

# Criterion for ultra-fast bubble walls: the impact of hydrodynamic obstruction (2401.05911)

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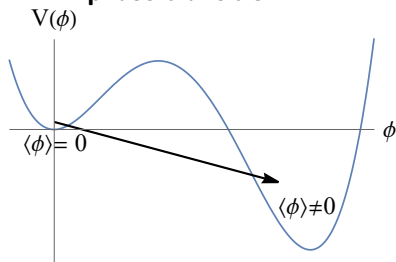
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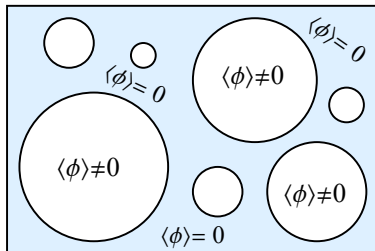
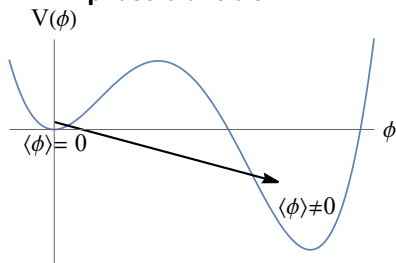
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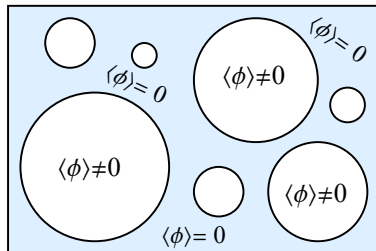
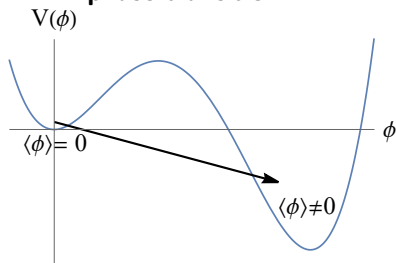
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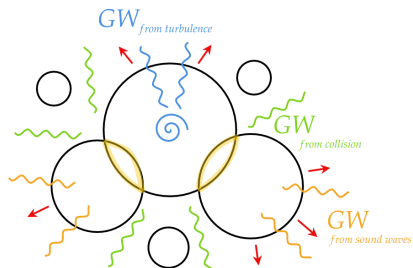
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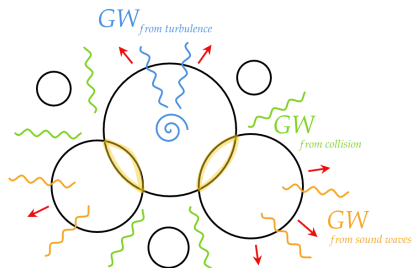
- Arises naturally in many BSM models
  - baryo/leptogenesis, dark matter, primordial black holes, ...

# Parameters determining GW spectrum



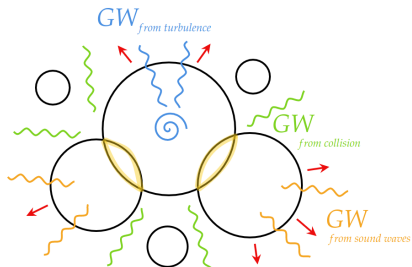


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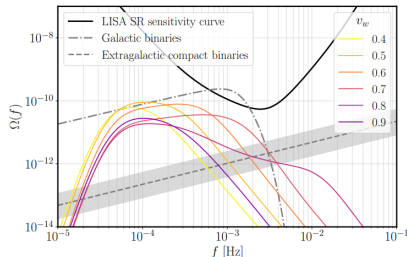


- Why determine  $\nu_w$ ?
  - Baryogenesis, DM
  - Predicting GW spectrum

# Parameters determining GW spectrum



- Why determine  $v_w$ ?
  - Baryogenesis, DM
  - Predicting GW spectrum



Gowling & Hindmarsh (2106.05984)


- Integrating EoM of scalar field coupled to plasma over stationary wall profile:

$$\int_{-\delta}^{\delta} dz \left( \underbrace{\frac{dV_0(\phi)}{dz}}_{\mathcal{P}_{\text{driving}}} + \underbrace{(\partial_z \phi) \sum_i \frac{\partial m_i^2(\phi)}{\partial \phi} \int \frac{\partial^3 \mathbf{p}}{(2\pi)^3 2E_i} (f_i^{\text{eq}} + \delta f_i)}_{-\mathcal{P}_{\text{friction}} > 0} \right) = 0 \quad (1)$$

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- Ballistic limit ( $L_w \ll \gamma_w L_{\text{MFP}}$ ),  $\delta f_i \rightarrow f_i^{\text{eq}}(p, T_n) - f_i^{\text{eq}}(z, p, T(z))$ :

$$\Delta V_0 = \sum_i \frac{a_i n_i m_i^2 T_n^2}{48} = \mathcal{P}_{BM}$$

$a_i$  : DoF of species  $i$

$$n_i = \begin{cases} 1 & \text{fermions} \\ 2 & \text{bosons} \end{cases}$$

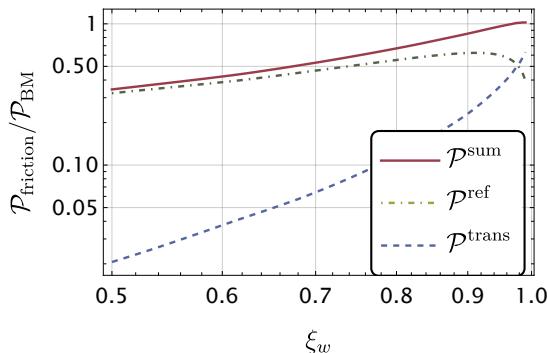
# Bödeker & Moore criterion

- If pressure increases **monotonously** with wall velocity

$$\Delta V_0 > \mathcal{P}_{BM}$$

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determines if bubble is runaway (Dine et al. 9203203)



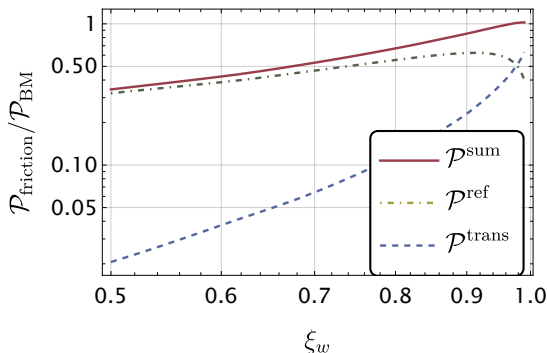
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- Pressure from *hydrodynamical* effects will not be monotonously and gives extra criterion

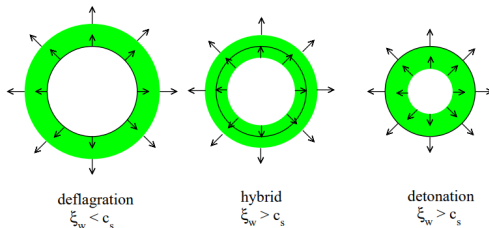
# Expansion modes

- No out-of-equilibrium effects:  $\delta f_i = 0 \Rightarrow$  lower bound on pressure

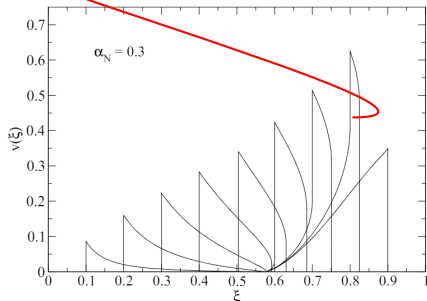


# Expansion modes

- No out-of-equilibrium effects:  $\delta f_i = 0 \Rightarrow$  lower bound on pressure
- Different expansion modes:



Jouguet velocity



$$\xi \equiv r/t$$

Espinosa, Konstandin,  
No & Servant (1004.4187)

# Hydrodynamic obstruction: A brief history

- Konstandin & No (1011.3735): Heating of plasma in front of wall  
⇒ hydrodynamic obstruction

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- Balaji, Spannowsky & Tamarit (2010.08013):

$$\partial_{\mu}(su^{\mu}) = 0 \quad (3)$$

→ Equivalent to total entropy conservation = LTE

# Hydrodynamic obstruction: A brief history

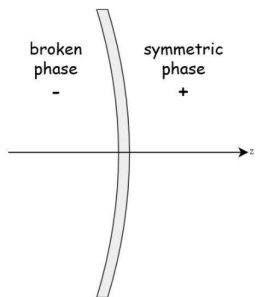
- Ai, Garbrecht & Tamarit (2109.13710): Assuming LTE

$$\begin{cases} \nabla_\mu T^{\mu\nu} = 0 \\ \partial_\mu (su^\mu) = 0 \end{cases} \Leftrightarrow \begin{cases} \omega_+ \gamma_+^2 v_+ = \omega_- \gamma_-^2 v_- \\ \omega_+ \gamma_+^2 v_+^2 + p_+ = \omega_- \gamma_-^2 v_-^2 + p_- \\ \gamma_+ T_+ = \gamma_- T_- \end{cases} \quad (4)$$

$$T^{\mu\nu} = T_{\text{plasma}}^{\mu\nu} + T_\phi^{\mu\nu}$$

$\omega$  : enthalpy

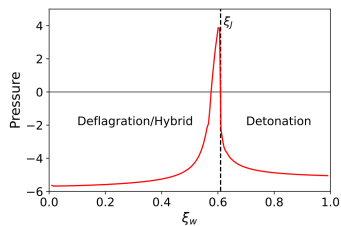
$p$  : pressure



- 3 unknowns ( $v_\pm$  &  $T_-$ )  $\Rightarrow$  Solve numerically

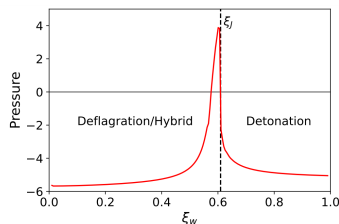
- Pressure numerically :

$$\mathcal{P}_{\text{net}} = -\Delta V_0 - \Delta V_T| + \bar{\mathcal{P}}_{\text{LTE}}$$



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- However, at Jouguet, where pressure is maximal:

$$\bar{\mathcal{P}}_{\text{LTE}}^{\text{max}} = a_+ T_+^4 \left[ \frac{4}{3} [\gamma(v_+)]^2 v_+ \left( v_+ - \frac{1}{\sqrt{3}} \right) \right]$$

$$(-\Delta V_T)|_{\xi_w = \xi_J} = \frac{a_+ T_+^4}{3} \left[ 1 - \frac{4}{9} \gamma(v_+)^4 b \right]$$

$$\Delta V_0 = a_+ T_+^4 \alpha_+$$

$a$  : Total DoF

$$b \equiv \frac{a_-}{a_+}$$

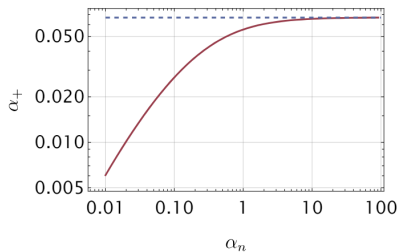
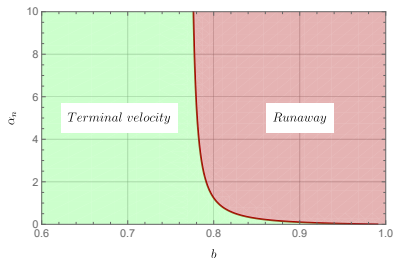
$$\alpha_+ \equiv \frac{\Delta V_0}{a_+ T_+^4}$$

$$\alpha_n \equiv \frac{\Delta V_0}{a_n T_n^4}$$

- $v_+(\alpha_+(\alpha_n)) \Rightarrow$  pressure depends only on  $\alpha_n$  &  $b$

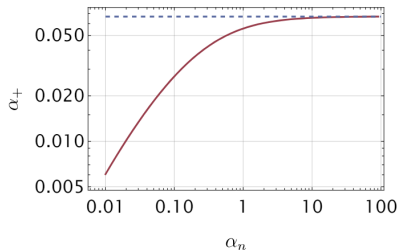
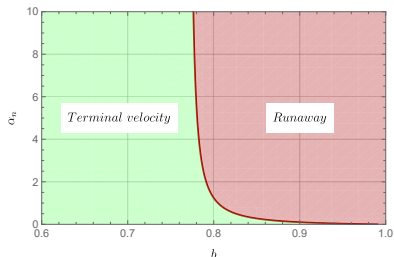
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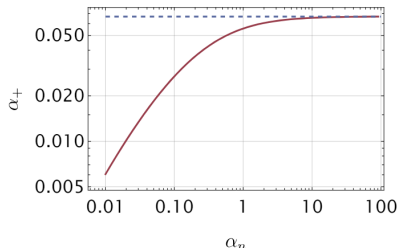
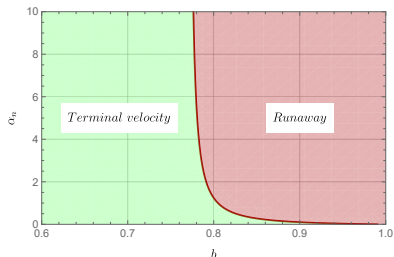


- Divergence at  $b \approx 0.77$  when  $\alpha_n \rightarrow \infty$ 
  - Boundary depends heavily on DoF in and outside



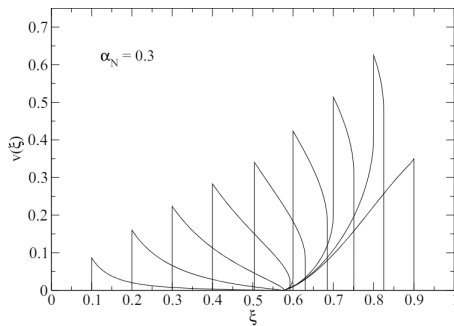
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- Divergence at  $b \approx 0.77$  when  $\alpha_n \rightarrow \infty$ 
  - Boundary depends heavily on DoF in and outside
- Saturation of  $\alpha_+$ 
  - Pressure independent of  $\alpha_n$  for supercooled PT

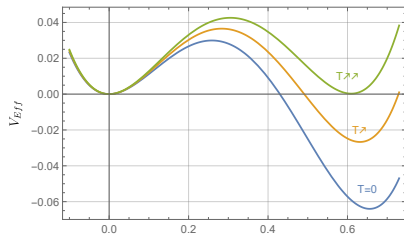
# Unlimited efficient heating



- $\xi_w \rightarrow \xi_J \Rightarrow |\xi_{sw} - \xi_w| \searrow$   
 $\Rightarrow T_+ \nearrow \Rightarrow \alpha_+ \searrow$

- Unphysical result  
→ Need for microphysics

Espinosa, Konstandin,  
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$$\left. \begin{array}{l} R_{\text{initial}} \sim \frac{1}{T_n} \\ \gamma_w \sim \frac{R}{R_{\text{initial}}} \end{array} \right\} \Rightarrow \text{Jouguet speed reached when } R \sim \text{few} \times R_{\text{initial}}$$

→ Strongly interacting particles such that  $R \gg L_{\text{MFP}}$

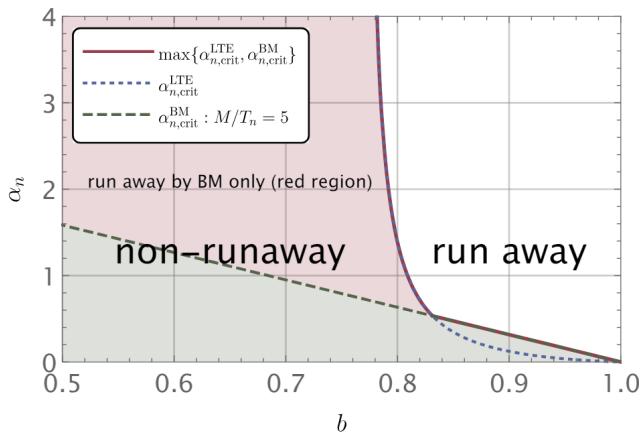
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- Strongly interacting particles such that  $R \gg L_{\text{MFP}}$
- Validity of expansion modes obtained in static picture
  - Numerical simulations: Is adaptation to profile with new velocity faster than bubble acceleration?

# LTE vs. B & M



- If LTE picture holds, much more parameter space excluded to be runaway!

# Conclusions

- Analytic expressions for hydrodynamic quantities when  $\xi_w \rightarrow \xi_J$

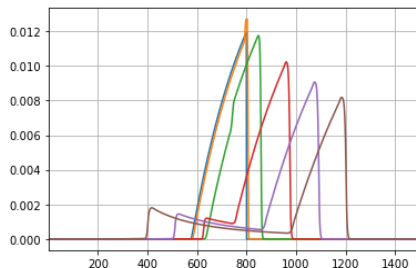
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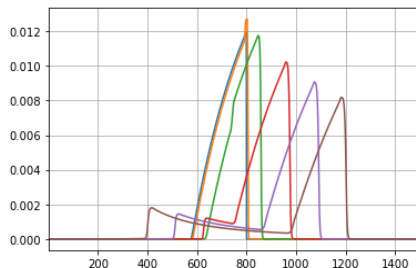
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- Extra criterion on top of B & M criterion

# Bibliography

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