



# Why is Parity *RESTORED* ?

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## Why is parity restored ? ...

a somewhat provocative title, since we would more usually ask  
« Why is parity violated? »

In this talk, I plan to show that

- Parity violation is the **default expectation** in gauge theories
- Parity violation has nothing to do with the presence/absence of right-handed neutrinos

and ask the question:

- Since *we were fooled for centuries* to think parity was a good symmetry, **why is it indeed restored at the large distances then accessible?**
- And then...is some *fairly long distance* P violation possible?

## Parity violation

For centuries, getting to the root of physical law has let us to assume that Parity was a law of nature, with **exceptions linked to biological life (seen as boundary conditions)**.

This abstraction proved right for **gravitation, for mechanics**.

It also proved correct **for electromagnetism** and later for **nuclear forces**.

*Remark : Electromagnetism is a bit tricky, since we seem to introduce a «right-hand rule » to define  $B$ , but this is only an intermediate construction, the convention applies twice to calculate any physical process ...*





# Parity is Broken

## The discovery of P violation was a real shock

It was first met with disbelief, in a **purely hadronic context**

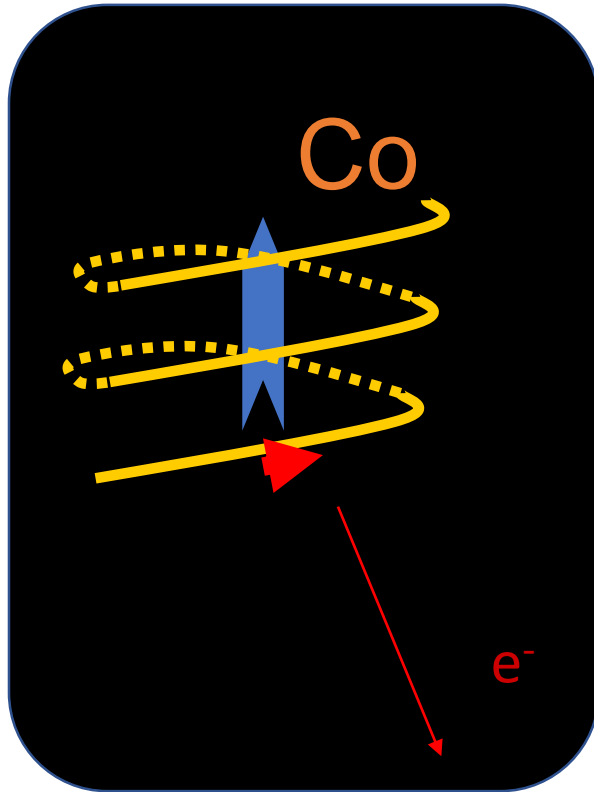
2 particles, then called Theta and Tau (nothing to do with the lepton) were observed with similar masses....close to 500 MeV

The  $\tau \rightarrow 3 \pi$  and the  $\Theta \rightarrow 2 \pi$  ,

Since the decays were in S wave and the p parity was known to be (-) there were 2 possibilities

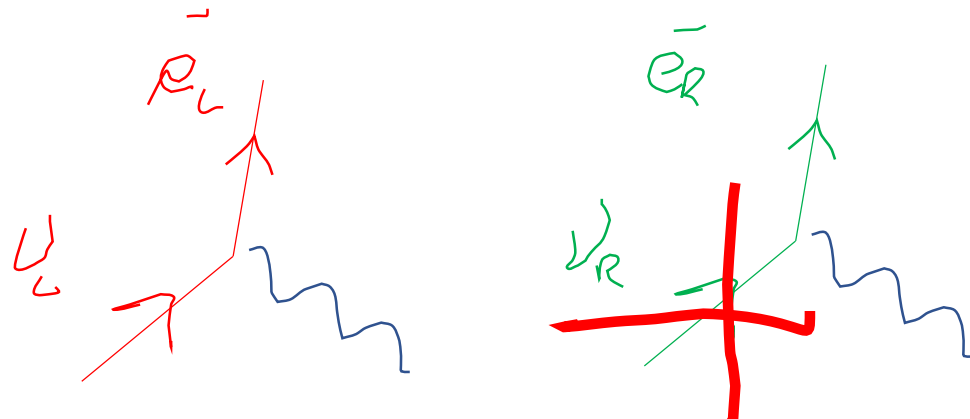
- **$\Theta$  and  $\tau$**  were 2 distinct particles, with identical mass and production (parity doubling)
- OR they were the same particle ( $K^+$ ) **with broken parity?**

It was so hard to accept the breaking of Parity (Lee and Yang) ,  
that a “demonstration” experiment was conceived , the famous Wu experiment.



P violation was clearly demonstrated  
in the Wu experiment ..

It is easy to explain if only left-handed electrons  
are produced in a charged vector current.



*Killing the right-handed neutrino allows for parity violation in charged  
currents, even if the coupling is pure vector*

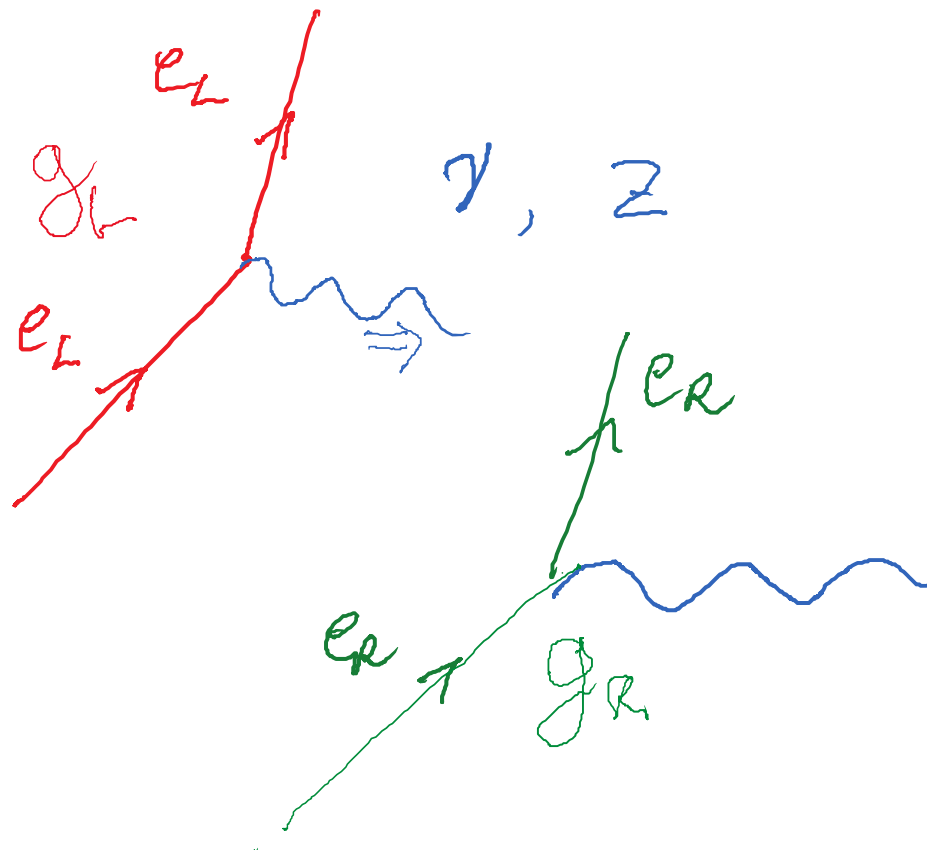
The WU experiment was convinced ...  
.....but led to a wrong track!

By focusing on neutrinos, and thus on leptons,  
it probably led to the often-encountered *folklore*  
that “Parity violation is due to the absence of the right-handed neutrino”,  
and indirectly to the **artificial exclusion of the  $\nu_R$**  from the Standard Model.

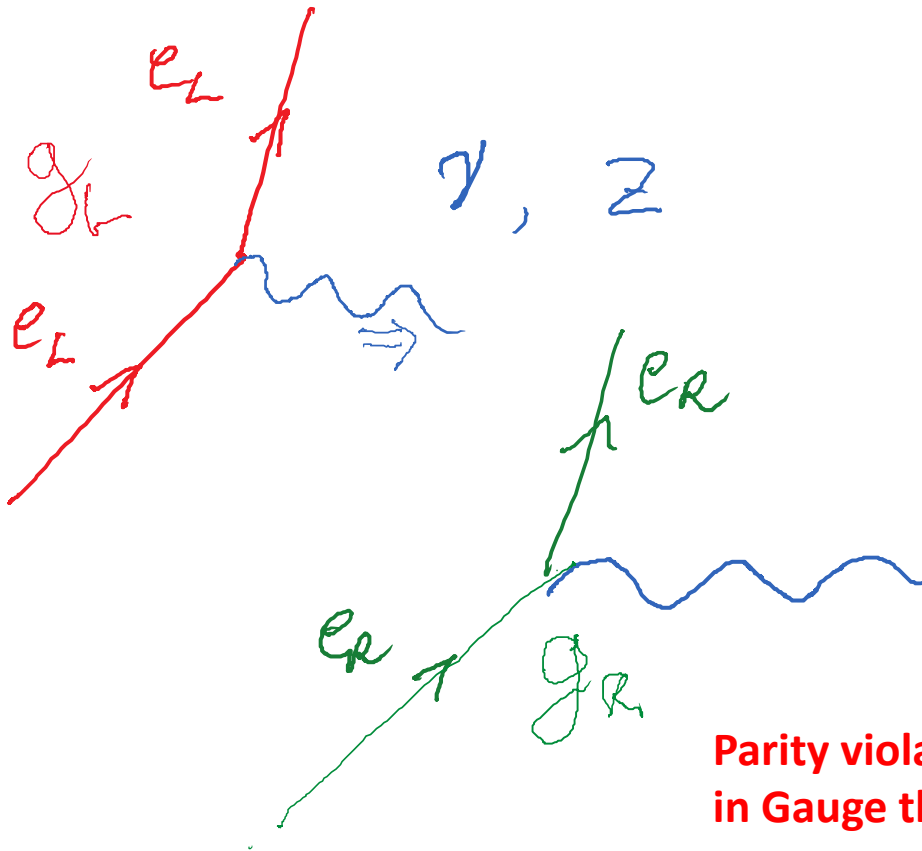
**Of course, this was in contradiction to the initial observation of the  
 $K \rightarrow 2 \pi$  vs  $K \rightarrow 3 \pi$  Parity violation!**

Soon, the experiments establishing the Standard Model proved

- The existence of neutral currents (they could have been included in the Fermi Lagrangian, but were still considered proof of the SM)
- The violation of parity in neutral current interactions (atomic parity violation).







**Parity violation is indeed the EXPECTED SITUATION in Gauge theories ! (In 3+1 dimensions)**

**They are purely chiral, with L spinors speaking only to L spinors and R to R**

$$g_L \neq g_R$$

In fact, the mystery would rather be ....  
Why is Parity respected around us?

Take for example the SU(5) unification ... (or any SuSy approach)

All fermions are re-written in terms of the Left-Handed spinors

e.g;  $((u_R)^c)_L$  .... In 10 and  $\widetilde{5}$

Is it an accident that after breaking, the « long-distance »  
gauge interactions (in which I would include U(1)<sub>em</sub> but  
also the unbroken SU(3)<sub>color</sub> ) are parity invariant ???

*...with the result that we have been fooled for many centuries in believing  
in Parity as an exact symmetry?*

**I have no (complete, satisfactory) answer ! ...**

The mathematical coherence of the theory may give some hints.

- Anomaly conservation
- Gauge invariance of mass term for long-distance interactions
- Singularities in massless gauge bosons coupling with massless fermions prevent the existence of the latter.

- Anomaly conservation

**The discussion must involve the  $U(1)$  em and the  $SU(3)$  long distance forces (unbroken symmetries)**

Let us assume that we start from a grand-unified theory, say  $SU(5)$  (not parity-conserving) or  $SO(10)$

For all the gauged currents, the quantum anomalies must cancel, and the same must remain true after symmetry breaking.

**IF we have only relatively small representations**, this matching can only be done for  $SU(3)$  by compensating 3 by  $\bar{3}$  ... (no other representation like 6 present) ... which would lead to Parity restoration .

- Only massive fermions have long-distance interactions  
(one neutrino could be massless, but it does not have long-distance interactions)

Hence a mass term in the Lagrangian must be invariant under the corresponding gauge transformation (rotation by  $\varphi$ )

$$m \bar{\psi}_L \psi_R + h.c.$$

Diagram illustrating the gauge transformation of the mass term:

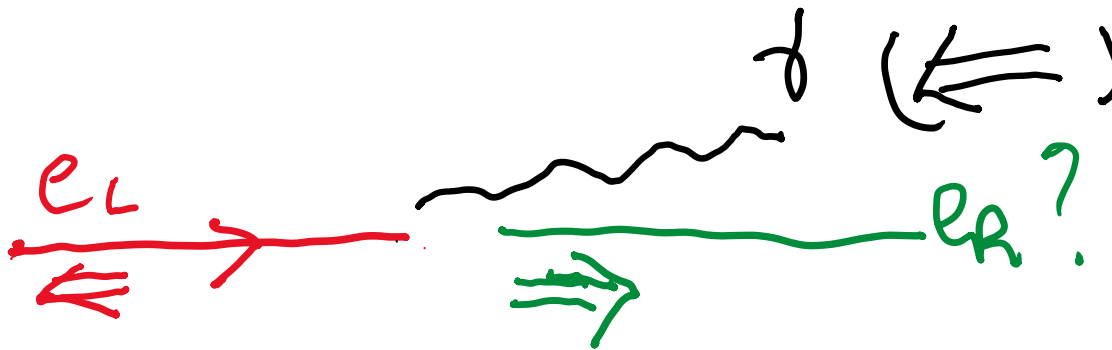
- The left fermion field  $\psi_L$  (red) transforms as  $\psi_L \rightarrow e^{i\alpha_L} \psi_L$  (red).
- The right fermion field  $\psi_R$  (green) transforms as  $\psi_R \rightarrow e^{i\alpha_R} \psi_R$  (green).
- The mass term  $m \bar{\psi}_L \psi_R$  is invariant only if  $\alpha_R = \alpha_L$  (green).

This brings us to an old question :

The problem of parity restoration is solved if only massive fermions can have charges under long-distance interactions (massless gauge bosons)

**Indeed, there are singularities,**

**For instance consider the longitudinal emission,** either from a fermion with  $m \rightarrow 0$   
or a massless fermion



The L R transition  
is only possible with  
a small mass term ...



Could we still have some « residual » *fairly long-distance* P violation?

## Magnetic Moments

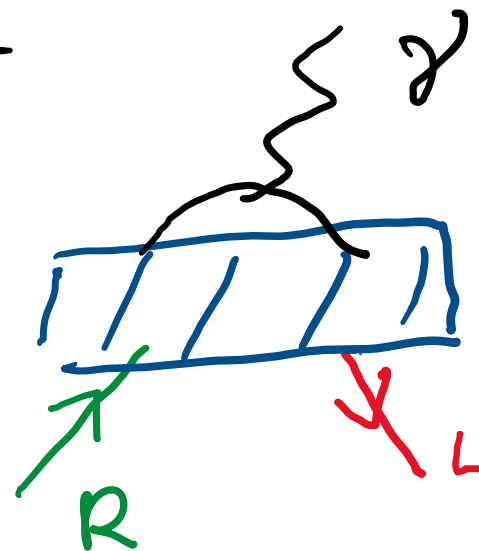
$$\mu \vec{S} \cdot \vec{B}$$

As LR transitions, diagonal magnetic moments are intrinsically P conserving.

$$\bar{p}_L \sigma^\mu p_R F_{\mu\nu} + h.c$$

## Transition Magnetic Moments

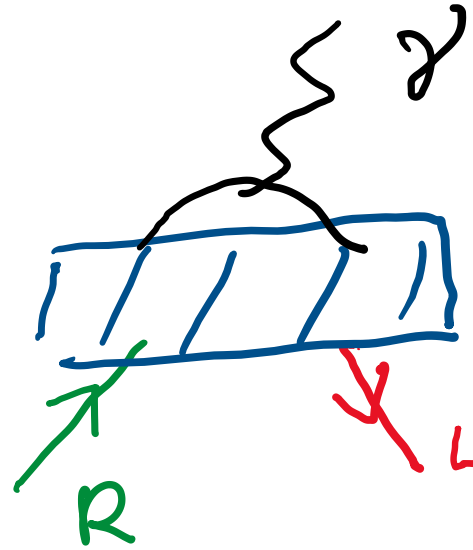
$$\mu \bar{e}_L \sigma^\mu s_R + h.c$$



## Electric Dipole Moments

EDMs are intrinsically P violating for an *elementary particle*, as they must align with the spin\* but can be induced by local forces.

$$\vec{S} \cdot \vec{E}$$

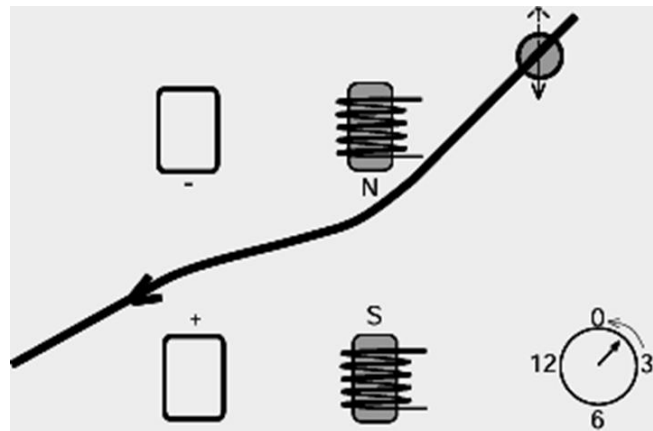
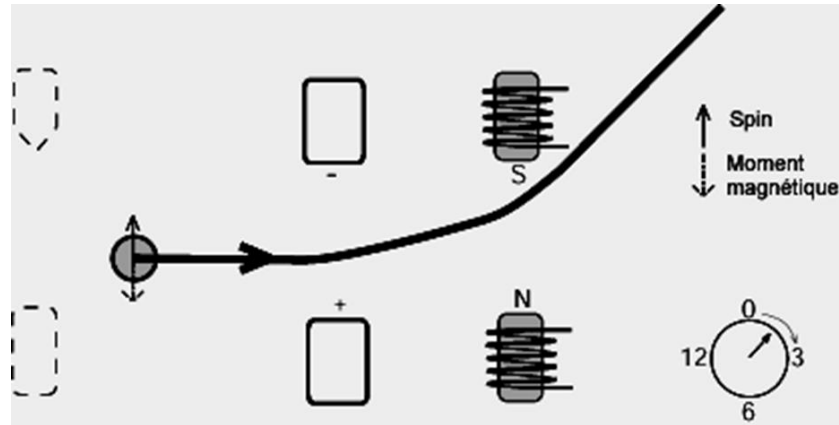


**They violate P, CP and T .**

**They are expected to be tiny in SM, but are predicted (and currently bound by experiment) in most extensions (SuSy, LR, ...)**

*\*except maybe in some non-commutative geometries.*

## T-violating effect of hypothetical electric dipole moment (gedanken experiment)





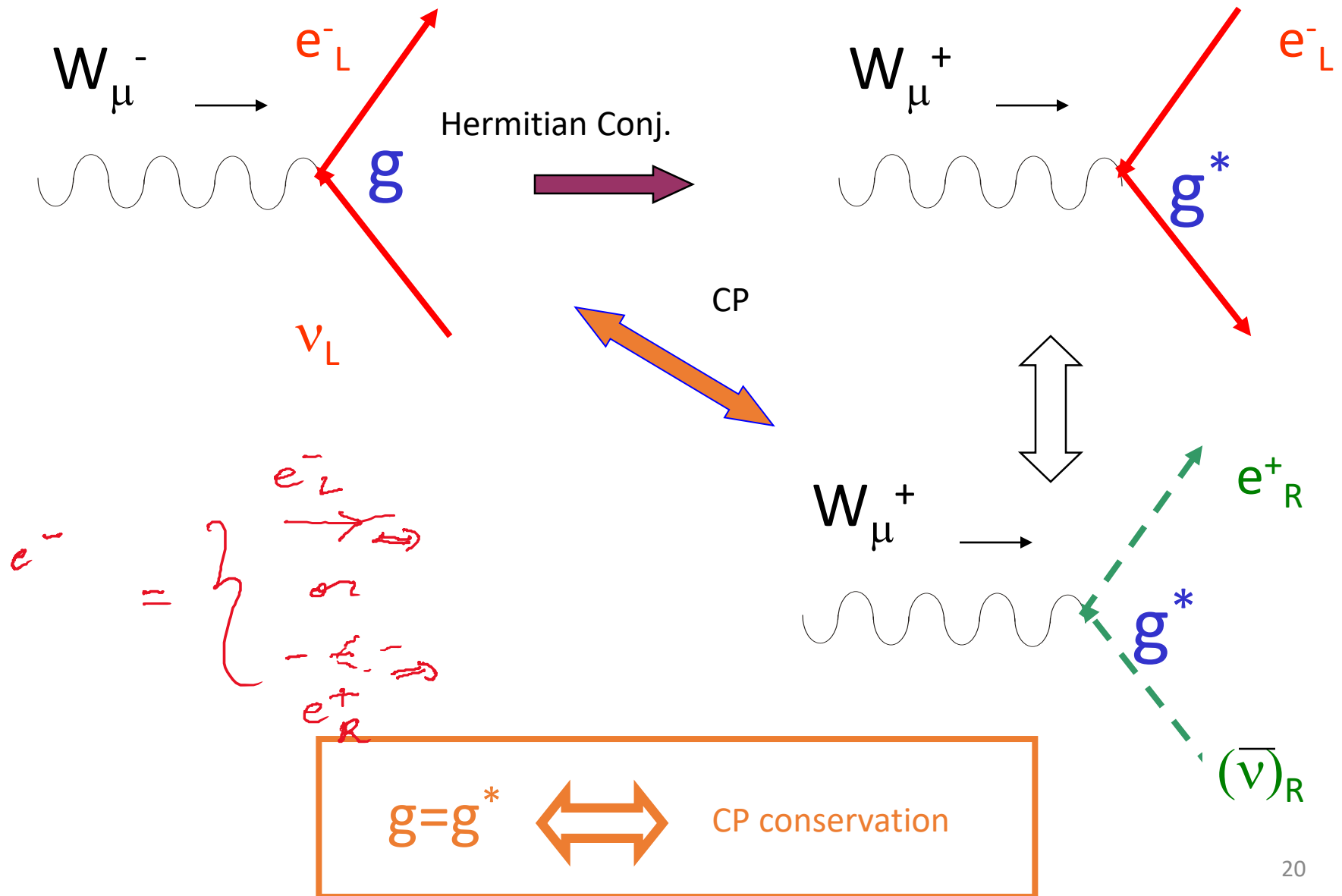
Parity restoration  
due  
to fermion masses?

Some fairly long-distance  
P violation  
could persist through  
dipole moments.

If time allows ... a few words about CP violation

- CP is the natural symmetry of pure gauge interactions of fermions (no scalar or mass terms)
- After the discovery of P violation, CP violation was expected (predicted) by theorists (in particular Lev B. Okun, who had to counter Landau's opposition) and the K decay experiments were suggested
- CP violation is due to something different: the feeble Scalar interaction (there are no weak interactions, but still feeble ones)

CP is the intrinsic symmetry of gauge interactions in 3+1 dimensions  
see Basic Building Bloc : chiral fermion and gauge boson





# The feeble interaction (H)

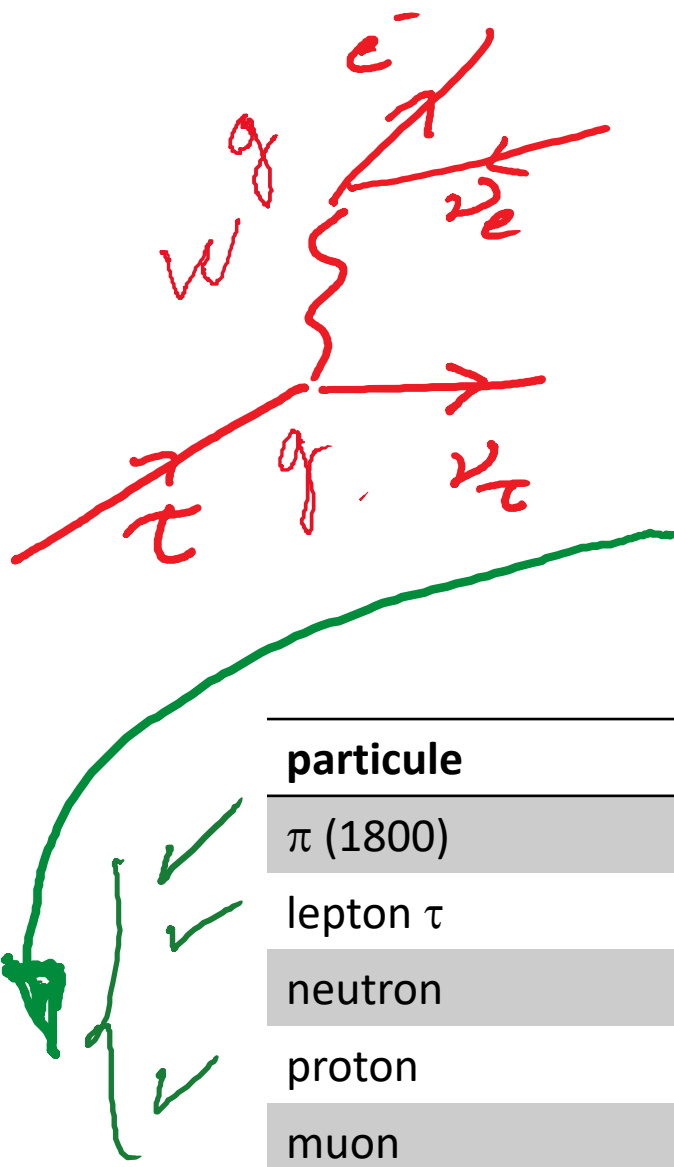
Naïve question: where is the border between « weak » and « strong »?

Answer (ca 1970) ...didn't you realize they are 10 orders of magnitude apart ???

particule	Mass (GeV/c <sup>2</sup> )	« lifetime »
$\pi$ (1800)	1.8	$3.3 \cdot 10^{-24} \text{ s}$
lepton $\tau$	1.777	$2.9 \cdot 10^{-13} \text{ s}$
neutron	0.9396	880 s
proton	0.938	$> 10^{+31} \text{ years}$
muon	0.113	$2.2 \cdot 10^{-6} \text{ s}$

« Dimensional analysis » expectation ....

$$1 \text{ GeV} \quad T = 6.58 \cdot 10^{-25} \text{ s}$$



$$T \sim \frac{1}{m_\tau} * (1/g)^4 * \left(\frac{M_W}{m_\tau}\right)^4$$

Playing either on MW or g allows to fit lifetimes, but to fit all, a combination of M of order 100GeV and **g of order of the electric charge** is needed....

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$$g \sin \theta = e$$

# ENTERS THE *feeble* FORCE



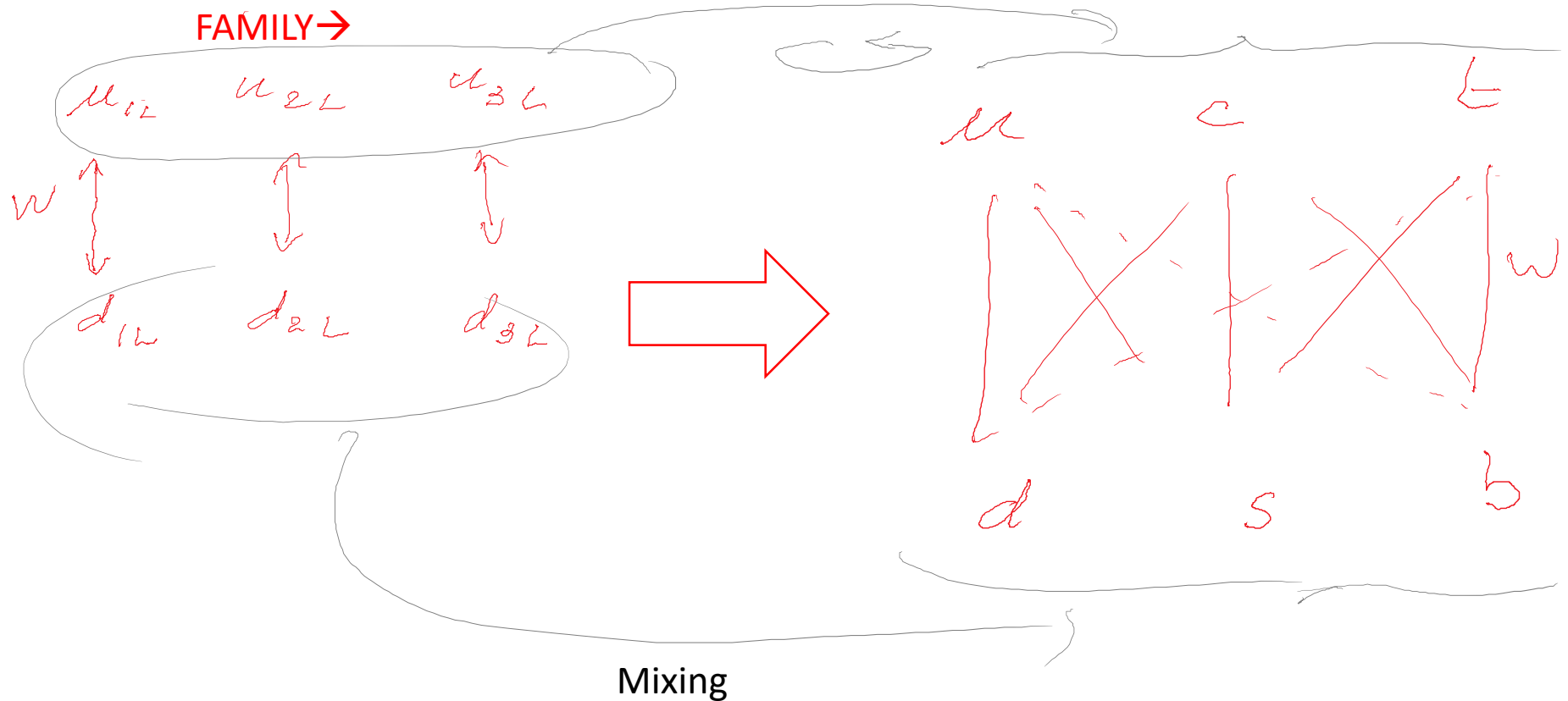
The Scalar Boson (B-E-H) which was introduced to give mass to the gauge bosons ... finds a new use. It is also responsible for the fermion masses!

A priori, 2 different roles, BUT

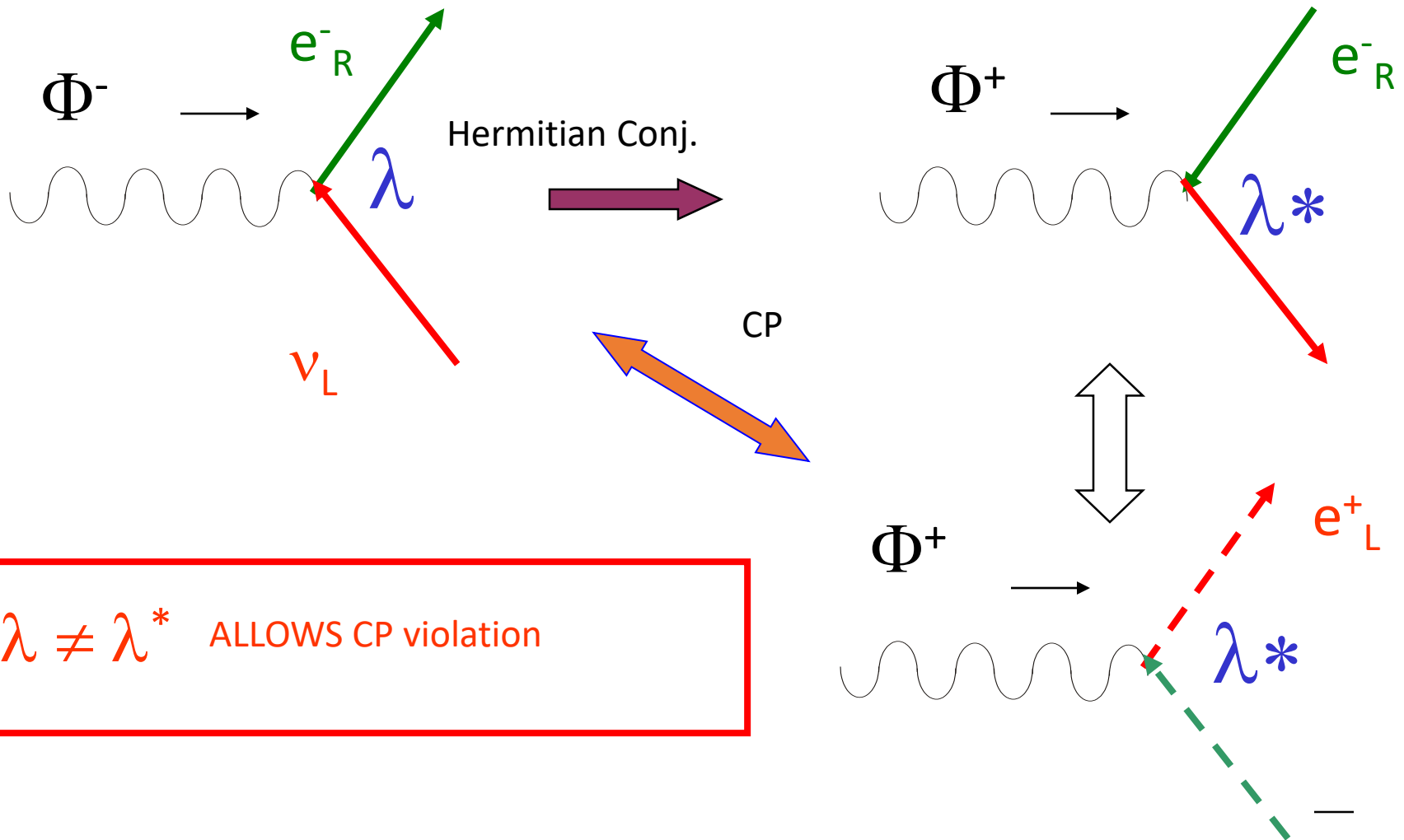
- SU(2) breaking is needed to split masses in a multiplet
- Such splitting necessarily contributes to the W and Z masses

$$\lambda_e = \frac{g}{\sqrt{2}} \frac{m_e}{M_W} \simeq 6 \cdot 10^{-6} \frac{g}{\sqrt{2}}$$

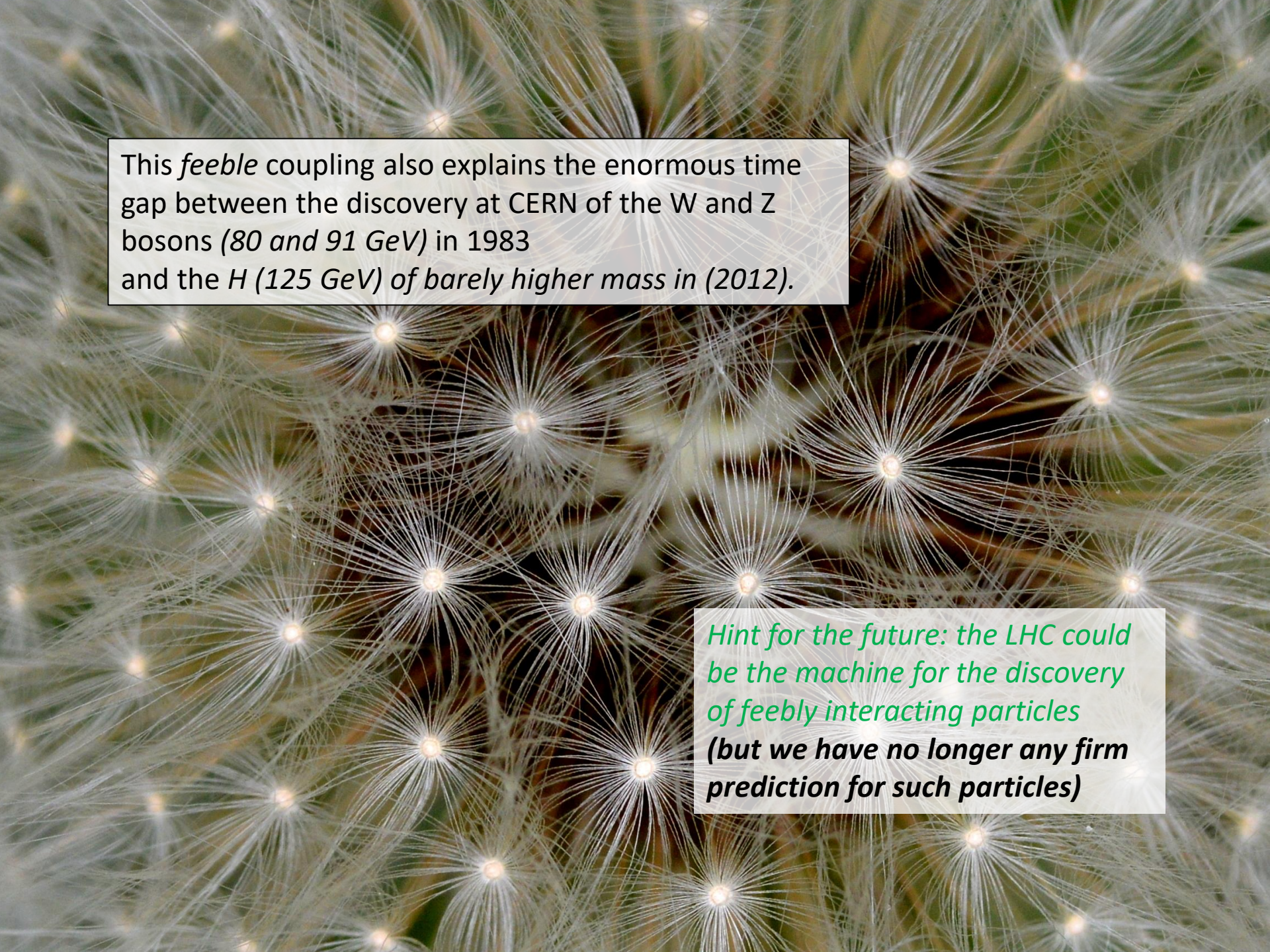
Moving from a « current » basis to a « mass basis »,  
the Scalar (feeble) interaction accounts for the instability of the heavy flavours ..



But CP violation is easily introduced by arbitrary, complex couplings







This *feeble* coupling also explains the enormous time gap between the discovery at CERN of the W and Z bosons (*80 and 91 GeV*) in 1983 and the *H (125 GeV)* of *barely higher mass* in (2012).

*Hint for the future: the LHC could be the machine for the discovery of feebly interacting particles (but we have no longer any firm prediction for such particles)*