

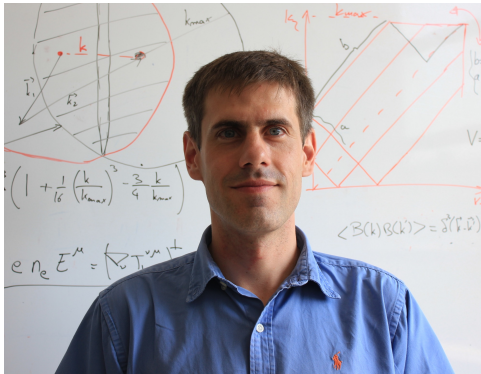
# Primordial nucleosynthesis (with Pierre)

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# The three observational evidences for the big bang model

## 1. The expansion of the Universe

Galaxies move away from each other according to Hubble's law:

$V = H_0 \times D$  with  $H_0 \approx 70$  km/s/Mpc, the Hubble parameter (or "constant").

More precisely distances  $\propto a(t)$ , the cosmological scale factor

## 2. The Cosmic Microwave Background radiation (CMB)

A black body radiation at 2.7 K corresponding to the redshifted spectrum emitted when the universe became transparent

## 3. Primordial nucleosynthesis

Reproduces the "light-elements" ( $^4\text{He}$ ,  $^2\text{H}$  or  $\text{D}$ ,  $^3\text{He}$  and  $^7\text{Li}$ ) primordial abundances over a range of nine orders of magnitudes (predictions vs observations).

# Determination of primordial abundances

Primordial abundances :

1) Observe a set of primitive objects born when the Universe was young

- $^4\text{He}$  in H II (ionized H) regions of blue compact galaxies
- $^3\text{He}$  in H II regions of *our* Galaxy (not very useful)
- $\text{D}$  in remote **cosmological clouds** (i.e. at high redshift) on the line of sight of quasars
- $^7\text{Li}$  at the surface of low metallicity\* stars in the halo of our Galaxy

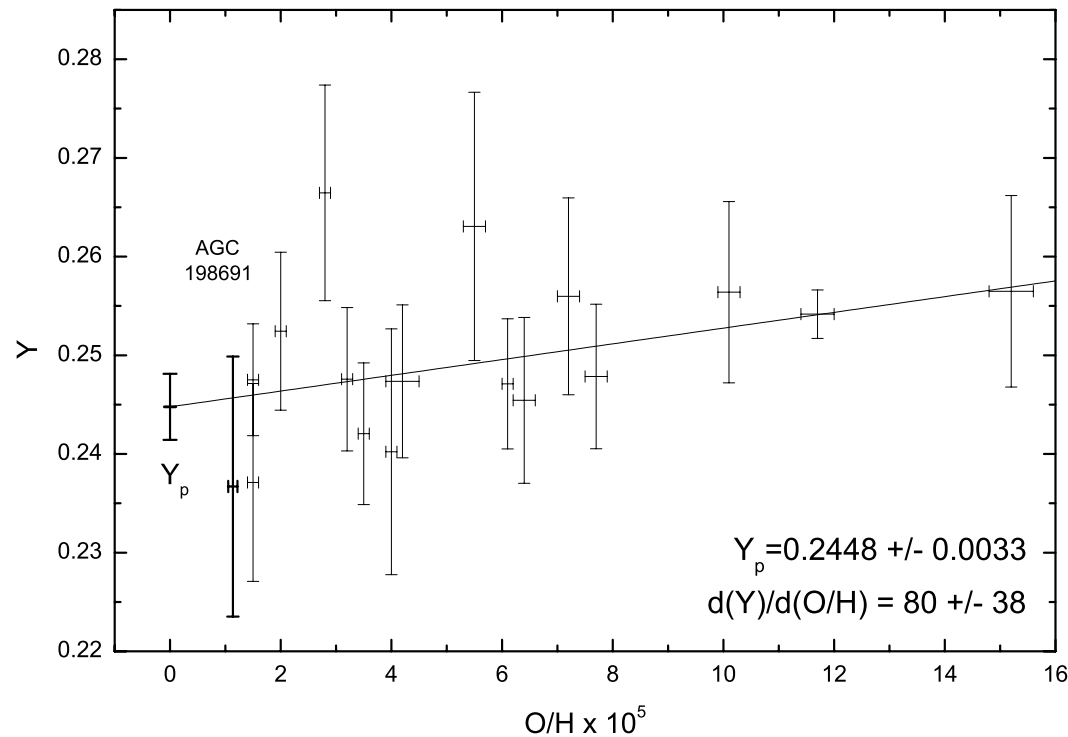
2) Possibly extrapolate to zero metallicity\* :  $\text{Fe}/\text{H}, \text{O}/\text{H}, \dots \rightarrow 0$

\*In astrophysics: “metals” = everything beyond helium

Notation :  $[\text{X}/\text{H}] \equiv \log(\text{X}/\text{H}) - \log(\text{X}_{\odot}/\text{H}_{\odot})$ ,  $\text{X}=\text{Fe}, \text{O}, \dots$

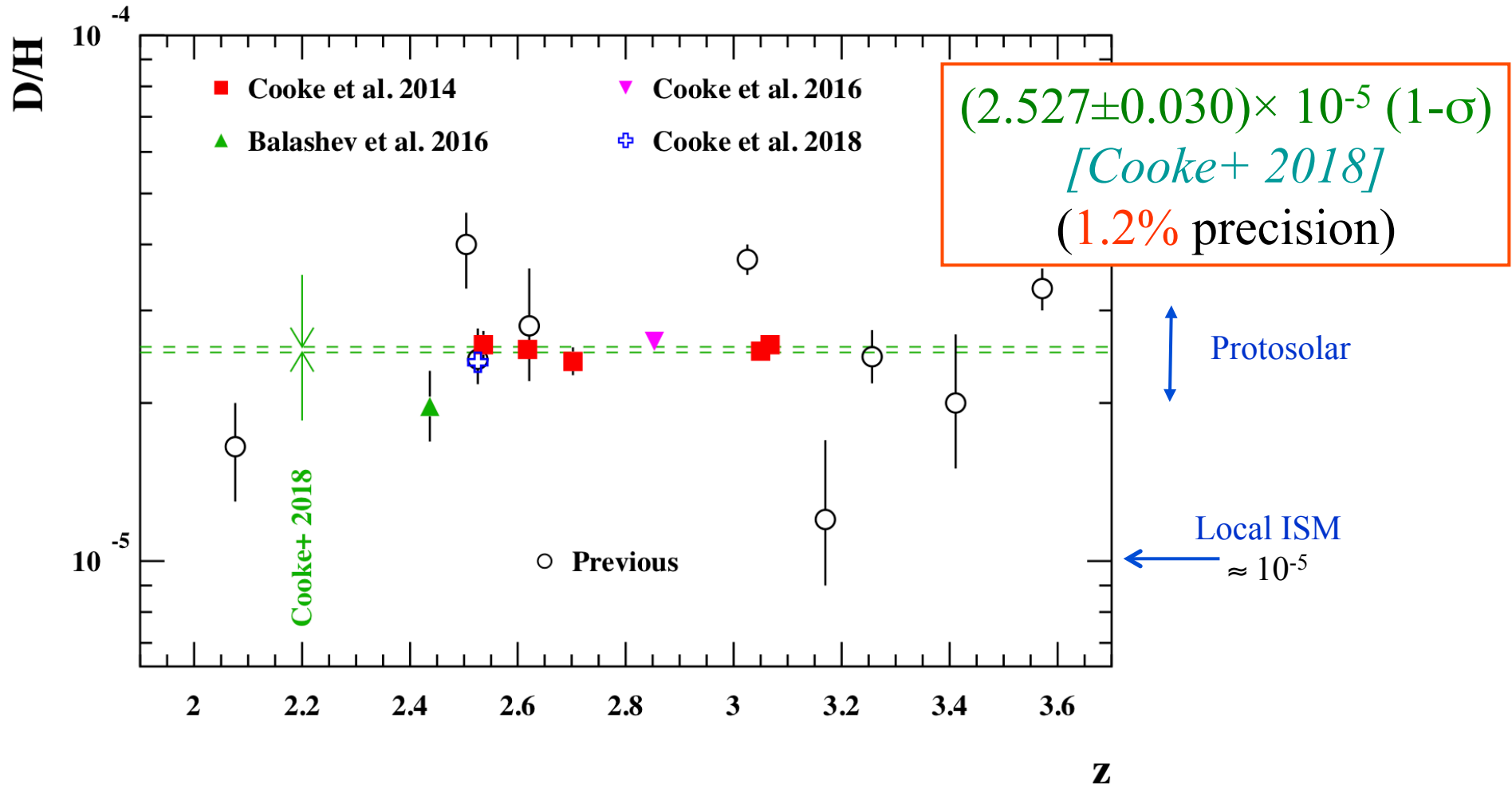
# $^4\text{He}$ observations in blue compact galaxies

Observations:  $^4\text{He}$  from a sample of H II regions in blue compact galaxies [Izotov+ 2007;2014] with additional extremely metal poor galaxies (AGC 198691) [Aver+ 2021;2022]



$Y_p = 0.2453 \pm 0.0034$  ( $^4\text{He}$  mass fraction) [Aver+ 2021;2022]  
(1.4% precision)

# D/H observations in cosmological clouds

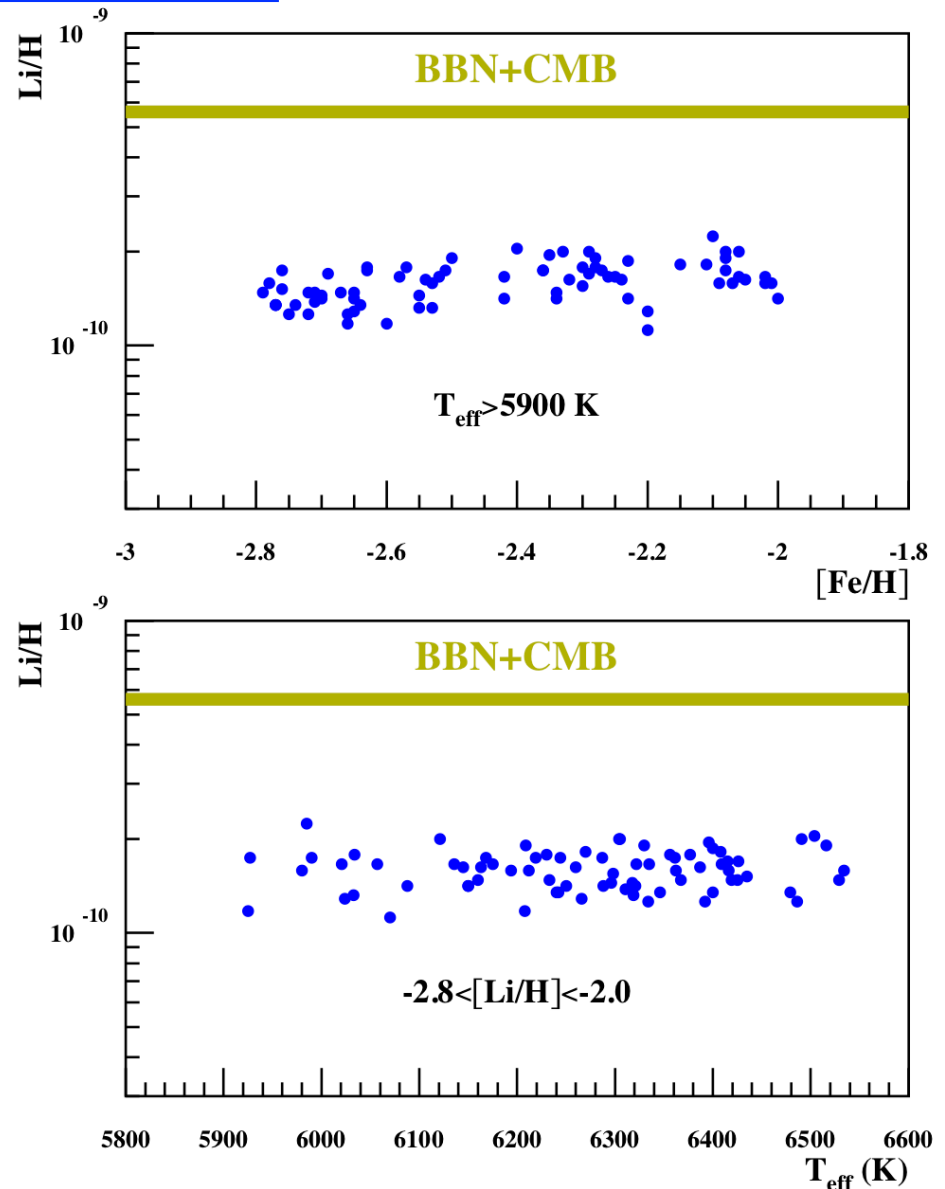


Burles & Tytler 1998; O’Meara+ 2001, 2006; Pettini+ 2001, 2008, 2012; Kirkman+ 2003, Crighton+ 2004; Srianand+ 2010; Cooke+ 2011; Fumagalli+ 2011; Cooke+ 2016; 2018

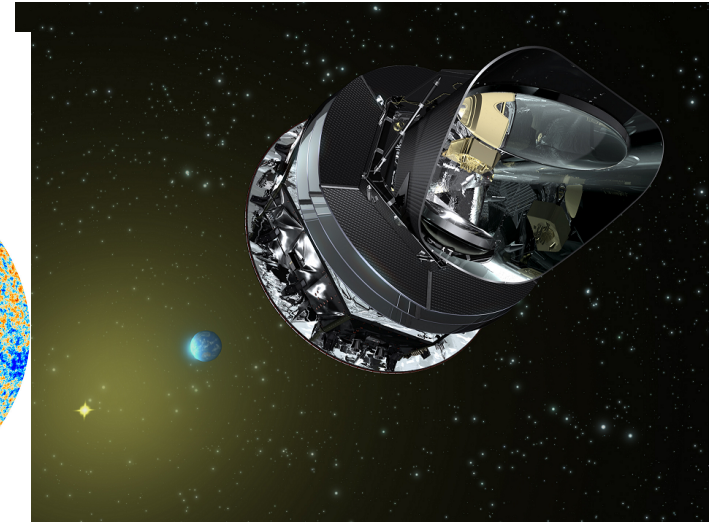
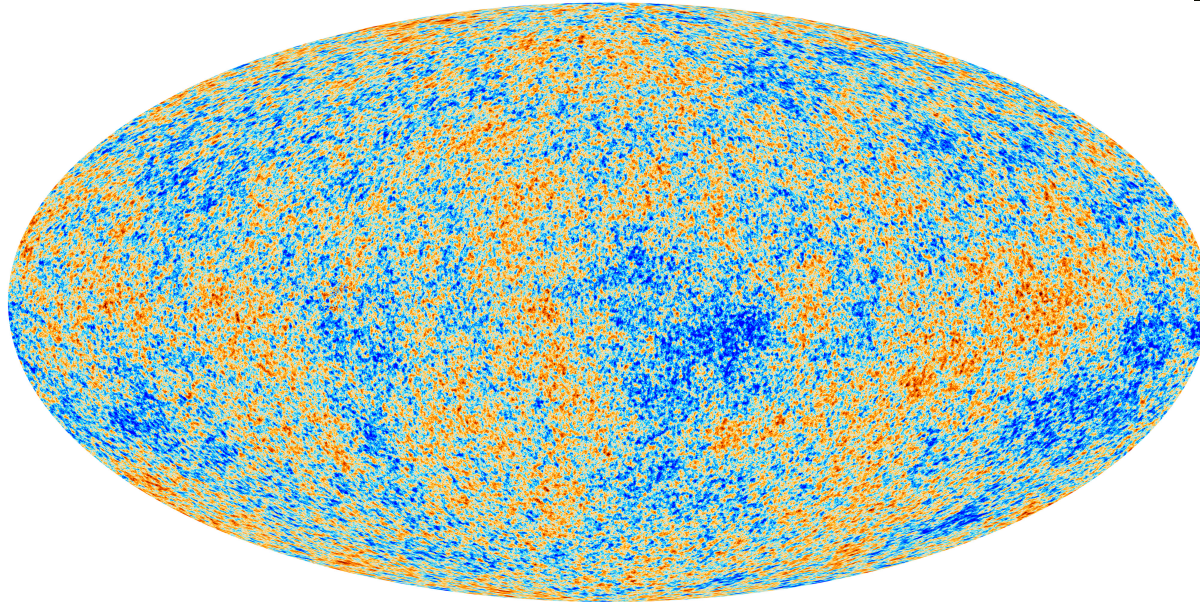
# Primordial Li from observations

[M. Spite, priv. comm.]

- Oldest (low metallicity) stars in galactic halo
- Lifetime of  $M < 0.9 M_{\odot}$  stars  $> 15$  Gy
- For  $T_{eff} > 5900$  K, no deep convection and no Li surface depletion (?)
- $\text{Li}/\text{H} = (1.58 \pm 0.35) \times 10^{-10}$   
[Sbordone et al. 2010]



# Density components of the Universe



$$\Omega \equiv \frac{\rho}{\rho_{\text{Critical}}}$$

Some  $\Omega$  values [Aghanim+ 2020 (Planck)]

Radiation (CMB)	$\Omega_R$	$5 \cdot 10^{-5}$
Visible matter	$\Omega_L$	$\approx 0.003$
Baryons	$\Omega_b$	0.049
Dark Matter	$\Omega_c$	0.262
Vacuum	$\Omega_\Lambda$	0.680
Total	$\Omega_T$	$\approx 1.0$

$$\Omega_b h^2 = 0.02242 \pm 0.00014$$

(0.6% precision)

[Alam+ 2017; Aghanim+ 2020 (Planck)]

$$h = H_0 / 100 \text{ km/s/Mpc}$$

$$\approx 0.69 - 0.73$$

# Dynamics of the expanding Universe

Einstein equation & Friedmann-Lemaître-Robertson-Walker metrics

$$ds^2 = g_{\alpha\beta} dx^\alpha dx^\beta = dt^2 - a^2(t) \left( \frac{dr^2}{1-kr^2} + r^2(d\theta^2 + \sin^2\theta d\phi^2) \right) \quad a(t): \text{scale factor}$$

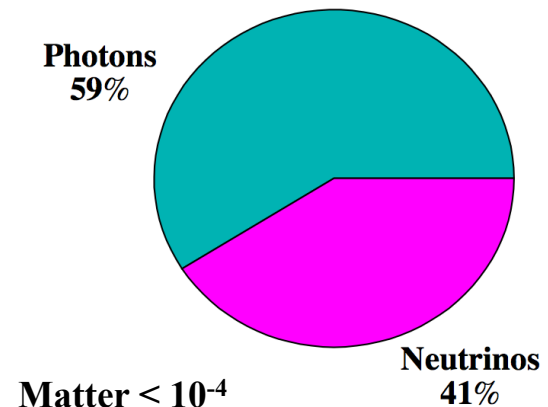
Friedmann equation :

$$\left( \frac{\dot{a}}{a} \right)^2 \equiv H^2 = \frac{8\pi G}{3} (\rho_R + \cancel{\rho_M} + \cancel{\rho_\Lambda}) - \cancel{\frac{k}{a^2}}$$

EoS:  $p = \text{pressure} \equiv w \times \rho \Rightarrow \boxed{\rho \propto a^{-3(1+w)}}$

$$w = \begin{cases} 0 \text{ (matter)} & \Rightarrow a^{-3} \\ 1/3 \text{ (radiation)} & \Rightarrow a^{-4} \\ -1 \text{ (\Lambda, dark energy)} & \Rightarrow a^0 \\ \text{Curvature} & \Rightarrow a^{-2} \end{cases}$$

....and then ( $a \approx 10^{-8}$ )





# Dynamics of the expanding Universe

Cosmological distances  $\propto a \equiv (1+z)^{-1}$  ( $z = \text{redshift}$ )

Rate of expansion  $\propto$   
(radiation energy density) $^{1/2}$

$$\textcircled{1} \quad \frac{1}{a} \frac{da}{dt} \propto \sqrt{\rho_{\text{rad}}^{e\gamma}(T)} \propto \sqrt{g_*^{e\gamma\nu}(T) T^2}$$

“Radiation”:  $\gamma$ ,  $(e^-)$ ,  $\nu_x$  and antiparticles

$T_\nu = T$  for  $T \gg 1 \text{ MeV}$

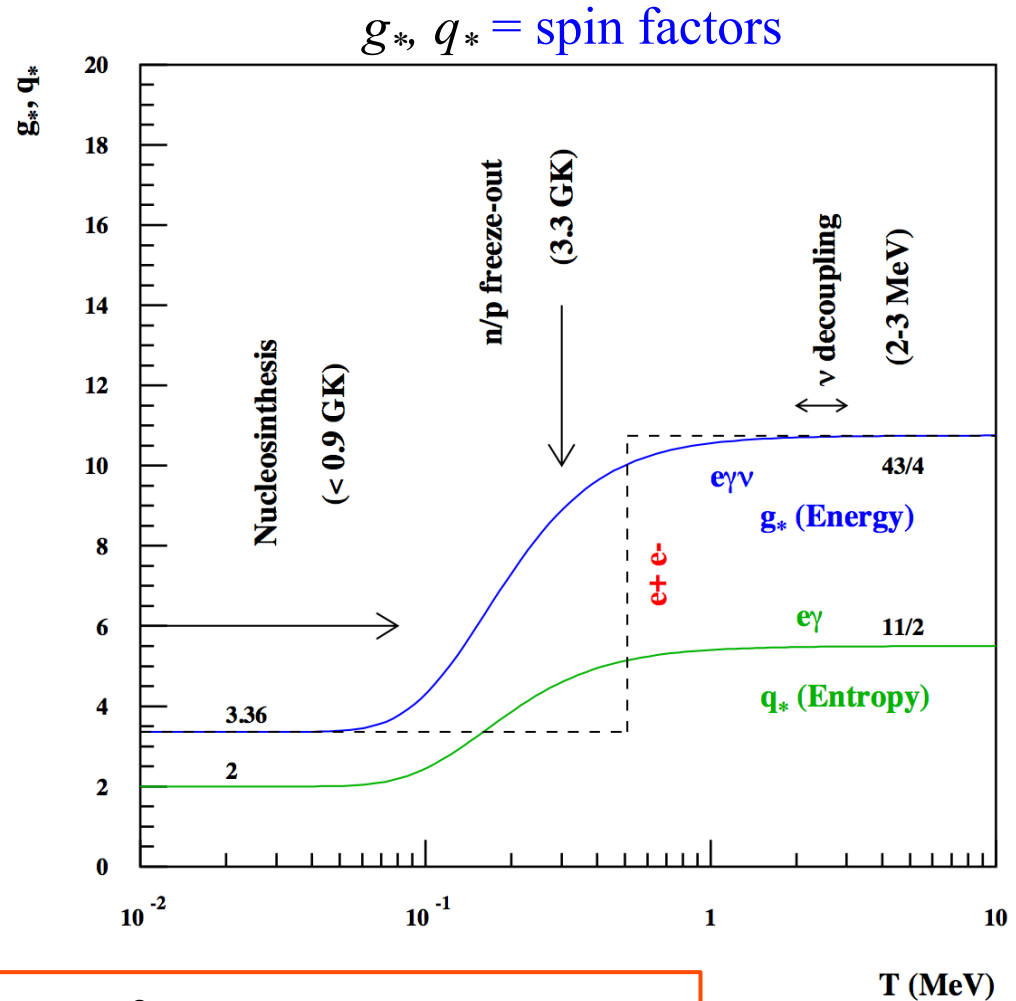
$$\textcircled{2} \quad a^3 T_\nu^3 = C_{\text{ste}}$$

Entropy constant

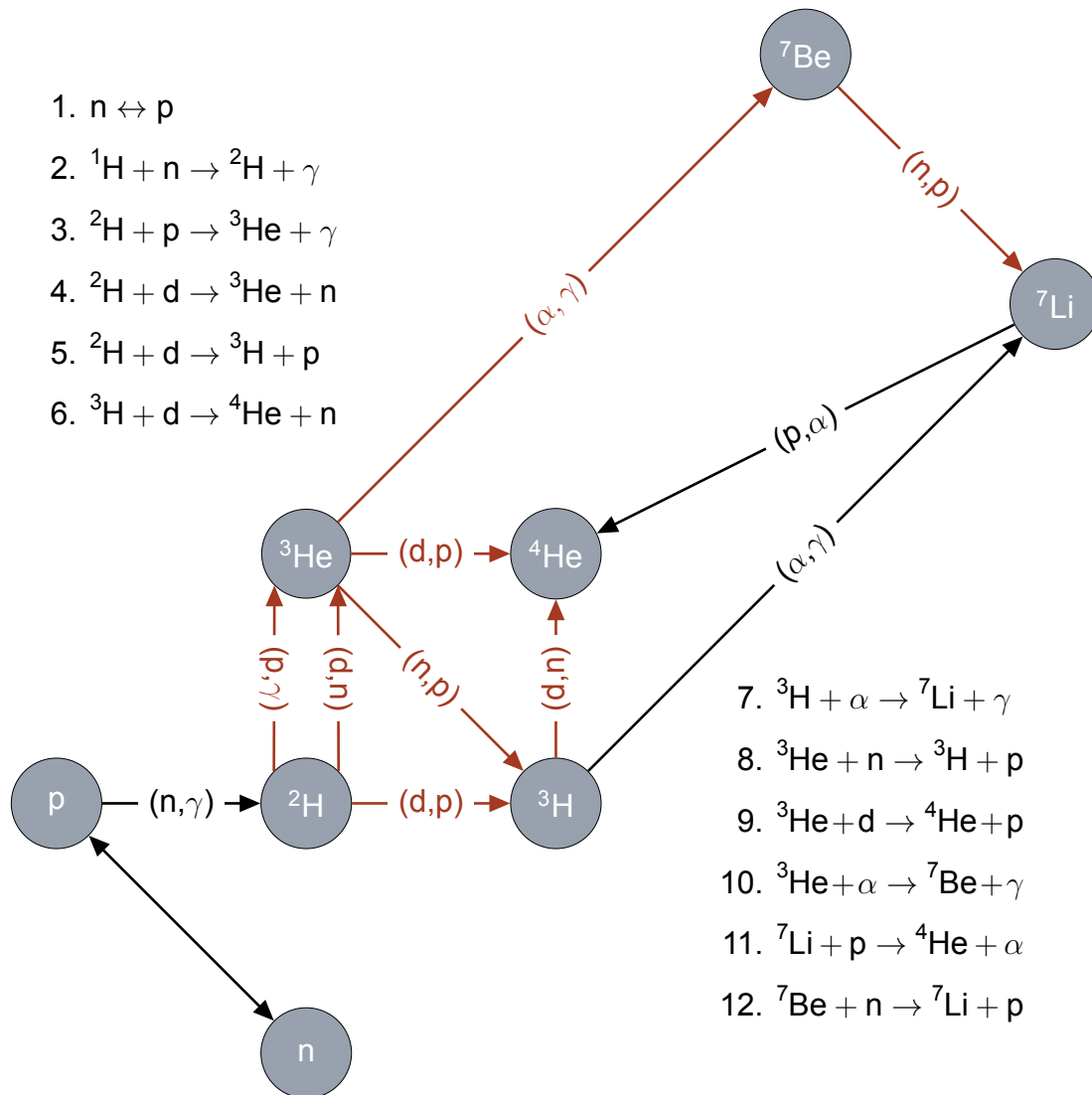
$$\textcircled{3} \quad a^3 q_*^{e\gamma}(T) T^3 = C_{\text{ste}}$$

+ a few small corrections

$$\textcircled{1} + \textcircled{2} + \textcircled{3} \Rightarrow \rho_b(t) \propto \Omega_b a^{-3}(t), T(t) \text{ and } T_\nu(t)$$



# Big Bang Nucleosynthesis calculations



Needs:

- Density  $\rho_b(t)$ , ions and photons  $T(t)$  and neutrino  $T_\nu(t)$  temperatures as a function of time
- Reaction rates ( $\sim 400$  to reach CNO)

New BBN public code:

- *PRIMAT* [Pitrou+ 2018]

Other codes:

- *PARthENoPE* [Gariazzo+ 2022]
- *AlterBBN* [Arbey+ 2020]
- *IL+MN* [Yeh, Olive & Fields]

give the same results when **same reaction rates** are used



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# Precision big bang nucleosynthesis with improved Helium-4 predictions



Cyril Pitrou<sup>a,b,\*</sup>, Alain Coc<sup>c</sup>, Jean-Philippe Uzan<sup>a,b</sup>, Elisabeth Vangioni<sup>a,b</sup>

Monthly Notices  
of the

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MNRAS **502**, 2474–2481 (2021)

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## A new tension in the cosmological model from primordial deuterium?

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## n ↔ p weak reaction rates (<sup>4</sup>He)

$$n \rightarrow p : n \rightarrow p + e^- + \bar{\nu}_e \quad n + e^+ \rightarrow p + \bar{\nu}_e \quad n + \nu_e \rightarrow p + e^-$$

$$p \rightarrow n : p + e^- + \bar{\nu}_e \rightarrow n \quad p + \bar{\nu}_e \rightarrow n + e^+ \quad p + e^- \rightarrow n + \nu_e$$

$$\lambda_{n \leftrightarrow p} = C \int \int (\text{phase space}) \times (\text{e distribution}) \times (\nu_e \text{ distribution}) \, dE$$

+ “some small corrections”

$$\lambda_{n \rightarrow pev} = C \int_1^q \frac{\varepsilon(\varepsilon - q)^2 (\varepsilon^2 - 1)^{1/2} \, d\varepsilon}{[1 + \exp(-\varepsilon z)] \{1 + \exp[(\varepsilon - q)z_\nu]\}} \quad T \rightarrow 0 \quad \boxed{\frac{1}{\tau_n} = C \int_1^q \varepsilon(\varepsilon - q)^2 (\varepsilon^2 - 1)^{1/2} \, d\varepsilon}$$

$$(q \equiv Q_{np}/m_e, \varepsilon \equiv E_e/m_e, z \equiv m_e/T_p, z_\nu \equiv m_e/T_\nu)$$

- Experimental neutron lifetime?  $\Delta Y_p = +0.0002 \times \Delta \tau_n$  (s)  $\tau_n = 879.4 \pm 0.6$  s [Zyla+ 2020]
- Calculation of the “small corrections” (radiative corrections, finite nucleon mass corrections, finite temperature radiative corrections, weak-magnetism, QED plasma effects, incomplete neutrino decoupling) [Pitrou+ 2018]  $\Delta Y_p \approx 1.8\%$ ,  $\Delta(D/H) \approx 1.5\%$
- Precision e.g.  $N_{\text{eff}} = 3.044$  (effective number of neutrino families)

# Evaluation of reaction rates for BBN















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A compilation of charged-particle induced thermonuclear reaction rates  
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- 2  2004ApJ...600..544C 2004/01 cited: 310     
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- 3  2004ADNDT..88..203D 2004/09 cited: 252     
Compilation and R-matrix analysis of Big Bang nuclear reaction rates  
Descouvemont, Pierre; Adahchour, Abderrahim; Angulo, Carmen *and 2 more*
- 4  2018PhR...754...1P 2018/09 cited: 245     
Precision big bang nucleosynthesis with improved Helium-4 predictions  
Pitrou, Cyril; Coc, Alain; Uzan, Jean-Philippe *and 1 more*

## Compilation and **R-matrix** analysis of Big Bang nuclear reaction rates

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THE ASTROPHYSICAL JOURNAL, 600:544–552, 2004 January 10

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## UPDATED BIG BANG NUCLEOSYNTHESIS COMPARED WITH *WILKINSON MICROWAVE ANISOTROPY PROBE* OBSERVATIONS AND THE ABUNDANCE OF LIGHT ELEMENTS

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AND

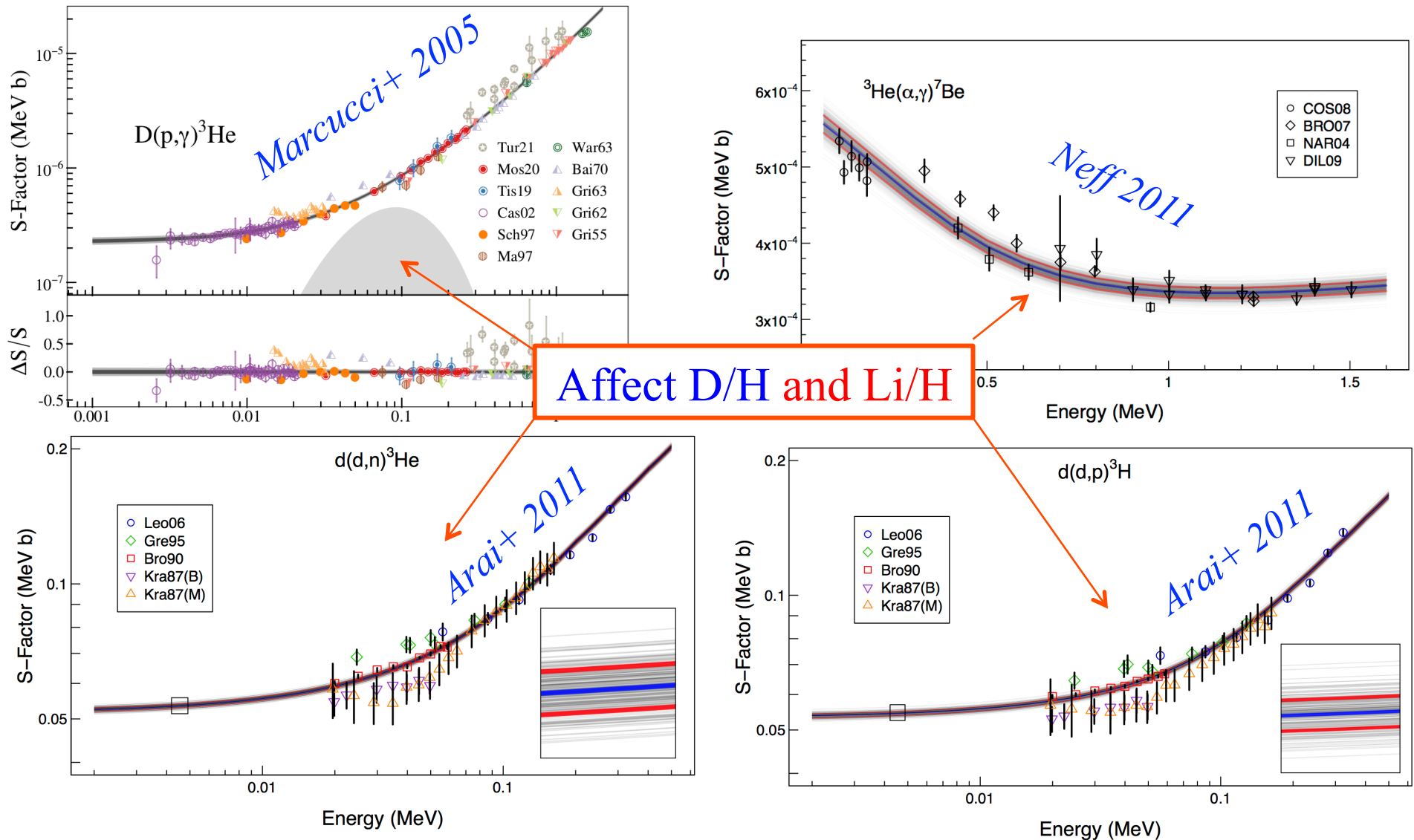
CARMEN ANGULO

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*Received 2003 July 23; accepted 2003 September 24*

# Re-evaluation of main reaction rates

Bayesian analyses of reaction rates for BBN :  $D(p,\gamma)^3\text{He}$ ,  $D(d,n)^3\text{He}$ ,  $D(d,p)^3\text{H}$  and  $^3\text{He}(\alpha,\gamma)^7\text{Be}$  [*Moscoso+ 2021; Gómez Iñesta+2017; Iliadis+ 2016*]

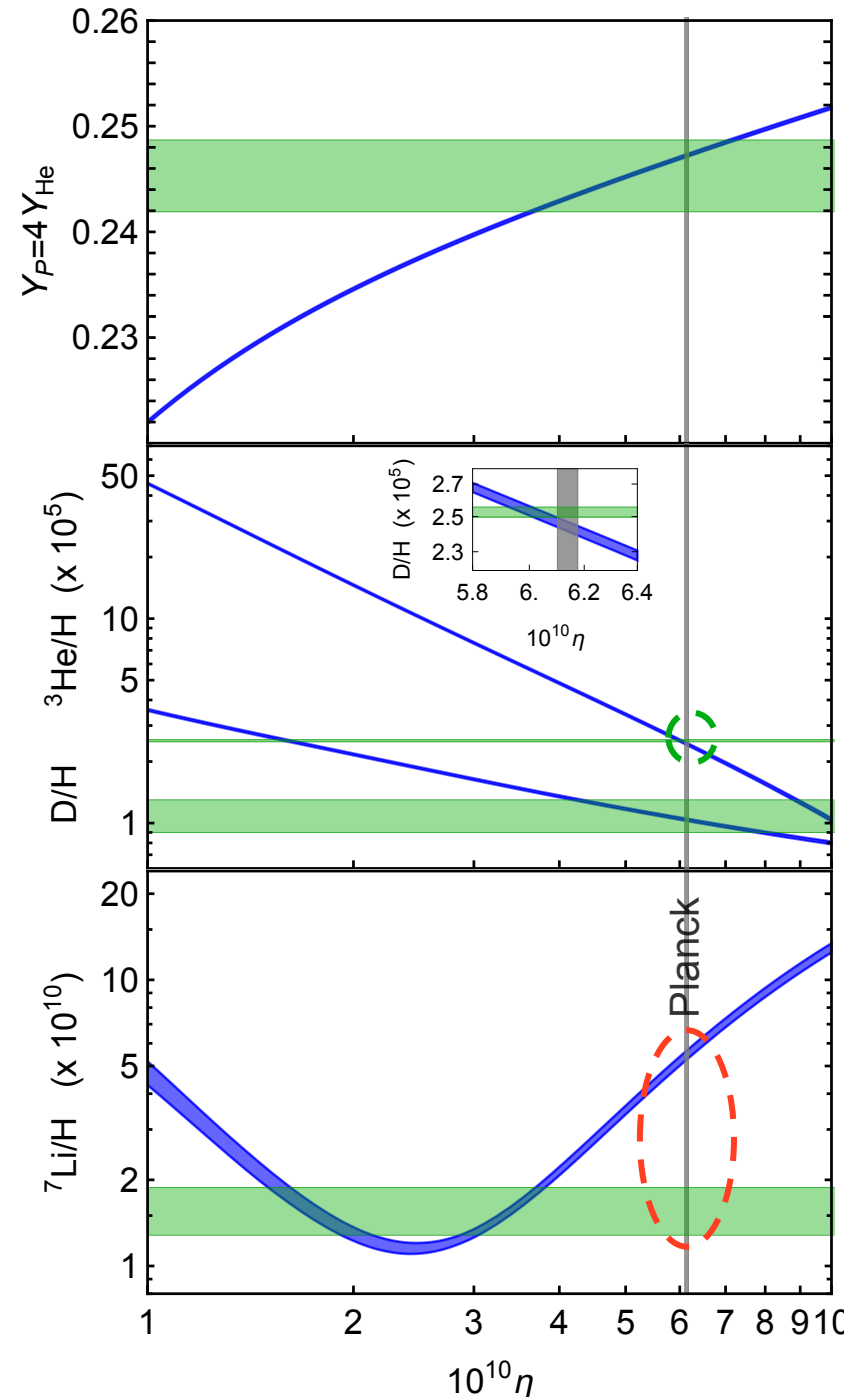


## Comparison between observed and calculated abundances

Limits ( $1-\sigma$ ) obtained by Monte-Carlo from reaction rate uncertainties [Pitrou + 2018;2021]

Concordance (?) BBN, spectroscopy and CMB

- $\Omega_B h^2$  [Planck: Alam+ 2017; Aghanim+ 2020 ]
- $^4\text{He}$  [Aver+ 2021; 2022]
- $\text{D}$  [Cooke+ 2018]
- $^3\text{He}$  [Bania et al. (2002)]
- $^7\text{Li}$  [Sbordone+ 2010] : difference of a factor of  $\approx 3$  between calculated (BBN +CMB) and observed (Spite plateau) primordial lithium

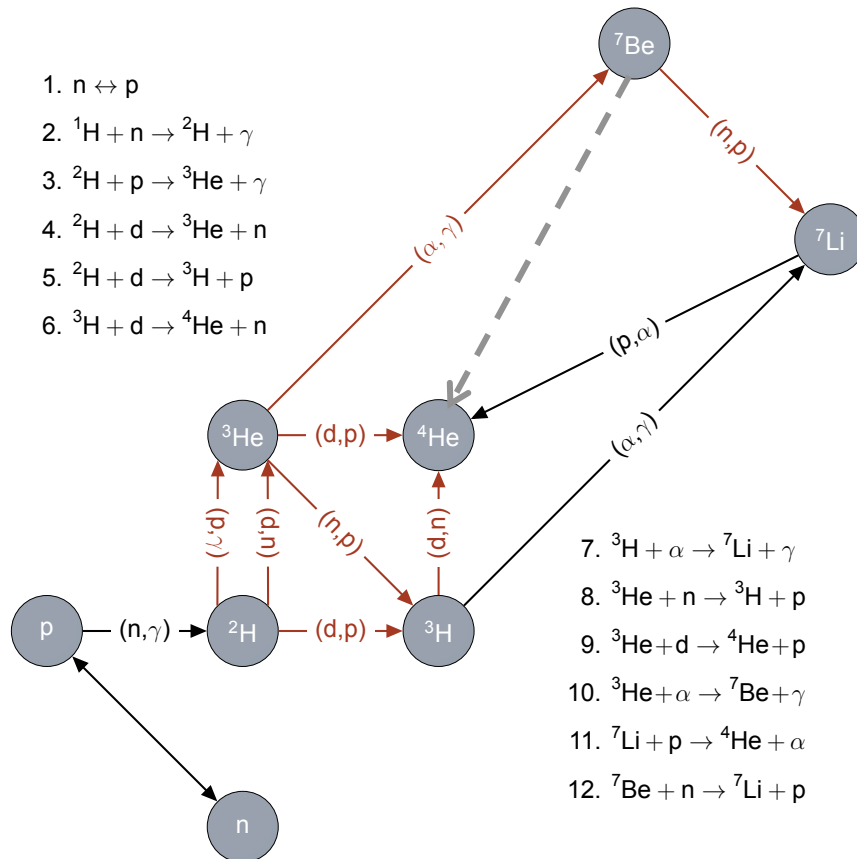


# Nuclear solution to the Li problem ?

At  $\eta_{\text{CMB}}$   ${}^7\text{Li}$  from  ${}^7\text{Be}$  post BBN decay

Tentatives nuclear solutions:  
 ${}^7\text{Be}$  destruction by:

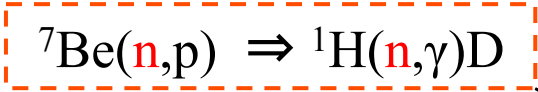
- ~~Supplementary reactions  
 e.g.  ${}^7\text{Be}(d,p){}^8\text{Be}^* \rightarrow 2\alpha$~~
- Increased neutron destruction efficiency by  ${}^7\text{Be}(n,p){}^7\text{Li}(p,\alpha){}^4\text{He}$  from exotic neutron sources



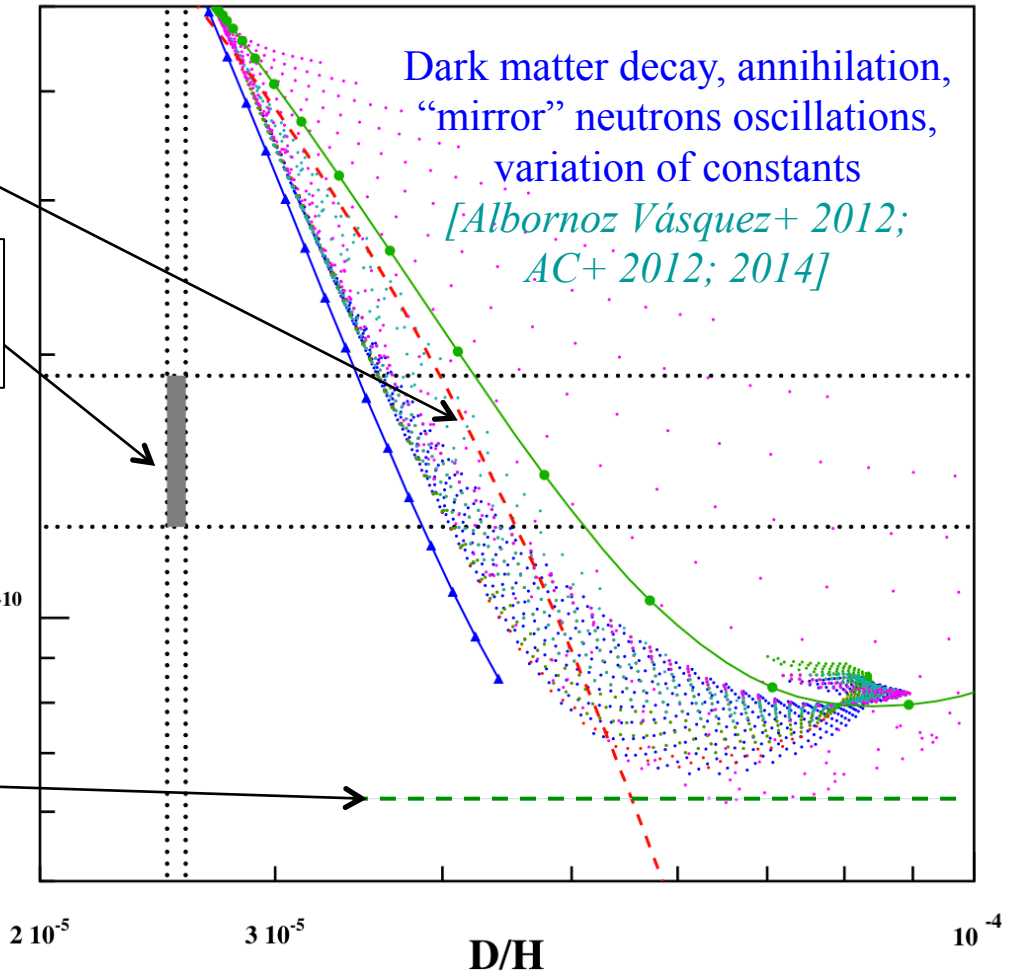
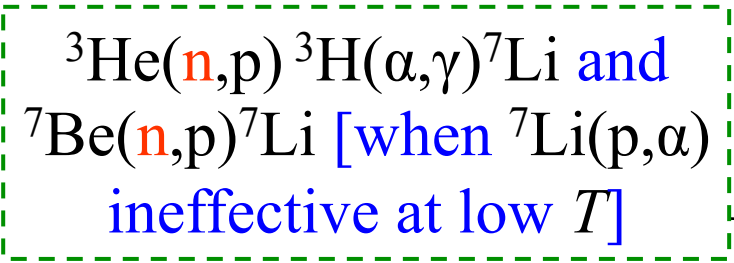


# The limits to ${}^7\text{Li}+{}^7\text{Be}$ destruction by extra neutrons

See also e.g. *Olive+ 2012; Kusakabe+ 2014; .....*



Li/D observational limits  
*[Sbordone+ 2010 × Cooke+ 2014]*

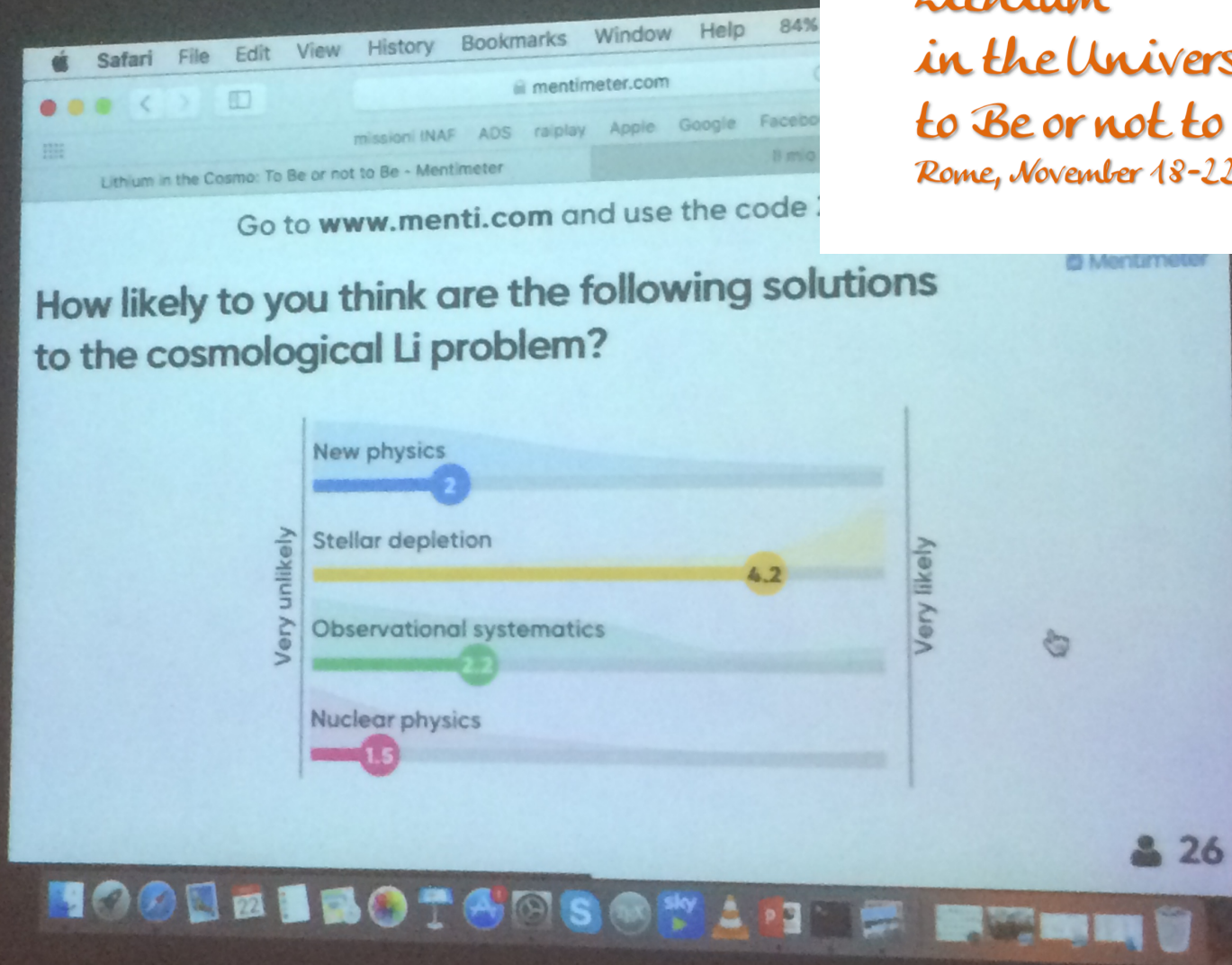


Even worse for (non-thermal) high energy neutrons *[Kusakabe+ 2004]*

Lower Li/H  $\Rightarrow$  higher D/H

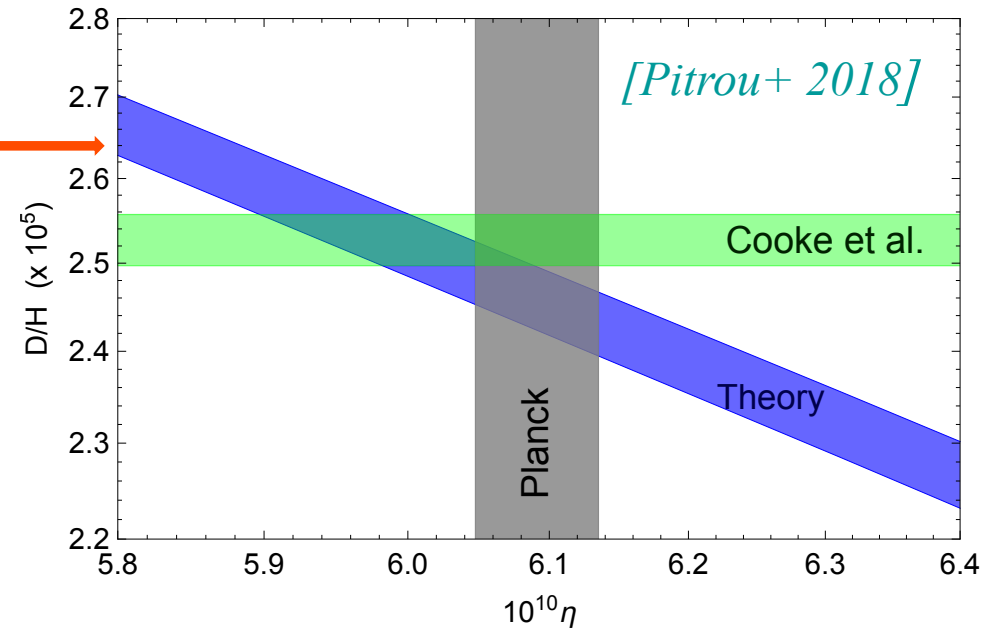


Lithium  
in the Universe:  
to Be or not to Be?  
Rome, November 18-22, 2019



# BBN tension on D/H

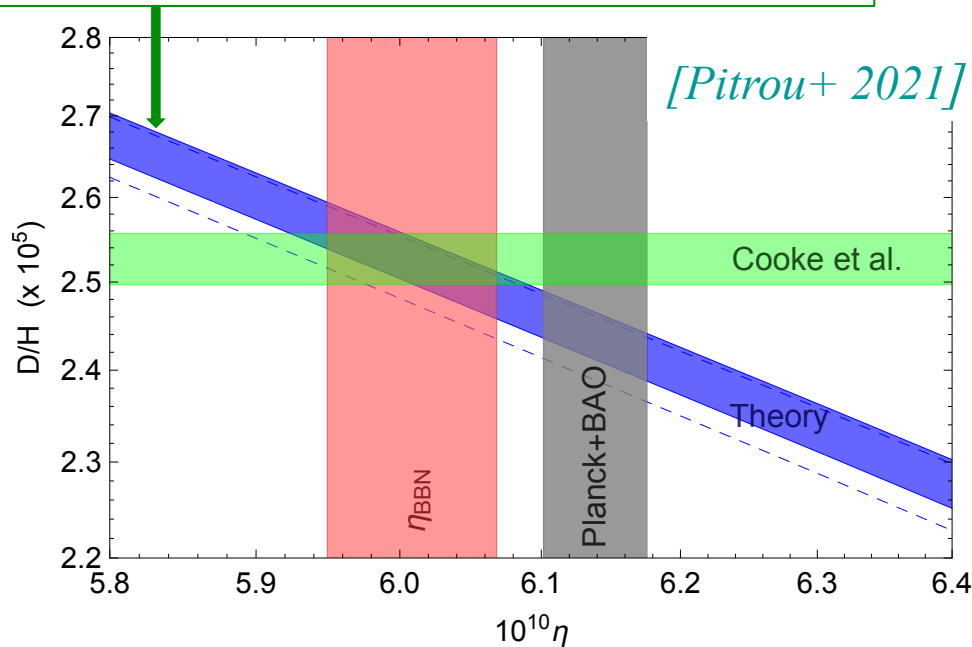
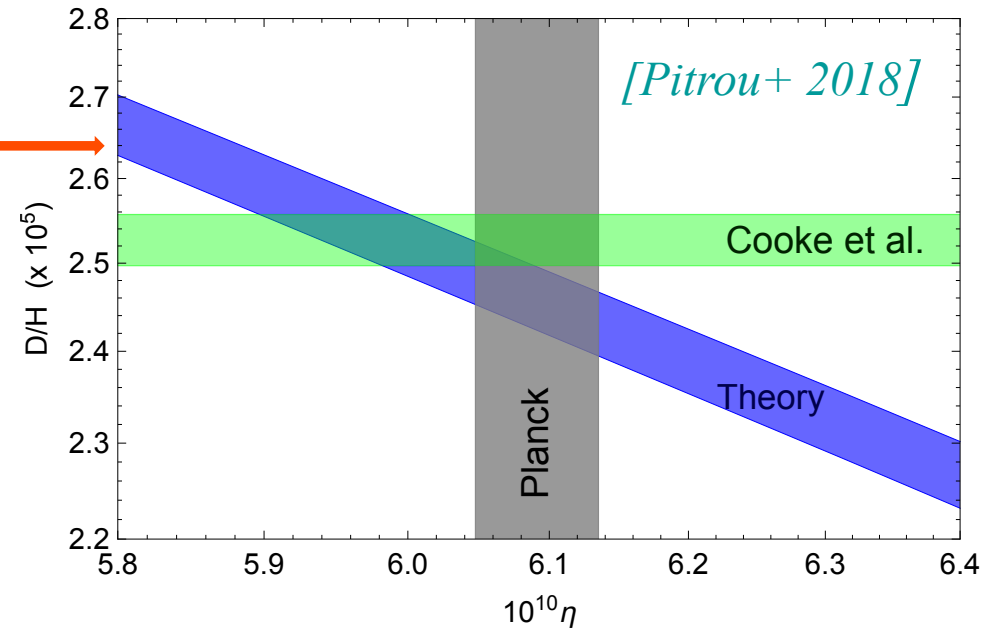
$D(p,\gamma)^3\text{He}$  [Iliadis+2016]  
 $D(d,n)$  &  $D(d,p)$  [Gómez Iñesta+ 2017]  
 $\Omega_b h^2 = 0.02225 \pm 0.00016$  [CMB: Ade+ 2016 (Planck)]



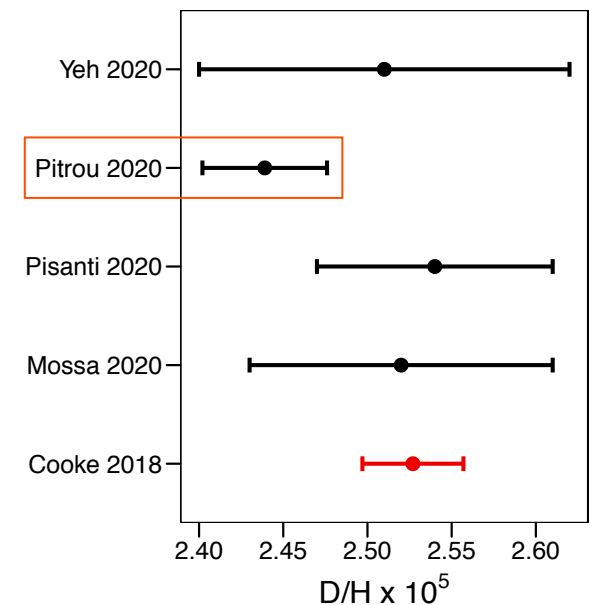
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$D(p,\gamma)^3\text{He}$  [Mossa+ 2020 (LUNA)]  
 $D(d,n)$  &  $D(d,p)$  [Gómez Iñesta+ 2017]  
 $\Omega_b h^2 = 0.02242 \pm 0.00014$  [CMB+BAO: Alam+ 2017; Aghanim+ 2020]

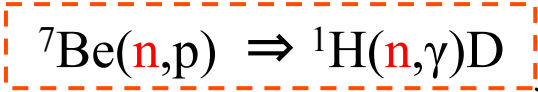


1.8- $\sigma$   
tension

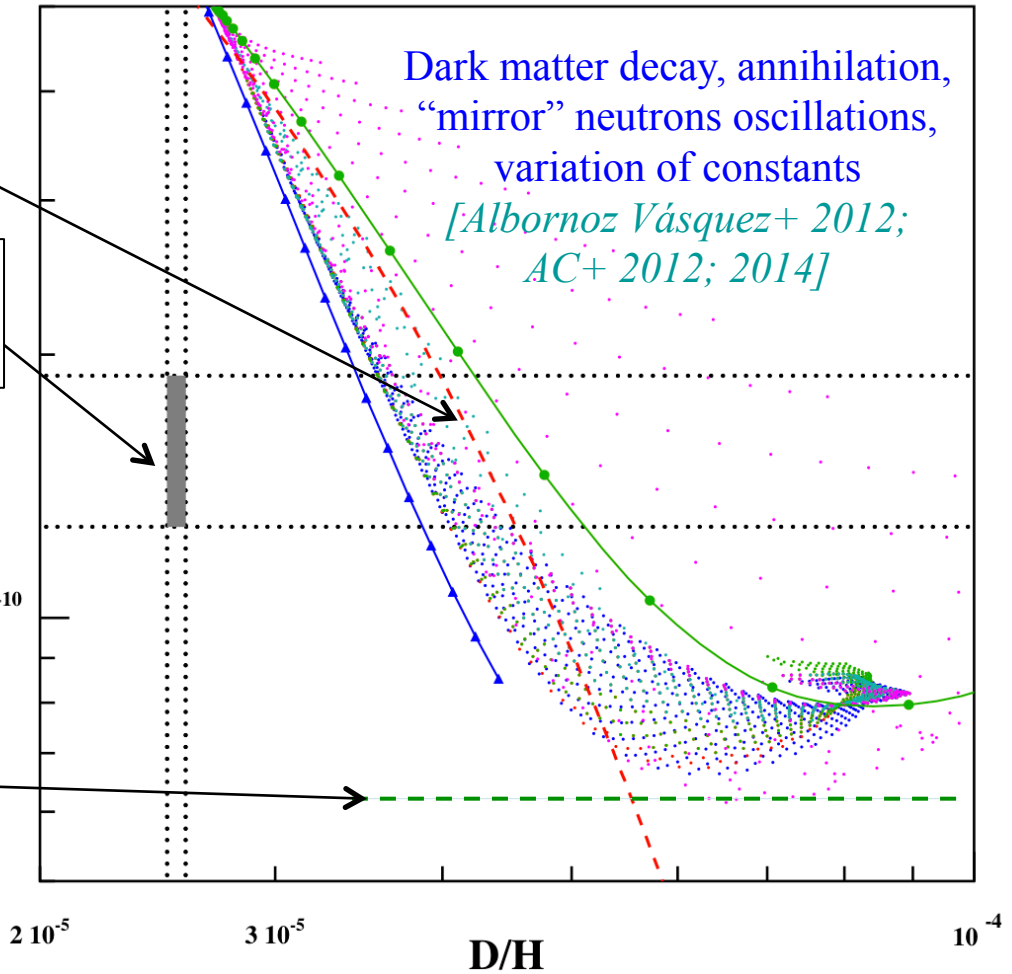
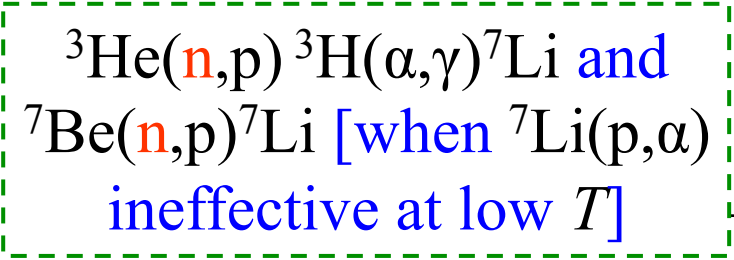


# The limits to ${}^7\text{Li}+{}^7\text{Be}$ destruction by extra neutrons

See also e.g. *Olive+ 2012; Kusakabe+ 2014; .....*



Li/D observational limits  
*[Sbordone+ 2010 × Cooke+ 2014]*



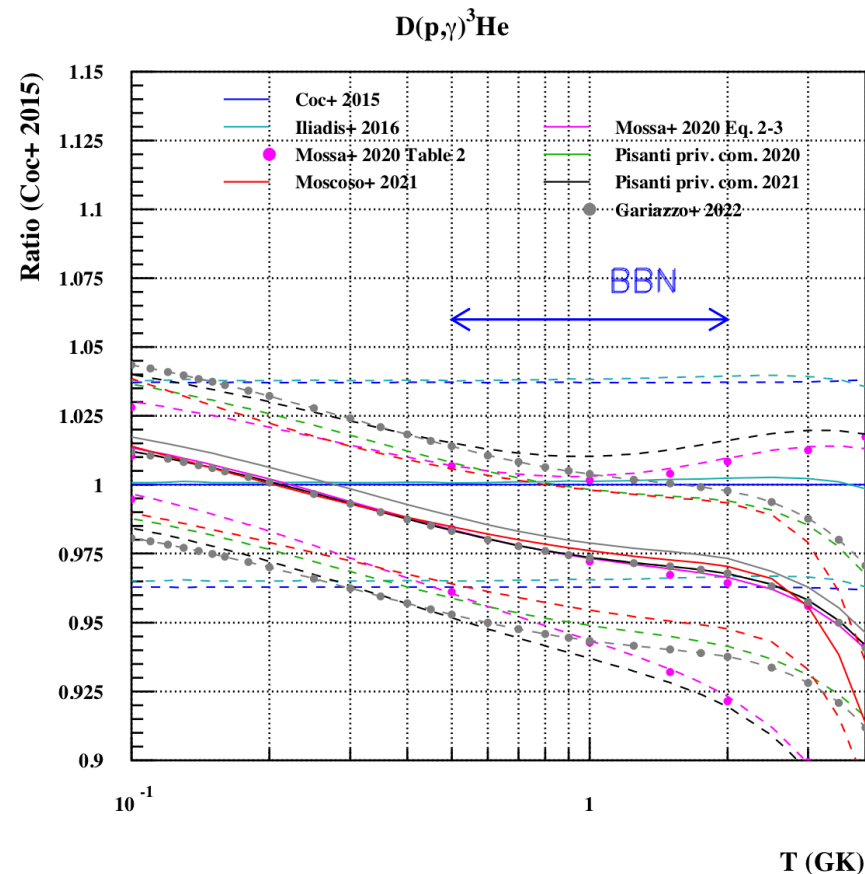
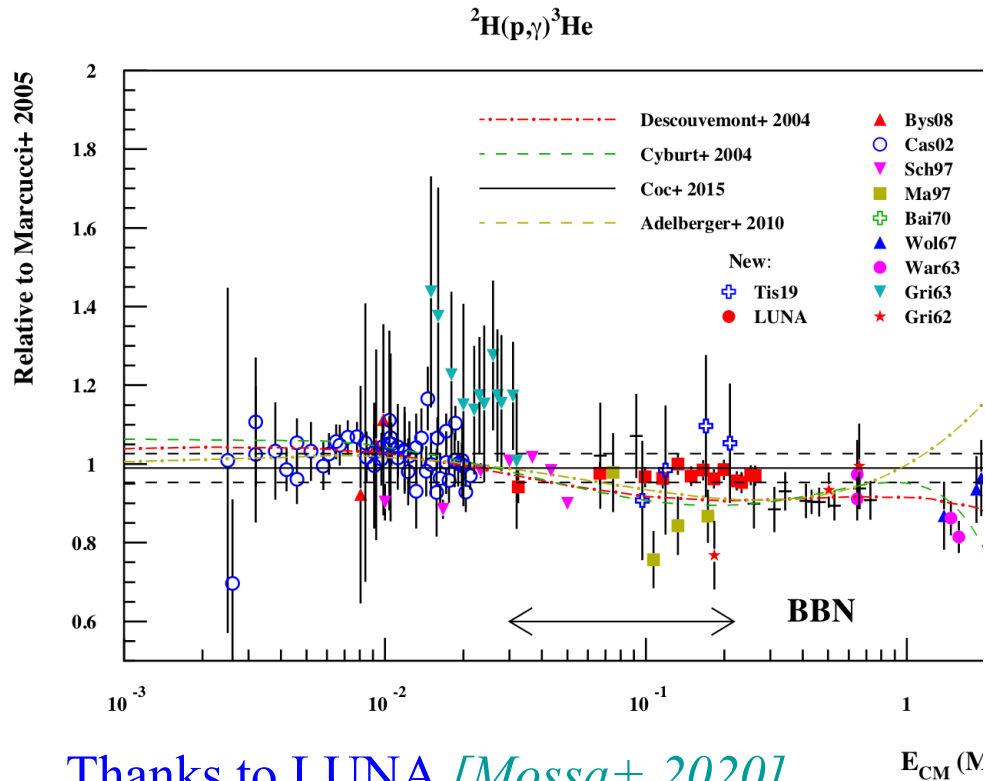
Even worse for (non-thermal) high energy neutrons *[Kusakabe+ 2004]*

Lower Li/H  $\Rightarrow$  higher D/H

# Comparison of reaction rates : $D(p,\gamma)^3\text{He}$

$$\frac{\Delta(D/H)}{D/H} = -0.32 \times \frac{\Delta\langle\sigma v\rangle_{d(p,\gamma)^3\text{He}}}{\langle\sigma v\rangle_{d(p,\gamma)^3\text{He}}} - 0.54 \times \frac{\Delta\langle\sigma v\rangle_{d(d,n)^3\text{He}}}{\langle\sigma v\rangle_{d(d,n)^3\text{He}}} - 0.46 \times \frac{\Delta\langle\sigma v\rangle_{d(d,p)^3\text{H}}}{\langle\sigma v\rangle_{d(d,p)^3\text{H}}}$$

$D(p,\gamma)^3\text{He}$ ,  $D(d,n)^3\text{He}$  and  $D(d,p)^3\text{H}$  reaction rates need to be known at a % level to match the 1.2% precision on observations!

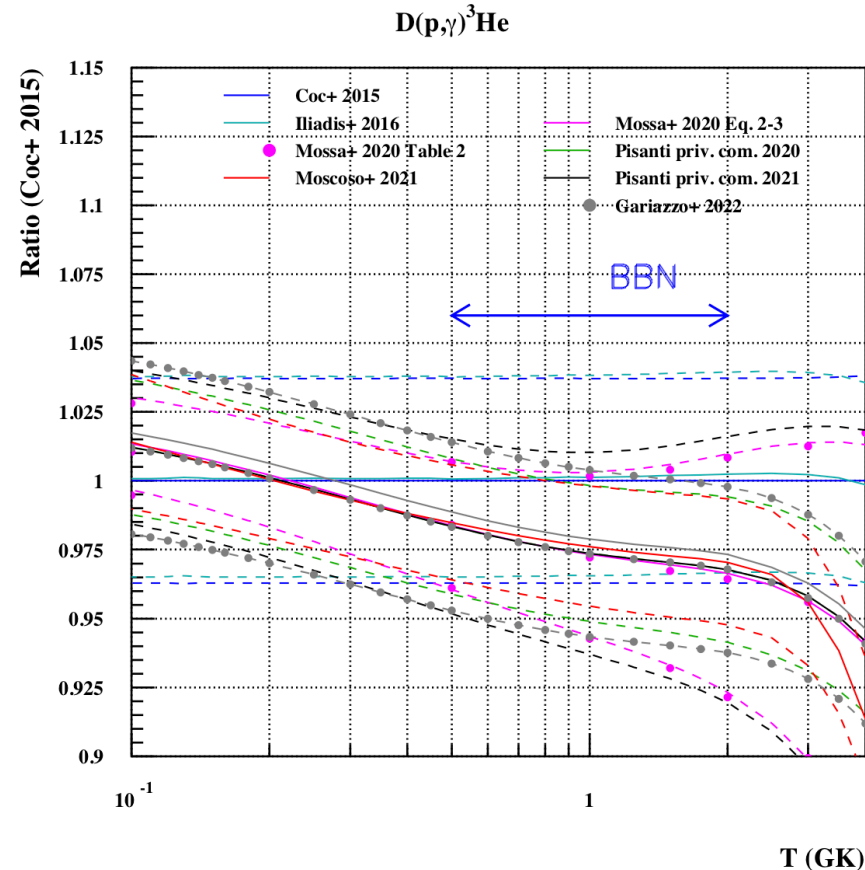
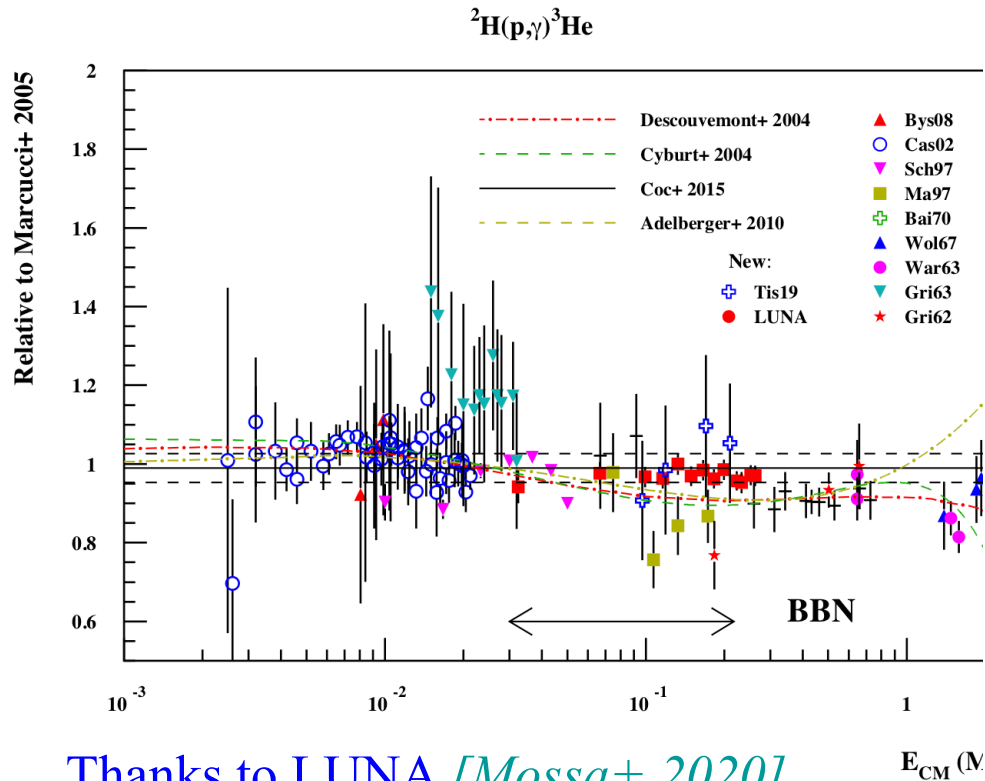


Thanks to LUNA [Mossa+ 2020]  
 $D(p,\gamma)^3\text{He}$  rates known to  $\sim 2\%$

# Comparison of reaction rates : $D(p,\gamma)^3\text{He}$

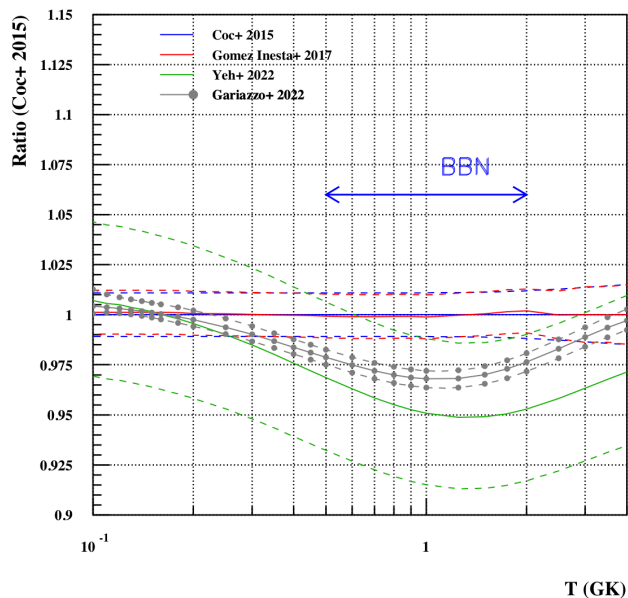
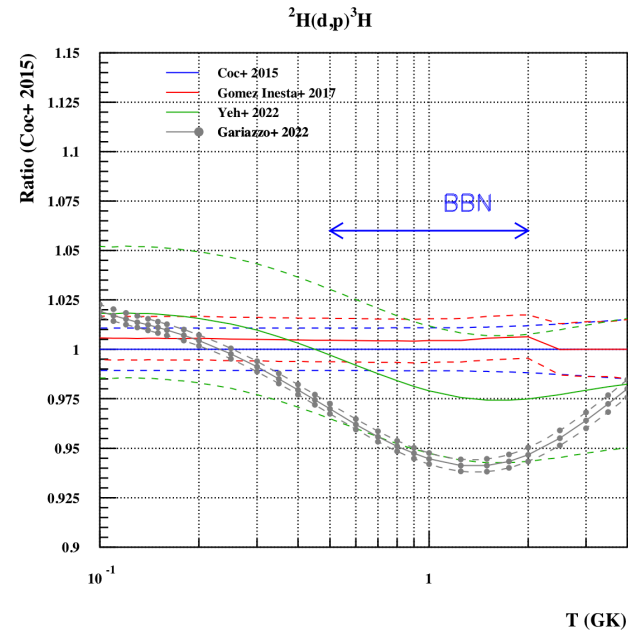
$$\frac{\Delta(D/H)}{D/H} = -0.32 \times \frac{\Delta\langle\sigma v\rangle_{d(p,\gamma)^3\text{He}}}{\langle\sigma v\rangle_{d(p,\gamma)^3\text{He}}} - 0.54 \times \frac{\Delta\langle\sigma v\rangle_{d(d,n)^3\text{He}}}{\langle\sigma v\rangle_{d(d,n)^3\text{He}}} - 0.46 \times \frac{\Delta\langle\sigma v\rangle_{d(d,p)^3\text{H}}}{\langle\sigma v\rangle_{d(d,p)^3\text{H}}}$$

$D(p,\gamma)^3\text{He}$ ,  $D(d,n)^3\text{He}$  and  $D(d,p)^3\text{H}$  reaction rates need to be known at a % level to match the 1.2% precision on observations!



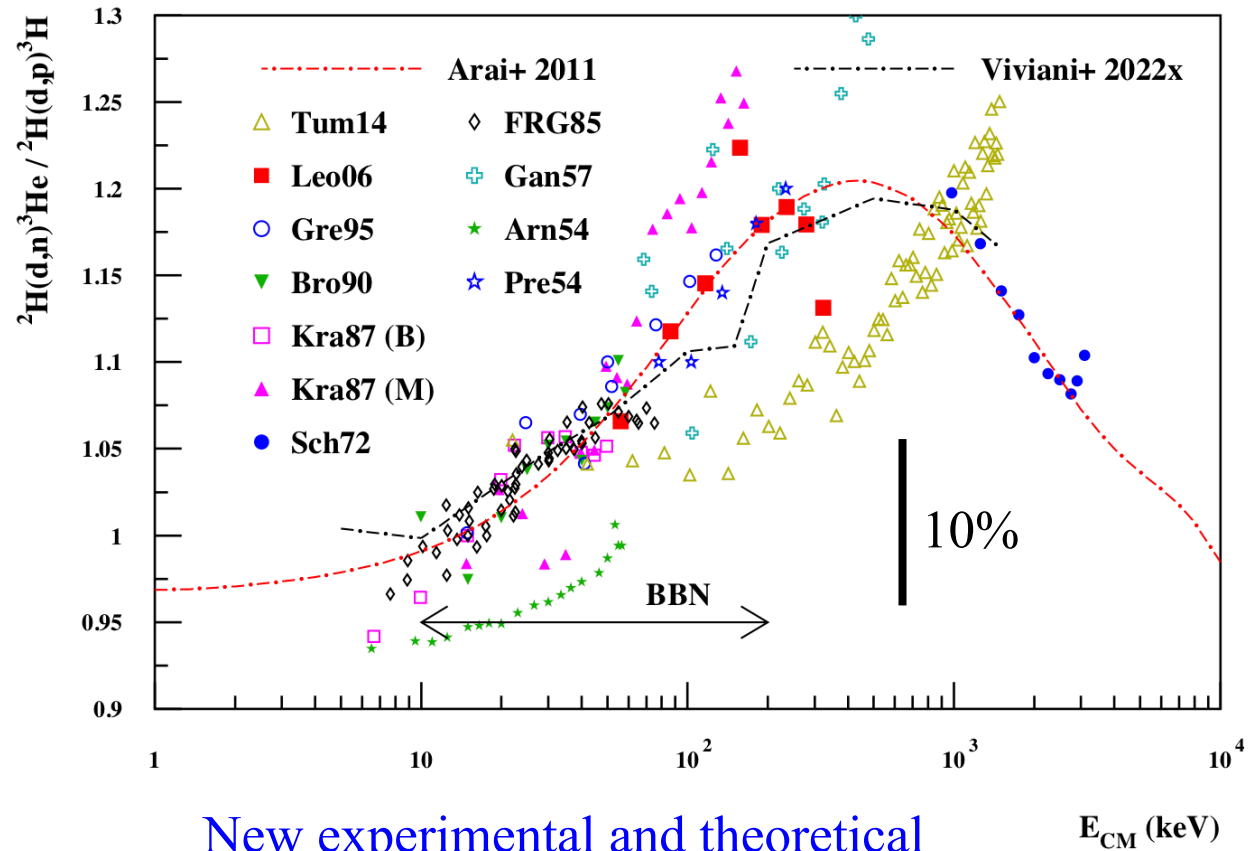
Thanks to LUNA [Mossa+ 2020]  
 $D(p,\gamma)^3\text{He}$  rates known to  $\sim 2\%$

# Comparison of reaction rates : $D(d,n)^3\text{He}$ and $D(d,p)^3\text{H}$



At the origin of discrepancies between groups  
*[Pitrou+ 2021, Pisanti+ 2021, Yeh+ 2021]*

Ratio of  $D(d,n)^3\text{He}$  and  $D(d,p)^3\text{H}$  cross sections  
 exp. and th. (to remove some of the systematics)



New experimental and theoretical  
 studies welcomed!



# Evaluation of main reaction rates

( $D(p,\gamma)^3\text{He}$ ,  $D(d,p)^3\text{H}$  and  $D(d,n)^3\text{He}$  )

## 1. Selection of experimental data sets

- Direct measurements (no model involved)
- Systematic uncertainties (normalization)
- Vintage experiments (50's,60's) with few details published!

## 2. Selection of fitting function

- From theory?
- No resonances in cross sections, no structure
- Polynomial, splines?

## 3. Statistical analyses

- “Classical”  $\chi^2$  minimisation
- Bayesian

## **Effects of the variation of fundamental constants on Population III stellar evolution**

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## **Variation of fundamental constants and the role of $A = 5$ and $A = 8$ nuclei on primordial nucleosynthesis**

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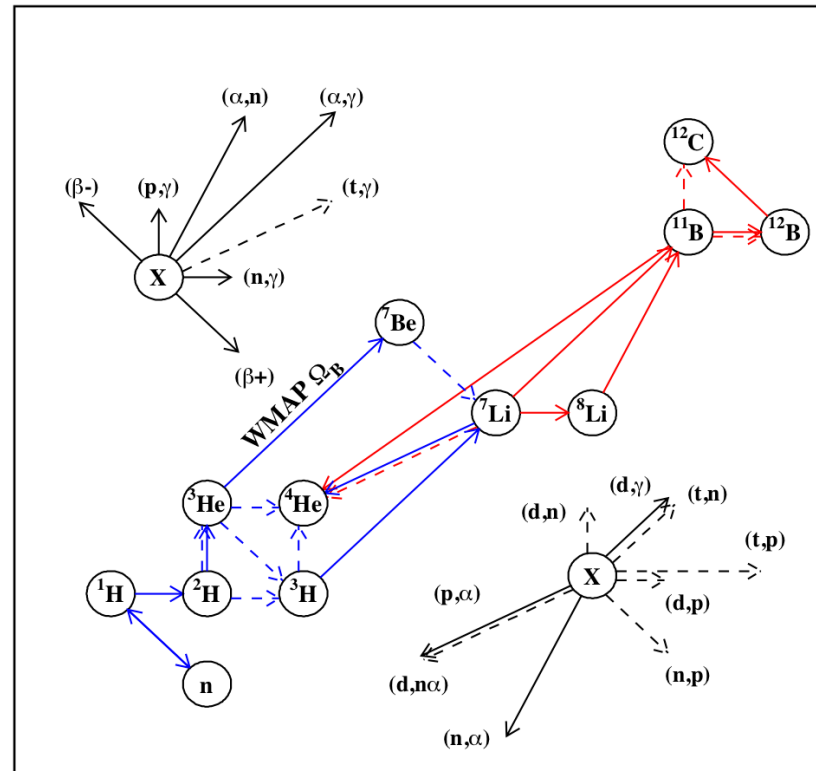
Jean-Philippe Uzan<sup>§</sup> and Elisabeth Vangioni<sup>||</sup>

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# CNO seeds for first (Population III) stars

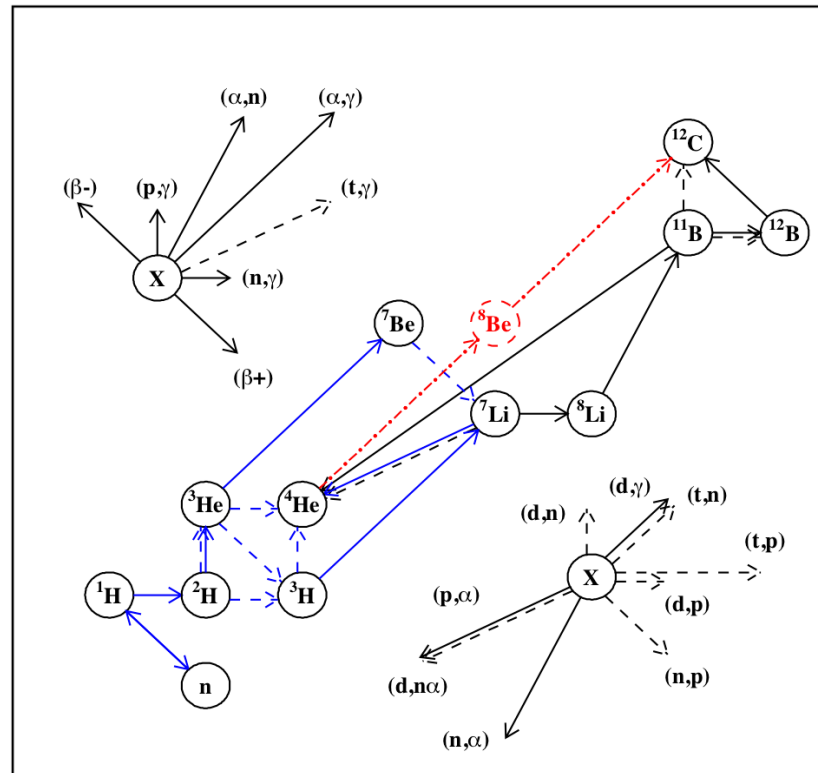
$\text{CNO}/\text{H} > 10^{-13}$  to affect Pop III stars [Ekström+ 2008]



$\text{CNO}/\text{H} = (0.5-3) \times 10^{-15}$  from BBN [Iocco  
+ 2017; Coc+ 2012; 2014]

# CNO seeds for first (Population III) stars

$\text{CNO}/\text{H} > 10^{-13}$  to affect Pop III stars [Ekström+ 2008]



However, “*well known*” that a stable  $^8\text{Be}$  would bridge the  $A=8$  gap!

# Variation of the fundamental constants

Physical theories involve constants

These parameters cannot be determined by the theory that introduces them; we can only measure them.

These arbitrary parameters have to be assumed constant:

- *experimental validation*
- *no evolution equation*

Very small variations, best studied on cosmological scales, from astronomical observations

See reviews : *J.-P. Uzan in Rev. Mod Phys. 2003, Living Rev. Relativity 2011; E. García-Berro, J. Isern & Y.A. Kubishin in Astron. Astrophys. Rev. 2007*

Assuming that coupled variations of physical constants ( $\alpha_{em}, \Lambda_{QCD}, G_F, \dots$ ) may lead to a *stable*  ${}^8\text{Be}$  *at the time of BBN* [*Coc, Descouvemont, Olive, Uzan & Vangioni 2012*]

# The triple-alpha reaction

1. Equilibrium between  ${}^4\text{He}$  and the short lived ( $\sim 10^{-16}$  s)  ${}^8\text{Be}$  :  $\alpha\alpha \leftrightarrow {}^8\text{Be}$
2. Resonant capture to the ( $l=0, J^\pi=0^+$ ) Hoyle state:  ${}^8\text{Be} + \alpha \rightarrow {}^{12}\text{C}^* (\rightarrow {}^{12}\text{C} + \gamma)$

Simple formula used in previous studies

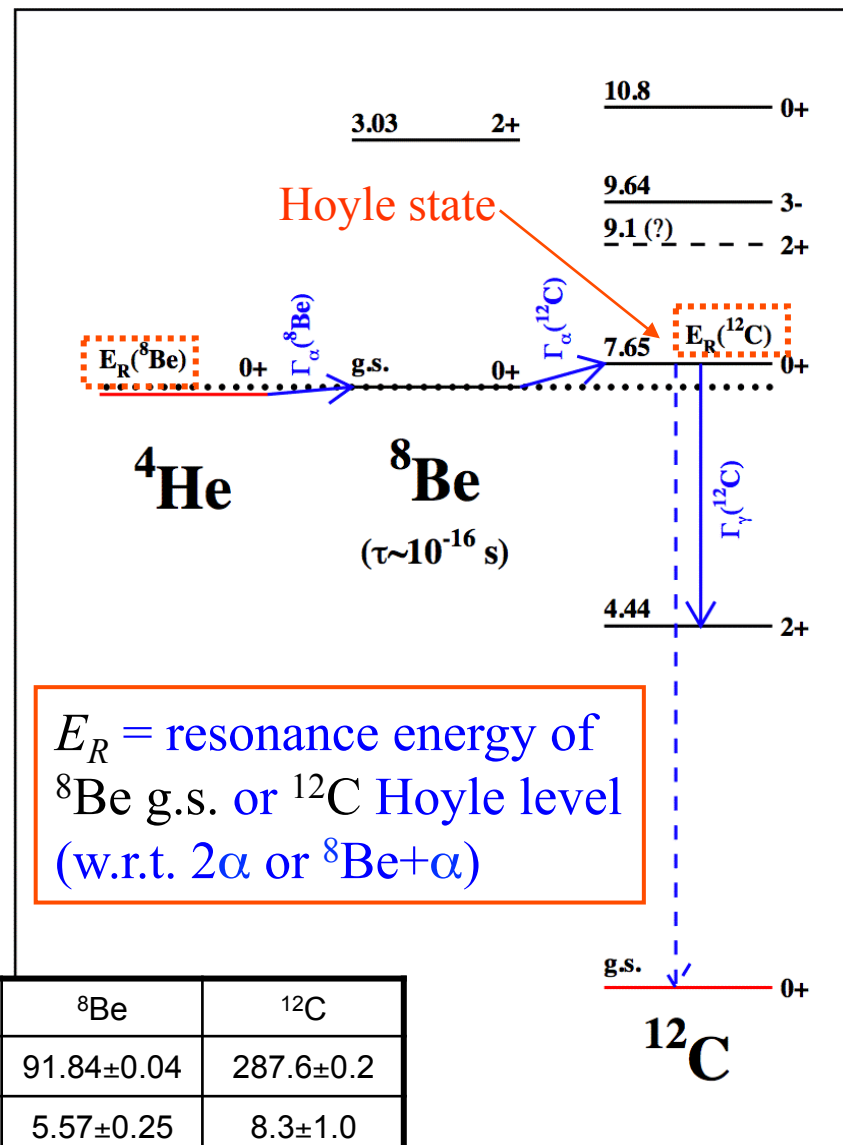
1. Saha equation (thermal equilibrium)
2. Sharp resonance analytic expression:

$$N_A^2 \langle \sigma v \rangle^{\alpha\alpha\alpha} = 3^{3/2} 6 N_A^2 \left( \frac{2\pi}{M_\alpha k_B T} \right)^3 \hbar^5 \gamma \exp\left( \frac{-Q_{\alpha\alpha\alpha}}{k_B T} \right)$$

with  $Q_{\alpha\alpha\alpha} = E_R({}^8\text{Be}) + E_R({}^{12}\text{C})$  and  $\gamma \approx \Gamma_\gamma$

Approximations

1. Thermal equilibrium
2. Sharp resonance
3.  ${}^8\text{Be}$  decay faster than  $\alpha$  capture



Nucleus	${}^8\text{Be}$	${}^{12}\text{C}$
$E_R$ (keV)	$91.84 \pm 0.04$	$287.6 \pm 0.2$
$\Gamma_\alpha$ (eV)	$5.57 \pm 0.25$	$8.3 \pm 1.0$
$\Gamma_\gamma$ (meV)	-	$3.7 \pm 0.5$

# Nuclear microscopic calculations

□ Hamiltonian:

$$H = \sum_{i=1}^A T(r_i) + \sum_{i<j=1}^A (V_{Coul.}(r_{ij}) + V_{Nucl.}(r_{ij}))$$

Where  $V_{Nucl.}(r_{ij})$  is an effective Nucleon-Nucleon interaction

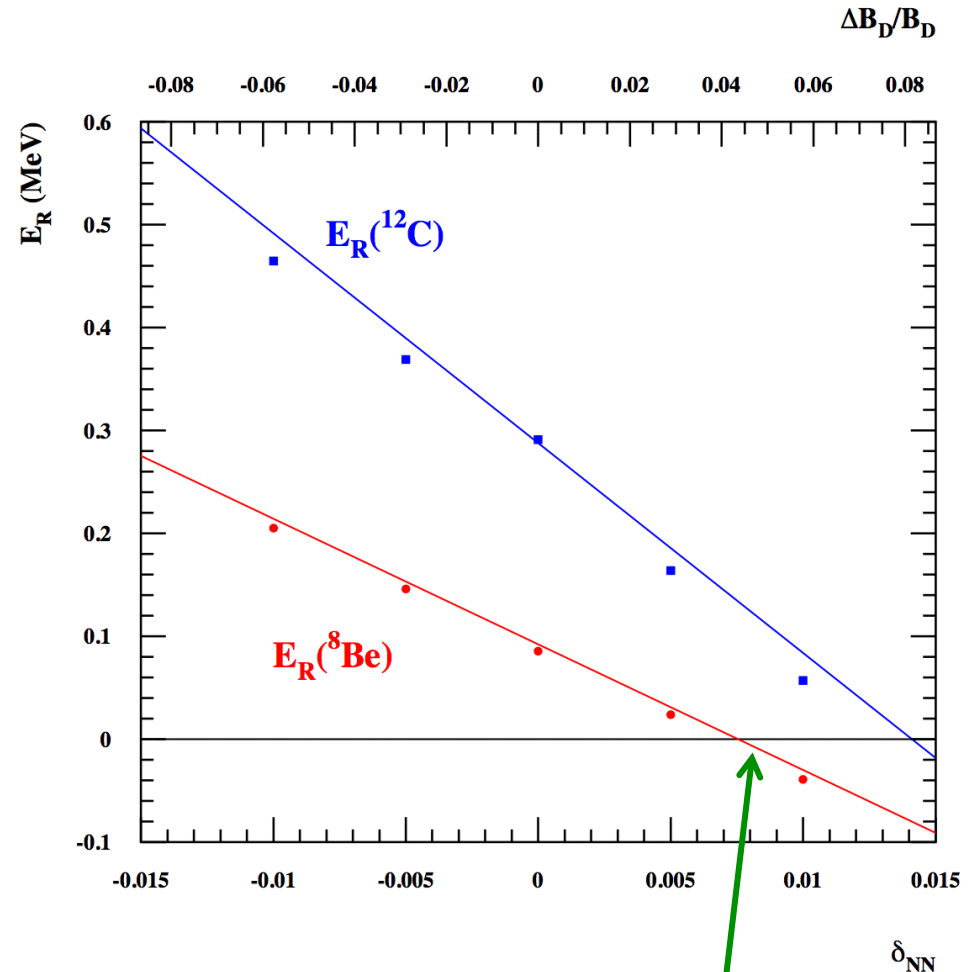
□ Minnesota N-N force [Thompson *et al.* 1977] optimized to reproduce low energy N-N scattering data and  $B_D$  (deuterium binding energy)

□  $\alpha$ -cluster approximation for  ${}^8\text{Be}^{\text{g.s.}}$  ( $2\alpha$ ) and the Hoyle state ( $3\alpha$ ) [Kamimura 1981]

□ Scaling of the N-N interaction

$$V_{Nucl.}(r_{ij}) \rightarrow (1 + \delta_{NN}) \times V_{Nucl.}(r_{ij})$$

to obtain  $B_D$ ,  $E_R({}^8\text{Be})$ ,  $E_R({}^{12}\text{C})$  as a function of  $\delta_{NN}$ :



${}^8\text{Be}$  stable for N-N interaction higher by 0.75%

# Numerical rate calculation

□ At “low temperatures”, capture via resonance tails [*Nomoto et al. 1985*] requires numerical integration

➤ Even more important when resonances are shifted upwards with larger widths

- Radiative widths :  $\Gamma_\gamma(E) \propto E^{2L+1}$  (with  $L=2$  here)

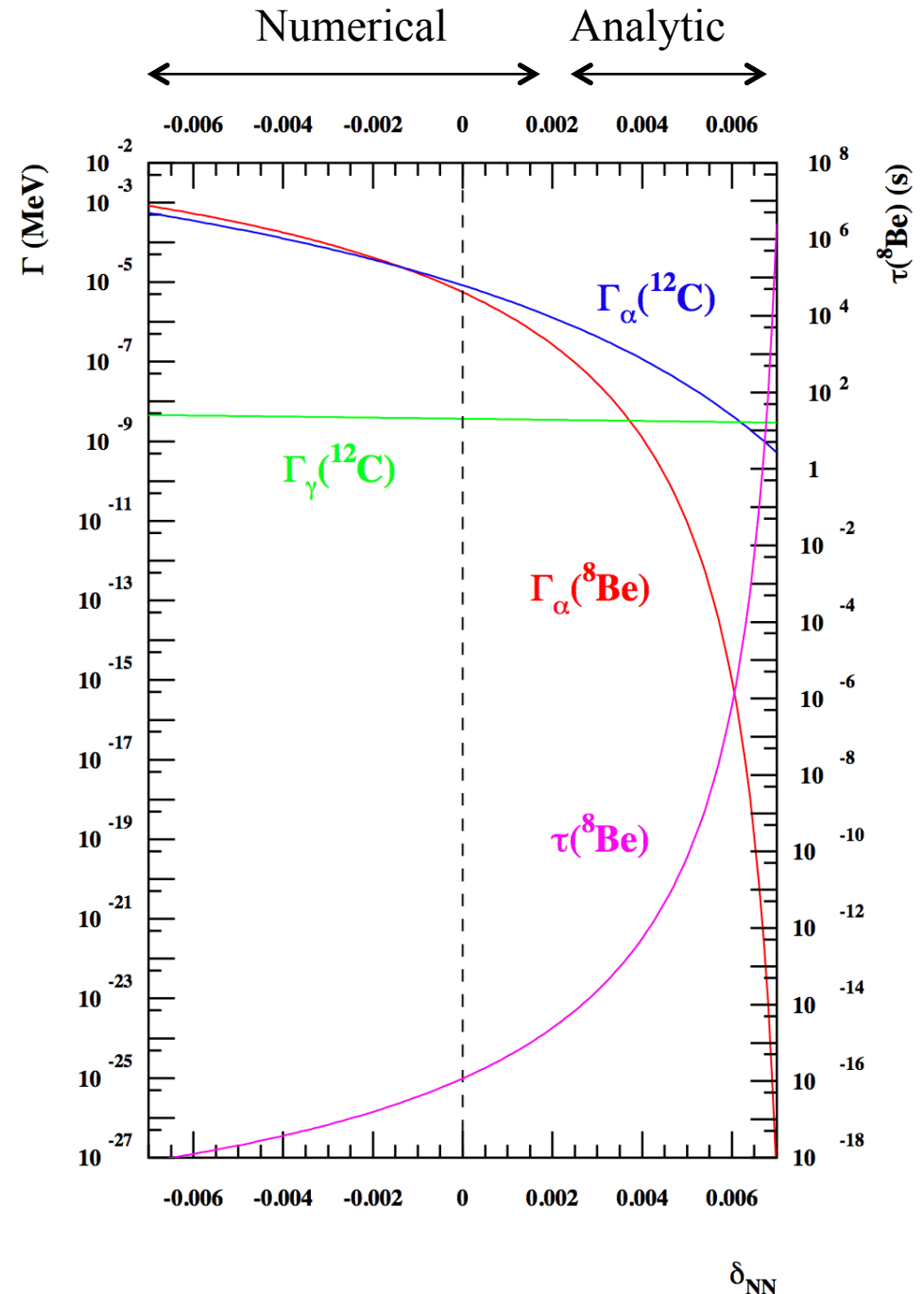
- Charged particle widths :

$$\Gamma_\alpha(E) = \Gamma_\alpha(E_R) P_L(E, R_C) / P_L(E_R, R_C) \text{ with}$$

$$P_L(E, R_C) = \propto (F_L^2(\eta, kR_C) + G_L^2(\eta, kR_C))^{-1}$$

the penetrability

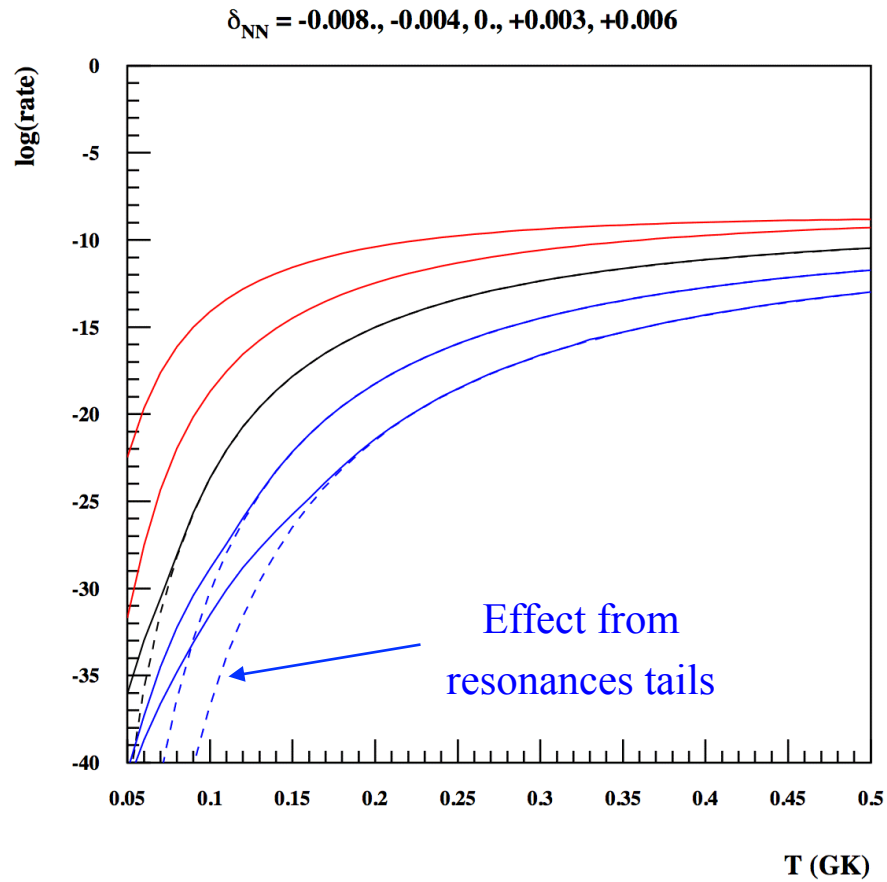
$\gamma \equiv \Gamma_\alpha(E) \Gamma_\gamma(E) / (\Gamma_\alpha(E) + \Gamma_\gamma(E)) \approx \Gamma_\gamma(E)$   
if  $\Gamma_\gamma(E) \ll \Gamma_\alpha(E)$  in analytic expression



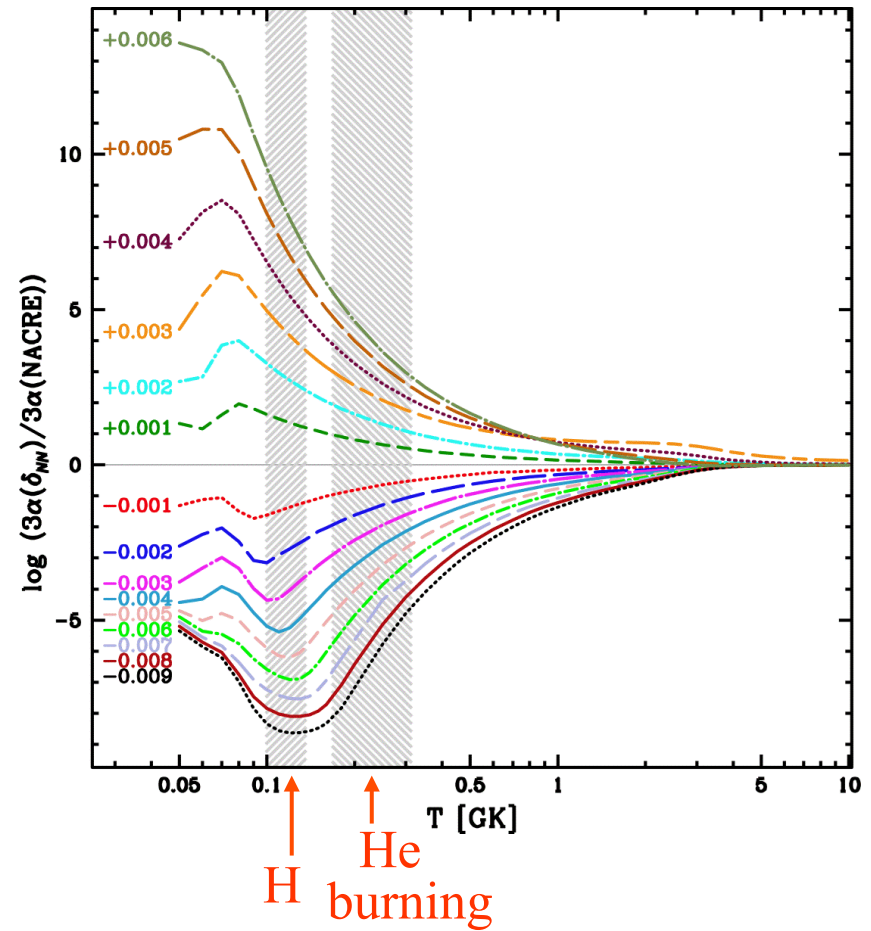


# The triple-alpha with unstable $^8\text{Be}$

Rates



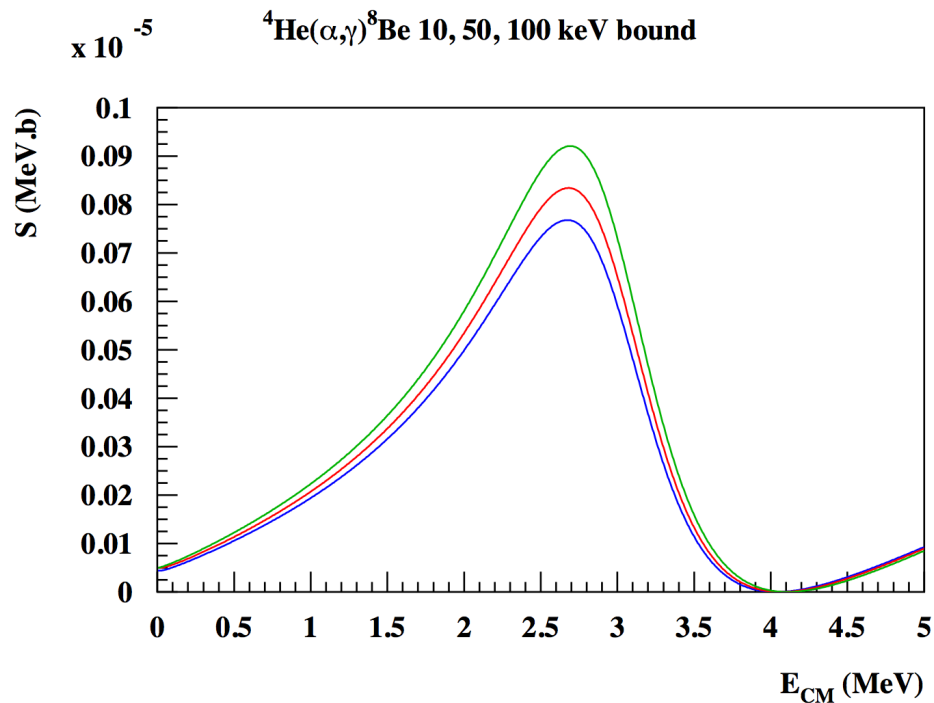
Rates / NACRE



# The triple-alpha with a stable $^8\text{Be}$

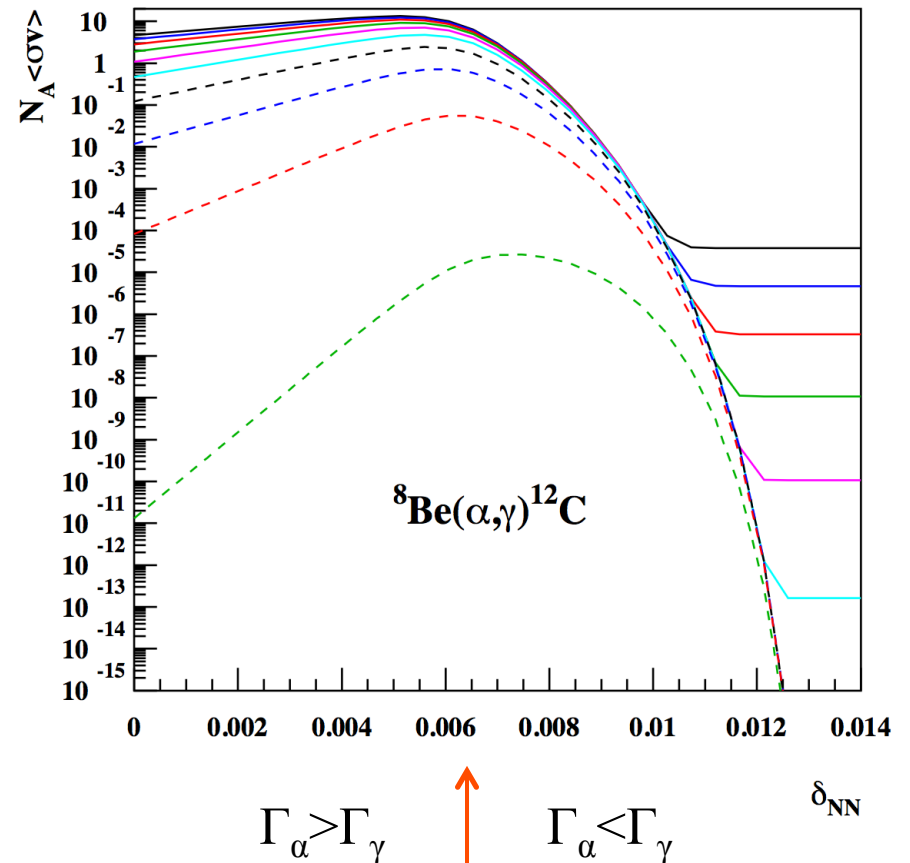
$^8\text{Be}$  bound by 10, 50 or 100 keV  
 ( $\delta_{\text{NN}} = 0.0083, 0.0116, 0.0156$ )

$^4\text{He}(\alpha, \gamma)^8\text{Be}^{\text{bound}}$  cross-section in  
 continuity with unbound one  
*[Baye & Descouvemont 1985]*



$^8\text{Be}^{\text{bound}}(\alpha, \gamma)^{12}\text{C}$  from  
 sharp resonance formula

$T = 0.1, 0.2, 0.3, 0.4, 0.5, 0.6, 0.7, 0.8, 0.9, 1.0$  GK

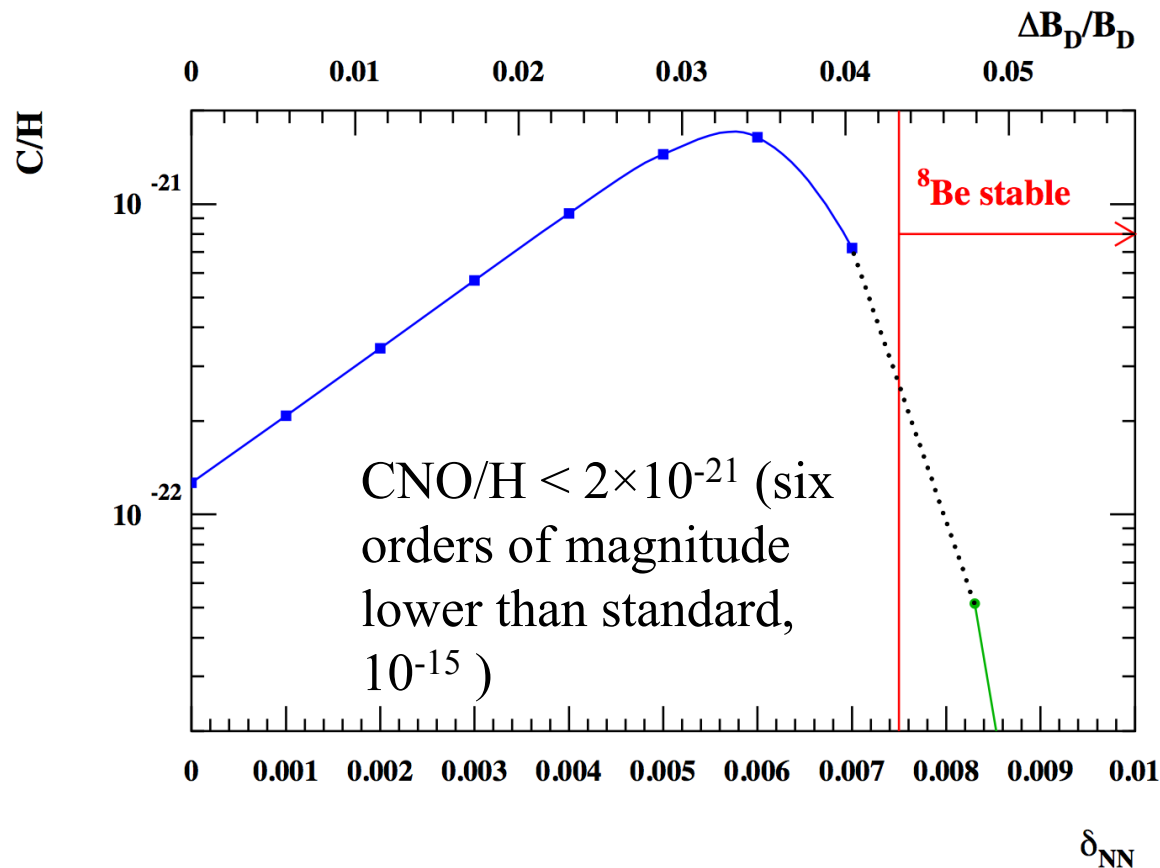


*[Newton et al. 2007]*

# CNO production with un/stable $^8\text{Be}$

BBN calculations [*Coc+ 2012*] with:

- ❑ Triple-alpha with shifted resonance energies ( $\delta_{NN} < 0.0075$ )
- ❑  $^4\text{He}(\alpha, \gamma)^8\text{Be}^{\text{bound}}(\alpha, \gamma)^{12}\text{C}$  [*Baye & Descouvemont 1985*] ( $\delta_{NN} < 0.0075$ )
- ❑ Other CNO producing reactions removed



- Unbound: density lower (3 body reaction) and timescale shorter than in stars!
- Bound:  $^4\text{He}(\alpha, \gamma)^8\text{Be}(\alpha, \gamma)^{12}\text{C}$  too slow (two  $Z=2$  captures)

# Conclusions

- ❑ Convergence of BBN codes when same nuclear reaction rates are used
  - Mathematica code with >400 reaction network publicly available at <http://www2.iap.fr/users/pitrou/primat.htm>
- ❑ Standard BBN is now in the (1%) precision era for D and  $^4\text{He}$ 
  - We confirm a  $1.8\sigma$  tension between observations/predictions of D/H with our evaluated rates + LUNA (2 codes  $\times$  2 rate analyses)
  - Calls for even better precision on  ~~$D(p,\gamma)^3\text{He}$~~ ,  $D(d,p)^3\text{H}$  and  $D(d,n)^3\text{He}$  cross sections. Experimentally very challenging.
  - New physics? Plenty of models overproduce D to reduce Li!
- ❑ The Cosmological Lithium Problem is worse than ever!
  - Disagreement (factor of 3) with Li observations
    - Nuclear : excluded by experiments
    - No  $^6\text{Li}$  plateau [*Lind+ 2013*] from BBN!
  - BBN: a tool to explore BSM physics in the early stages of the Universe