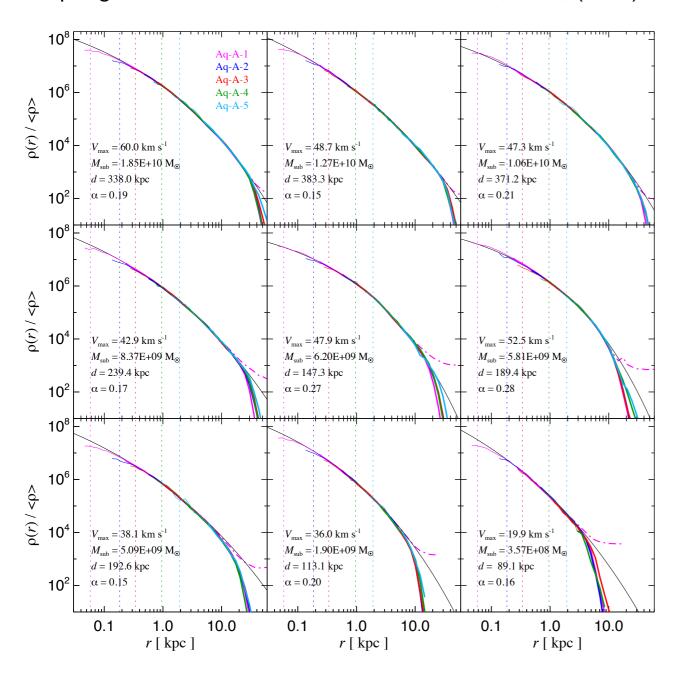
Springel et al., Mon. Not. R. Astron. Soc. **391**, 1685, (2008)



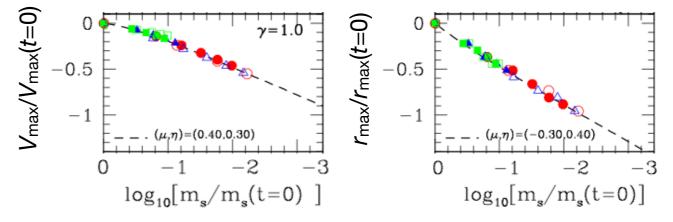
Procedure

- 1. Solve the differential equation from z_{acc} to z to get m
- 2. Calculate ps and rs following Penarrubia et al. (2010)

Springel et al., Mon. Not. R. Astron. Soc. **391**, 1685, (2008)



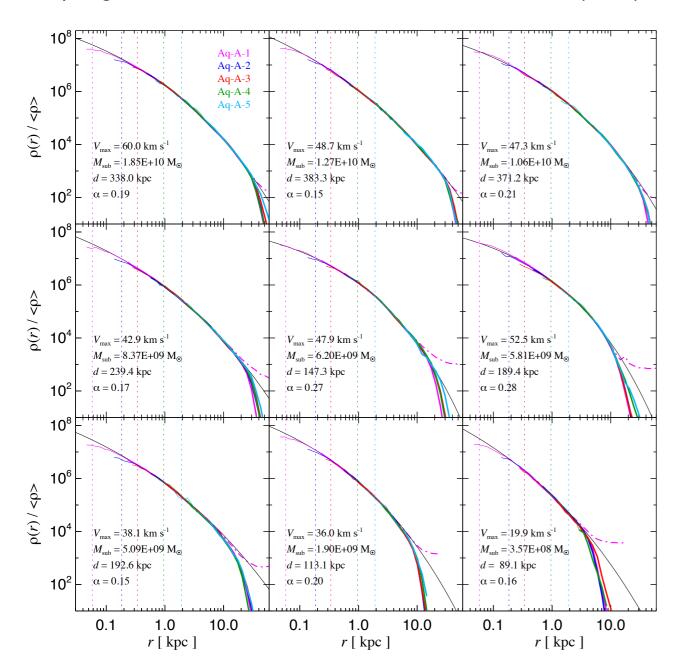
Penarrubia et al., Mon. Not. R. Astron. Soc. 406, 1290, (2010)



Procedure

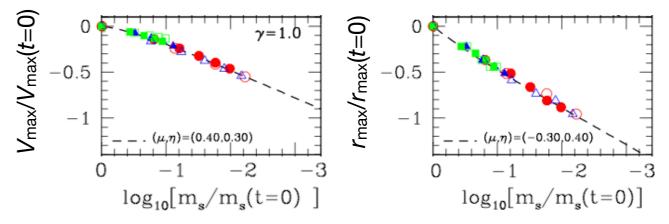
- 1. Solve the differential equation from z_{acc} to z to get m
- 2. Calculate ps and rs following Penarrubia et al. (2010)

Springel et al., Mon. Not. R. Astron. Soc. 391, 1685, (2008)



Truncated NFW

Penarrubia et al., Mon. Not. R. Astron. Soc. 406, 1290, (2010)



Procedure

- 1. Solve the differential equation from z_{acc} to z to get m
- 2. Calculate ps and rs following Penarrubia et al. (2010)
- 3. Obtain truncation radius *r*_t by solving

$$m = \int_0^{r_t} dr \ 4\pi r^2 \rho(r)$$

Evolution

$$L_{\rm sh}^{\rm total}(M,z) = \int dm_{\rm acc} \int dz_{\rm acc} \frac{d^2N_{\rm sh}}{dm_{\rm acc}dz_{\rm acc}} L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc})$$

Number of subhalos accreted at $z_{\rm acc}$ with mass $m_{\rm acc}$

Luminosity of the subhalo at z

$$L_{\rm sh}(z \mid m_{\rm acc}, z_{\rm acc}) \propto \rho_s^2(z \mid m_{\rm acc}, z_{\rm acc}) r_s^3(z \mid m_{\rm acc}, z_{\rm acc}) \left\{ 1 - \frac{1}{[1 + r_t(z \mid m_{\rm acc}, z_{\rm acc})/r_s(z \mid m_{\rm acc}, z_{\rm acc})]^3} \right\}$$

Parameters subhalo density profile after tidal mass loss

Results

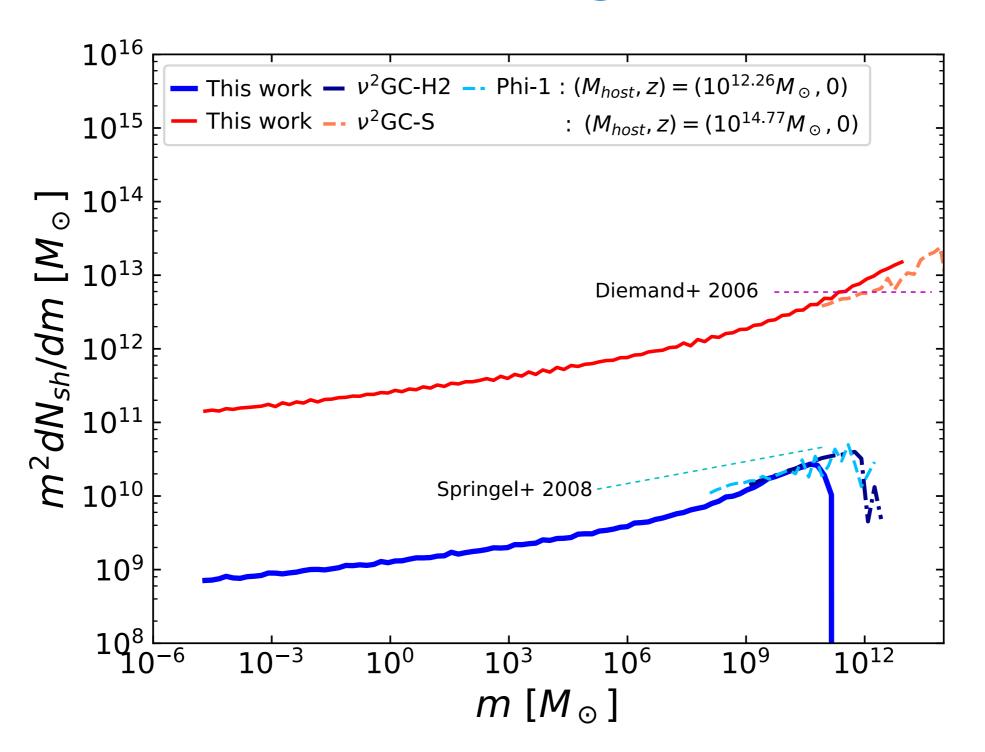
Comparison with simulations

Name		N	L	Softening	$m_{ m p}~({ m M}_{\odot})$	Reference
ν^2 GC-S	Cluster	2048^{3}	$411.8~\mathrm{Mpc}$	$6.28~{\rm kpc}$	3.2×10^8	[38, 44]
ν^2 GC-H2	Galaxy	2048^{3}	$102.9~\mathrm{Mpc}$	$1.57~\mathrm{kpc}$	5.1×10^6	[38, 44]
Phi-1	Dwarf	2048^{3}	$47.1~\mathrm{Mpc}$	$706~\mathrm{pc}$	4.8×10^{5}	Ishiyama et al. (in prep)
Phi-2	Dwarf	2048^{3}	$1.47~\mathrm{Mpc}$	11 pc	14.7	Ishiyama et al. (in prep)
A_N8192	L800 Micro	8192^{3}	$800.0 \ \mathrm{pc}$	$2.0 \times 10^{-4} \text{ pc}$	3.7×10^{-11}	Ishiyama et al. (in prep)

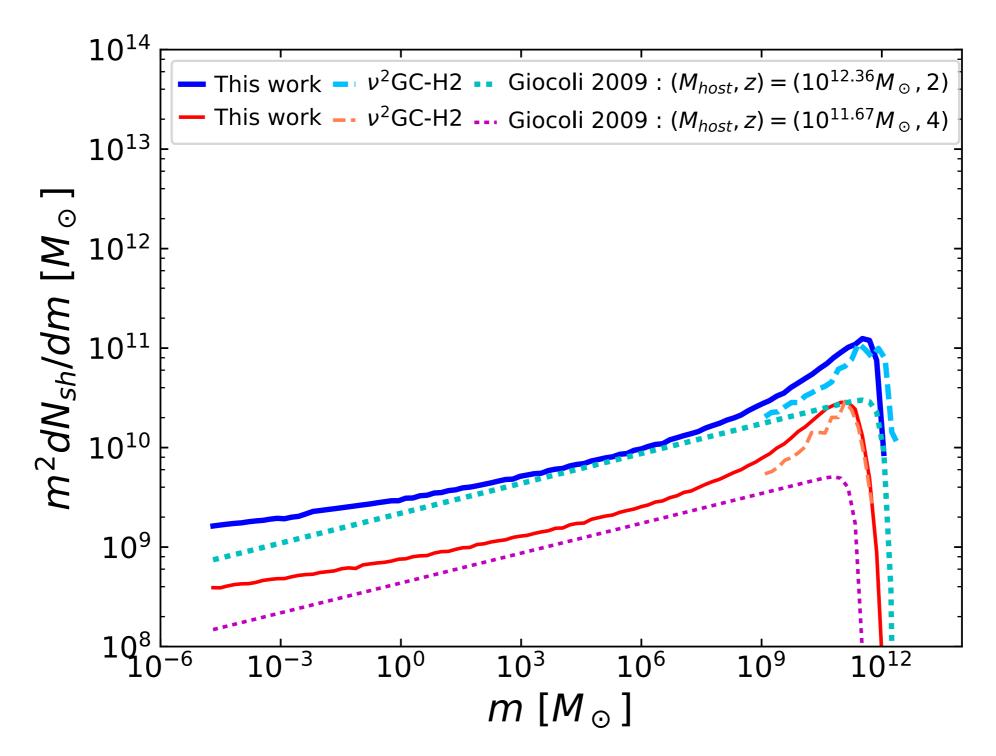
[38] Ishiyama et al., Pulb. Astron. Soc. Jap. 67, 61 (2015)

[44] Makiya et al., Pulb. Astron. Soc. Jap. 68, 25 (2016)

Subhalo mass function: Clusters and galaxies

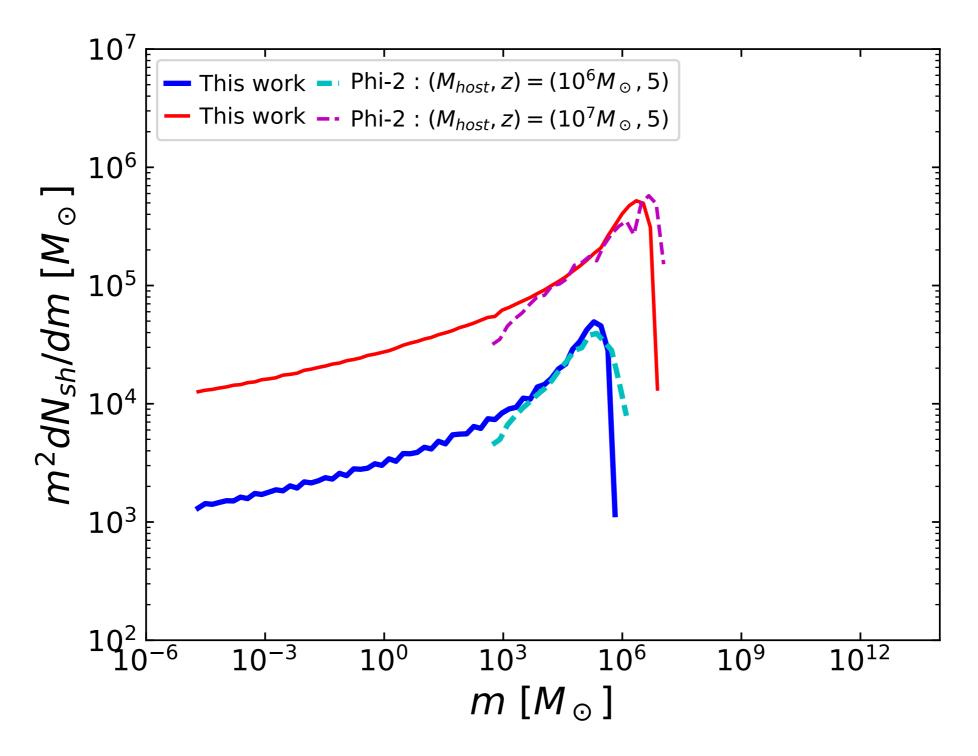


Subhalo mass function: Galaxies at z=2,4



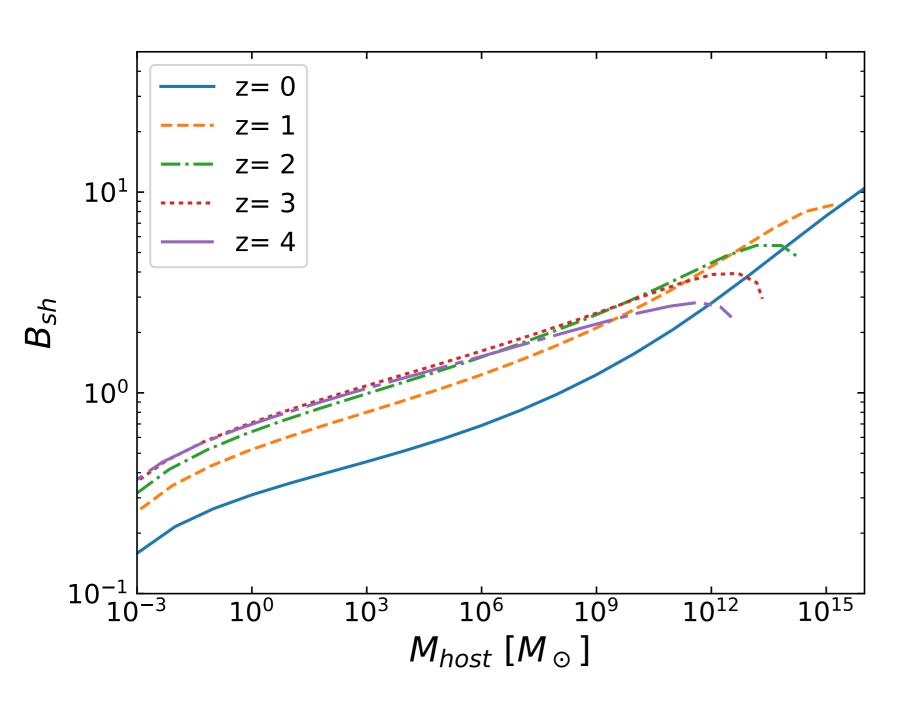
Hiroshima, Ando, Ishiyama, Phys. Rev. D 97, 123002 (2018)

Subhalo mass function: Dwarfs at z=5



Annihilation boost

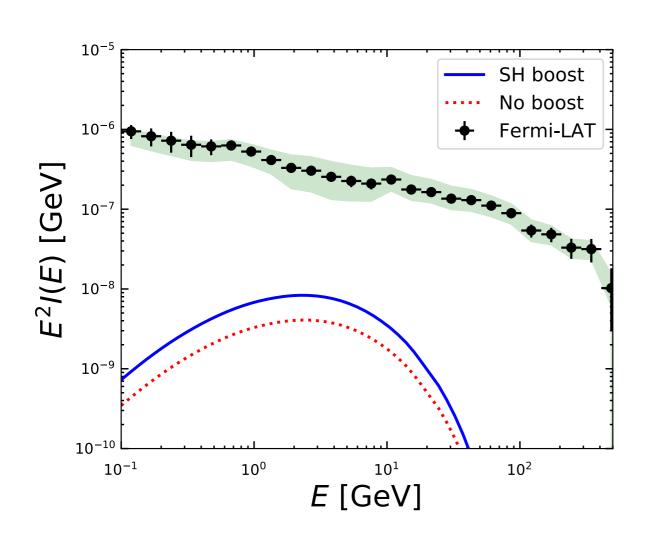
Hiroshima, Ando, Ishiyama, Phys. Rev. D 97, 123002 (2018)

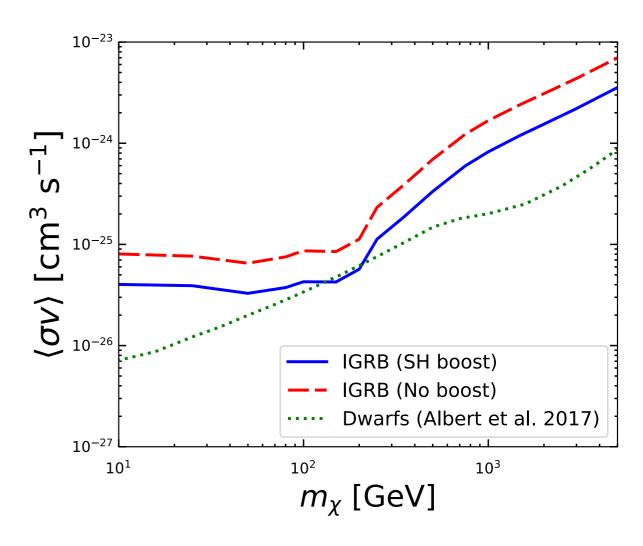


- Boost factors are higher at larger redshifts, but saturates after z = 1
- Boost can be as large as ~3 (10) for galaxies (clusters)
- For one combination of host mass and redshifts (M, z), the code takes only ~O(1) min to calculate the boost on a laptop computer

Application: IGRB

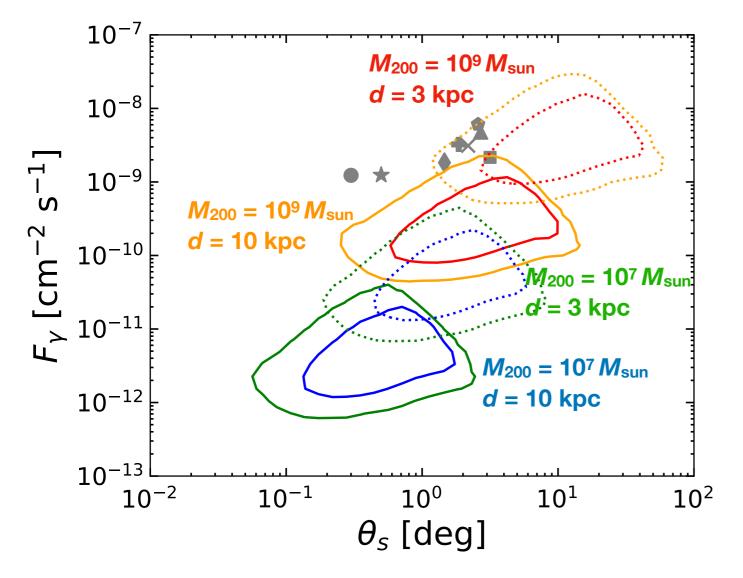
Hiroshima, Ando, Ishiyama, Phys. Rev. D 97, 123002 (2018)





Application: Fermi unassociated sources

Ciuca, Kawata, Ando, Calore, Read, Mateu, Mon. Not. R. Astron. Soc. 480, 2284 (2018)



3FGL J2212.5+0703 (star), 3FGL J1924.8–1034 (circle), FHES J1501.0–6310 (pentagon), FHES J1723.5–0501 (diamond), FHES J1741.6–3917 (square), FHES J2129.9+5833 (cross), FHES J2208.4+6443 (plus), FHES J2304.0+5406 (square)

- Test of Fermi unassociated sources in light of *Gaia* non-detection: upper limit 10⁹ M_{sun} within 20 kpc
- Analytic subhalo model enables to compute PDF of source extension and gamma-ray flux (for a fixed distance)

$$\langle \sigma v \rangle = 2 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

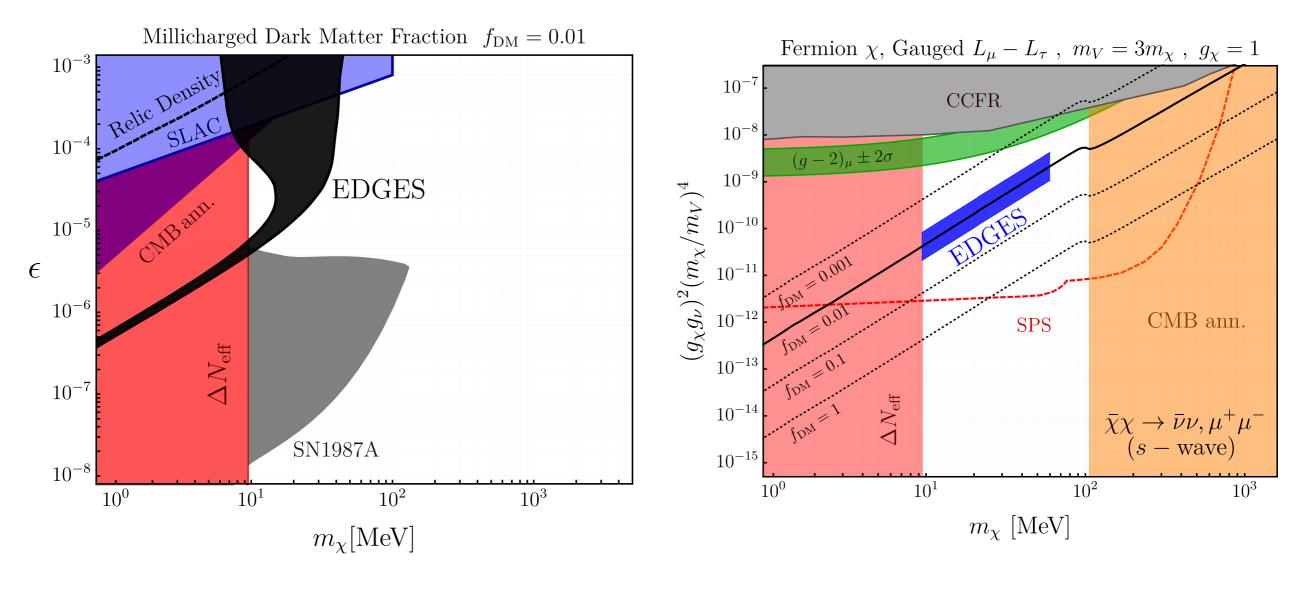
 $m_{\gamma} = 25 \text{ GeV}$

- Only they can be dark matter annihilation for $10^9 M_{sun}$ at d = 3 kpc
- This is unlikely because (1) probability is very small and (2) it will be depleted by the disk
- Conclusion: no Fermi unassociated sources are subhalos

Application: MeV neutrinos and implications for 21-cm measurements

Model to explain EDGES 21-cm obs

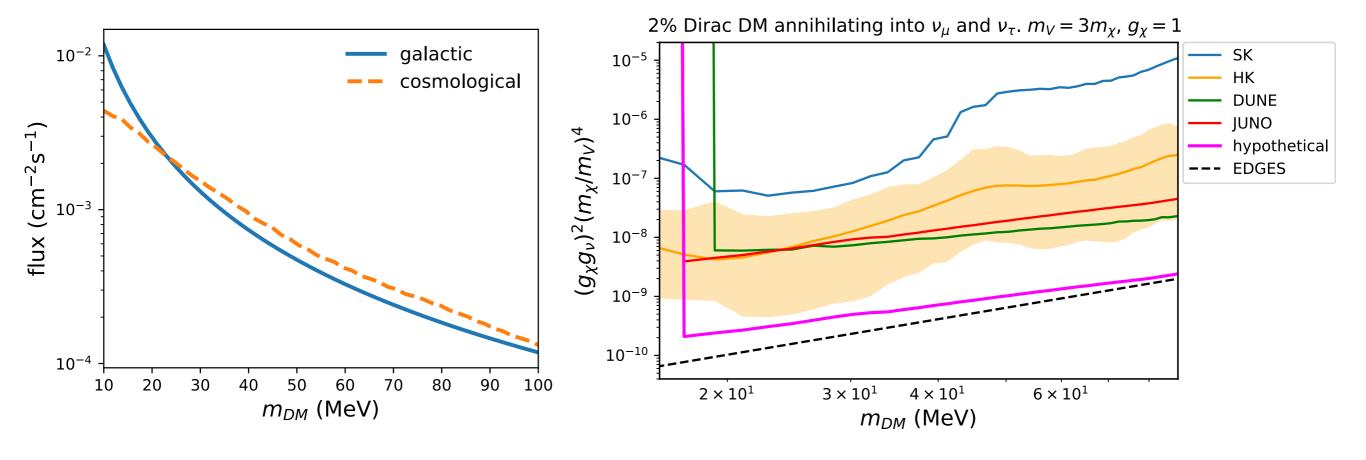
Berlin et al., *Phys. Rev. Lett.* **121**, 011102 (2018)



- $L_{\mu} L_{\tau}$ gauge symmetry to explain relic density
- MeV dark matter annihilating into neutrino-antineutrino pairs still allowed

Neutrino constraints on MeV dark matter

Klop, Ando, *Phys. Rev. D* **98**, 103004 (2018)



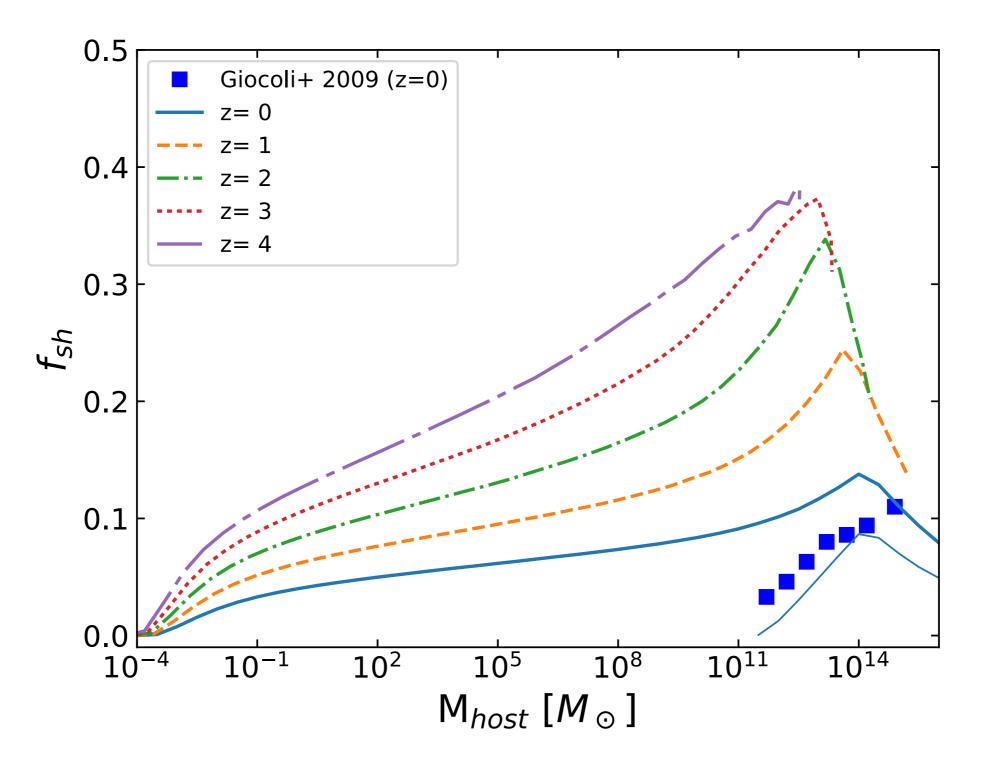
- Similar contribution to neutrino flux from Galactic halo and cosmological halos (including substructure boost)
- EDGES regions might be reached with future neutrino telescopes

Conclusions

- Combining the distribution of subhalo accretion with the evolution afterwards, we can analytically model various subhalo quantities such as mass function and annihilation boost factor
- The subhalo mass function appears to be in good agreement with results of numerical simulations for wide range of masses and redshifts
- The annihilation boost factors are predicted to be ~3 (10) for galaxy (cluster) halos
- The boost is not as uncertain as has been considered

Backup slides

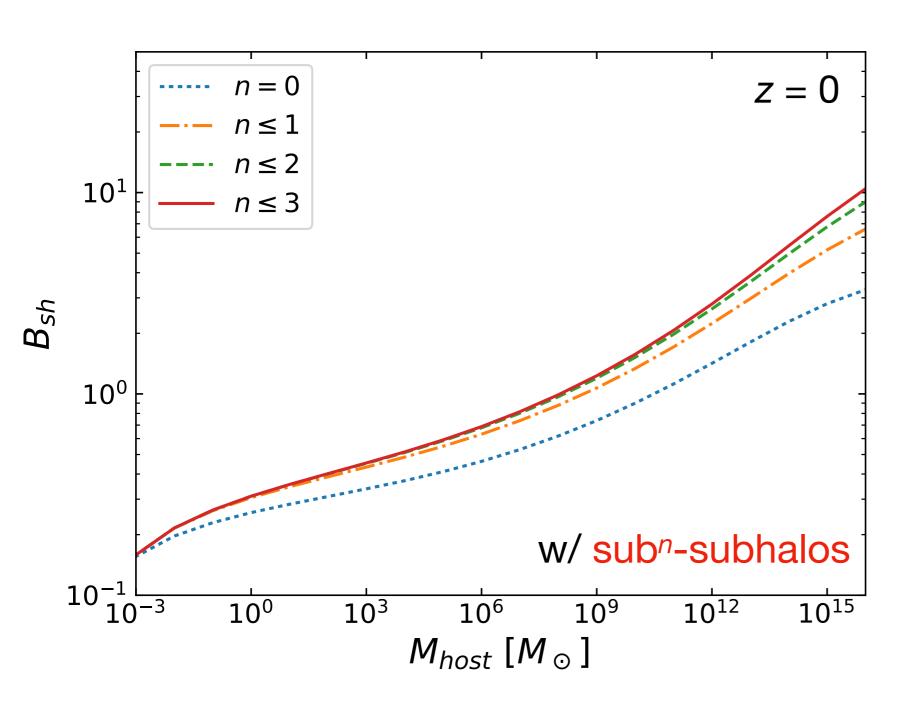
Subhalo mass function: Mass fraction in the subhalos



Hiroshima, Ando, Ishiyama, *Phys. Rev. D* **97**, 123002 (2018)

Annihilation boost

Hiroshima, Ando, Ishiyama, Phys. Rev. D 97, 123002 (2018)



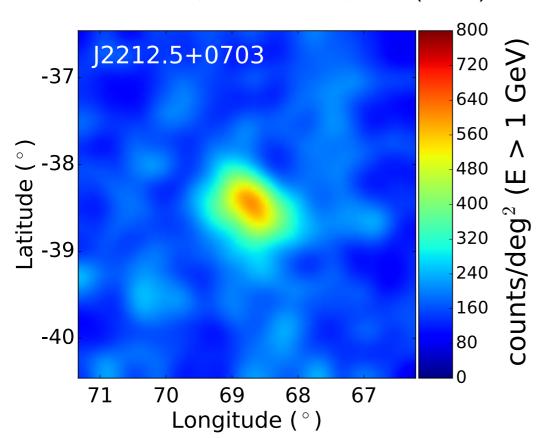
- Include effect of subⁿsubhalos iteratively
- They are assumed to be distributed following

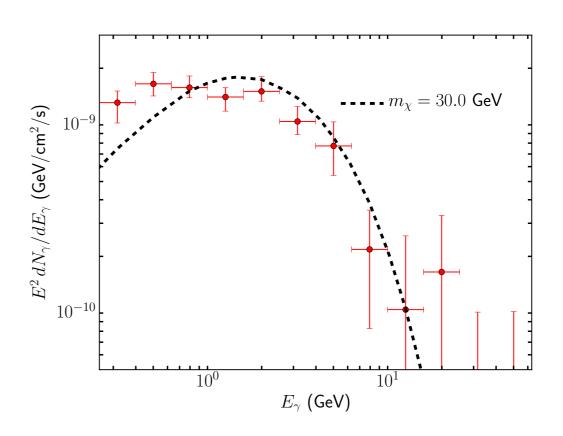
$$\propto [1 + (r/r_s)^2]^{-3/2}$$

- All the sub-subhalos outside of the tidal radius is assumed lost
- Important to include up to sub²substructures

Fermi unassociated sources

Bertoni et al., *JCAP* **1605**, 049 (2016)

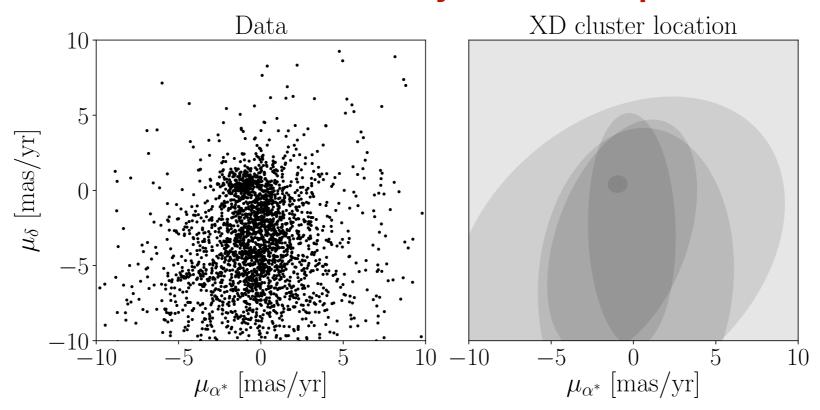




- There are several extended unassociated sources that might be compatible with dark matter annihilation from subhalos
- E.g., 3FGL J2212.5+0703 (Bertoni et al. 2016); 3FGL J1924+1034 (Xia et al. 2017)

Gaia DR2 search for subhalos

Simulation of 5000 M_{sun} stellar system at 10 kpc



Ciuca, Kawata, Ando, Calore, Read, Mateu, Mon. Not. R. Astron. Soc. 480, 2284 (2018)

- No detection of dwarfs (subhalos) towards any of the 8 unassociated sources
- Gaia DR2 should be sensitive to subhalos with pre-infall mass of >10⁹ M_{sun} within 20 kpc