



IN2P3
Les deux infinis



(α,γ) & (α,n) key reaction studies using ($^7\text{Li},t$) alpha-transfer reactions

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Why studying (α,γ) & (α,n) reactions via α -transfer reactions?

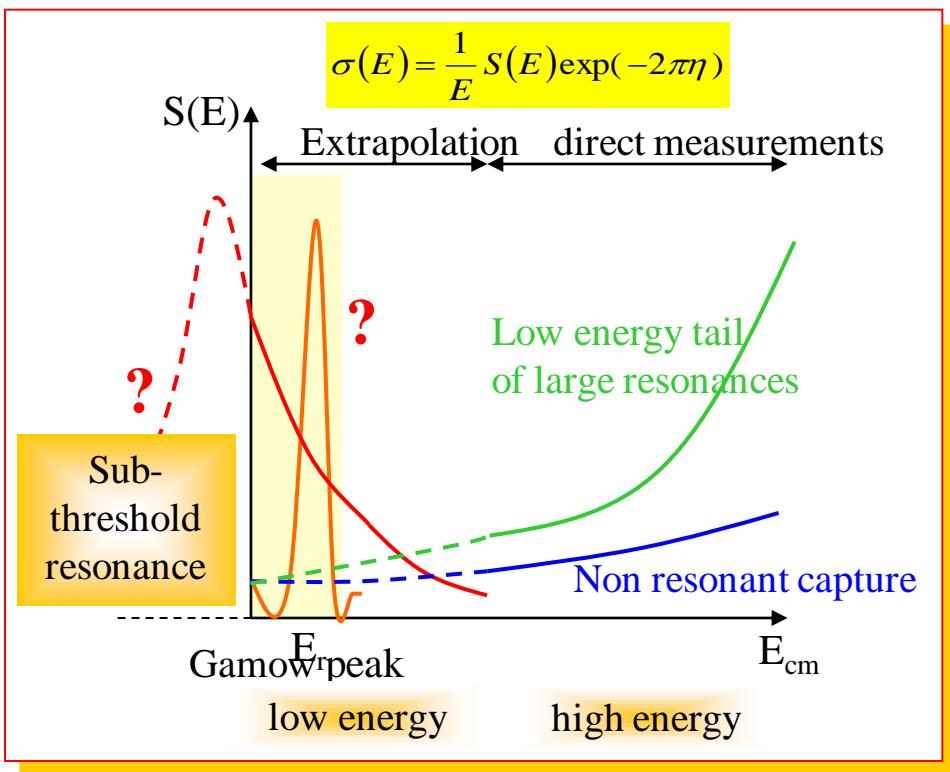
- $^{12}\text{C}(\alpha,\gamma)^{16}\text{O}$ (massive stars), $^{14}\text{C}(\alpha,\gamma)^{18}\text{O}$ (AGB stars & ^{19}F)
 - $^{13}\text{C}(\alpha,n)^{16}\text{O}$
 - $^{22}\text{Ne}(\alpha,n)^{25}\text{Mg}$, $^{22}\text{Ne}(\alpha,\gamma)^{25}\text{Mg}$
 - $^{17}\text{O}(\alpha,n)^{20}\text{Ne}$, $^{17}\text{O}(\alpha,\gamma)^{21}\text{Ne}$
 - $^{15}\text{O}(\alpha,\gamma)^{19}\text{Ne}$ (Type I X-ray bursts)
- $\left. \begin{array}{l} \text{s-process in AGB} \\ \text{\& massive stars} \end{array} \right\}$

T=0.1-1 GK

\rightarrow hundreds keV- MeV $\ll E_C$

$\rightarrow \sigma(E)$ very weak (0.01 fb- \sim 100 pb)

\rightarrow Direct measurements are very challenging or impossible



➤ In case of stable nuclei:
Direct measurements of $\sigma(E)$ at high energies
then extrapolation at stellar energies

But:

Problems with extrapolation: resonances at very low energy, sub-threshold resonances

➤ In case of radioactive nuclei:
Low beam intensities ($\sim 10^5$ - 10^7 p/s)
 \rightarrow direct measurements challenging

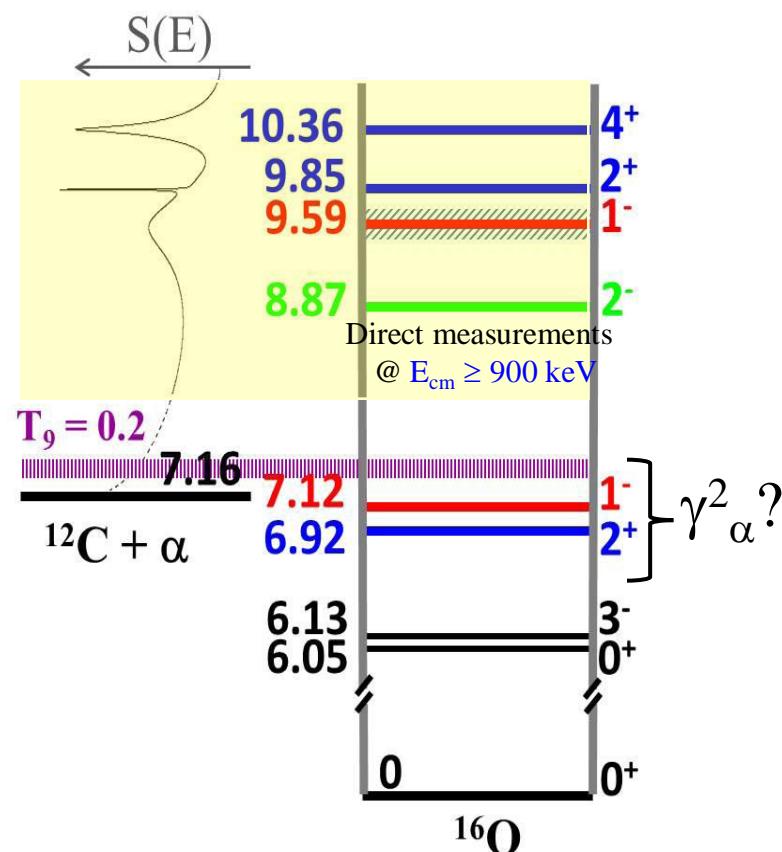
Resonant cross-sections & $N_a \langle \sigma v \rangle$ can be calculated if (E_R , J^π , $\Gamma_\alpha, \Gamma_n, \Gamma_\gamma, \Gamma_{tot}$) are known/constrained
 \rightarrow ($^7\text{Li},t$) or ($^6\text{Li},d$) α -transfer reactions $\rightarrow E_R, l_\alpha, S_\alpha \Rightarrow \gamma^2_\alpha, \Gamma_\alpha$

$^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$: the challenge of Nuclear Astrophysics

→ Crucial for energetics, nucleosynthesis, final fate of the massive stars,...

- $^{12}\text{C}(\alpha, \gamma)^{16}\text{O} \rightarrow 40\%$ uncertainty

- $T=0.2 \text{ GK}$ Gamow peak $\sim 300 \text{ keV}$, $\sigma(E_0) \sim 10^{-8} \text{ nb}$



Direct measurements @ 300 keV → impossible



Need of precise data at high energies & extrapolation at 300 keV (R-matrix formalism)

BUT

Overlap of various contributions:

- E1 & E2 transitions to the gs:
Effect of the high energy tail of the 1^- & 2^+ sub-threshold resonances ($S_\alpha, \gamma^2_\alpha$?)

$$S_\alpha(1^-) \rightarrow 0.02-1.08 !?$$

$$S_\alpha(2^+) \rightarrow 0.13-1.35 !?$$

&
Interference effect & cascades ?

⇒ Study of 6.92 & 7.12 MeV states ^{16}O via the transfect reaction $^{12}\text{C}(^{7}\text{Li}, t)^{16}\text{O}$

Study of ^{16}O states via $^{12}\text{C}(^{7}\text{Li},\text{t})^{16}\text{O}$ α -transfer reaction

Advantages

- multi-step effects less marked than in ($^{6}\text{Li},\text{d}$) transfer reaction

Becchetti et al. (78), Cobern et al.(81),
Keeley et al (06)

- less momentum mismatch
($L_\alpha=1$ in ^{7}Li)

$$\left. \frac{d\sigma}{d\Omega} \right|_{\text{exp}} = C^2 S'_{\alpha} S_{\alpha} \left. \frac{d\sigma}{d\Omega} \right|_{FR-DWBA}$$

$$S'_{\alpha} = \langle ^7\text{Li} | t \otimes \alpha \rangle = 1 \quad (\text{Kubo et al})$$

$$S_{\alpha} = \langle ^{16}\text{O} | ^{12}\text{C} \otimes \alpha \rangle$$

ANC

$$\tilde{C}^2 = S_{\alpha} \frac{r^2 |\varphi(r)|^2}{W(r)^2})$$

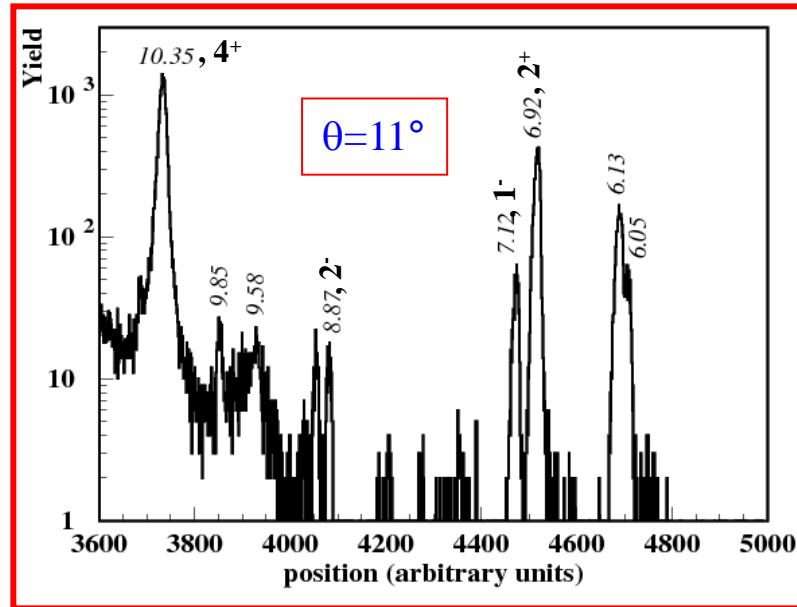
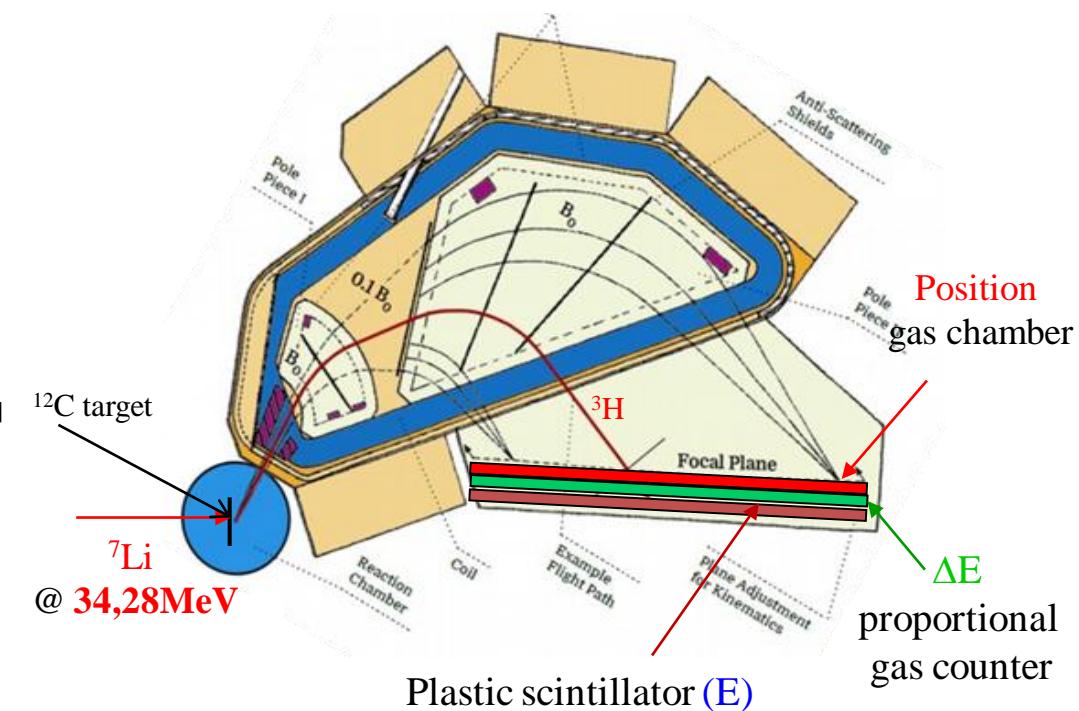
$$\gamma_{\alpha}^2 = \frac{\hbar^2 r}{2\mu} S_{\alpha} |\varphi(r)|^2$$

Detailed and elaborate finite-range DWBA analysis of the data is needed

Experimental Set-up

Split-Pole spectrometer (ALTO-Orsay)

- $\Delta\Omega \sim 1.7 \text{ msr}$
- $\Delta E/E \sim 5 \times 10^{-4}$

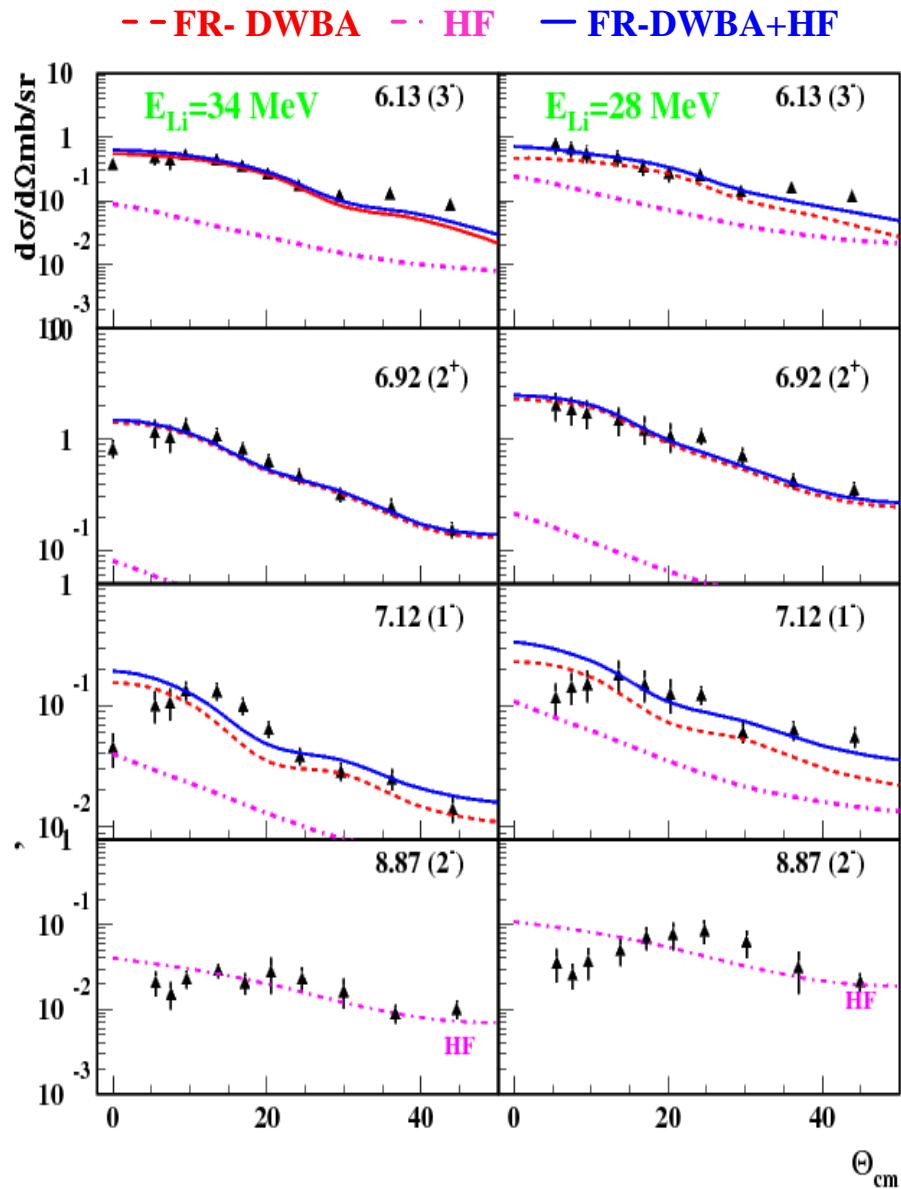


- $d\sigma/d\Omega$ measurements up to 45°

- Strong population of the α cluster 2^+ & 4^+ state \rightarrow direct transfer mechanism
- Population (weak) of the 2^- , 8.87 MeV state \rightarrow compound nucleus?

} FR-DWBA
& HF calculations

Results: Comparison exp & calculations



→ Good description of the data by DWBA
 (6.05, 6.13, 6.92, 7.12, 9.58 et 10.35 MeV)



Direct transfer mechanism

→ Disagreement at $\theta < 10^\circ$ for the 7.12 MeV



Multi-step transfer? $\rightarrow d\sigma/d\Omega \searrow$



No (CDCC calculations of Keeley)

$$S_\alpha(2^+) = 0.15 \pm 0.05$$

$$S_\alpha(1^-) = 0.08 \pm 0.03$$

Main source of error



DWBA input parameters

R matrix calculations—E2 & E1 components

$$S_\alpha(6.92) = 0.15 \pm 0.05 \rightarrow \gamma_\alpha^2 = 27 \pm 10 \text{ keV}$$

$$S_\alpha(7.12) = 0.08 \pm 0.03 \rightarrow \gamma_\alpha^2 = 8 \pm 3 \text{ keV}$$

} $r=6.5 \text{ fm}$ Asymptotic region



$$\tilde{C}^2(2^+) = (2.07 \pm 0.80) \times 10^{10} \text{ fm}^{-1}$$
$$\tilde{C}^2(1^-) = (4.00 \pm 1.38) \times 10^{28} \text{ fm}^{-1}$$

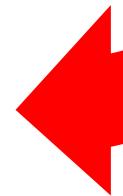
In agreement with
ANC sub-coulomb
experiments

Brune+2001, Avila+2014

Multi-level R-matrix analysis

P. Descouvemont DREAM Code

$$R_{CC'} = \sum_{\lambda} \frac{\gamma_{\lambda C} \gamma_{\lambda C'}}{E_{\lambda} - E}$$



➤ Fit E2 & E1 components separately

➤ Fit

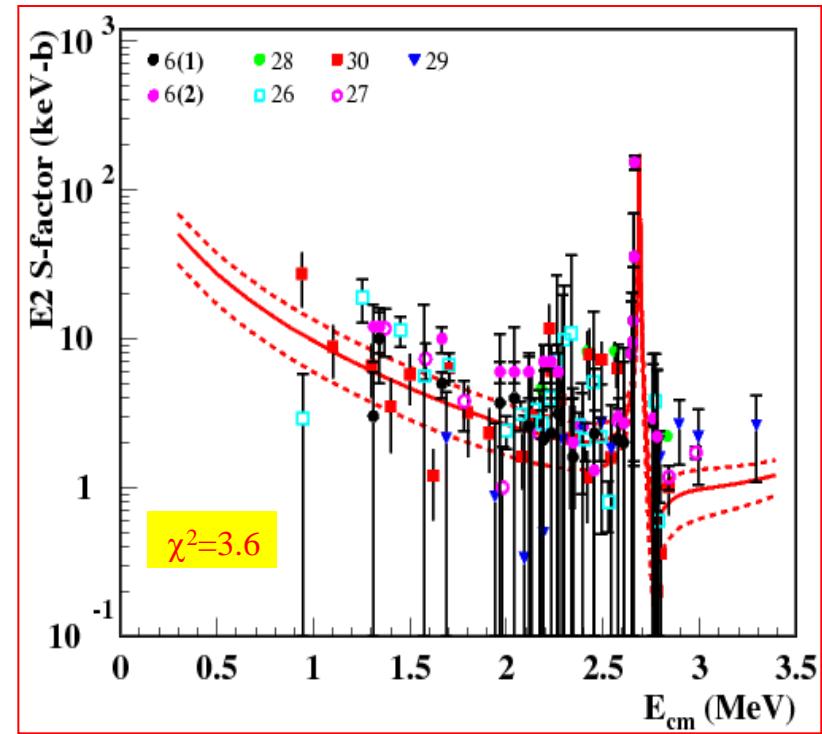
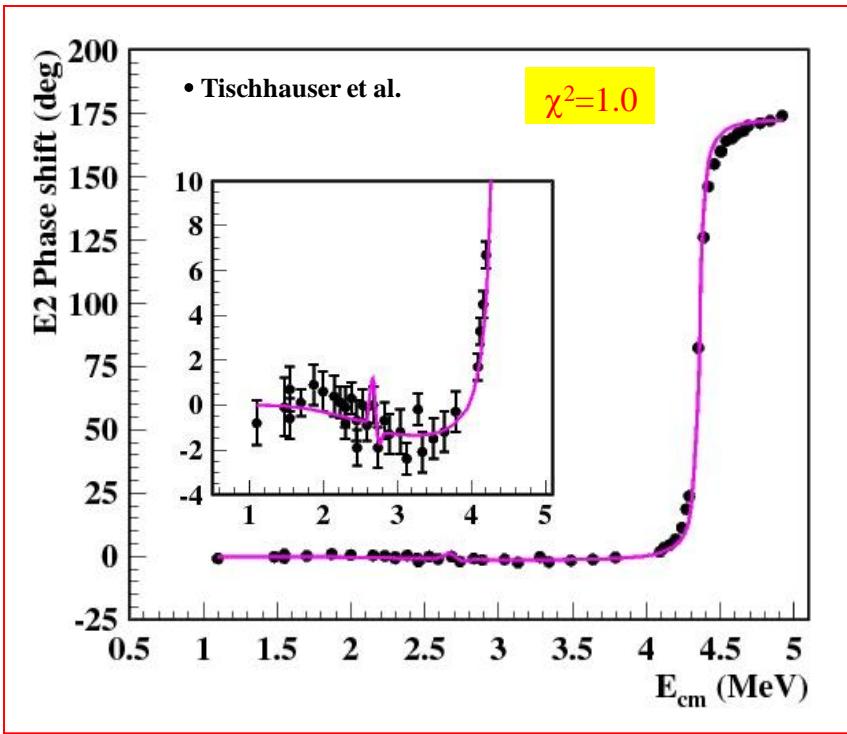
${}^{12}\text{C}(\alpha, \gamma){}^{16}\text{O}$ astrophysical S-factors (direct data)

phase shifts data $\rightarrow {}^{12}\text{C}(\alpha, \alpha){}^{12}\text{C}$ measurements

} $S_{\text{E2}}(300 \text{ keV})$
 $S_{\text{E1}}(300 \text{ keV})$

R matrix calculations– E2 component

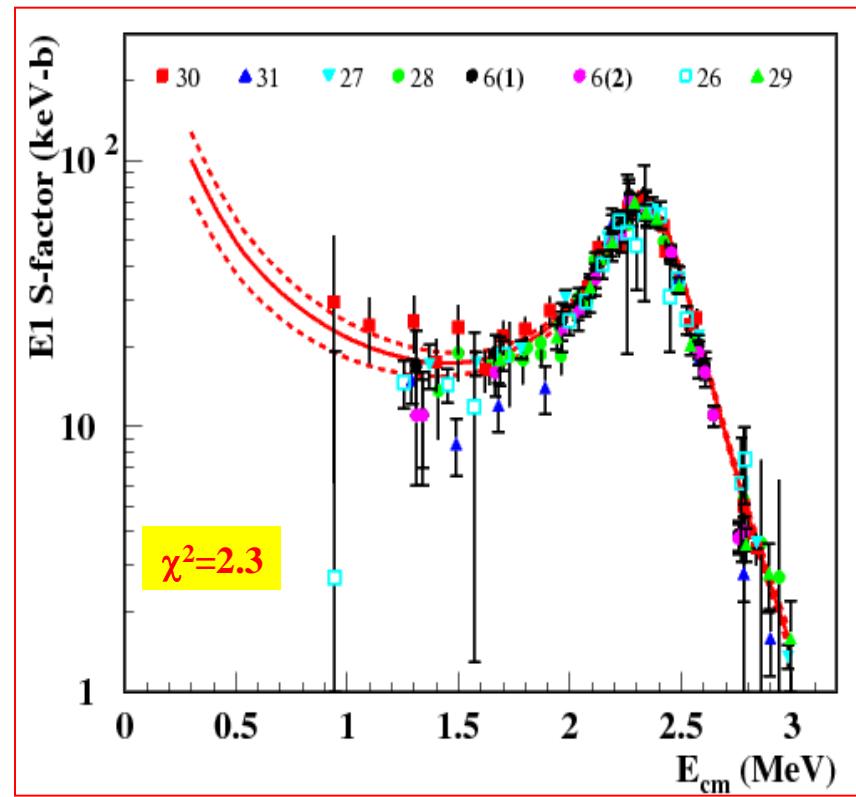
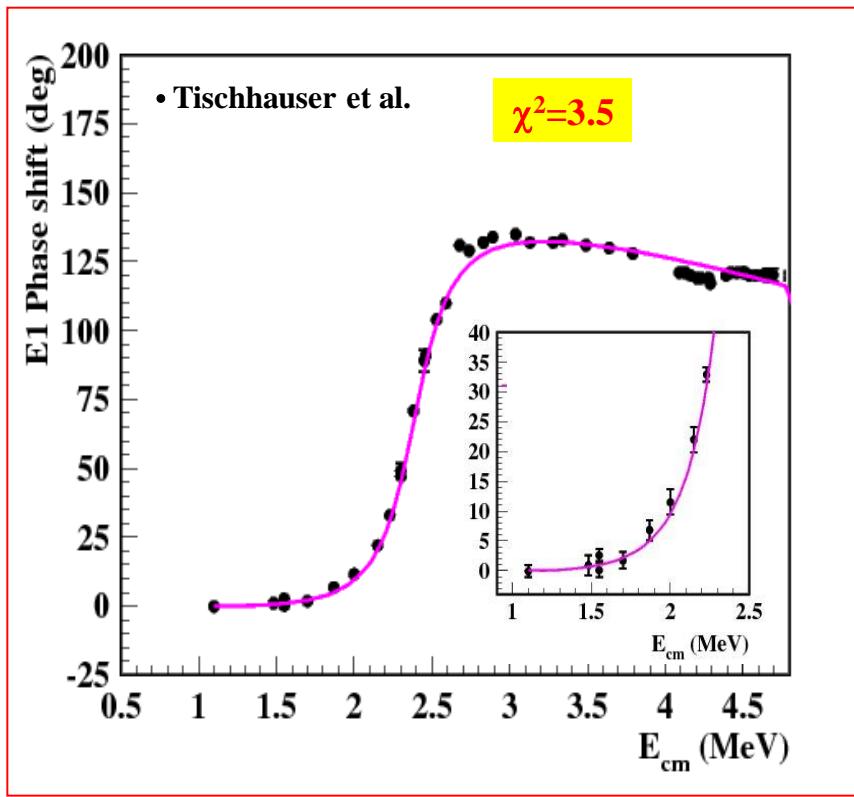
- E2 Component calculation → 4 states
 $\rightarrow 6.92, 9.85, 11.52 \text{ MeV} \rightarrow$ fixed resonance parameters
 \rightarrow Background equivalent state (E_{r4} , $\Gamma_{\alpha 4}$, $\Gamma_{\gamma 4}$)



$$S_{E2}(300 \text{ keV}) = 50 \pm 19 \text{ keV-barn}$$

R matrix calculations– E1 component

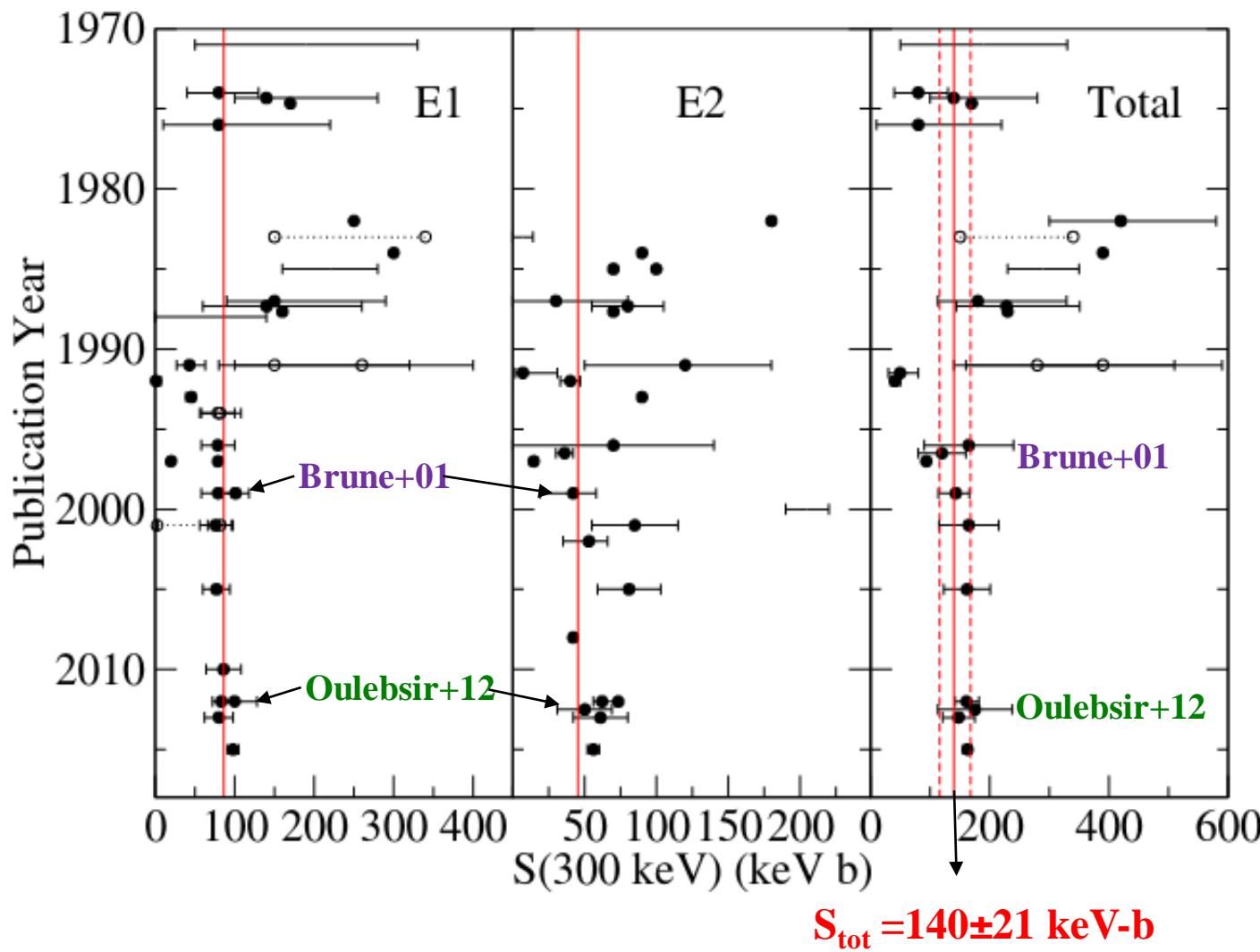
- E1 Component calculation → 3 states
→ 7.12, 9.58 → fixed resonance parameters
→ Background equivalent state (E_{r3} , $\Gamma_{\alpha 3}$, $\Gamma_{\gamma 3}$)



$$S_{E1}(300 \text{ keV}) = 100 \pm 28 \text{ keV-barn}$$

$S_{E1}(0.3 \text{ MeV})$, $S_{E2}(0.3 \text{ MeV})$ & $S_{\text{tot}}(0.3 \text{ MeV})$ over time

Credits: deBoer+17

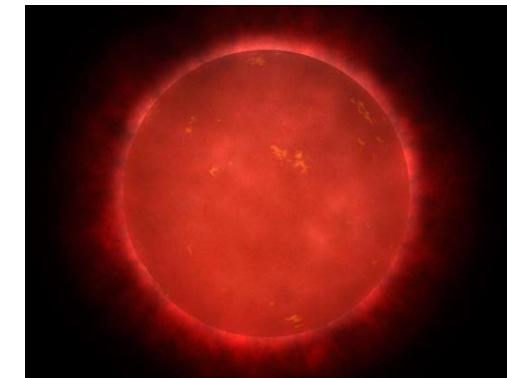


→ From MC analysis of all existing data including ANC's from transfer reactions (deBoer+17)

s-process in rotating metal-poor massive stars

- **s-process** nucleosynthesis → half of the abundance of heavy elements in Universe
- $60 < A < 90$ (weak s-process component) → massive stars $M > 8M_{\odot}$

Core He burning ($T \sim 3 \cdot 10^8$ K, $N_n = 10^6$ cm $^{-3}$)
&
shell Carbon burning ($T \sim 10^9$ K, $N_n = 10^{11}$ cm $^{-3}$)



- Metal-poor massive stars → negligible s-process production (low ^{22}Ne & Fe seed abundance)

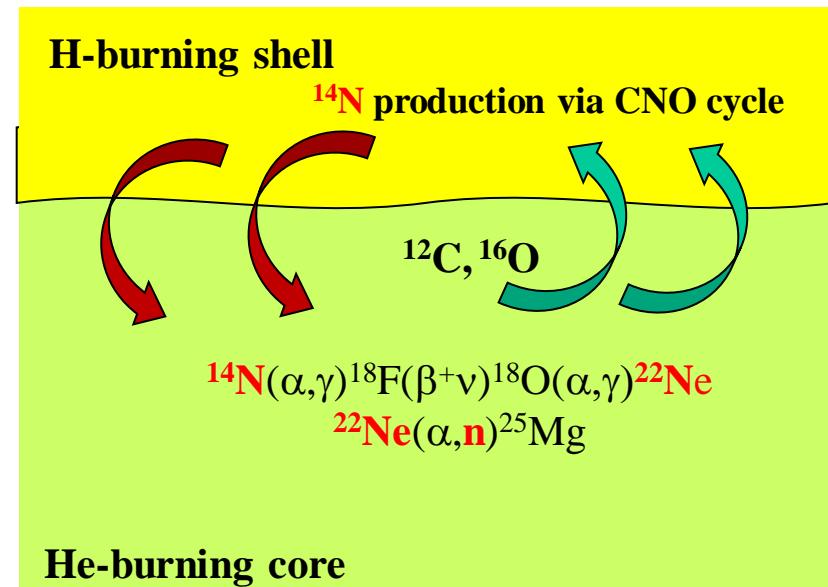
With fast rotation induced mixing
Nishimura+16, Chojalin+18

⇒ ^{22}Ne production in He core strongly enhanced



large production of s-elements between
Strontium & Barium

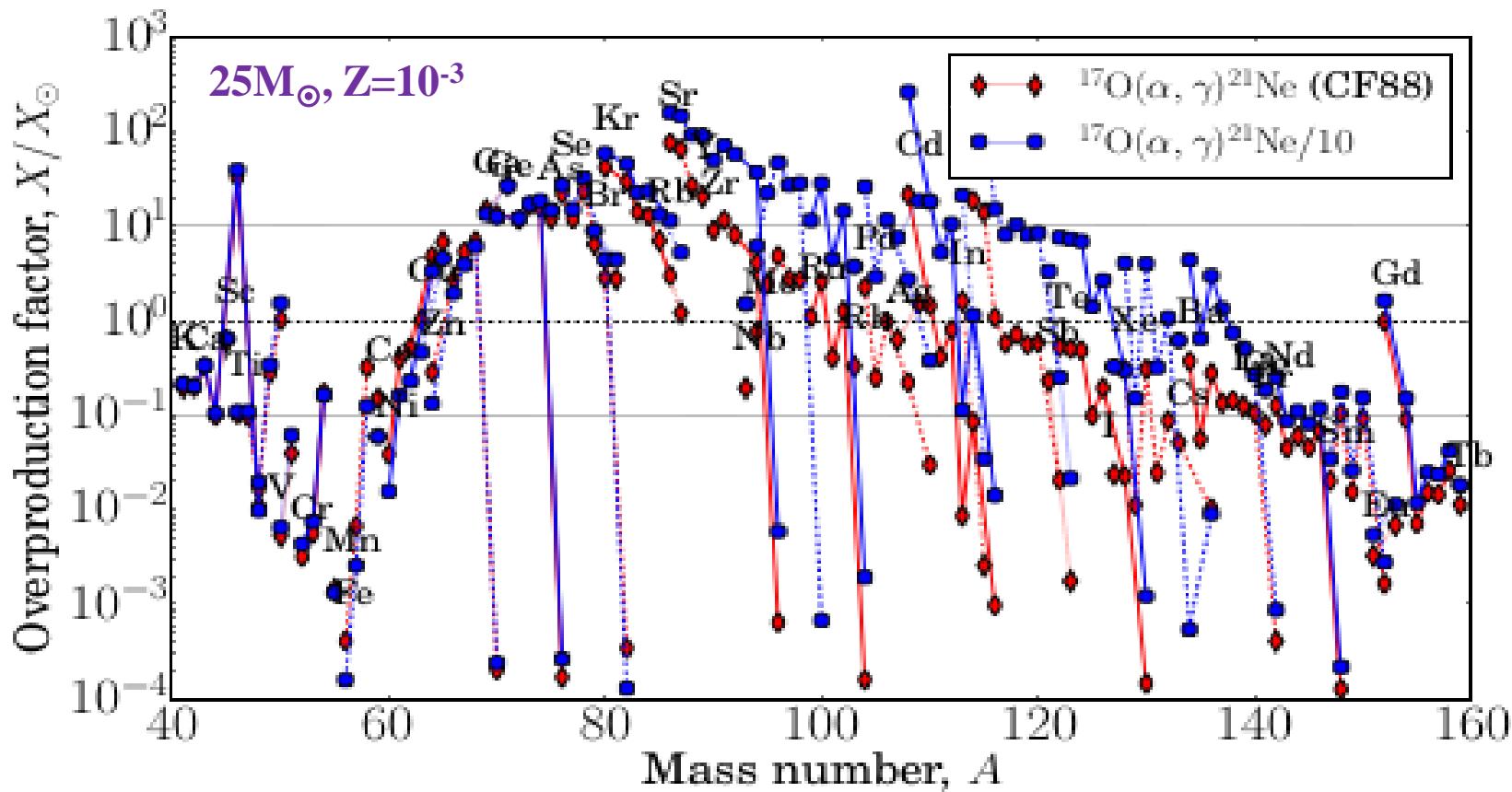
$90 < A < 140$



s-process in rotating metal-poor massive stars

But the final abundances of the enhanced weak s-process strongly depends on:

$^{16}\text{O}(\text{n},\gamma)^{17}\text{O}$ neutron poison effect & $^{17}\text{O}(\alpha,\text{n})/^{17}\text{O}(\alpha,\gamma)$ reaction rate ratio
→ neutron recycling efficiency



Calculation with $^{17}\text{O}(\alpha,\text{n})^{20}\text{Ne}$ Nacre adopted rate & $^{17}\text{O}(\alpha,\gamma)^{22}\text{Ne}$ CF88 rate

Present status on $^{17}\text{O}(\alpha, \text{n})^{20}\text{Ne}$ and $^{17}\text{O}(\alpha, \gamma)^{21}\text{Ne}$

- Core He burning: $T \sim 0.2\text{-}0.3 \text{ GK} \rightarrow E_{\text{c.m.}} \sim 0.297\text{-}0.646 \text{ MeV} \rightarrow E_x = 7.62\text{-}8.00 \text{ in } ^{21}\text{Ne}$
- Shell Carbon burning: $T \sim 1 \text{ GK} \rightarrow E_{\text{c.m.}} \sim 0.783\text{-}1.5 \text{ MeV} \rightarrow E_x = 8.13\text{-}8.85 \text{ in } ^{21}\text{Ne}$

$^{17}\text{O}(\alpha, \text{n})^{20}\text{Ne}$ & $^{17}\text{O}(\alpha, \gamma)^{21}\text{Ne}$ direct measurements:

- Denker+1994, Best+2013 $\Rightarrow 0.63 \leq E_{\text{cm}} \leq 1.8 \text{ MeV}$
- Best +2011, Taggart+2019 $\Rightarrow 0.63 \leq E_{\text{cm}} \leq 1.33 \text{ MeV}$
- Williams+2022

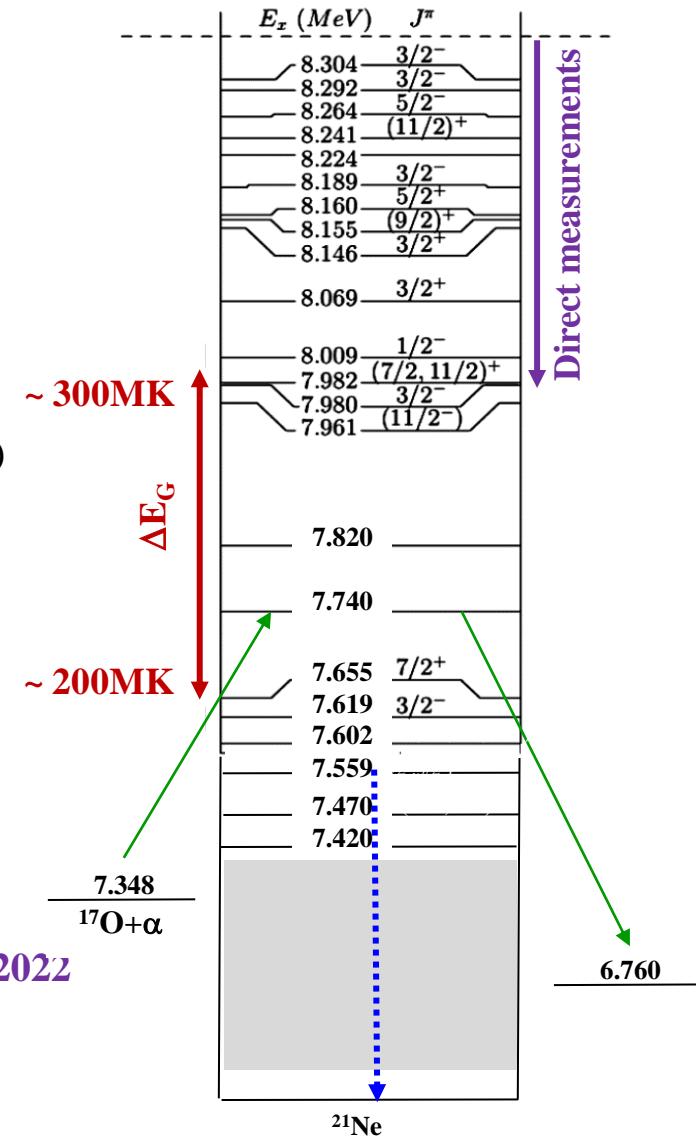
- No direct measurements @ $E_{\text{cm}} < 0.63 \text{ MeV}$ (Core He burning)

- Spectroscopy of ^{21}Ne : E_x , S_α or Γ_α , J^π , $\Gamma_\gamma/\Gamma_{\text{tot}}$, Γ_n

\rightarrow Unknown or poorly known $S_\alpha (\Gamma_\alpha)$ & Γ_n , $\Gamma_\gamma/\Gamma_{\text{tot}}$
 \rightarrow Few have spin-parity assignments

- Neutron transfer reaction $\rightarrow S_n \rightarrow \Gamma_n$ Frost-Schenk+MNRAS2022

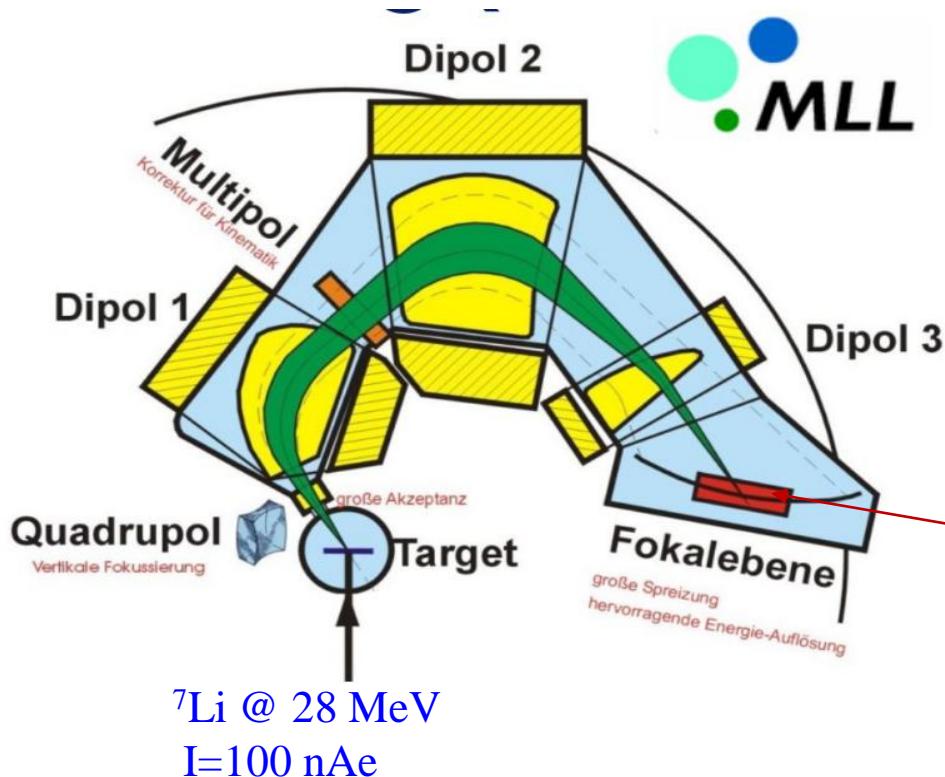
- α -transfer reaction $\rightarrow S_\alpha \rightarrow \Gamma_\alpha$ (present work/MLL-exp)



Study of ^{21}Ne states via $^{17}\text{O}(^{7}\text{Li},\text{t})^{21}\text{Ne}$ α -transfer reaction

Q3D spectrometer (MLL)

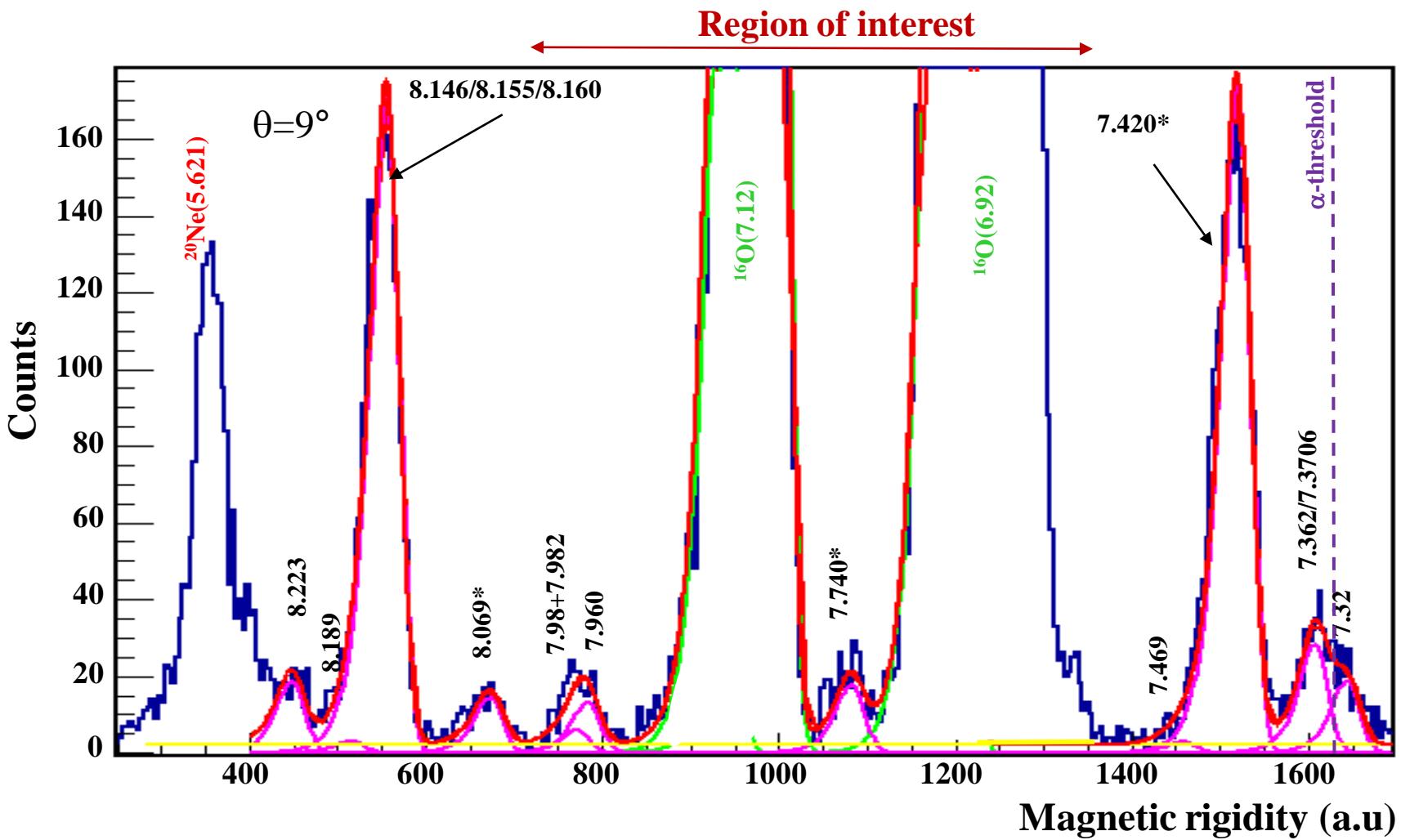
- $\Delta\Omega \sim 6$ to 12.4 msr
- $\Delta E/E \sim 2 \times 10^{-4}$



- Targets (manufactured at LNS Catania):
 - $41 \mu\text{g}/\text{cm}^2$ ^{17}O enriched (35%) WO_3 target on ^{nat}C
 - $39 \mu\text{g}/\text{cm}^2$ natural WO_3 target on ^{nat}C

- $d\sigma/d\Omega$ measurements @ $\theta_{lab} = 6^\circ - 36^\circ \Rightarrow \theta_{cm} \rightarrow 7.5^\circ - 45^\circ$ on enriched target & on natural target for calibration & background evaluation

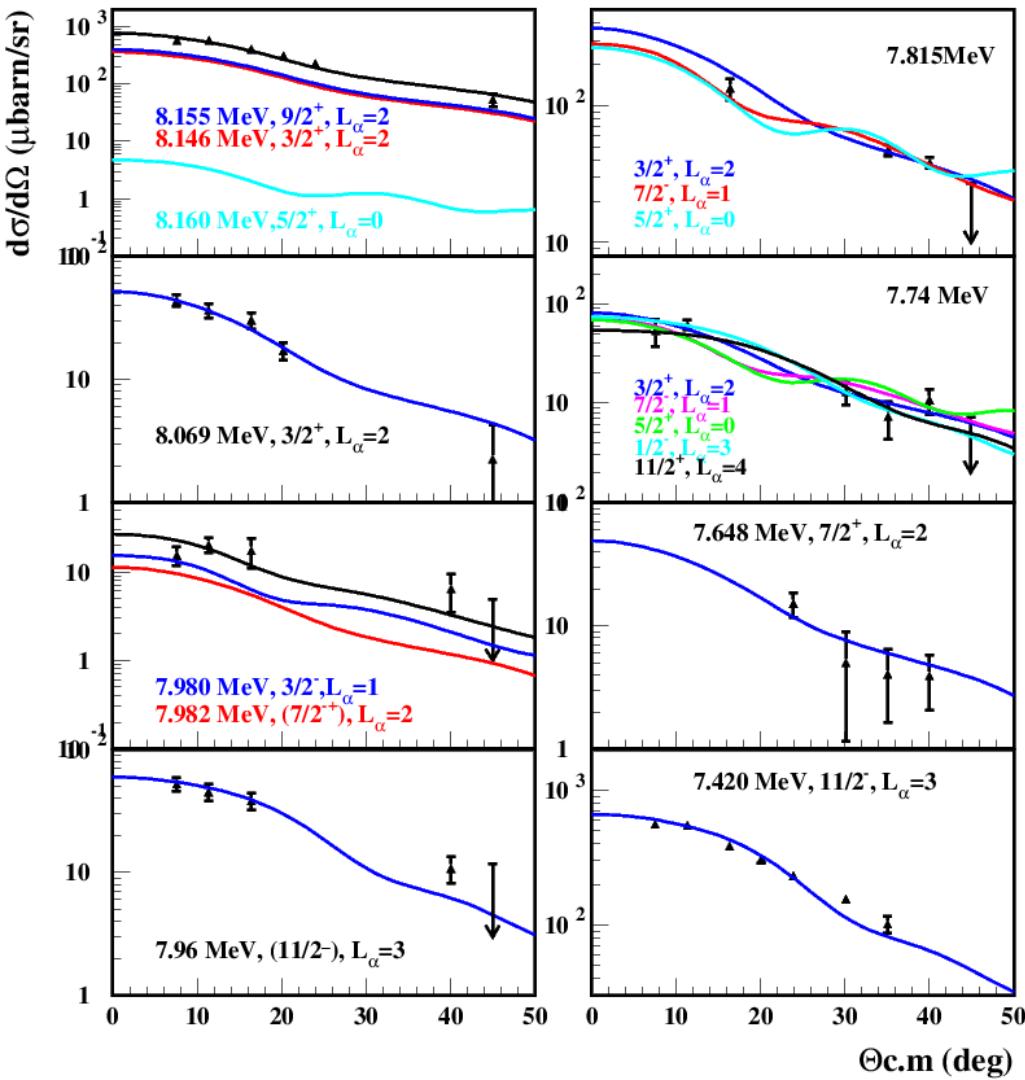
Excitation energy spectrum of ^{21}Ne



- Fit with multiple skewed gaussians with common width & exponential factor

Experimental energy resolution (FWHM) : ~ 30 keV (6°) - 71 keV (36°)

FR-DWBA calculations



- Γ_α uncertainty: 3- 40% (stat), 35% (optical pot)

- Good description of the data by DWBA
↓
Direct transfer mechanism

- Triplet 8.160/8.155/8.146:
Fit with 3 components $\rightarrow S_\alpha$ of 8.146 & 8.160 MeV derived from Γ_α Best+2013
 $\Rightarrow S_\alpha(8.155 \text{ MeV}) = 0.15$ (present work)

- Doublet 7.980/7.982 MeV
Fit with 2 components $\rightarrow S_\alpha$ of 7.98 MeV deduced using $\omega\gamma(\alpha, n)$ Denker+94
 $\Rightarrow S_\alpha(7.982 \text{ MeV}) = 0.005$ (present work)

- 7.815 MeV
 \rightarrow Best χ^2 for $L_\alpha=0,1$ & good for $L_\alpha=2$
 $\rightarrow L_\alpha=0 \rightarrow S_\alpha=0.66$ (unlikely)

$$S_\alpha \rightarrow \Gamma_\alpha = 2P_l \frac{\hbar^2 R}{2\mu} S_\alpha |\phi(R)|^2$$

$^{17}\text{O}(\alpha, \text{n})$ & $^{17}\text{O}(\alpha, \gamma)$ reaction rates $(\alpha, \text{n})/(\alpha, \gamma)$ rate ratio

Rates calculations:

RateMC code **Longland+13**

- For $E_\alpha < 721 \text{ keV}$ & $E_\alpha=807 \text{ keV}$:
 $\rightarrow \Gamma_\alpha$ (**present work**)

Γ_α (7.81 MeV) determined with $L_\alpha=1$

Γ_α (7.74 MeV) determined with $L_\alpha=0$

- $\rightarrow \Gamma_n$ **Frost-Schenk+2022**

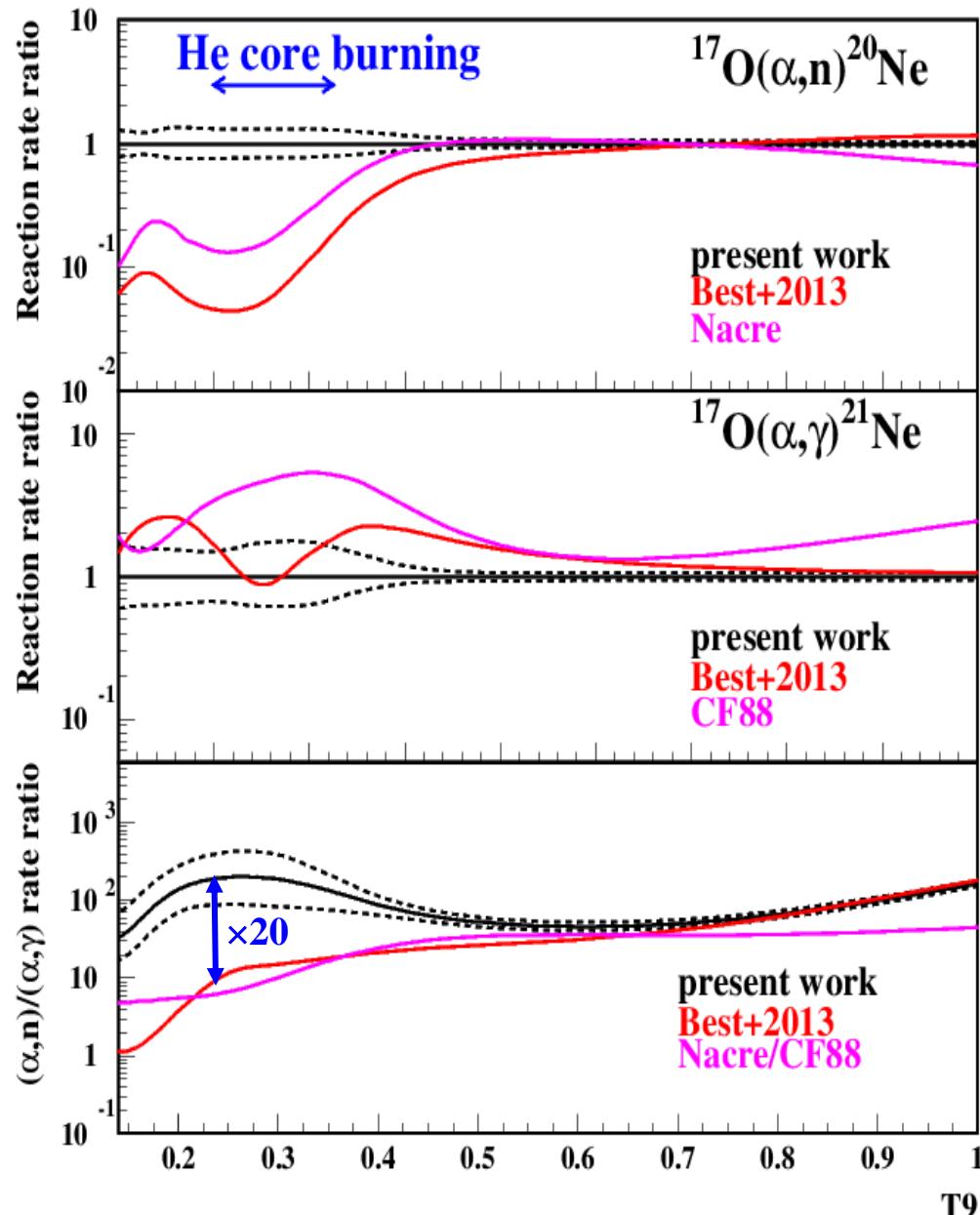
- For $E_\alpha \geq 721 \text{ keV}$: Γ_α & Γ_n (**Best+2013** direct measurement)

- Γ_γ from:

\rightarrow systematics of $\langle \tau \rangle_{\text{meas}}$ (**Rolfs+72**)

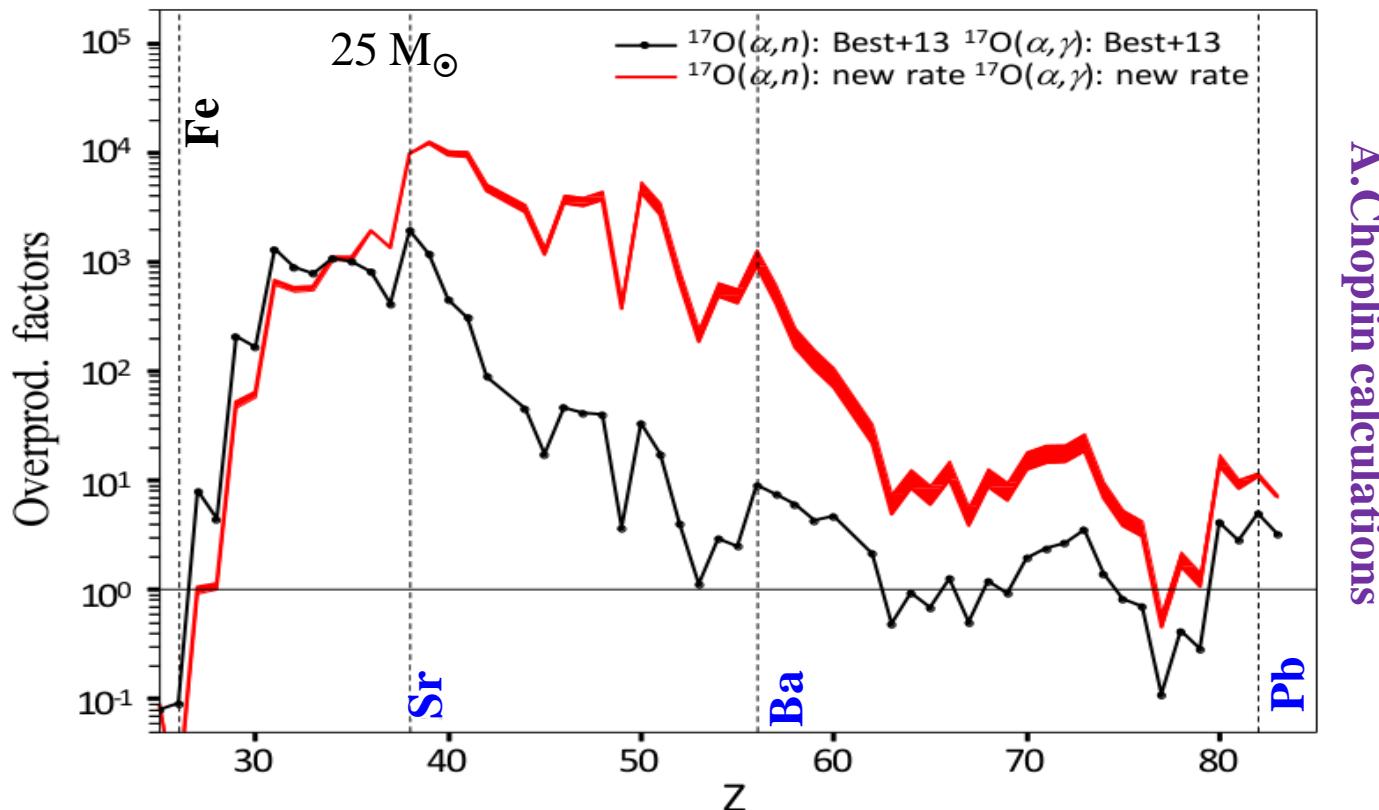
$\rightarrow \omega\gamma(\alpha, \gamma)$ **Williams+2022** combined with
present Γ_α & Γ_n (**Frost-Schenk+22**)

→ Better neutron efficiency recycling
with a factor of about **20** with the
present recommended rates than
Best+13 rates



Impact on the s-process in rotating poor-metal massive stars

- One-zone nucleosynthesis calculation mimicking the core He-burning phase of a low metallicity rotating massive star ($Z=0.001$)



- Large enhancement (>1.5 dex) of elements $40 < Z < 60$ with the present **new rates** in comparison to **Best+13** rates
- Two order of magnitude on **Barium**: largest effect

Collaboration

$^{12}\text{C}(^7\text{Li},\text{t})^{16}\text{O}$ experiment

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L. Gaudefroy **(GANIL-Caen)**

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$^{17}\text{O}(^7\text{Li},\text{t})^{21}\text{Ne}$ experiment

F.H., S. Harrouz, N. de Séréville, A. Meyer
IJCLab-Orsay

P. Adsley
(Ithemba)

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Romano, A. Tumino
LNS, Catania

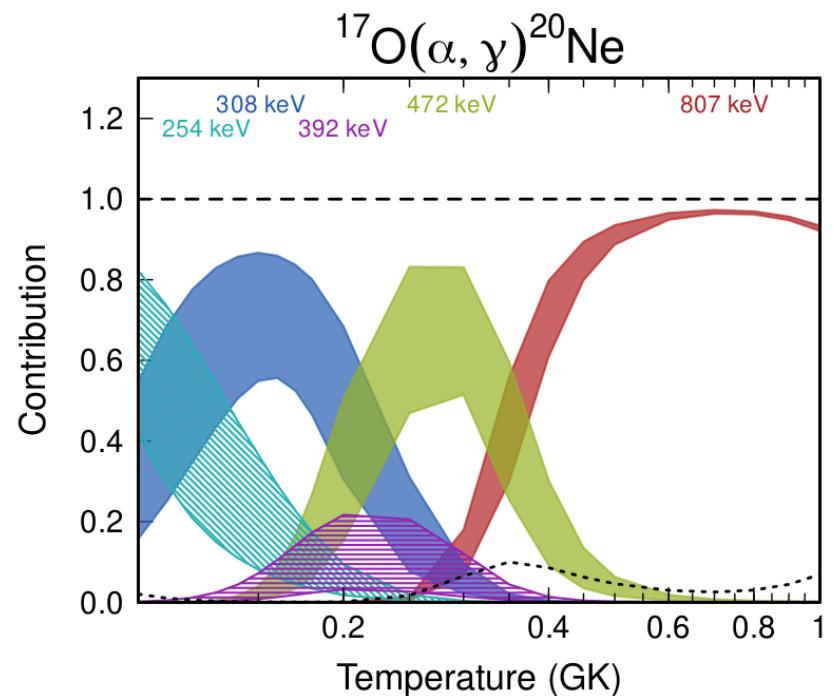
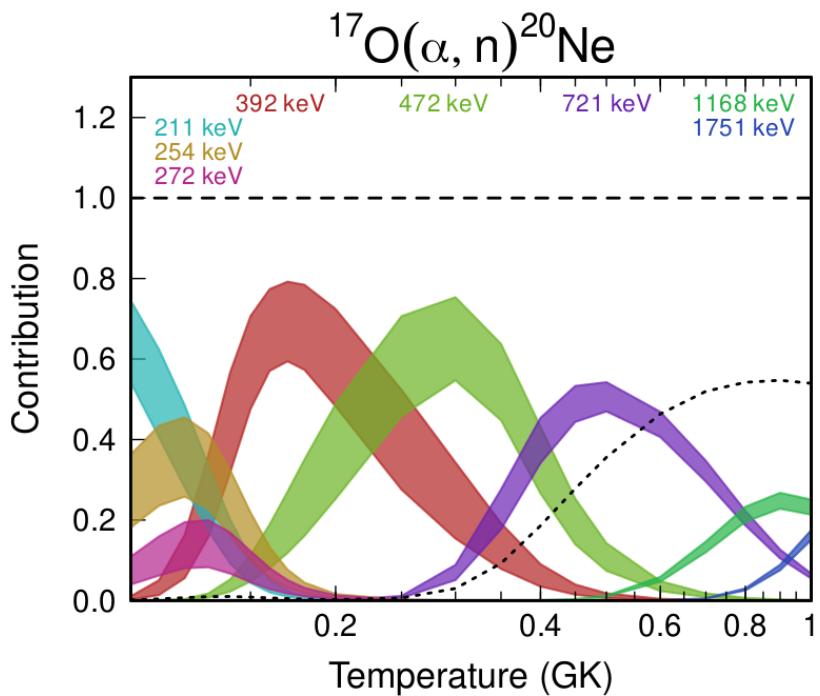
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Resonances contribution to the rates



- Er=392 (Ex=7.74 MeV) & 472 keV (Ex=7.82 MeV) contribute the most to the (α, n) rate
- Er=308 keV (Ex=7.655 MeV) & 472 keV (Ex=7.82 MeV) contribute the most to the (α, γ) rate

→ Ex=7.74 MeV unknown L_α & J^π

→ Ex=7.82 MeV $L_\alpha = 0, 1, 2$ & $L_n = 2, 3 \Rightarrow J^\pi = 5/2^+, 7/2^-, 3/2^+$

Key resonances

$^{17}\text{O}(\alpha, \text{n})$ & $^{17}\text{O}(\alpha, \gamma)$ reaction rates $(\alpha, \text{n})/(\alpha, \gamma)$ rate ratio

Worst case: 7.810 MeV $\mathbf{l}_\alpha=2 \rightarrow 3/2^+$ or $7/2^+$ & no 7.74 MeV

