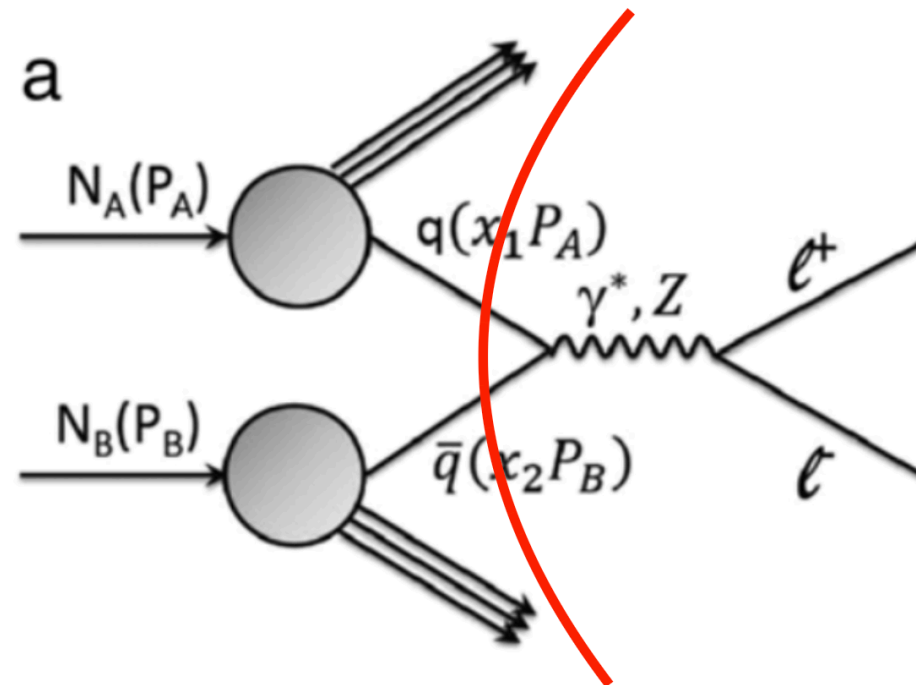


Proton structure information from A_{FB} of Drell-Yan process

Mingzhe Xie

2025.11.4

- QCD improved parton model
 - Interactions between partons can be ignored
 - Initial state parton information can be factorized into momentum-fraction x
 - Described by parton distribution function (PDF)



$$\sigma(p) = \sum_i \sigma(q_i) = \sum_i q_i(x) \otimes \text{HardProcess}(q_i; x)$$

Highly rely on the experimental measurements and QCD calculations

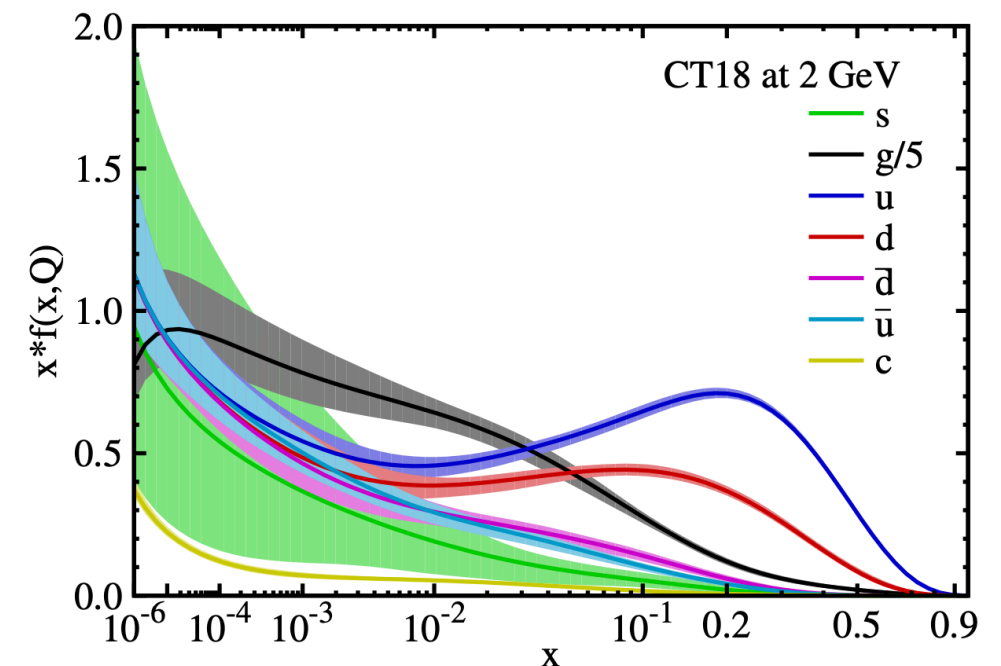
Low energy scale: non-perturbative quarks and gluon
 $u(x), d(x), \bar{u}(x), \bar{d}(x), s(x), \bar{s}(x), g(x)$
 i.e $q(x) = x^{a_0}(1-x)^{a_1} \cdot P(x; a_i)$



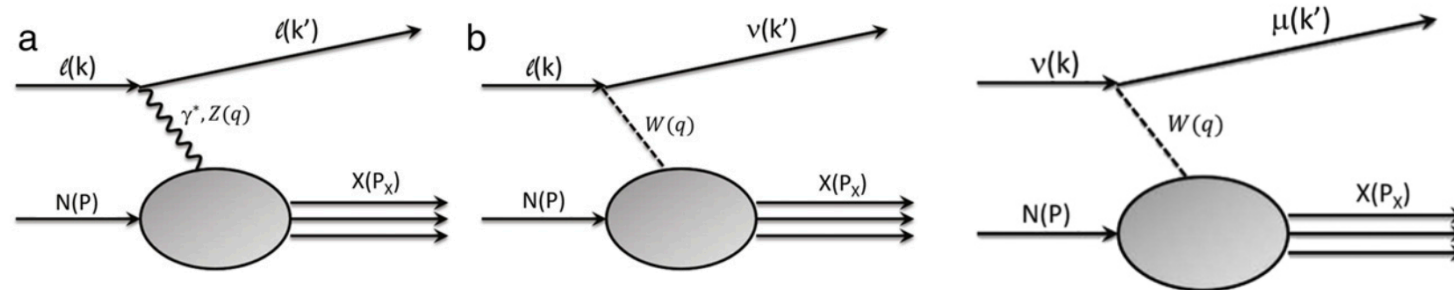
Compare with experimental measurement and fit the non-perturbative part



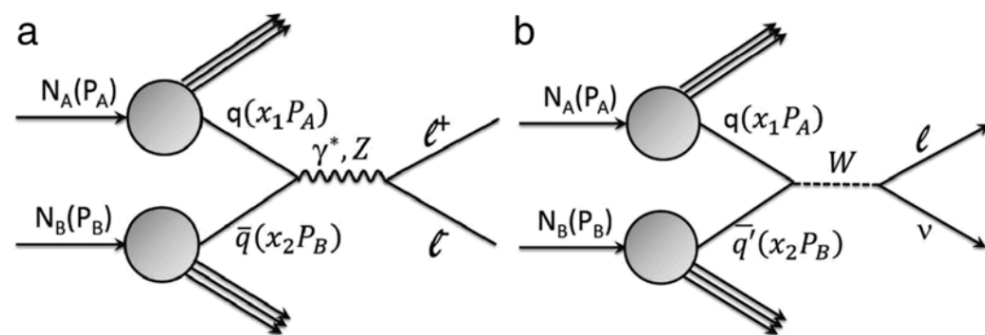
Perturbative QCD calculation
 Evolute to any energy scale (DGLAP)
 Heavy quarks $c(x), b(x)$
 Sum rule



- Mainly DIS and Drell-Yan measurements, plus a few jet and top measurements

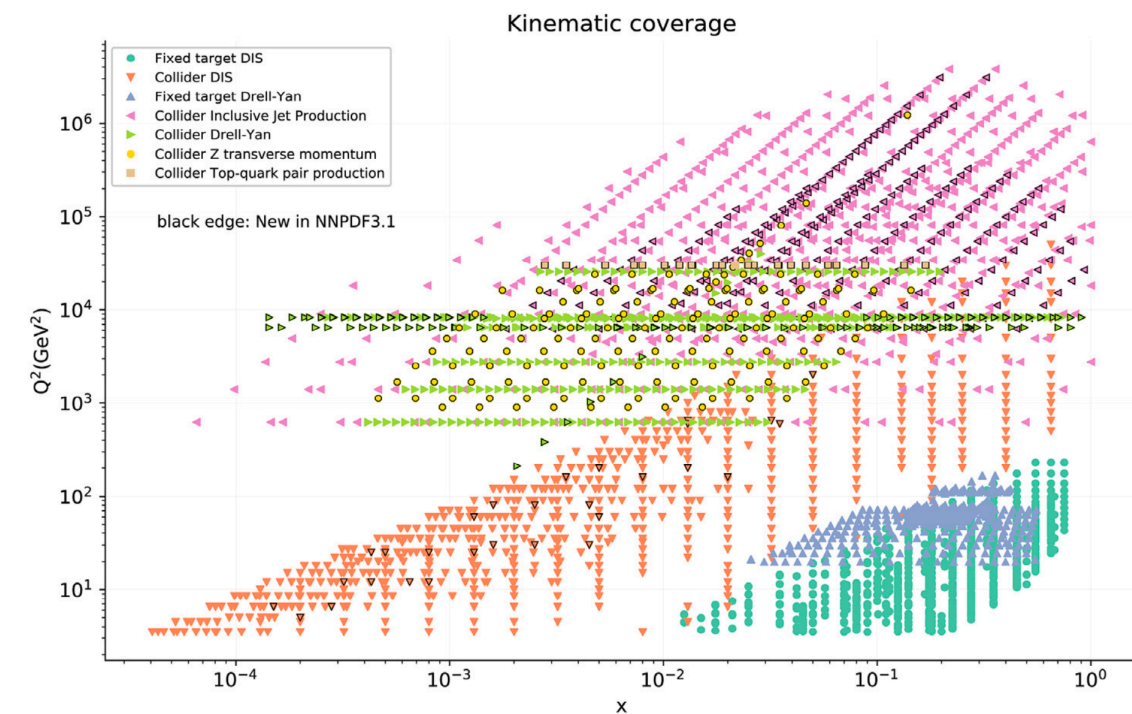


lepton(neutrino)-nucleon scattering



Neutral(charge)-current Drell-Yan process

- Experimental input covers large kinematic regions
 - Q : $O(1) \sim O(100)$ GeV
 - x : $O(10^{-4}) \sim 0.4$



Npt ~ 4000

- “Sum” quantity

- Cross section measurements are measuring sum of quark distributions
- In the past decades, “sum” quantity is measured from Drell-Yan and DIS cross sections, and then used to predict cross sections
- Well measured from the cross sections

$$\sigma_{DIS} \sim \sigma_u u(x) + \sigma_d d(x) + \sigma_{\bar{u}} \bar{u}(x) + \sigma_{\bar{d}} \bar{d}(x) + \dots$$

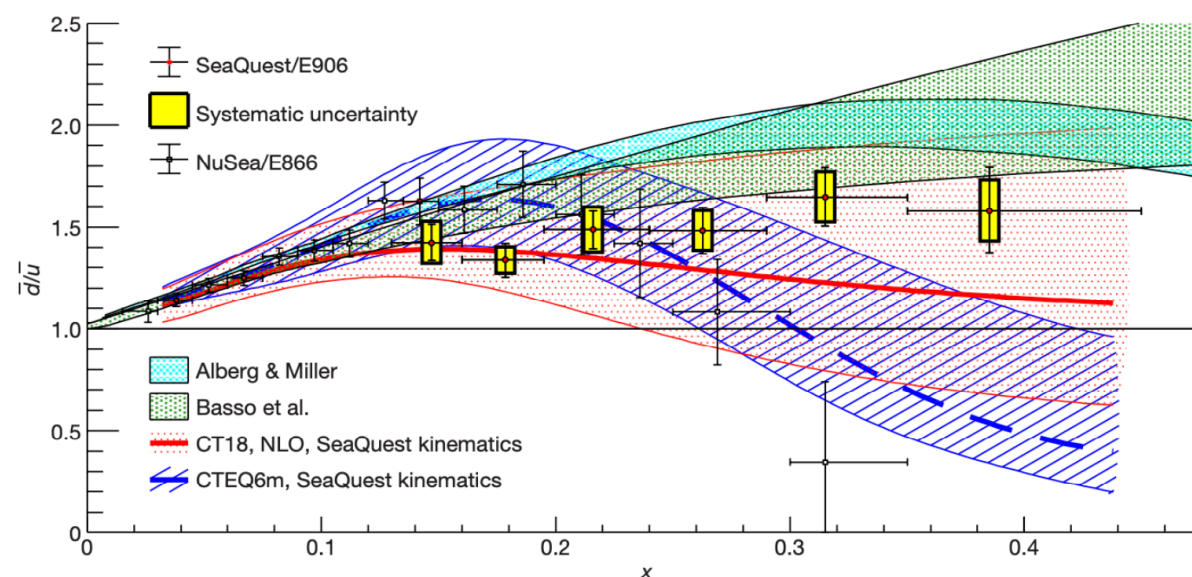
$$\sigma_{DY}^{pp} \sim \sigma_{u\bar{u}} u(x_1) \bar{u}(x_2) + \sigma_{d\bar{d}} d(x_1) \bar{d}(x_2) + \dots$$

- “Difference” quantity

- Relative difference between various quark distributions
- Directly related to proton structure and QCD theory
- Experimental measurement very limited

- Relative difference between different flavor quarks
 - Sea quark generation mechanism in proton: gluon splitting?
 - $\bar{d}/\bar{u} \neq 1$: SU(2) flavor asymmetry
 - $(s + \bar{s})/(\bar{d} + \bar{u}) \neq 1$: SU(3) flavor asymmetry

SU(2) flavor asymmetry: $\bar{d} \neq \bar{u}$



NuSea/SeaQuest pp/pD scattering

SU(3) flavor asymmetry: $(s + \bar{s}) \neq (\bar{d} + \bar{u})$

$$\kappa_s = \frac{\int_0^1 x[s(x, \mu^2) + \bar{s}(x, \mu^2)]dx}{\int_0^1 x[\bar{u}(x, \mu^2) + \bar{d}(x, \mu^2)]dx}$$

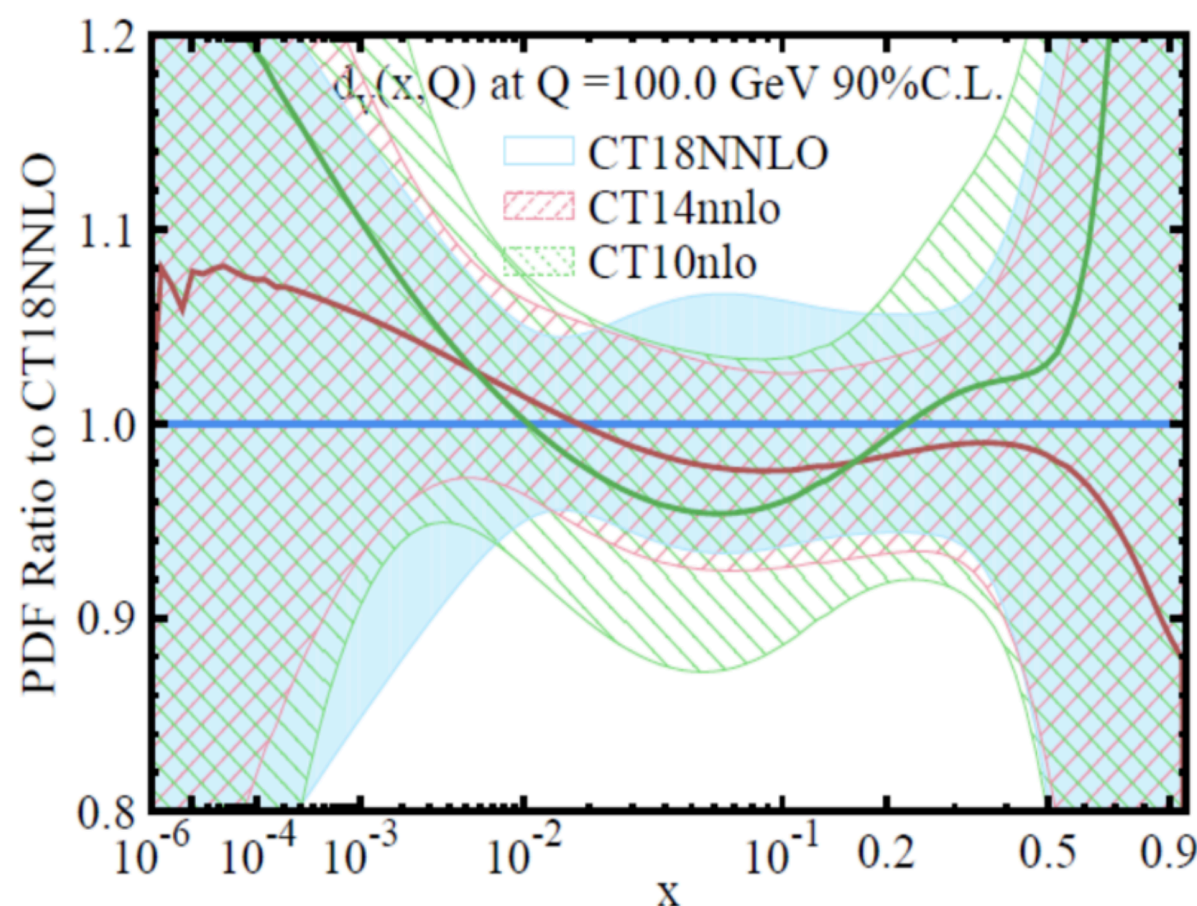
Experiment (year)	QCD order	κ_s
CDHS (1982)	LO	0.52 ± 0.09
CCFR (1993)	LO	$0.373^{+0.048}_{-0.041} \pm 0.018$
CCFR (1995)	NLO	$0.477^{+0.051-0.017}_{-0.050+0.036}$
CHARMII (1999)	LO	$0.388^{+0.074}_{-0.061} \pm 0.067$
NOMAD (2000)	LO	$0.48^{+0.09+0.17}_{-0.07-0.12}$
NuTeV (2001)	LO	$0.38 \pm 0.08 \pm 0.043$
NuTeV (2007)	NLO	
CHORUS (2008)	NLO	$0.33 \pm 0.05 \pm 0.05$
NOMAD (2013)	NNLO	0.591 ± 0.019

Some experiments showing less s quark

Obtain “difference” indirectly from “sum”?

- Difficult, as shown from valence d_V quark
 - From CT14 to CT18, lots of high precision LHC Drell-Yan measurements introduced, but the precision of d_V is not improved.
 - LHC Drell-Yan process contains large d quark contributions. However, these measurements are the sum of $u, d, \bar{u}, \bar{d} \dots$

$$\sigma_{DY}^{pp} \sim \sigma_{u\bar{u}} u(x_1) \bar{u}(x_2) + \sigma_{d\bar{d}} d(x_1) \bar{d}(x_2) + \dots$$

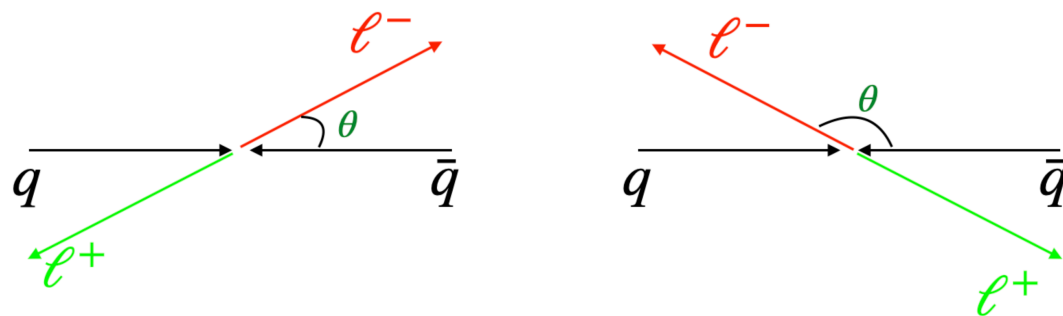


More direct measurement on the “difference” quantity is needed!

Precision of d_V is not improved!

- Originate from parity violation in weak interactions

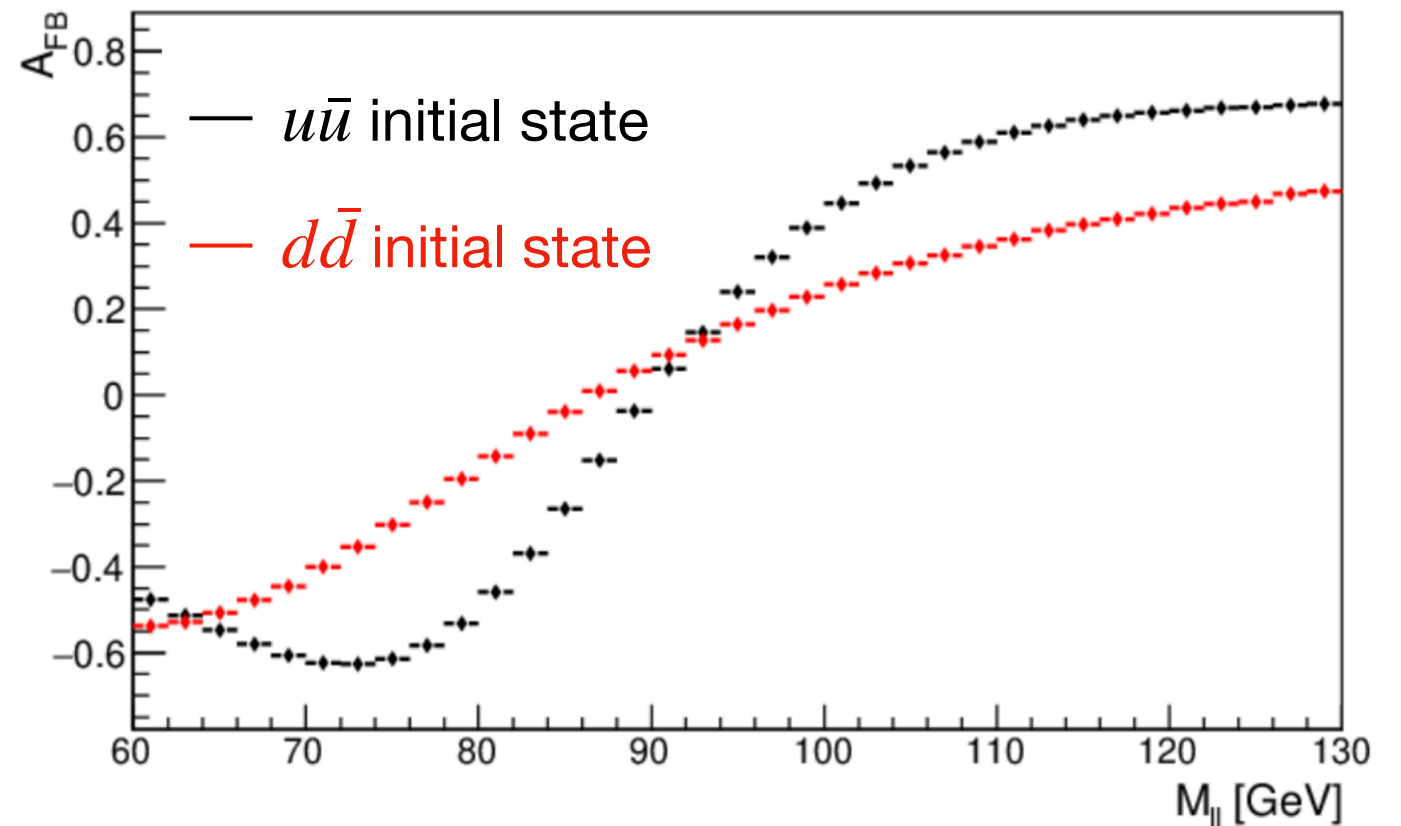
- $qh(q\bar{q}) \rightarrow Z/\gamma^* \rightarrow l^+l^-$



Forward event
 $\cos\theta > 0$

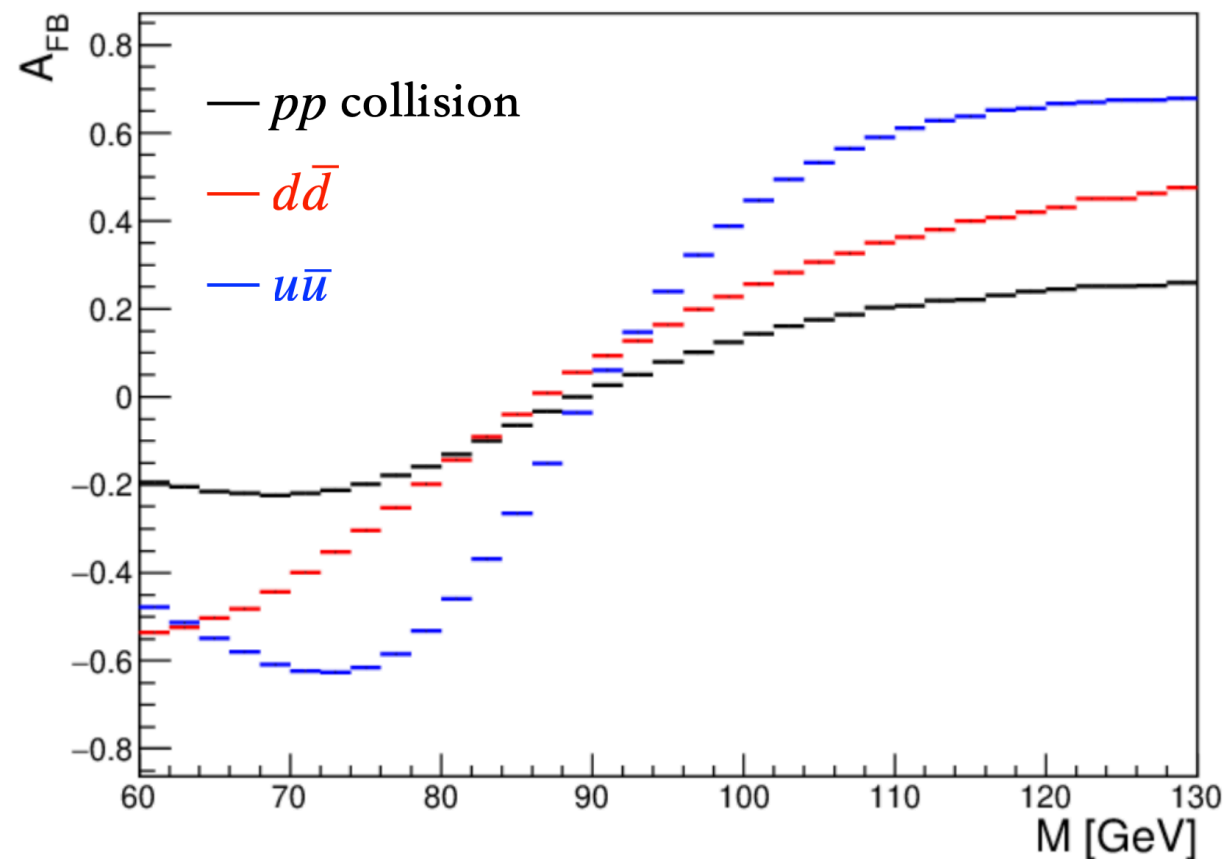
backward event
 $\cos\theta < 0$

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$

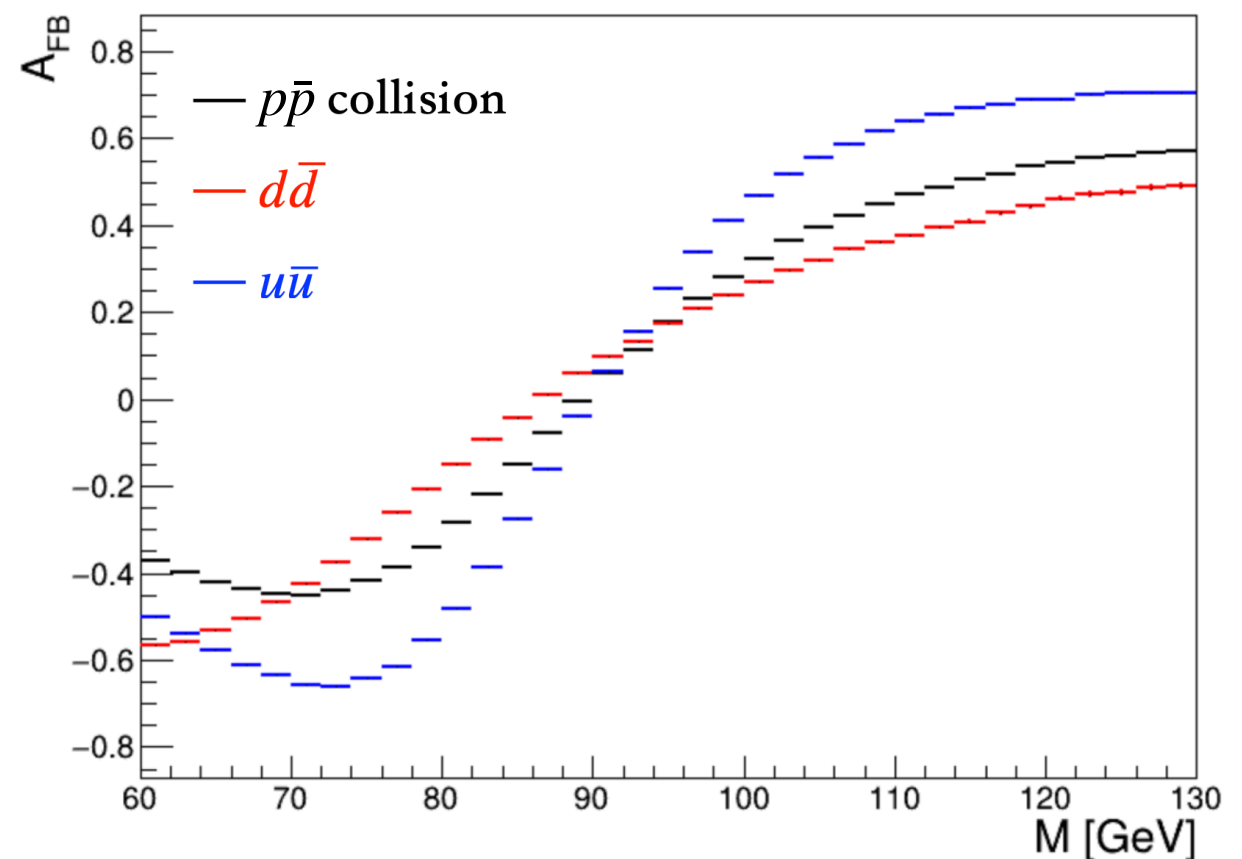


A_{FB} from different $q\bar{q} \rightarrow Z/\gamma^* \rightarrow l^+l^-$ hard process

- The observed A_{FB} is significantly affected by initial state quarks' PDF
 - At LHC, the observed A_{FB} is globally suppressed
 - At Tevatron, the observed A_{FB} is an average between $u\bar{u}$ and $d\bar{d}$ A_{FB}



LHC pp collision



Tevatron $p\bar{p}$ collision

- The observed A_{FB} at hadron collider can be factorized
 - A_{FB}^u : asymmetry of hard process $u\bar{u} \rightarrow Z/\gamma^* \rightarrow l^+l^-$
 - A_{FB}^d : asymmetry of hard process $d\bar{d} \rightarrow Z/\gamma^* \rightarrow l^+l^-$
 - C_u : structure parameter of u quark
 - C_d : structure parameter of d quark
- } Pure electroweak calculation, independent of proton structure and QCD calculation
 } Fully represent proton structure information

$$A_{FB}^h(Y, M, Q_T) = C_u(Y, M, Q_T) \times A_{FB}^u(Y, M, Q_T) + C_d(Y, M, Q_T) \times A_{FB}^d(Y, M, Q_T)$$

The factorization is valid to all orders.

Phys. Rev. D 106, 033001 (2022)

Eur. Phys. J. C 82:368 (2022)

Chin. Phys. C 45, 053001 (2021)

$$A_{FB}^h(Y, M, Q_T) = C_u(Y, M, Q_T) \times A_{FB}^u(Y, M, Q_T) + C_d(Y, M, Q_T) \times A_{FB}^d(Y, M, Q_T)$$

$$x_{1,2} = \frac{\sqrt{M^2 + Q_T^2}}{\sqrt{s}} \times e^{\pm Y}$$

$$C_u(x_1, x_2) = \frac{[u(x_1)\bar{u}(x_2) - \bar{u}(x_1)u(x_2)]\sigma_u}{\sum_{q=u,d,s,c,b} [q(x_1)\bar{q}(x_2) + \bar{q}(x_1)q(x_2)]\sigma_q}$$

$$C_u(x_1, x_2) = \frac{[u(x_1)u(x_2) - \bar{u}(x_1)\bar{u}(x_2)]\sigma_u}{\sum_{q=u,d,s,c,b} [q(x_1)q(x_2) + \bar{q}(x_1)\bar{q}(x_2)]\sigma_q}$$

$$C_d(x_1, x_2) = \frac{[d(x_1)\bar{d}(x_2) - \bar{d}(x_1)d(x_2)]\sigma_d}{\sum_{q=u,d,s,c,b} [q(x_1)\bar{q}(x_2) + \bar{q}(x_1)q(x_2)]\sigma_q}$$

$$C_d(x_1, x_2) = \frac{[d(x_1)d(x_2) - \bar{d}(x_1)\bar{d}(x_2)]\sigma_u}{\sum_{q=u,d,s,c,b} [q(x_1)q(x_2) + \bar{q}(x_1)\bar{q}(x_2)]\sigma_q}$$

C_u and C_d for LHC pp collision

C_u and C_d for Tevatron $p\bar{p}$ collision

- Small x probes relative difference between quark and antiquark
 - At LHC, small Y region corresponds to $x \sim 0.01$

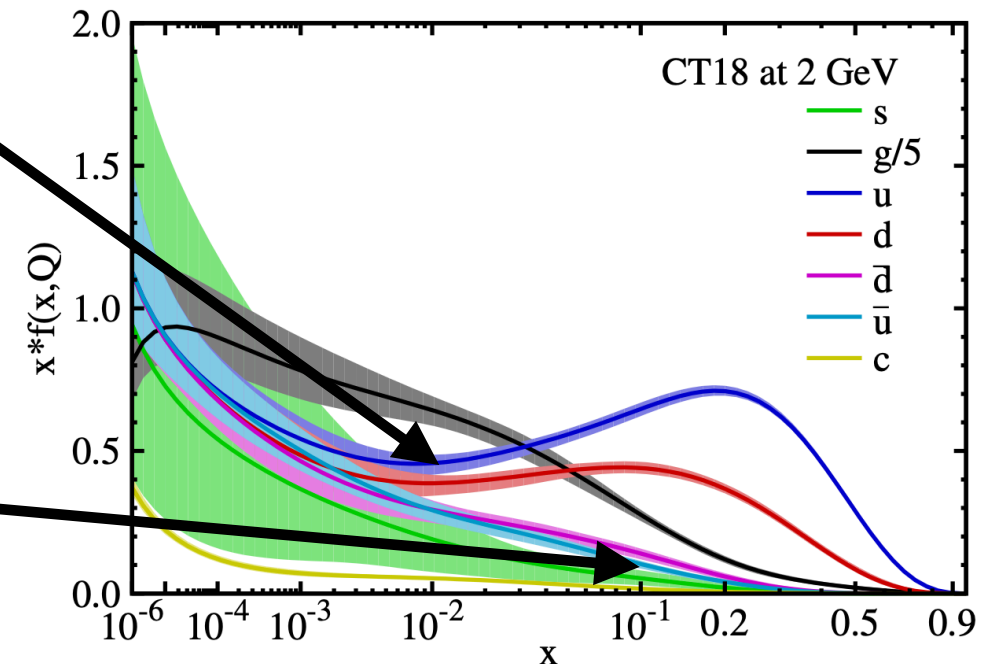
$$C_u \propto u(x_1)\bar{u}(x_2) - \bar{u}(x_1)u(x_2)$$

$$C_d \propto d(x_1)\bar{d}(x_2) - \bar{d}(x_1)d(x_2)$$

- Large x probes relative difference between u and d
 - At Tevatron and high Y region at LHC, $x \sim 0.1$

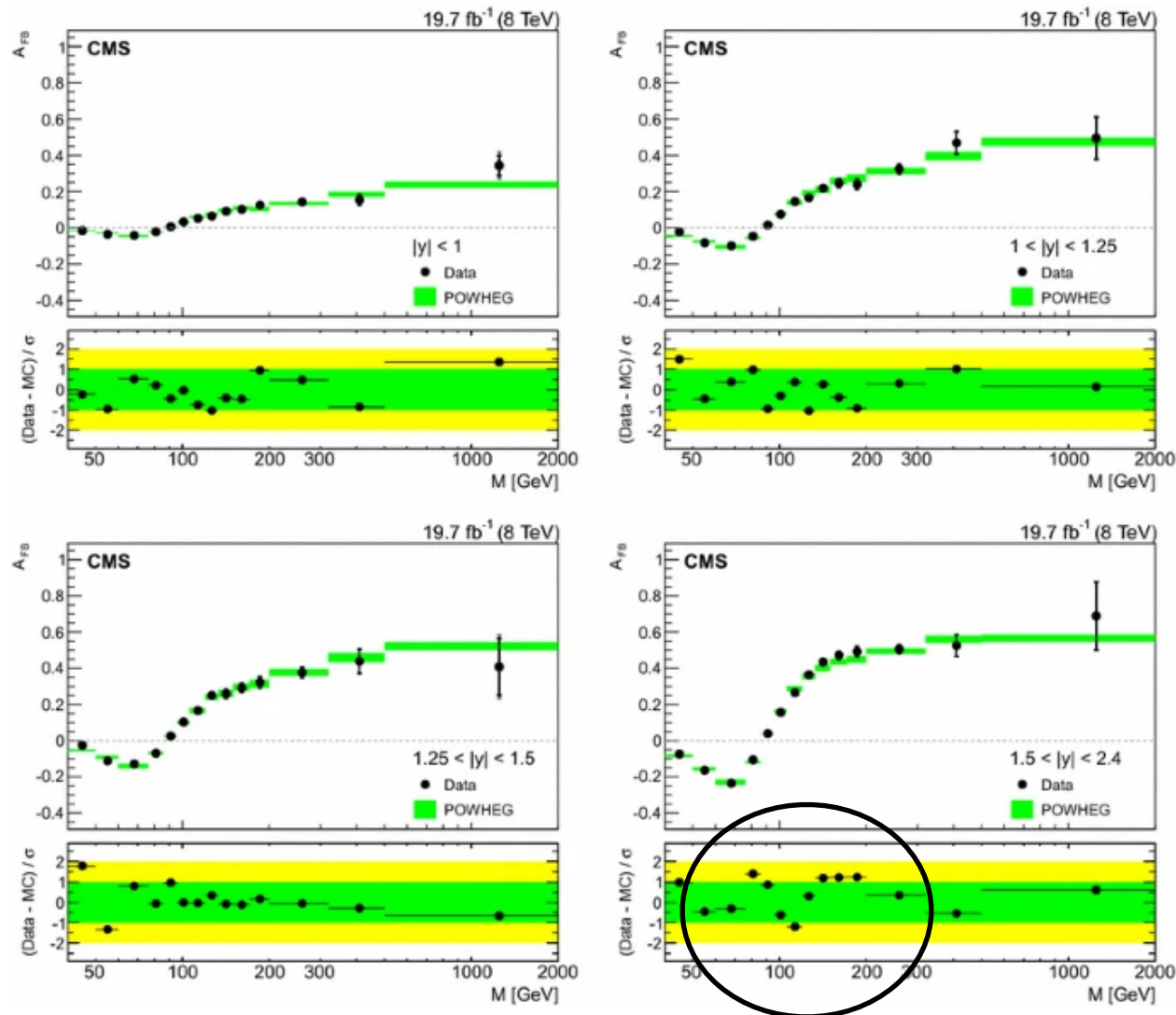
$$R_{\text{LHC}} = \frac{C_u}{C_d} \approx \frac{u}{d}$$

$$R_{\text{Tevatron}} = \frac{C_u}{C_d} \approx \frac{u^2}{d^2}$$

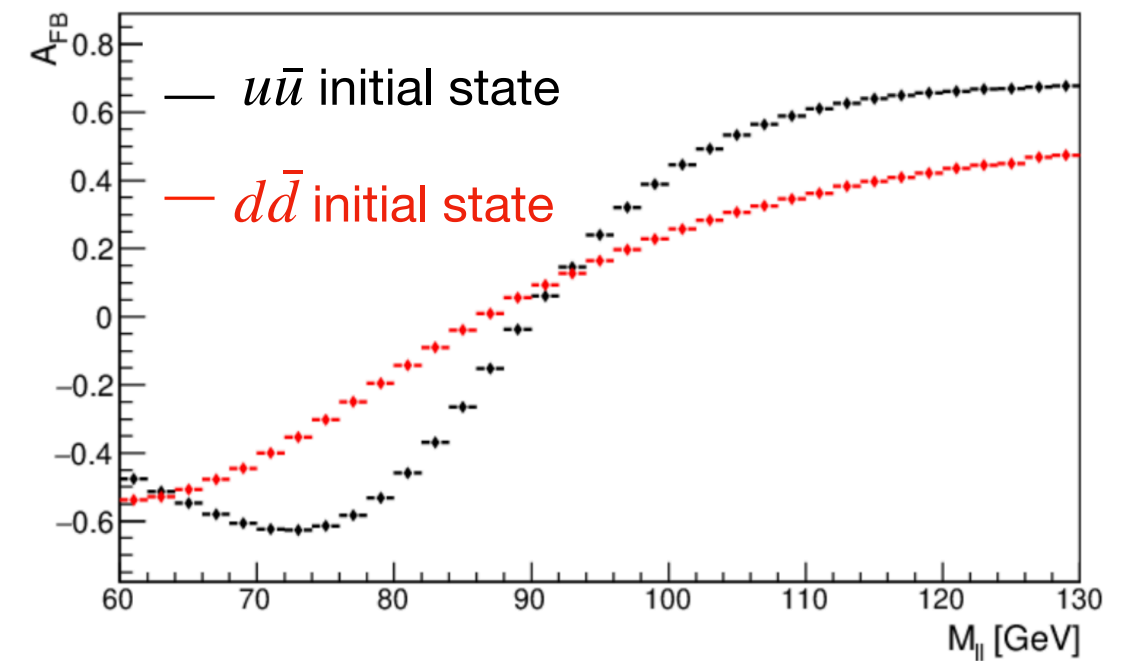


- Extract structure parameters from multiple experiments:
 - At LHC, from the published CMS 8 TeV and 13 TeV unfolded A_{FB}
 - At LHC, directly from ATLAS 13 TeV 2015+2016 data
 - At Tevatron, directly from D0 1.96 TeV RunII data

- CMS 8 TeV 19.6 fb^{-1} , both electron and muon channels used, 4 $|Y|$ bins



Theory prediction: Powheg+CT10

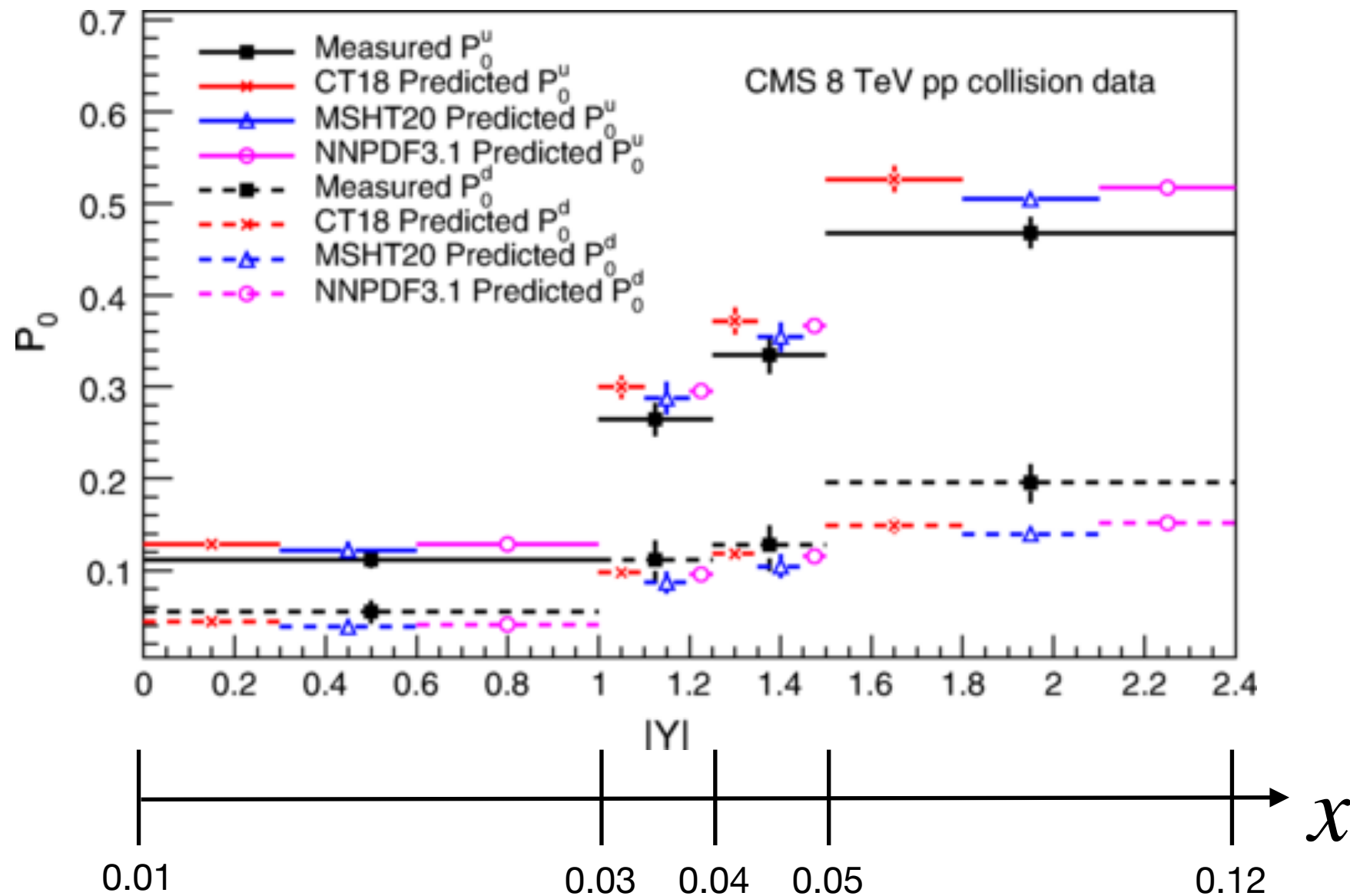


$$A_{FB}^h(Y, M, Q_T) = C_u(Y, M, Q_T) \times A_{FB}^u(Y, M, Q_T) + C_d(Y, M, Q_T) \times A_{FB}^d(Y, M, Q_T)$$

A_{FB} in data is smaller than theory prediction, suggesting more d quark contributions.

- The measured C_u is obviously lower than prediction, and measured C_d is higher
 - First evidence: more d quark contribution and less u quark contribution

Phys. Rev. D 107, 054008(2023)



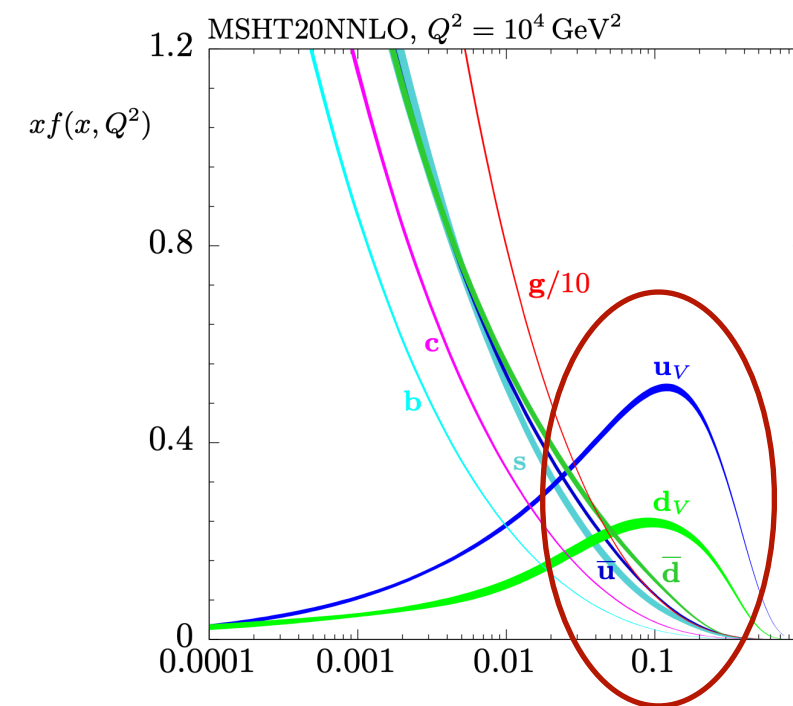
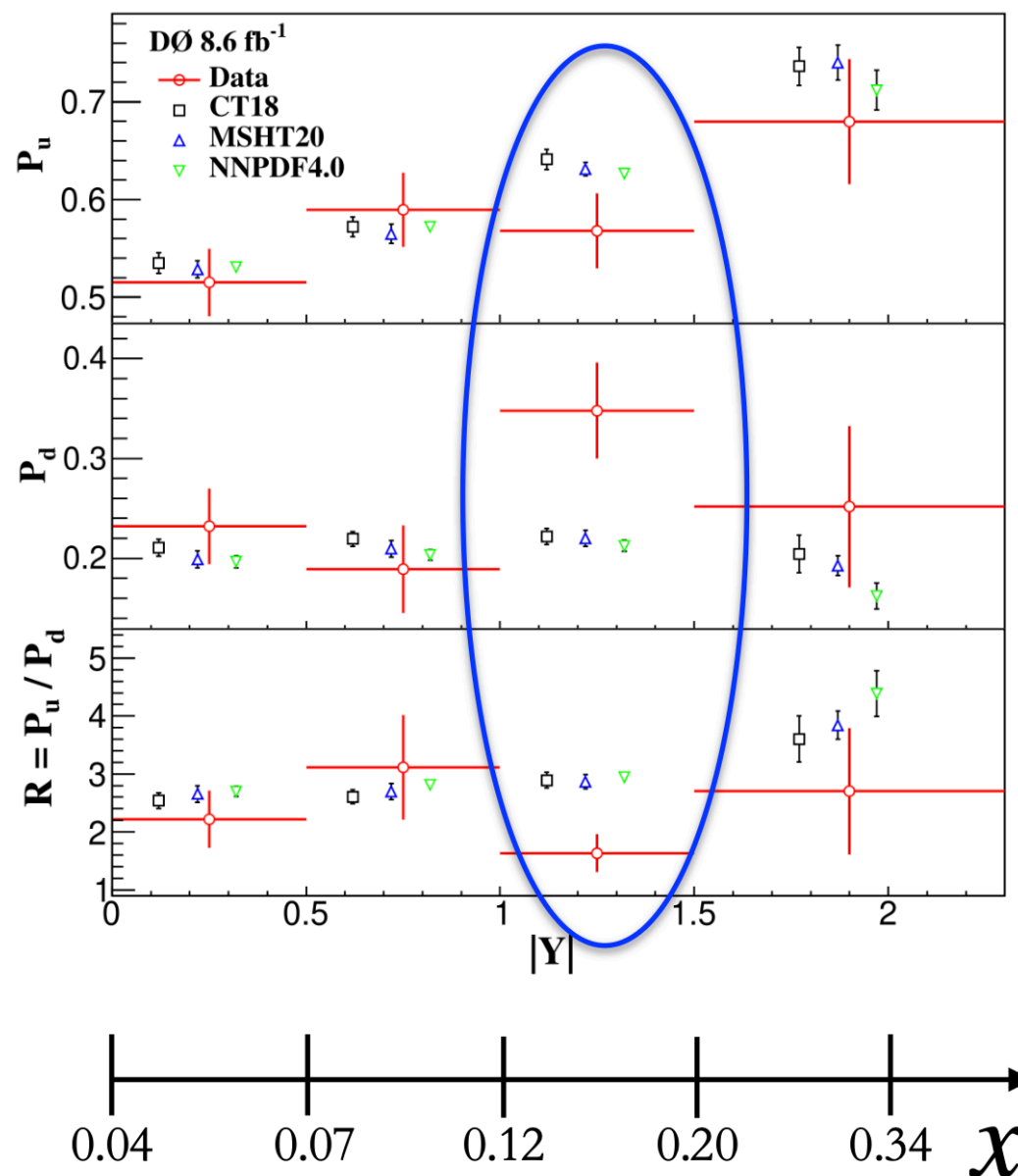
- 3.5σ deviation on u/d found in $x \sim [0.1, 0.2]$ (peak region of valence quark distribution)

Phys. Rev. D 110, L091101(2024)

$$C_u(x_1, x_2) \propto u(x_1)u(x_2) - \bar{u}(x_1)\bar{u}(x_2)$$

$$C_d(x_1, x_2) \propto d(x_1)d(x_2) - \bar{d}(x_1)\bar{d}(x_2)$$

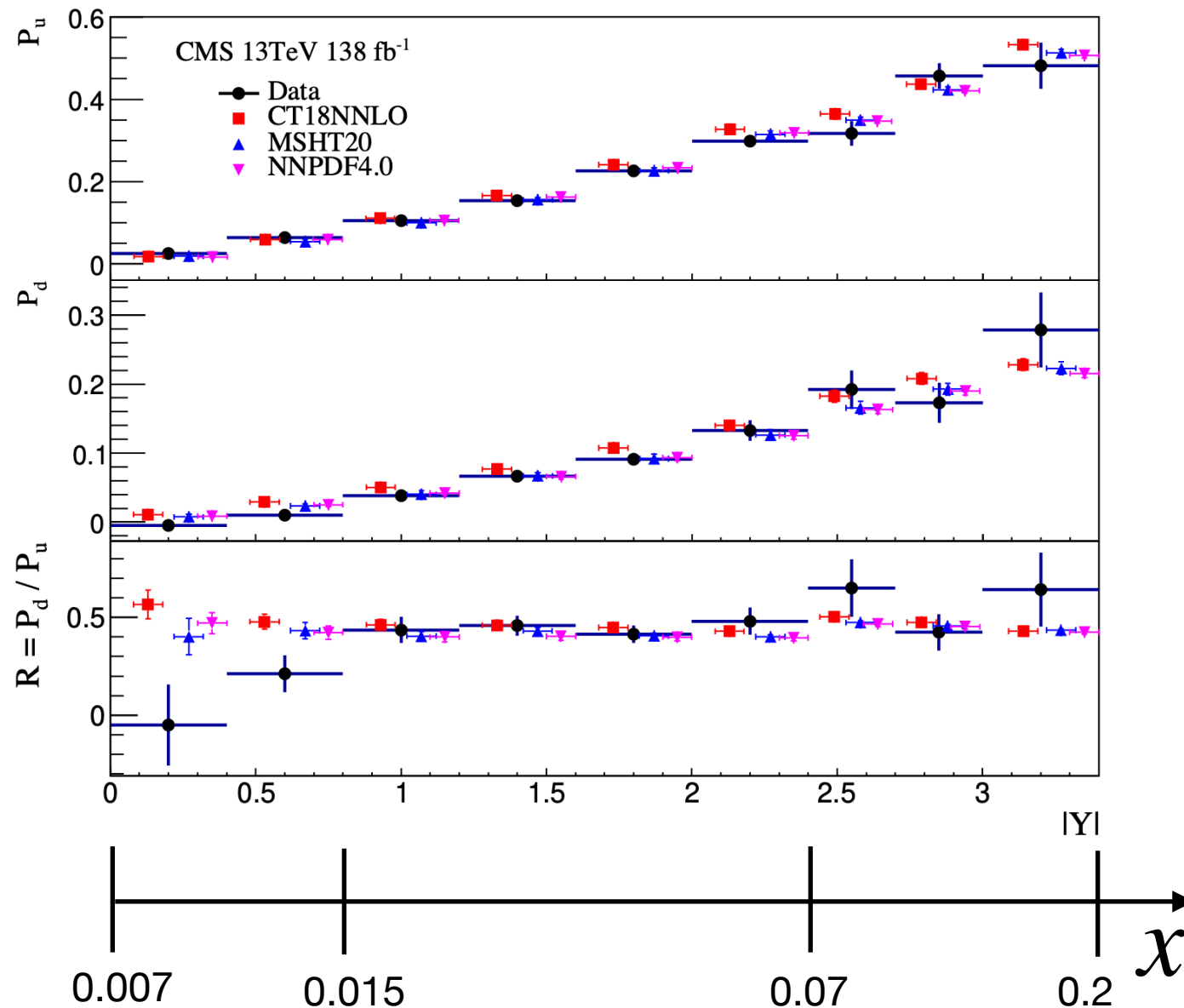
$$R_{\text{Tevatron}} = \frac{C_u}{C_d} \approx \frac{u^2}{d^2}$$



Further evidence that more d quark contribution and less u quark contribution

- Using 13 TeV unfolded A_{FB} , $|Y|$ to 3.4, corresponding x range [0.01, 0.2]
 - Provide a clear trend for the evolution of the u, d quark ratio with x

arXiv:2505.17608



- high $|Y|$ region: R is higher, but limited by statistics
- Middle $|Y|$: data and prediction agree well
- Low $|Y|$:
 - R is lower in data
 - C_d is closer to zero, means $d(x) \approx \bar{d}(x)$, perturbative QCD is more significant.

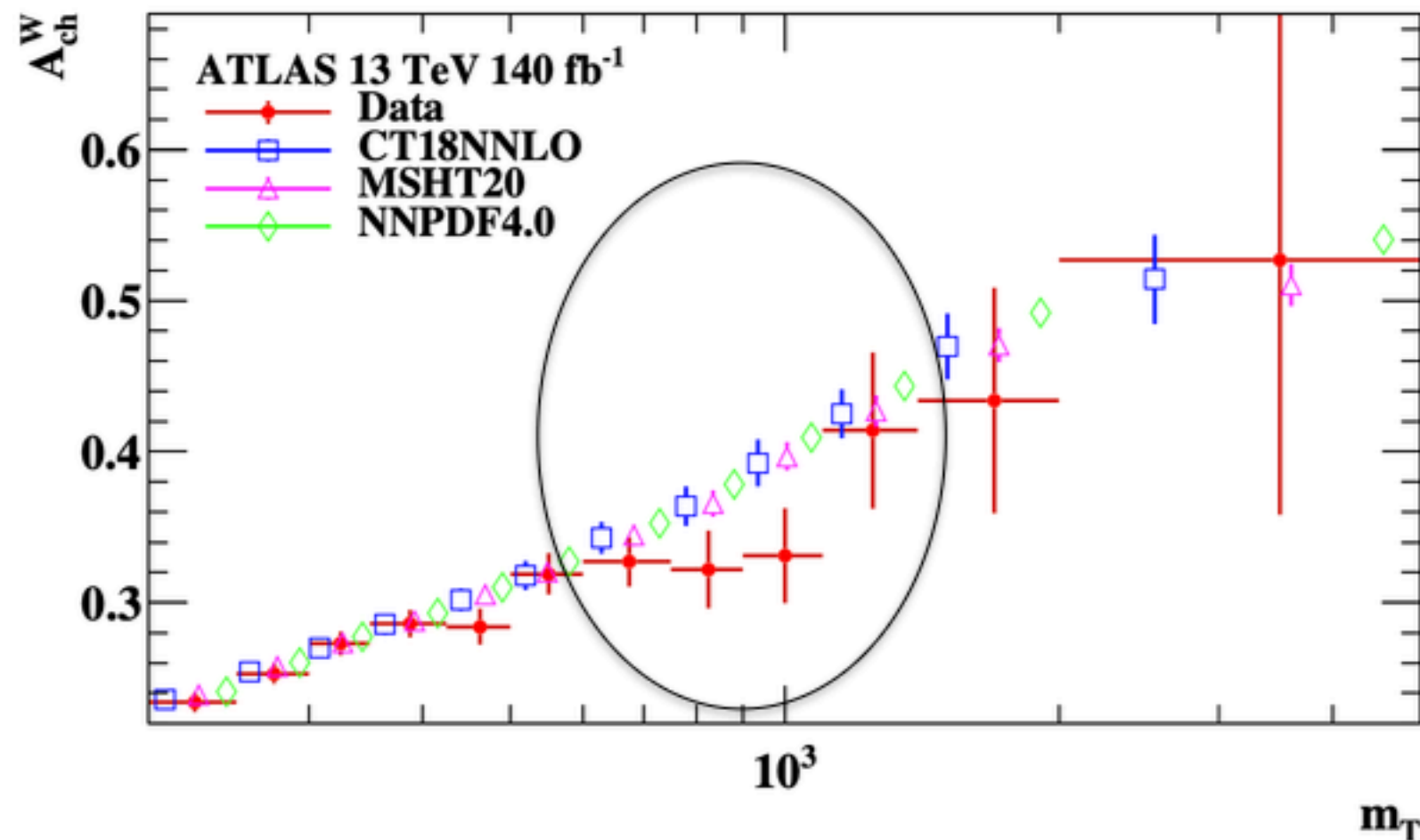
Extension

Same region covered by CMS 8TeV

Overlapping region with CMS 8TeV and D0

- High m_T measurement of the W charge asymmetry
 - With high m_T requirement, the high energy hadron collider data can reach a rather large x region around 0.1

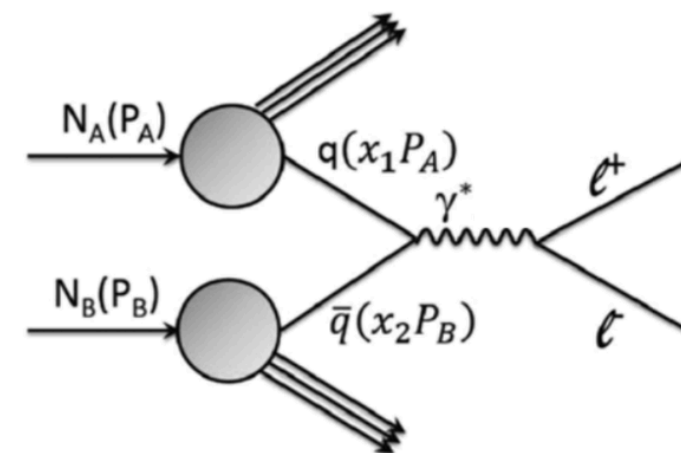
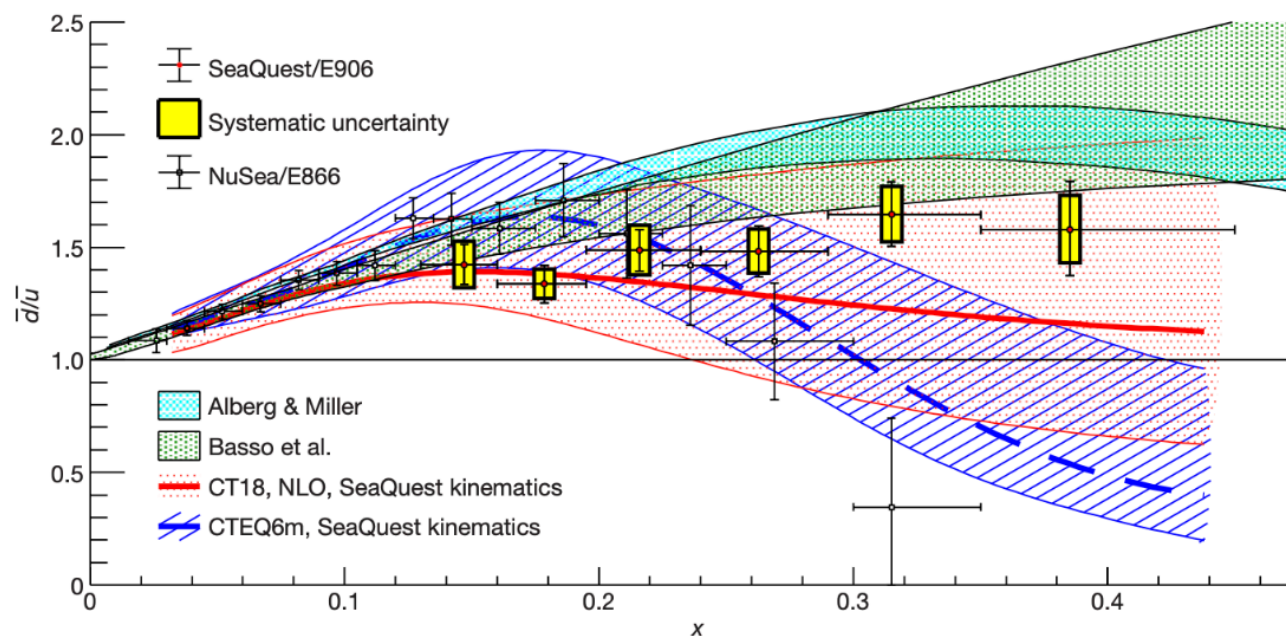
$$A_W \sim \frac{u(x_1)\bar{d}(x_2) + \bar{d}(x_1)u(x_2)}{d(x_1)\bar{u}(x_2) + \bar{u}(x_1)d(x_2)} = \frac{u(x_1) + \bar{d}(x_1)}{d(x_1) + \bar{u}(x_1)} \quad \text{JHEP 07 (2025) 026}$$



- The only evidence comes from the comparison between proton and deuteron interactions

NuSea: Phys. Rev. D 64, 052002 (2001)

SeaQuest: Nature Vol 590, 561-565, 2021



$$\frac{\sigma_{pD}}{\sigma_{pH}} = \frac{\sigma_{pp} + \sigma_{pn}}{\sigma_{pp}} = 1 + \frac{\sigma_{pn}}{\sigma_{pp}}$$

$$\begin{aligned} \frac{\sigma_{pn}}{\sigma_{pp}} &= \frac{u(x_1)\bar{d}(x_2) + \frac{1}{4}d(x_1)\bar{u}(x_2)}{u(x_1)\bar{u}(x_2) + \frac{1}{4}d(x_1)\bar{d}(x_2)} \\ &= \frac{\frac{\bar{d}(x_2)}{\bar{u}(x_2)} + \frac{1}{4}\frac{d(x_1)}{u(x_1)}}{1 + \frac{1}{4}\frac{d(x_1)}{u(x_1)}\frac{\bar{d}(x_2)}{\bar{u}(x_2)}} \approx \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \end{aligned}$$

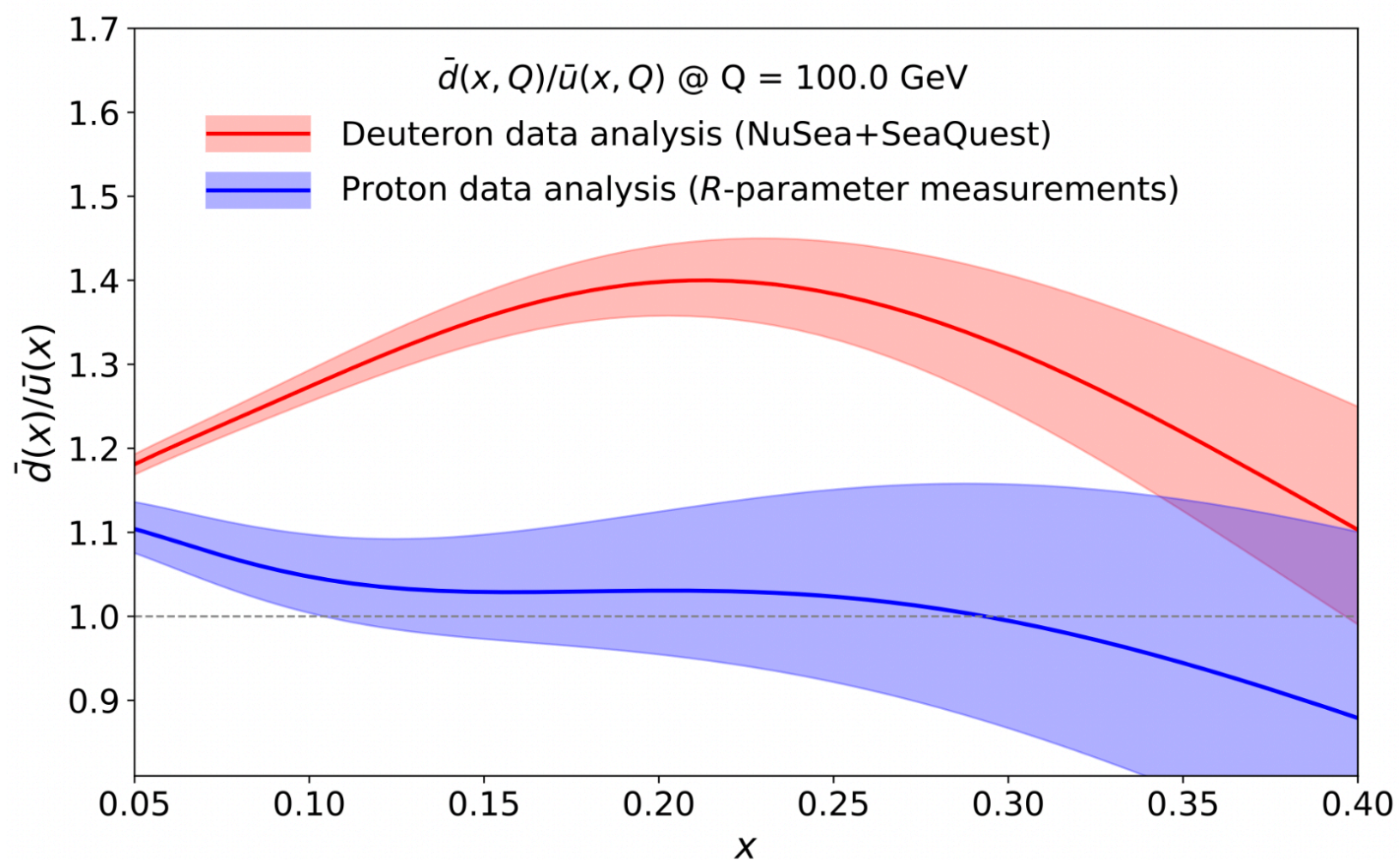
- Assumption 1: nuclear effect is small: $D = p+n$
- Assumption 2: isospin symmetry

$$\frac{\sigma_{pD}}{\sigma_{pp}} \approx 1 + \frac{\bar{d}}{\bar{u}}$$

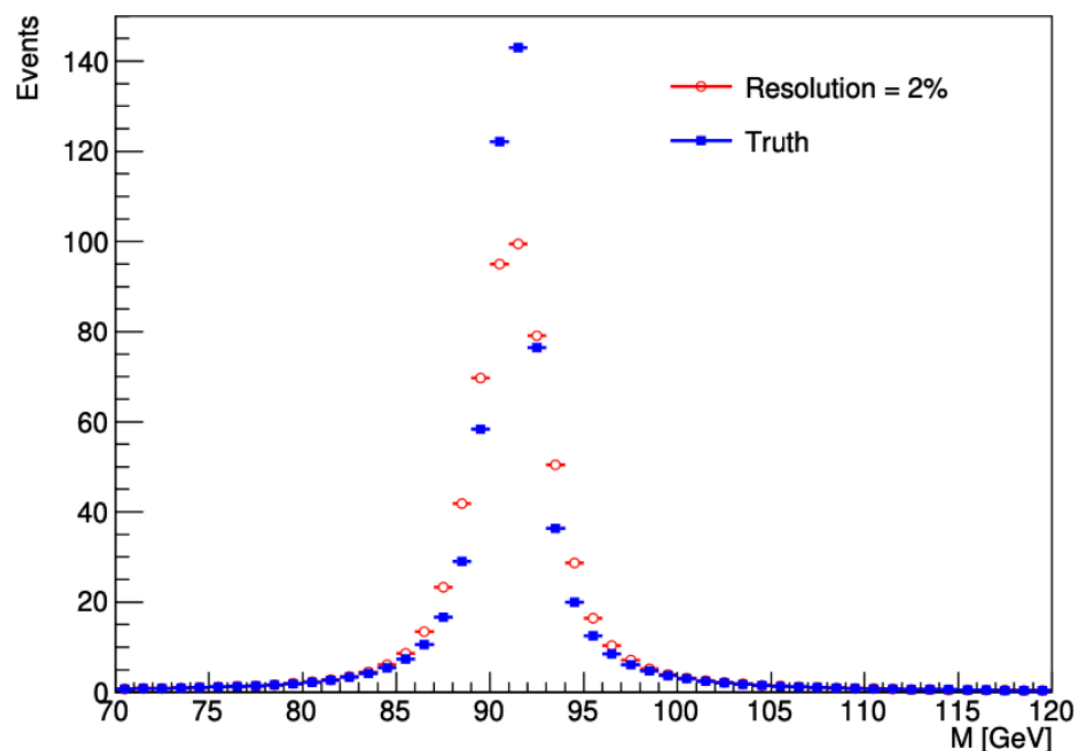
- PDF global fit with hadron collider data + specific data

arXiv:2510.08941

- Why do we see discrepancies between the pure proton data and deuteron data?
- What impact it would have on our knowledge of proton structure?



- Unfolding method introduce potential theoretical model uncertainties
- The purity is lower, the model uncertainties are larger
- Unfortunately, A_{FB} unfolding has low purity due to Z peak shape
- Unfortunately, the major model uncertainty in A_{FB} unfolding is PDF



- It's not reasonable to use a PDF-dependent result to measure PDF!

- Newly defined structure parameters from A_{FB}
- Measurement of structure parameters show deviation on u, d quark ratio
- Discrepancy observed from pure proton data and deuteron data
- Lack of direct measurement on CMS

