

Distinguishing Dark Matter Properties through Large Scale Structure

JCAP 01 (2024) 023

Yuan-Zhen Li

Oct 27, 2025
14th CosPa Meeting

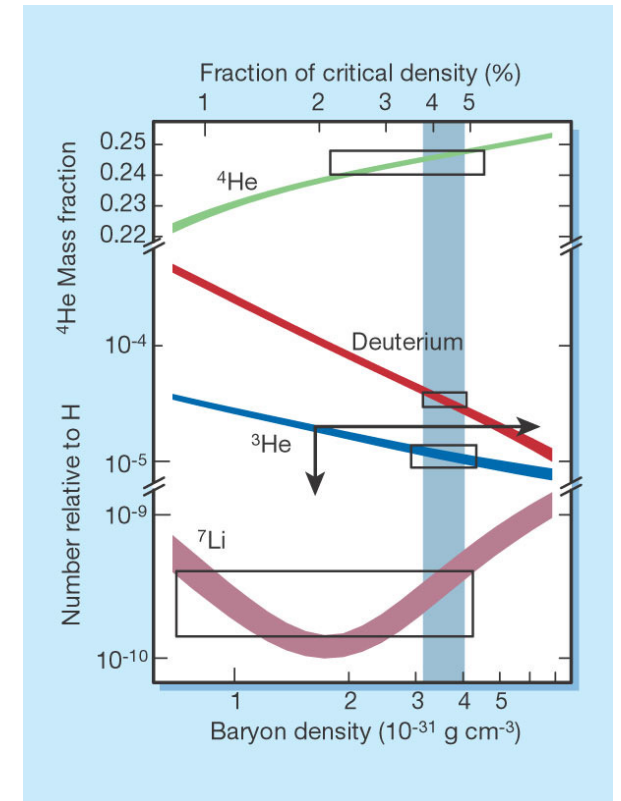
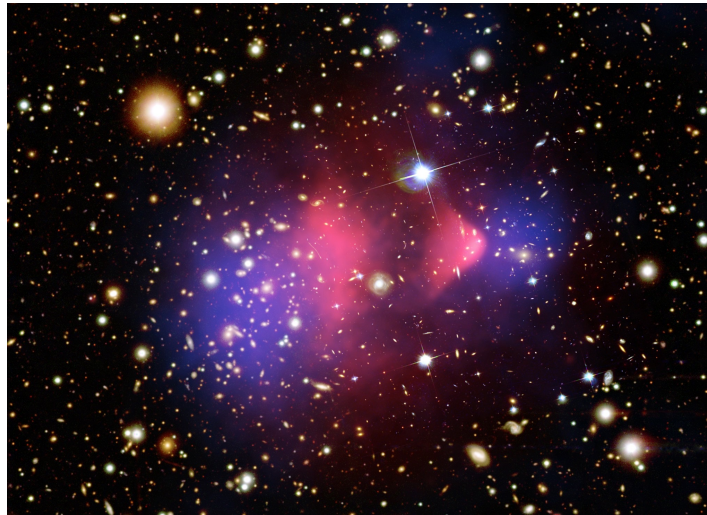
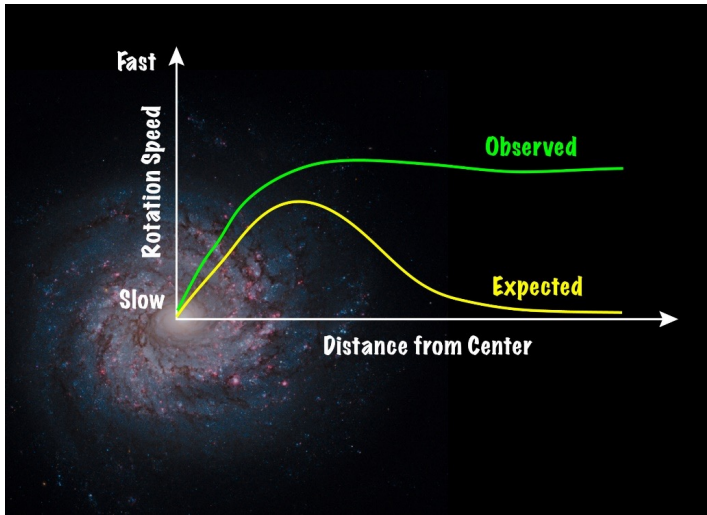
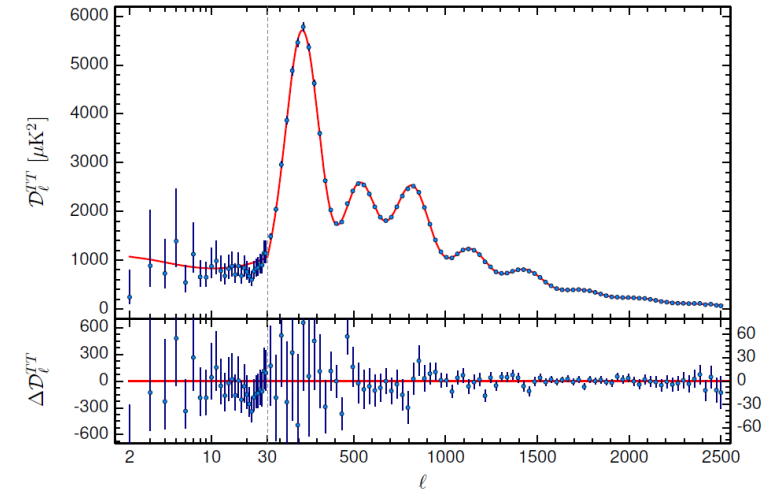


UCLouvain

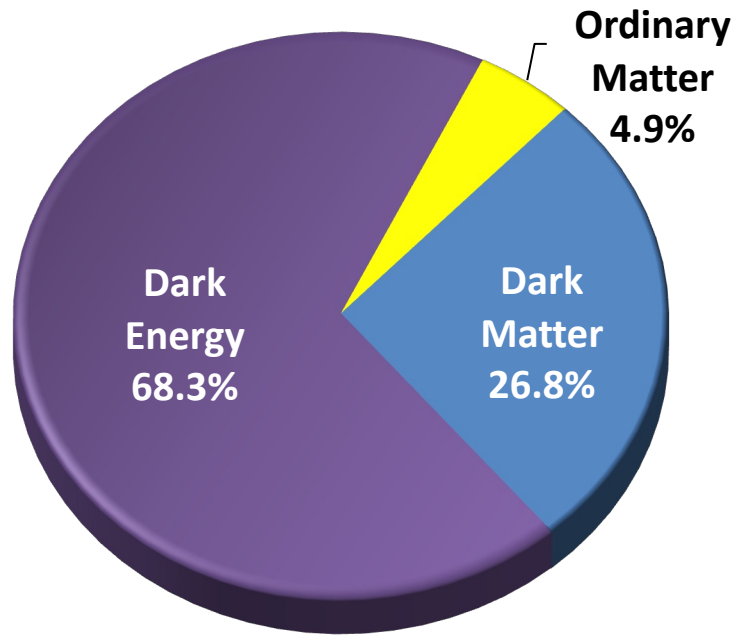


In collaboration with **Fei Huang** (Weizmann Institute, Israel)
and **Jiang-Hao Yu** (ITP, CAS, China)

The existence of dark matter is supported by observations across many different scales



What have we learnt?

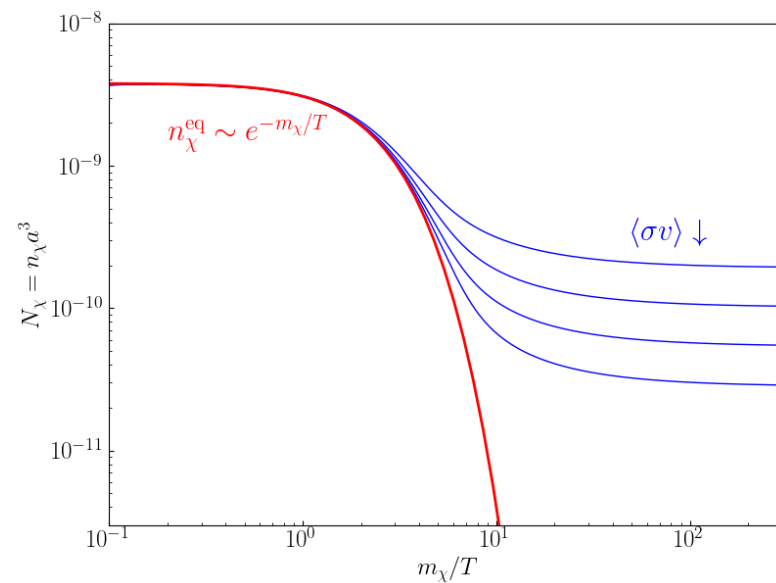


We still don't know:

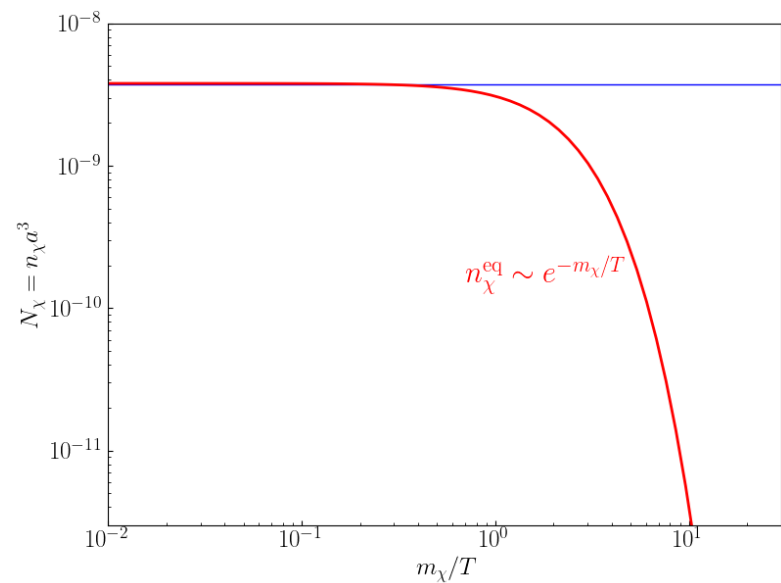
- **Abundance**: $\sim 26\%$ of the universe
- **Cold**: Non-relativistic, massive
- **Dark**: Negligible nongravitational interaction with Standard Model fields
- **Nonbaryonic**: Baryonic matter is simply not enough
- **BSM**: Not Standard Model particle

- **Particle properties**: mass, spin, fundamental or composite, single-component or multi-component, etc.
- **Production mechanism**: freeze-out, freeze-in, decays, misalignment?

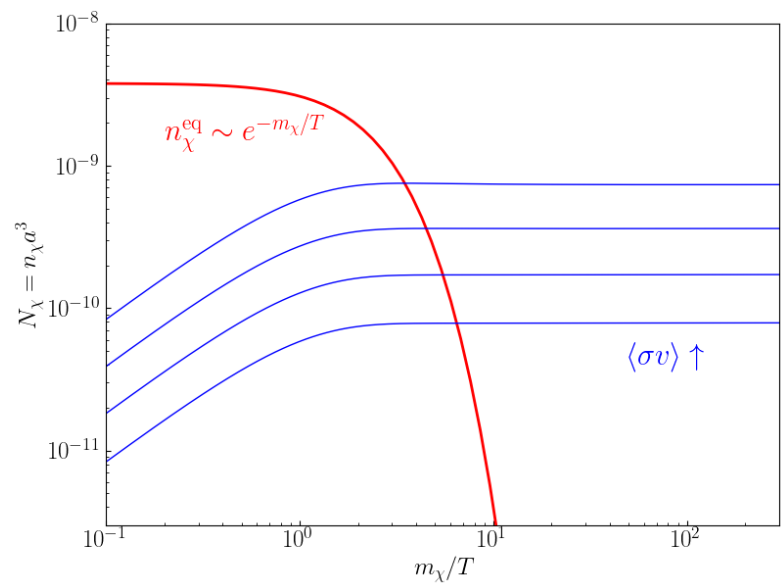
Possible Production mechanism of DM



Non-relativistic freeze-out
(WIMP)

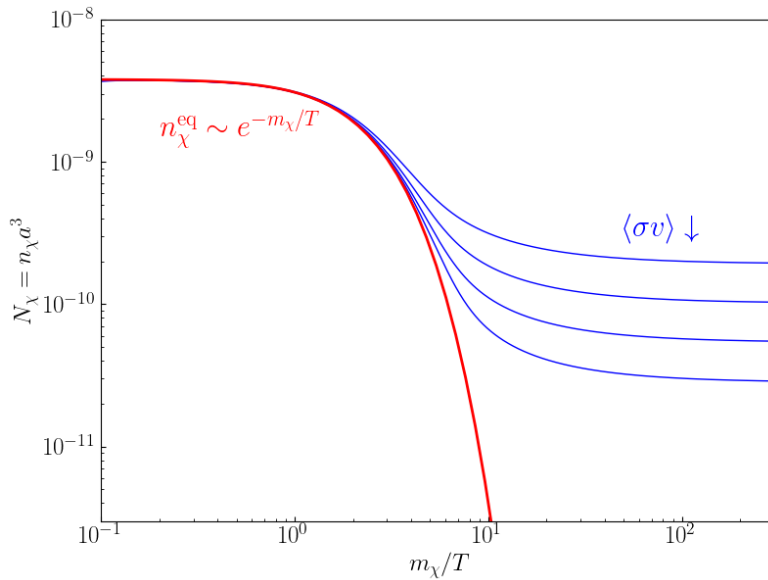


Relativistic freeze-out
(Warm Dark Matter)

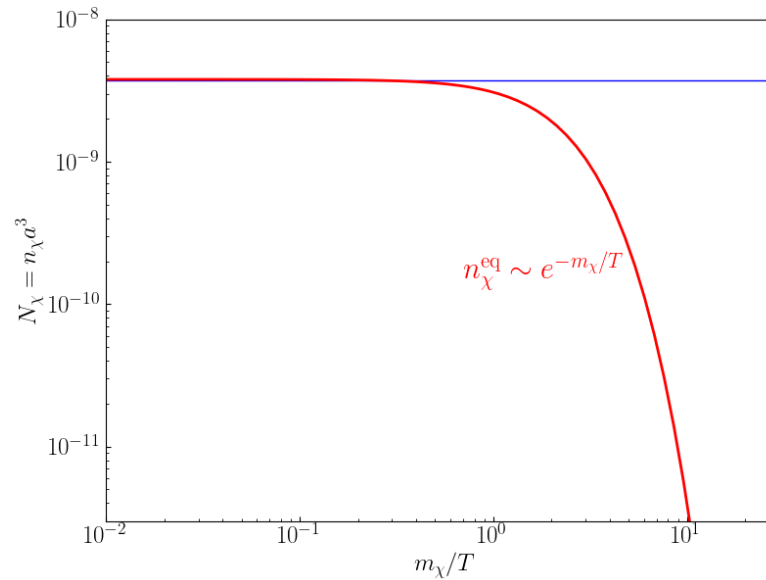


Freeze-in

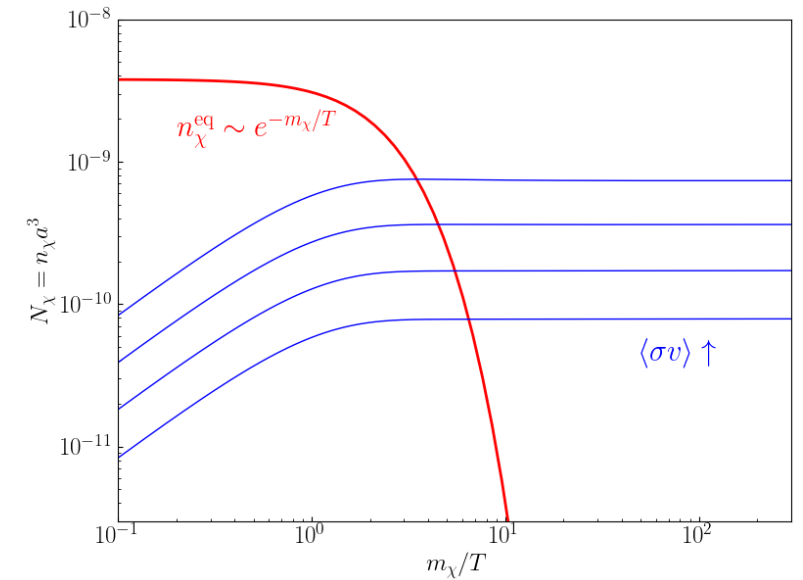
Possible Production mechanism of DM



Non-relativistic freeze-out
(WIMP)



Relativistic freeze-out
(Warm Dark Matter)



Freeze-in

Different thermal histories of DM result in different phase space distributions!

$$f_{\chi}(p) \approx \exp\left(-\frac{p^2}{2mT_{\chi}}\right)$$

$$f_{\chi}(p) \approx \exp\left(-\frac{p}{T_{\chi}}\right)$$

$$f_{\chi}(p) \approx C \frac{\exp(-p/T_{\chi})}{\sqrt{p/T_{\chi}}}$$

f_{χ} will then be encoded in the cosmic structures.

A little exercise with WDM

DM freezes-out relativistically: $f_\chi(p) \approx \exp(-p/T_\chi)$

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Energy density :	Average velocity:
$\rho_\chi \sim \int d^3p E f_\chi(p) \sim m_\chi T_\chi^3$	$\langle v \rangle \sim \frac{1}{m_\chi} \int d^3p E f_\chi(p) \sim \frac{T_\chi}{m_\chi}$

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well-measured

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Average velocity:

$$\langle v \rangle \sim \frac{1}{m_\chi} \int d^3p E f_\chi(p) \sim \frac{T_\chi}{m_\chi}$$

$$\langle v \rangle_0 \approx 1.1 \times 10^{-7} \times \left(\frac{2}{g_\chi} \right)^{1/3} \left(\frac{\Omega_\chi}{0.25} \right)^{1/3} \left(\frac{1 \text{ keV}}{m_\chi} \right)^{4/3}$$

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*Entropy conservation
+
cosmological redshift*

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*DM velocity is constrained
by structure formation!*

$$\begin{aligned} m_{\text{WDM}} &\geq 3.5 \text{ keV} \\ &\geq 5.3 \text{ keV} \end{aligned}$$

V. Iršič et al.
Phys. Rev. D 96 (2017) 023522

Constraints on decoupling temperatures

$$\begin{aligned} m_{\text{WDM}} &\geq 3.5 \text{ keV} \\ &\geq 5.3 \text{ keV} \end{aligned}$$



$$\begin{aligned} \langle v \rangle_0 &\lesssim 2.1 \times 10^{-8} \\ &\lesssim 1.2 \times 10^{-8} \end{aligned}$$

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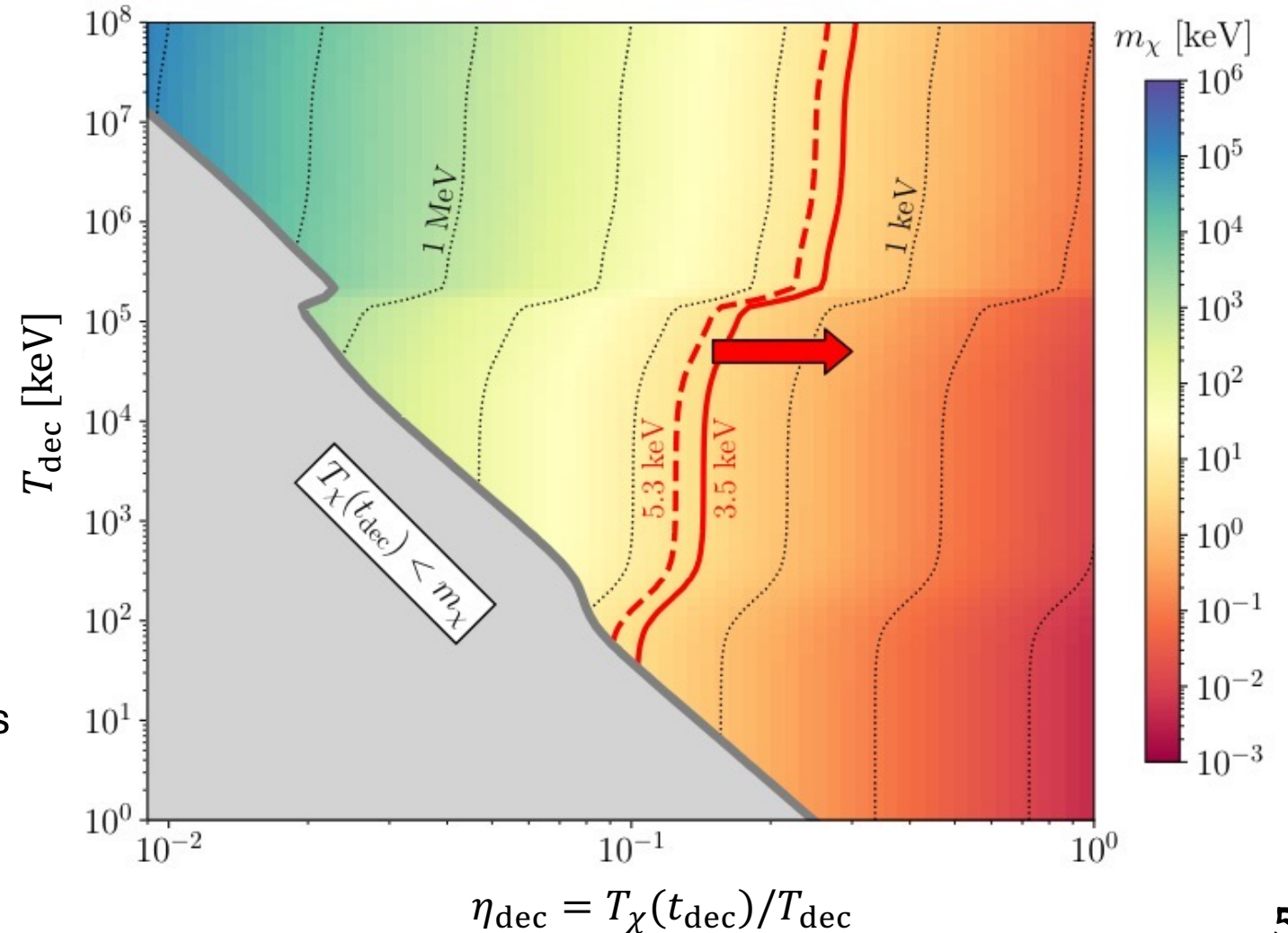
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$$\langle v \rangle_0 \lesssim 2.1 \times 10^{-8} \\ \lesssim 1.2 \times 10^{-8}$$

$$m_{\chi} \approx 1.9 \times 10^{-3} \text{ keV} \times \frac{\Omega_{\chi}}{0.25} \frac{g_{*,s}(T_{\text{dec}})}{g_{\chi}} \eta_{\text{dec}}^{-3}$$

❖ Mass contours are also velocity contours



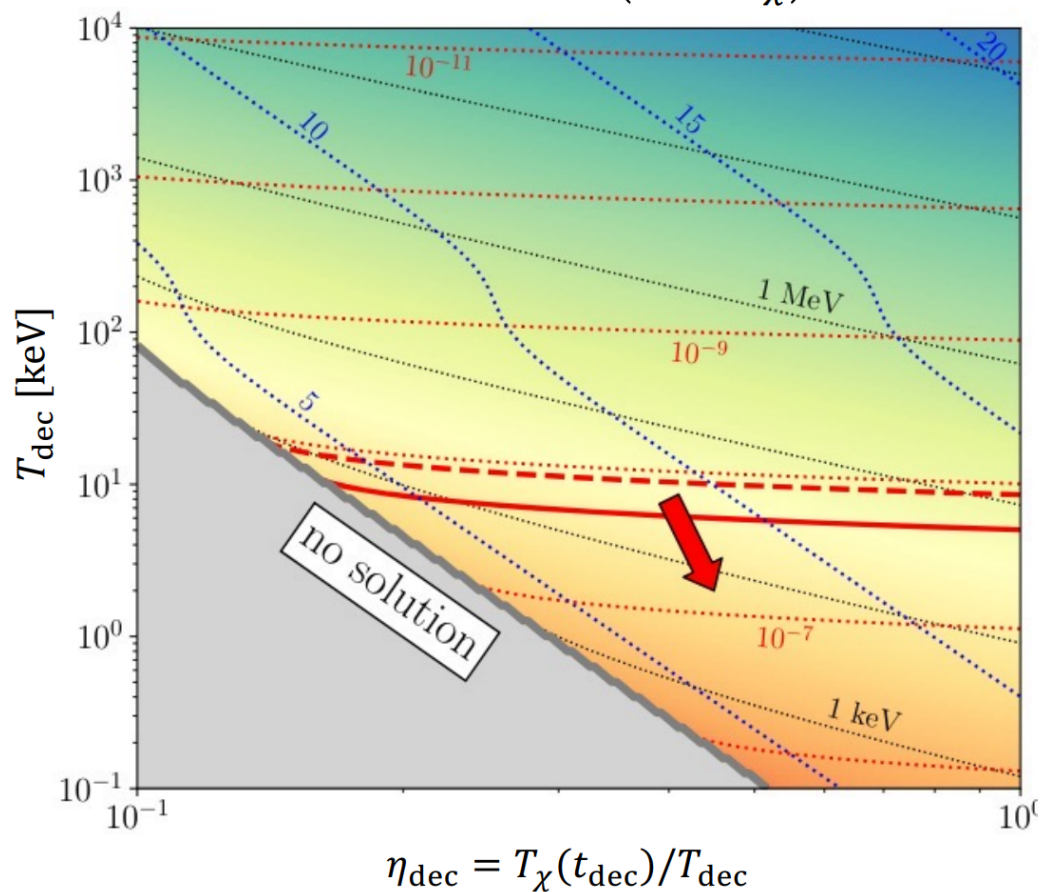
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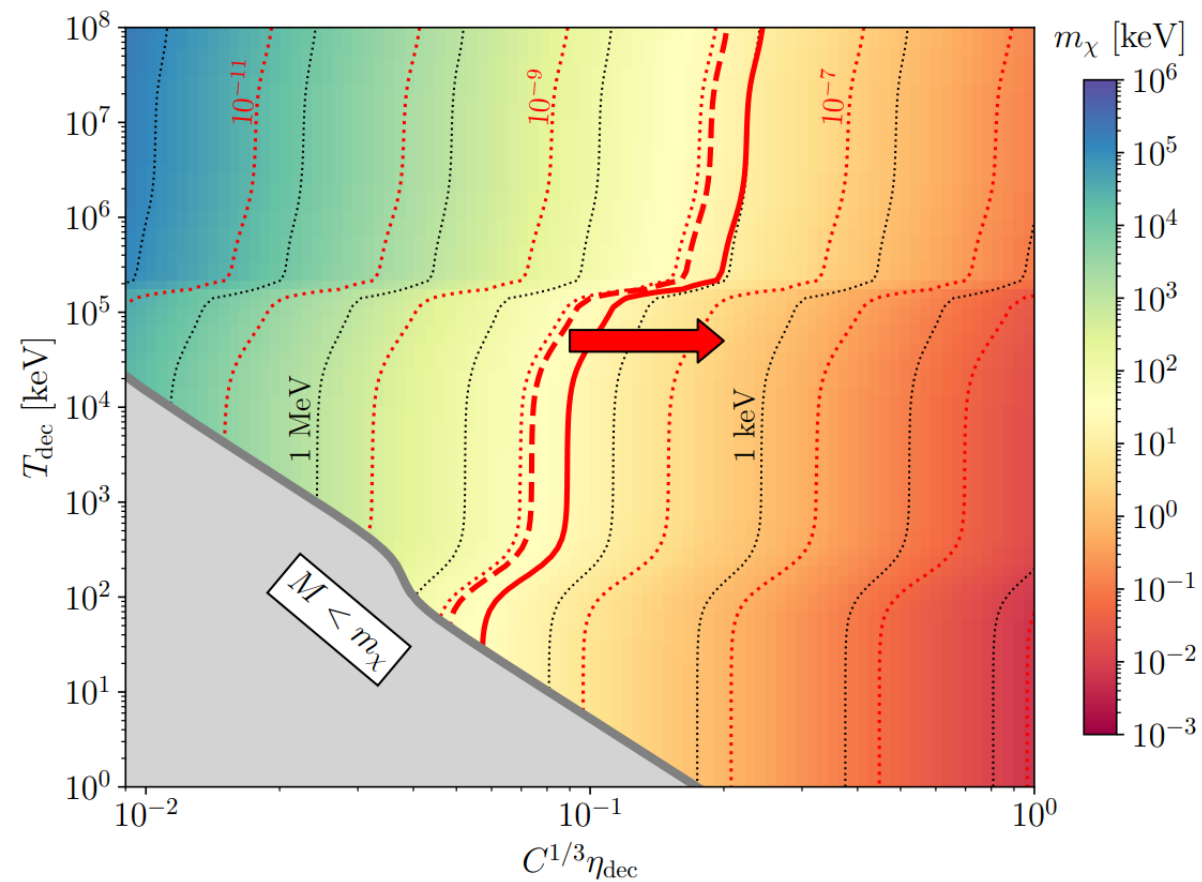
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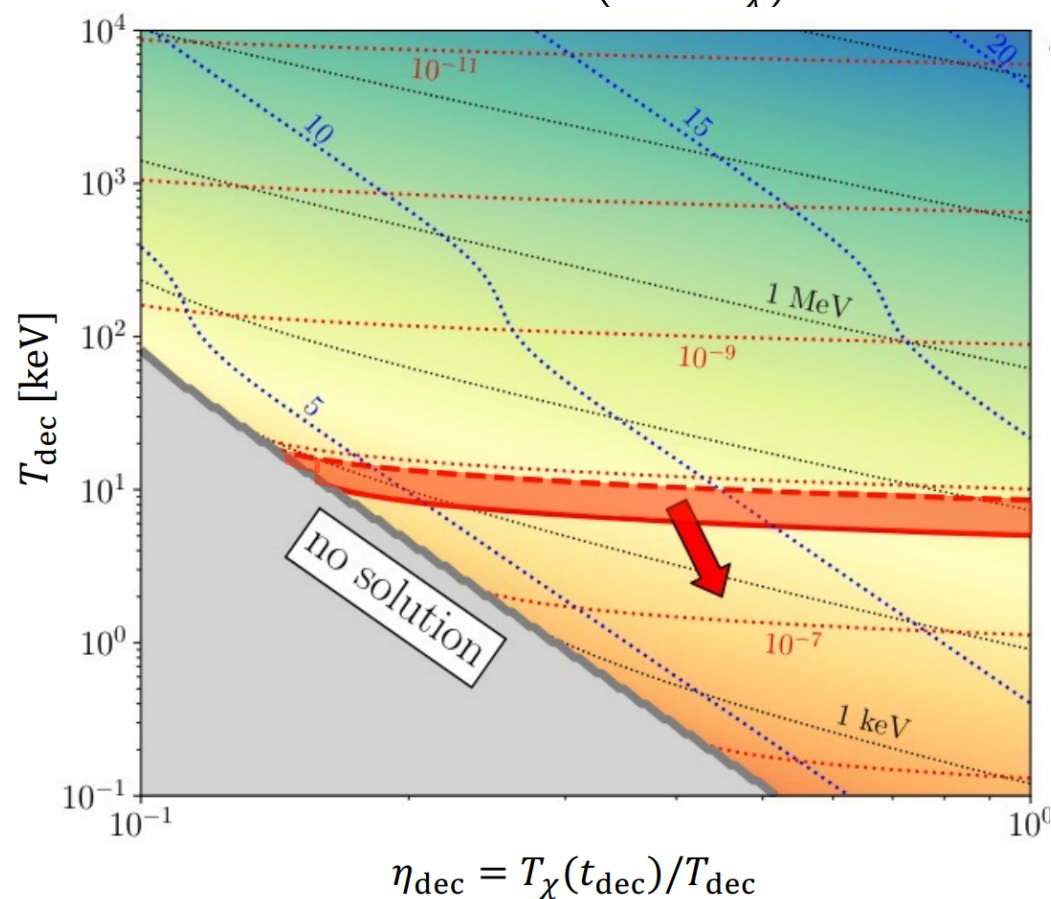
F. D'Eramo & A. Lenoci
JCAP 10 (2021) 045



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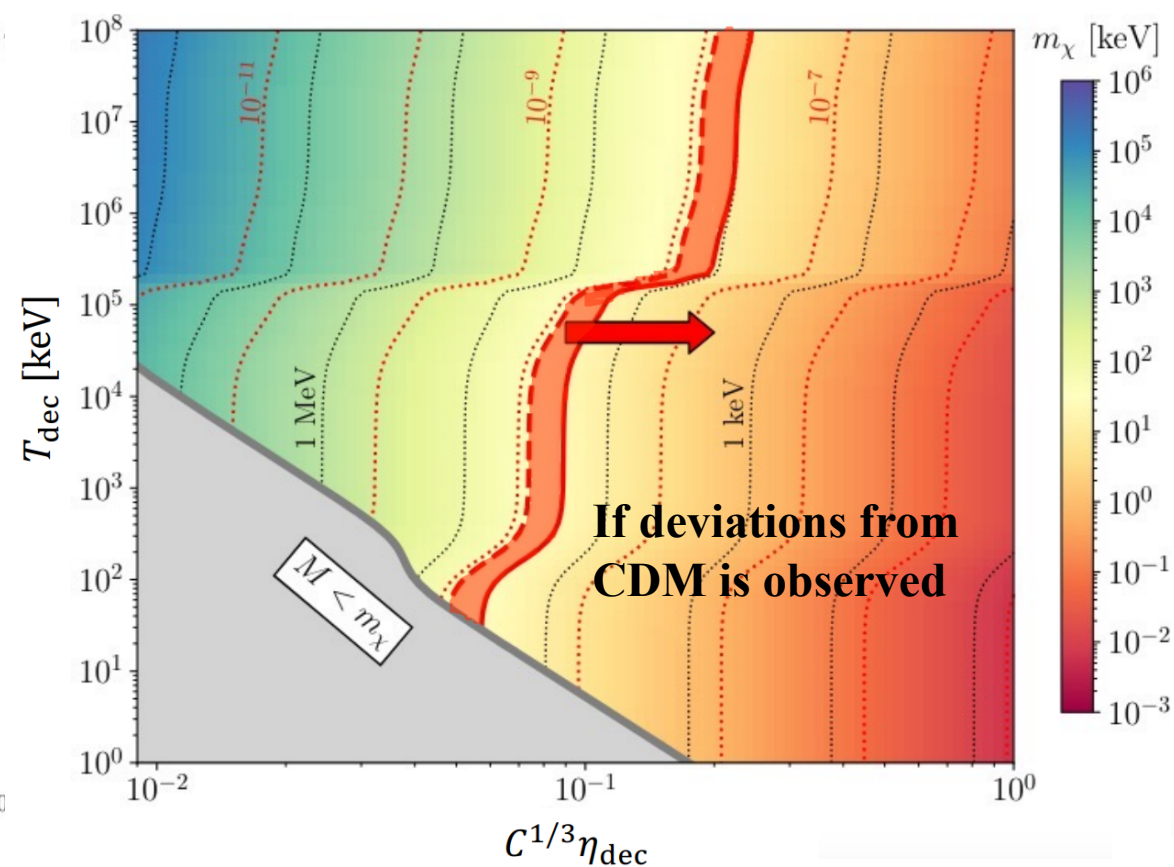
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Beyond the average velocity

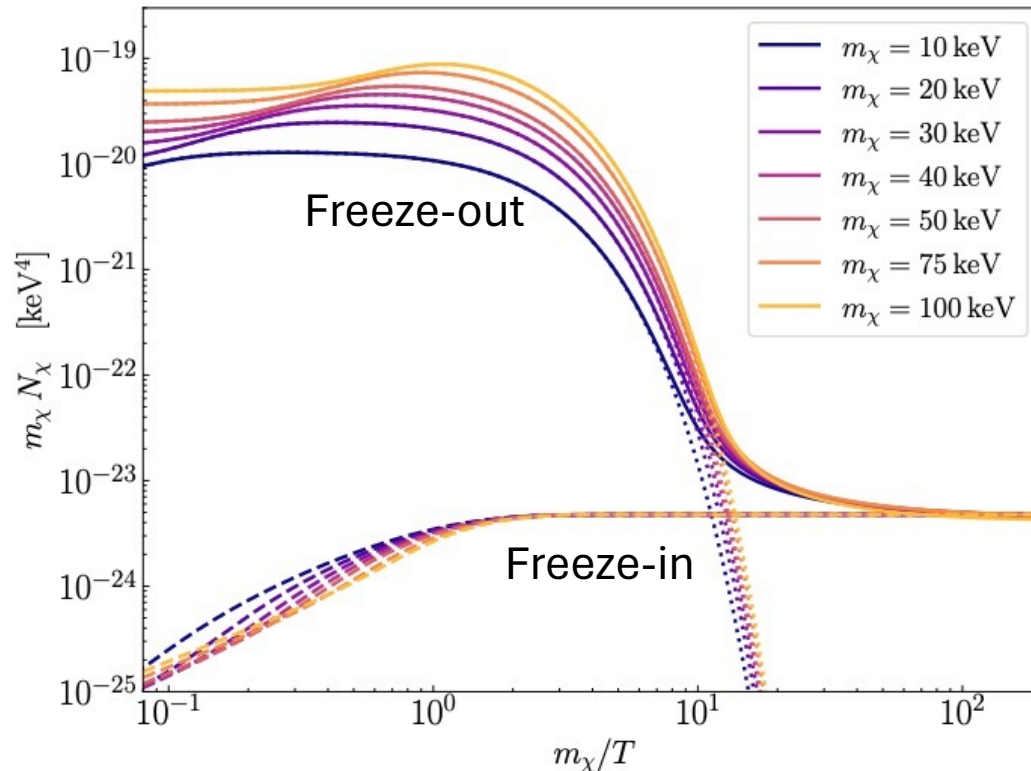
Solve Boltzmann equations numerically
to obtain the distribution function
(assuming $\eta_{\text{dec}} = 1$)

$$\left(\frac{\partial}{\partial t} - Hp \frac{\partial}{\partial p} \right) f_{\chi}(p, t) = \mathcal{C}[f]$$

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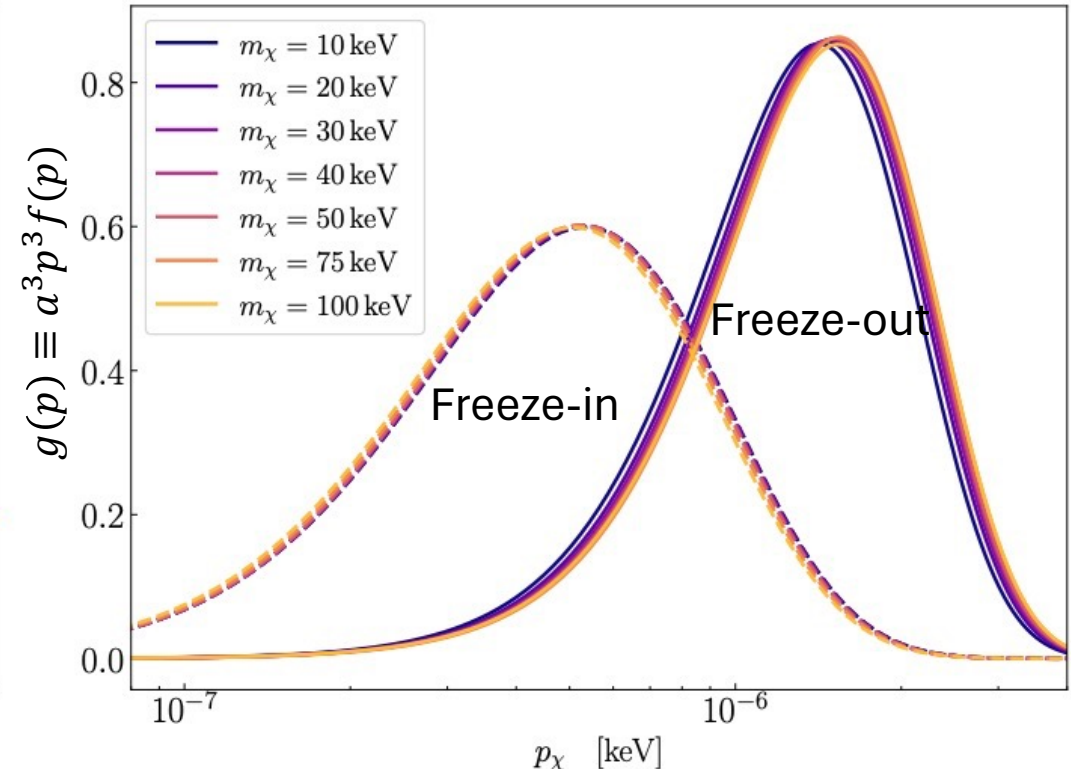
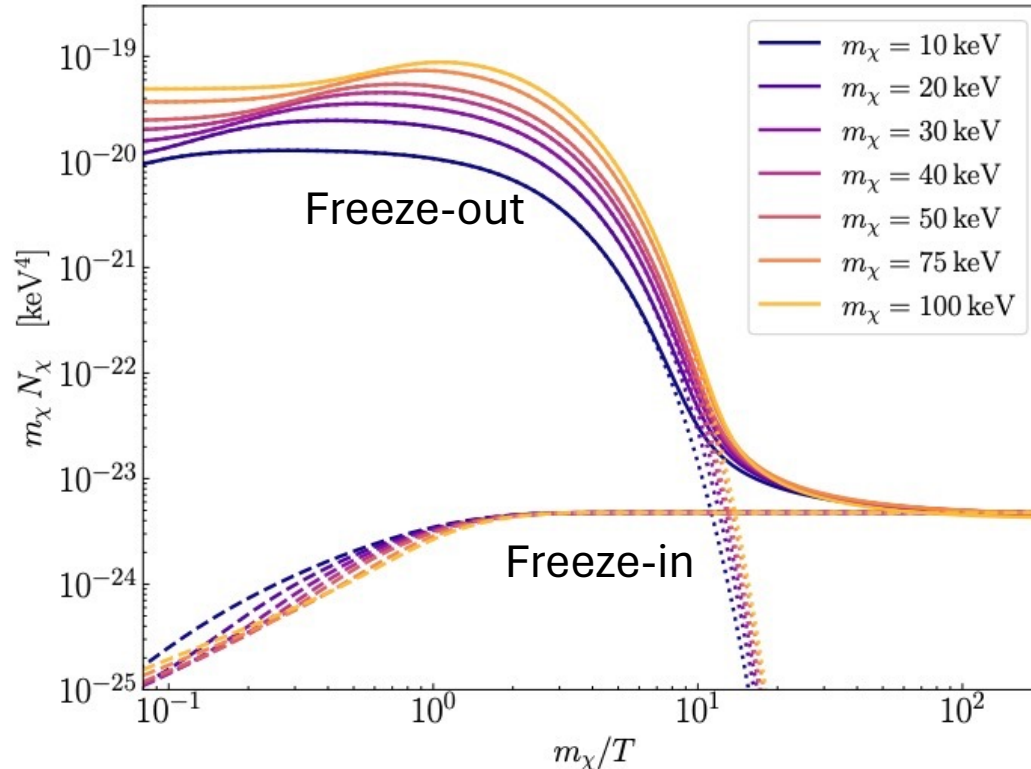
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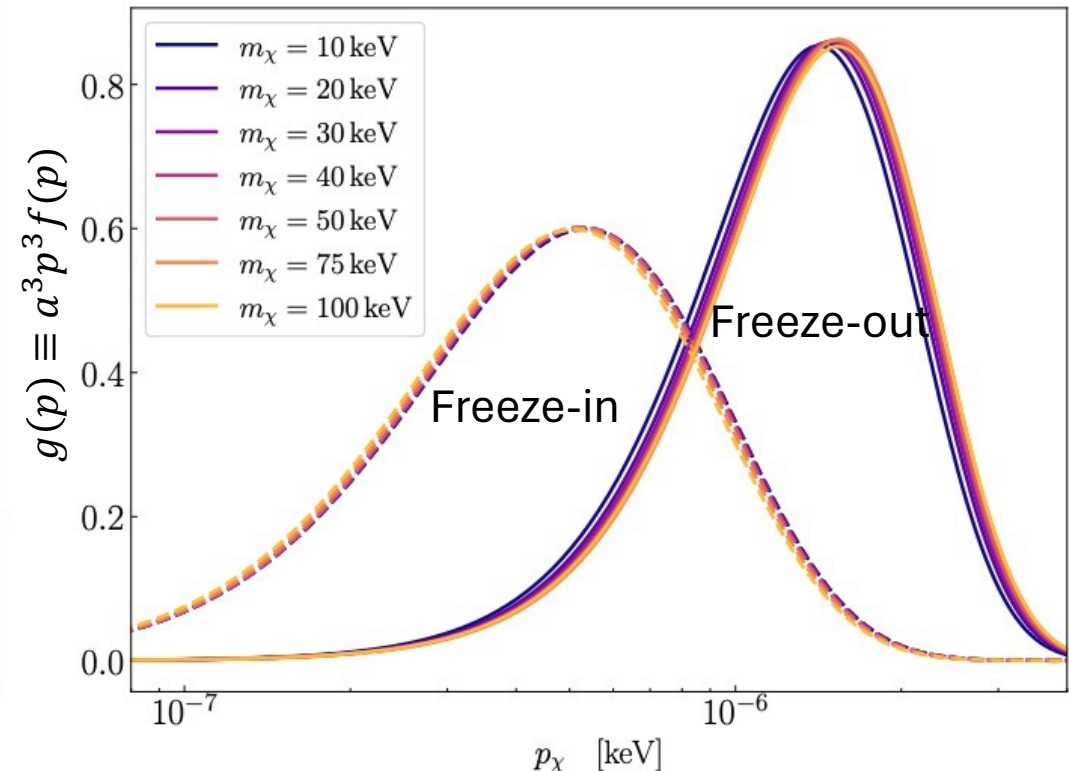
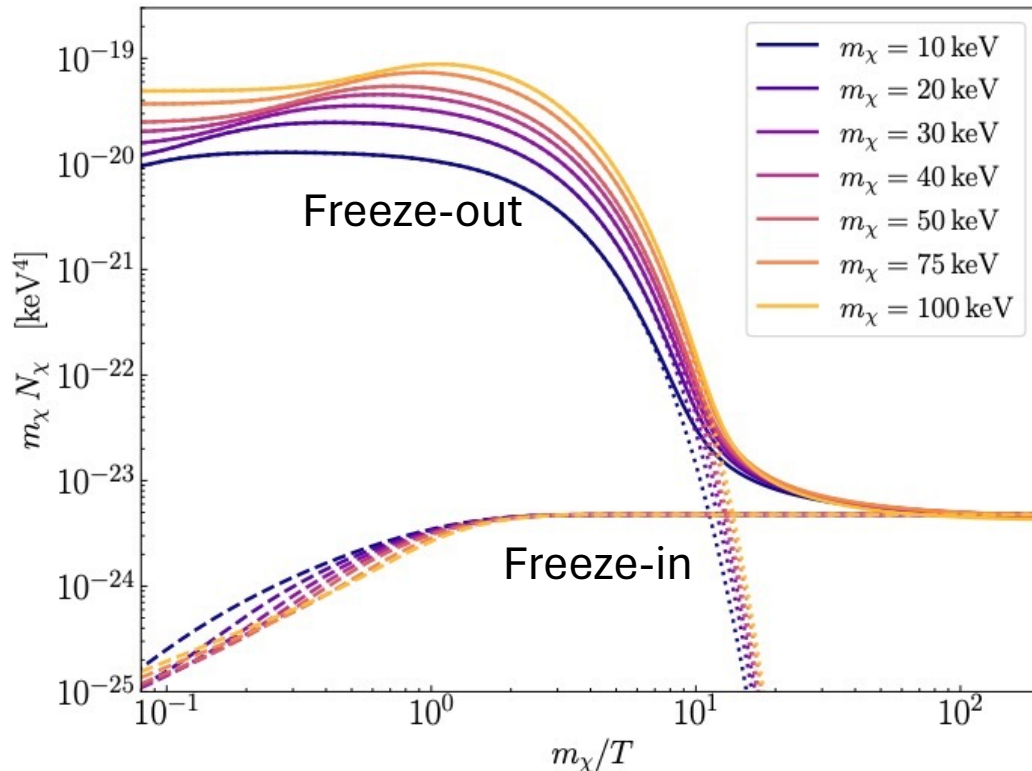


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- For the **same** production mechanism \rightarrow ***similar*** distributions
 - **Larger mass** \rightarrow **Smaller overall velocity**
- Distributions from freeze-in and freeze-out are **distinct**
 - **same mass** \rightarrow freeze-in distribution is **colder**
 - Even if same average velocity \rightarrow **different shape**

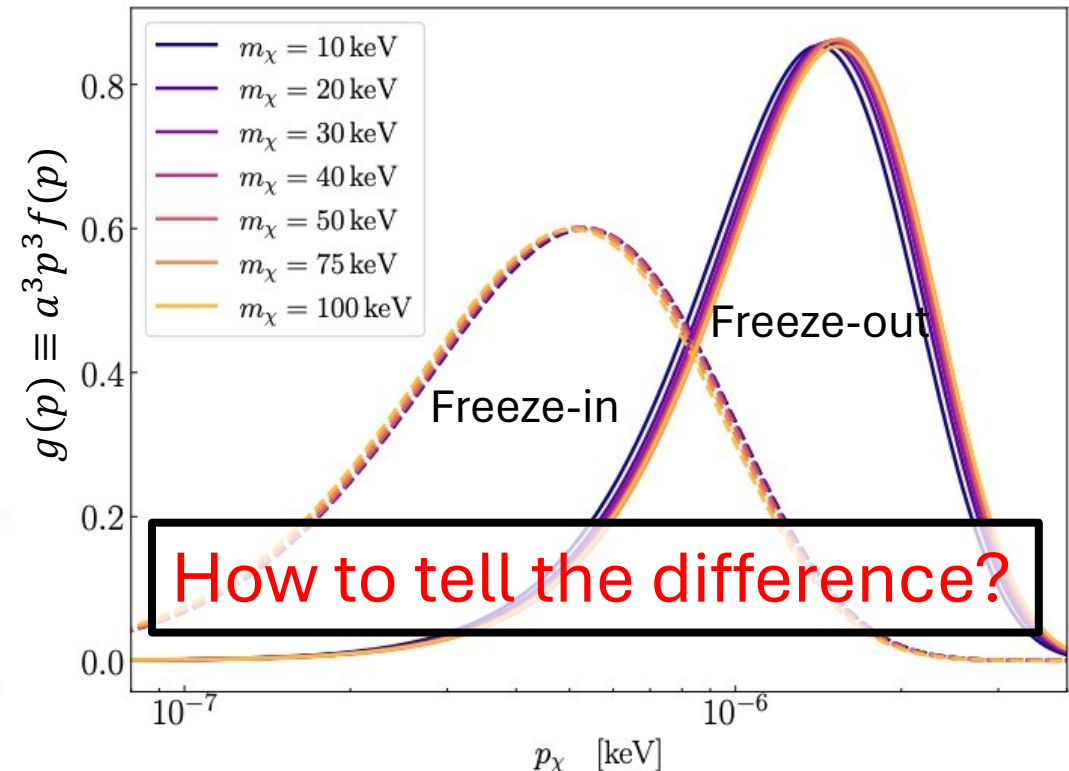
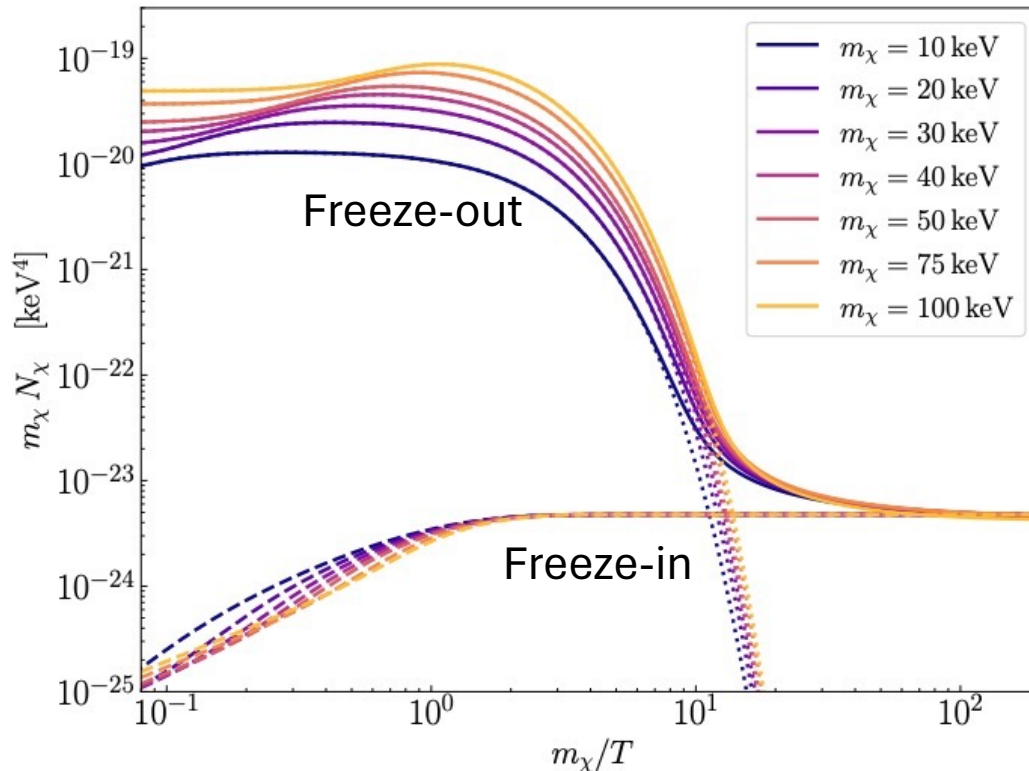


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Imprints on cosmic structure

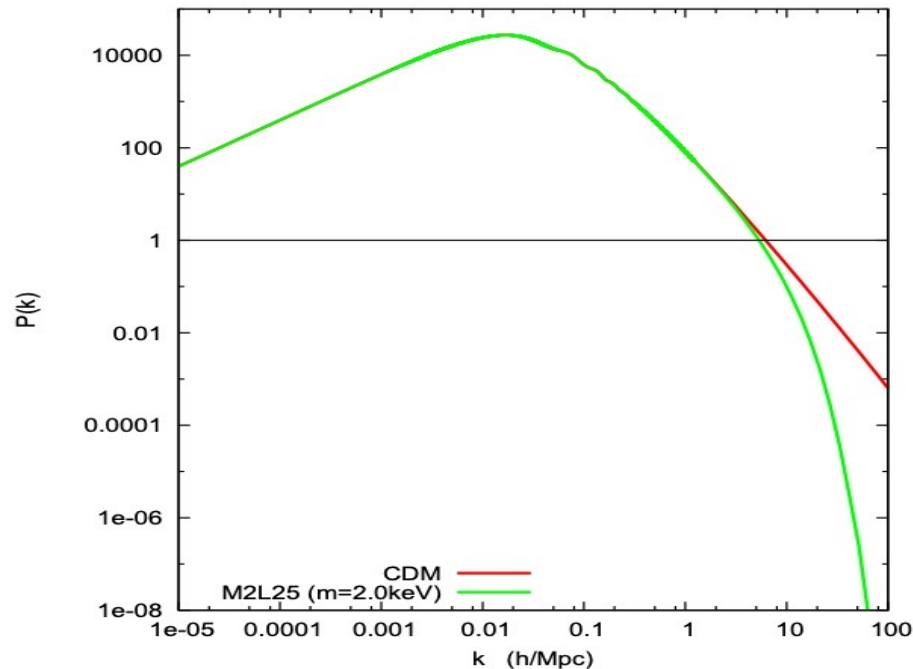
- Non-negligible velocities suppresses structure formation, reflected in the **matter power spectrum**

$$P(k)$$

- Often represented by the **squared transfer function**

$$T^2(k) \equiv P(k)/P_{\text{CDM}}(k)$$

- The shape of $T^2(k)$ **contains the information** of the distribution function



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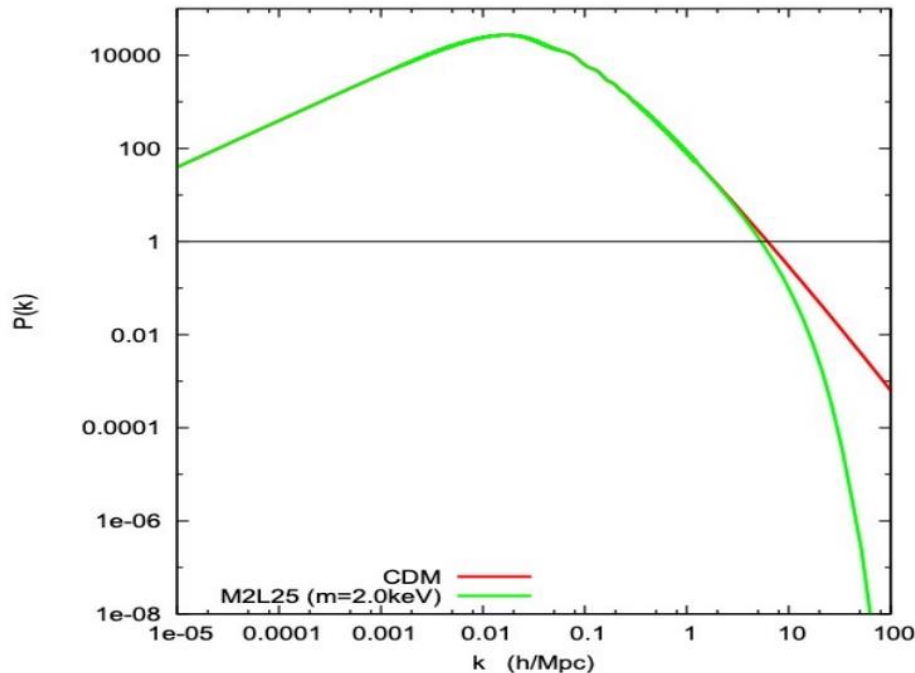
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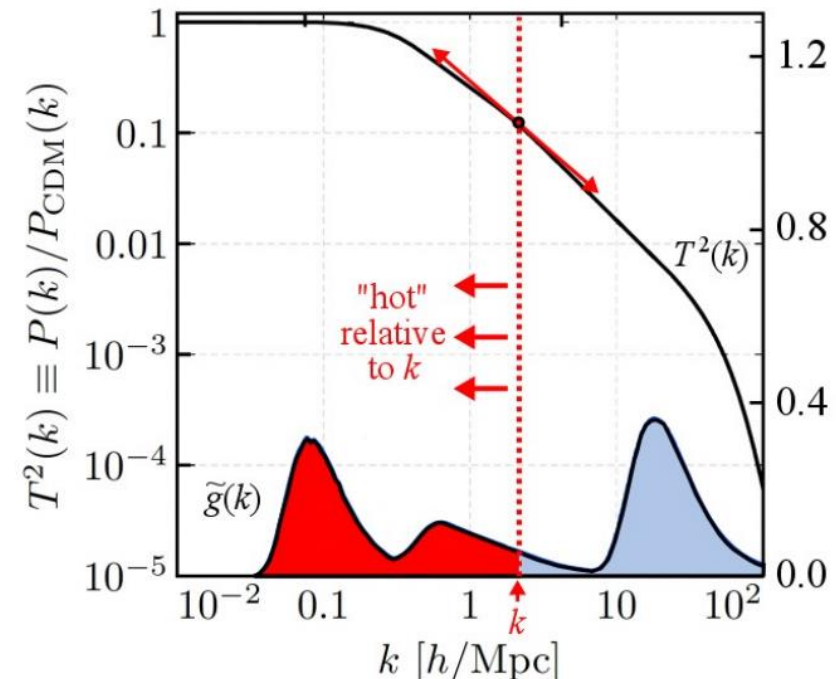
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J. Lesgourgues and T. Tram
JCAP 09 (2011) 032



K. Dienes, F. Huang, J. Kost, S. Su, B. Thomas
Phys.Rev.D 101 (2020) 12

Imprints on cosmic structure

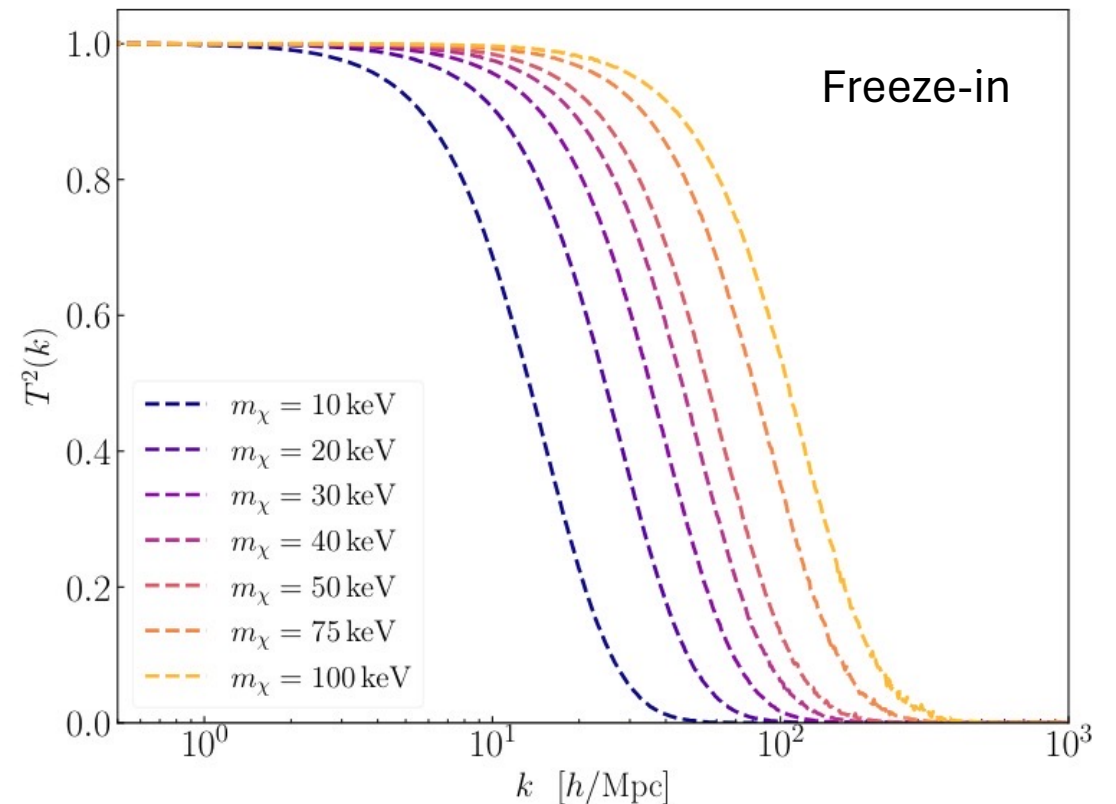
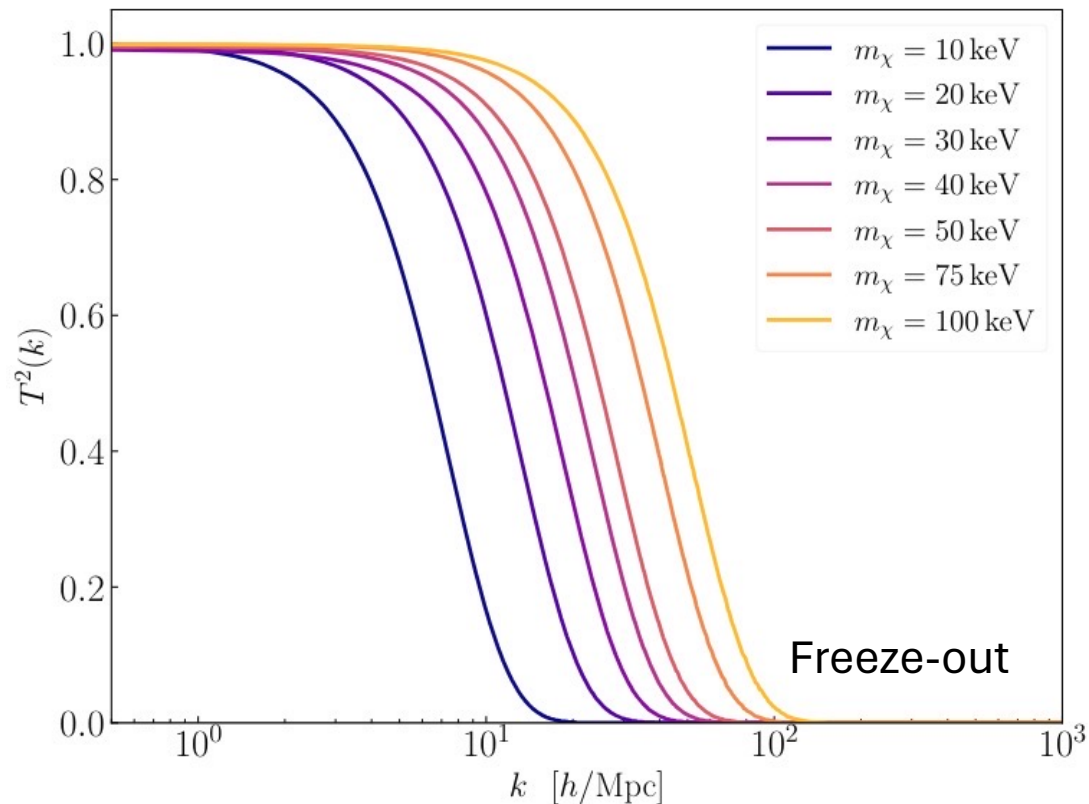
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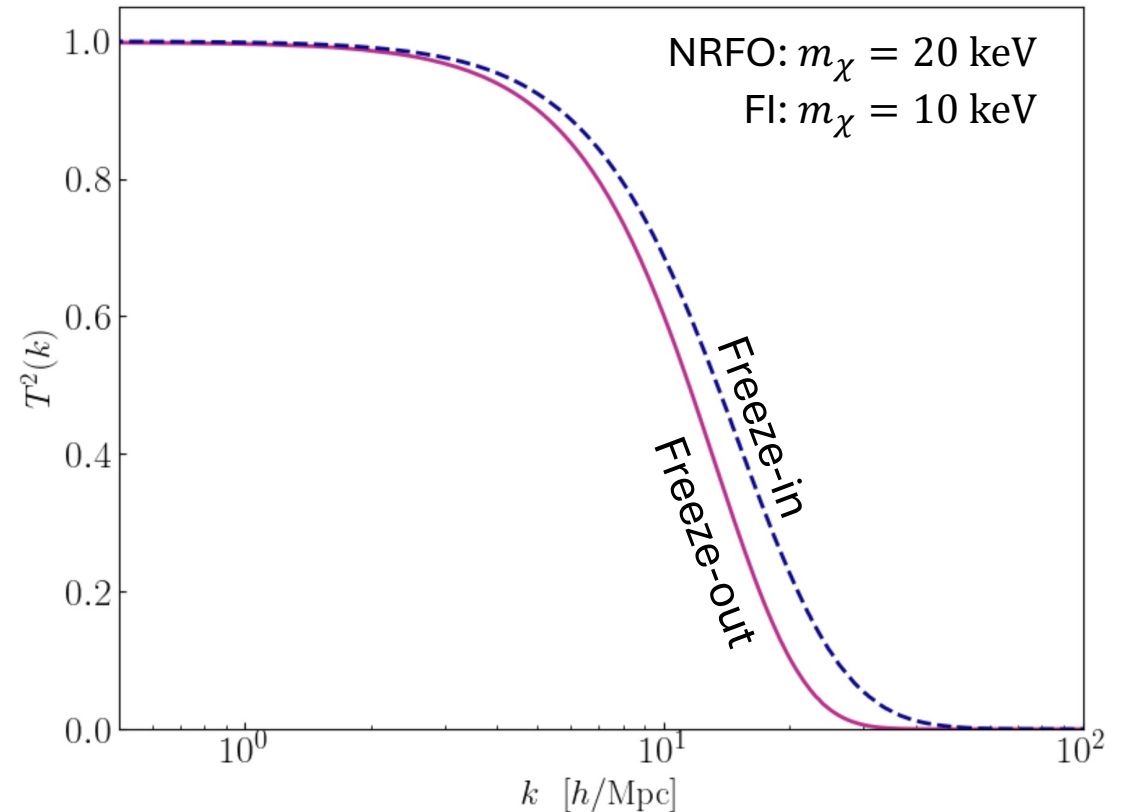
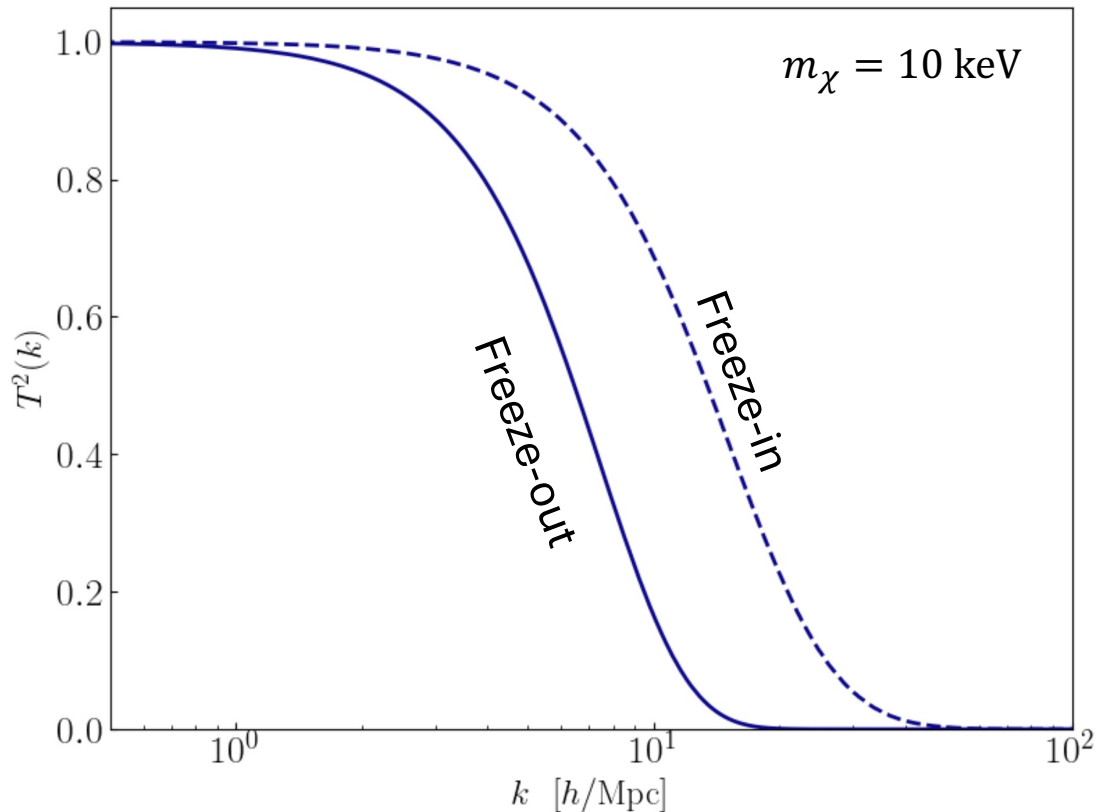
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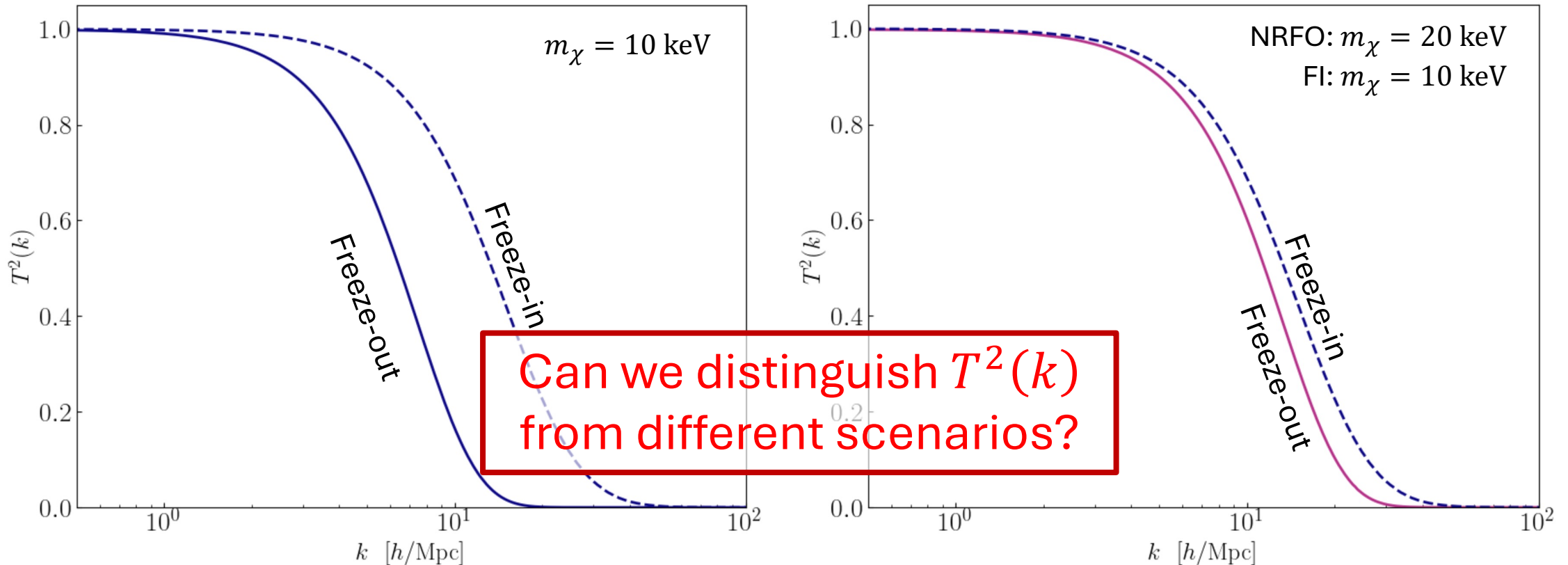
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Distinguishing different scenarios

- If future measurements on $P(k)$ finds deviation from CDM, use WDM as a baseline model

$$P_{\text{WDM}}(k)[1 + \sigma_P^+(k)] > P(k) > P_{\text{WDM}}(k)[1 - \sigma_P^-(k)]$$

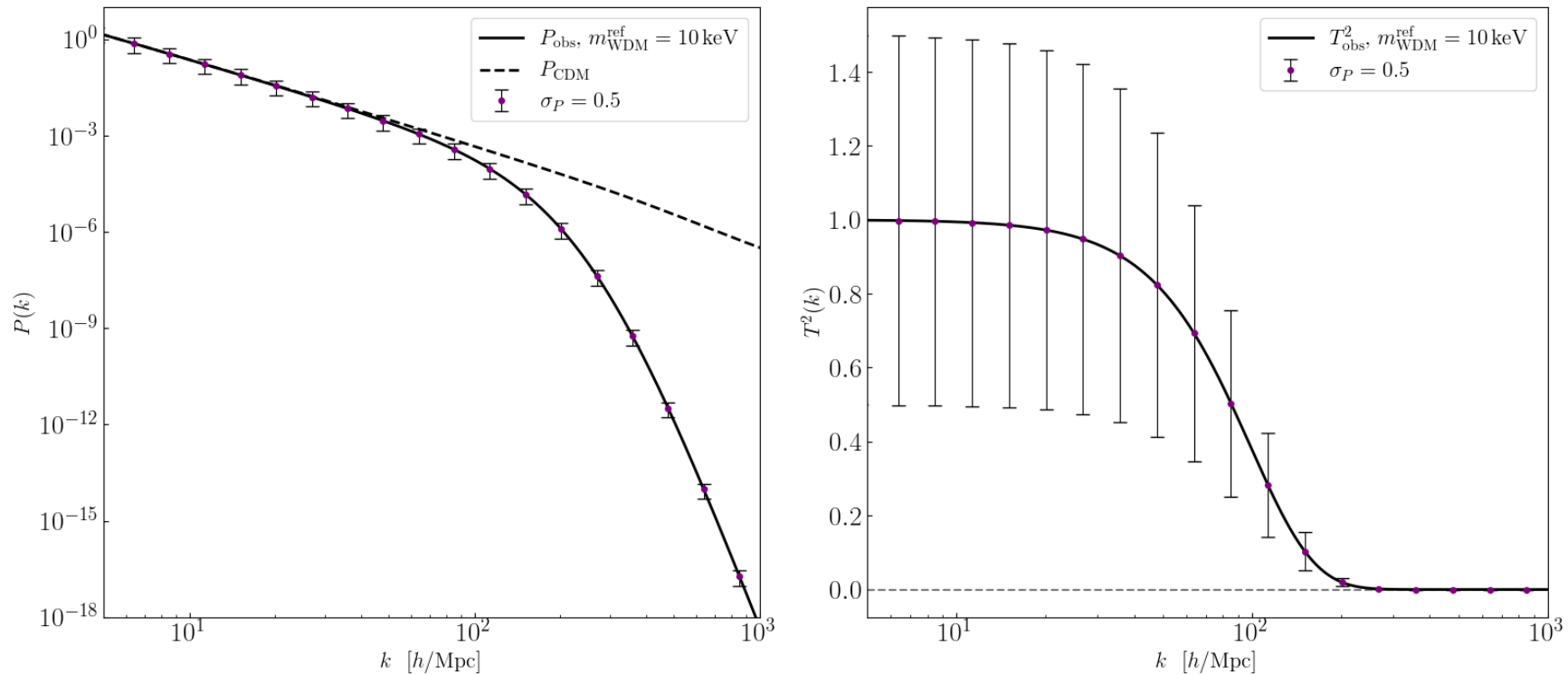
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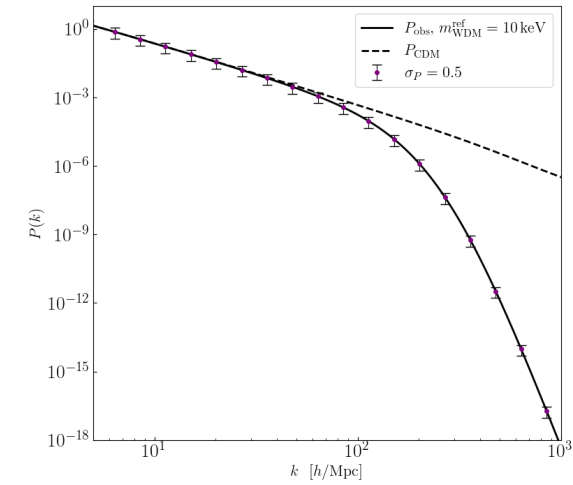
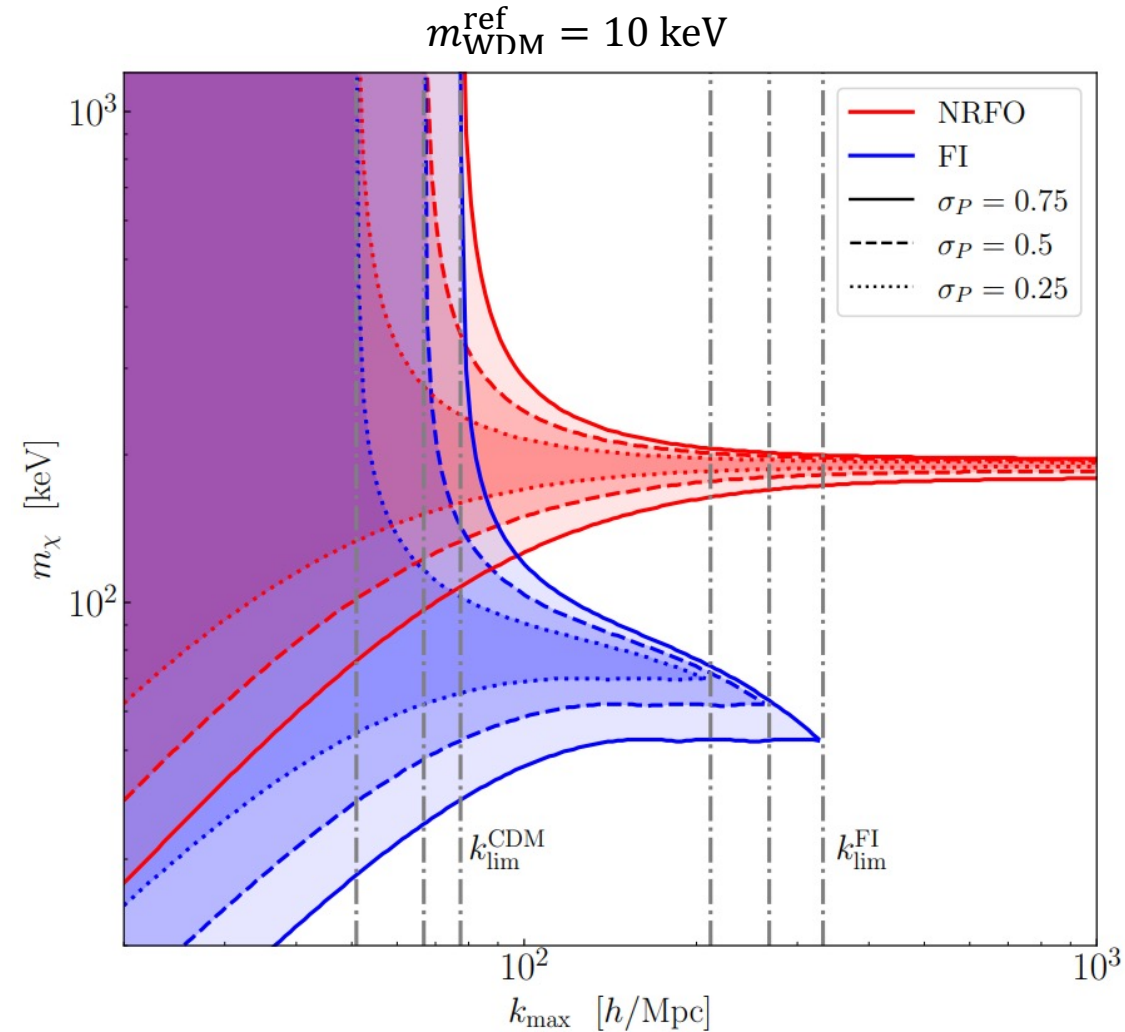
- Explore different possibilities: e.g., **constant symmetric relative errors** on $P(k)$

$$\sigma_P^\pm(k) = \sigma_P$$



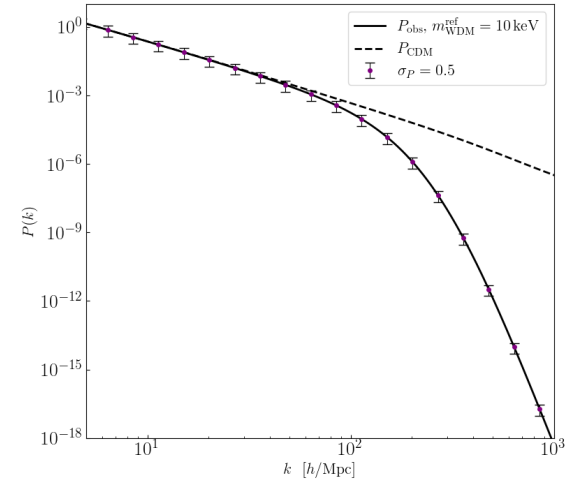
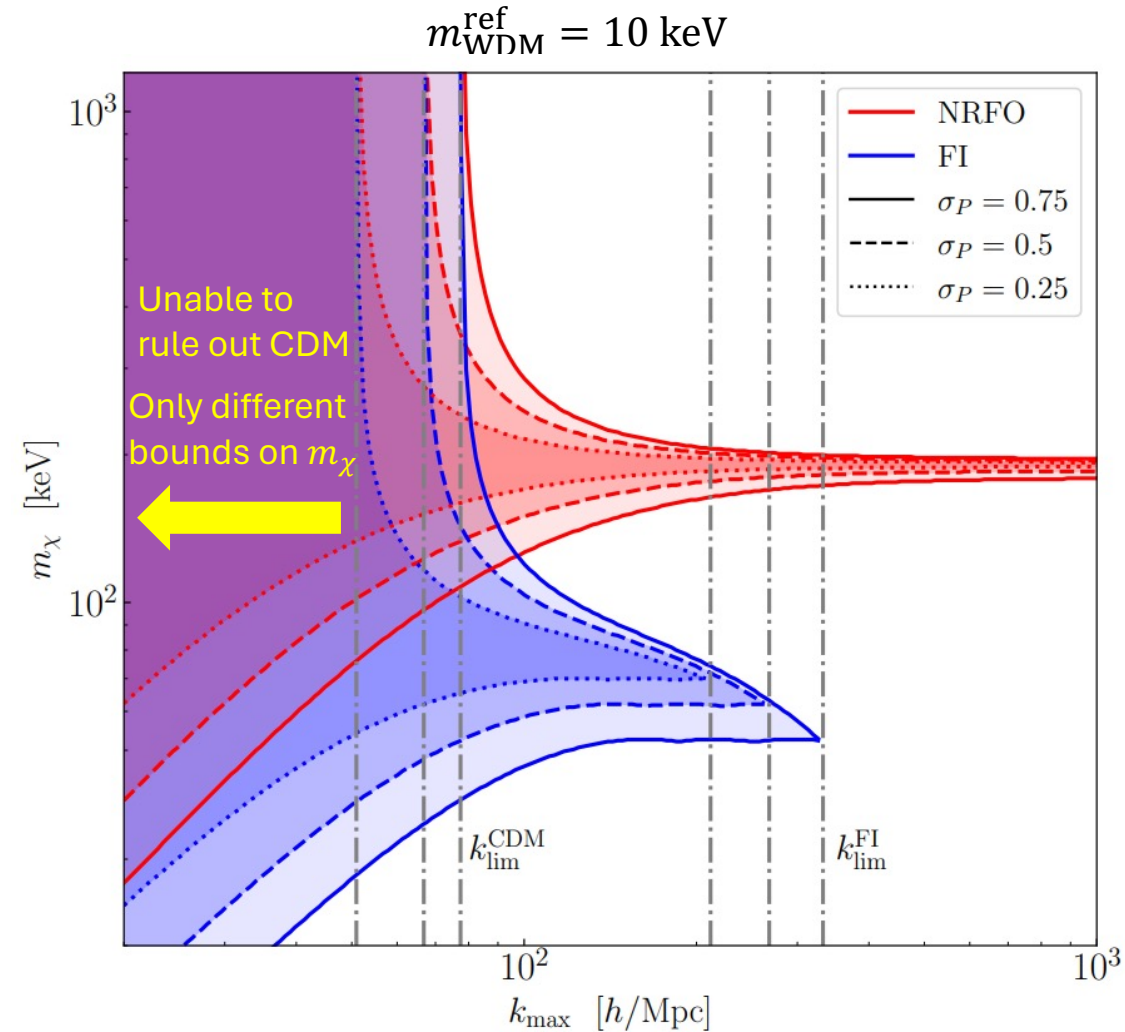
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- Assuming $N = 20$ data points evenly distributed on log-scale from $k_{\min} = 1 \text{ h/Mpc}$ to k_{\max}



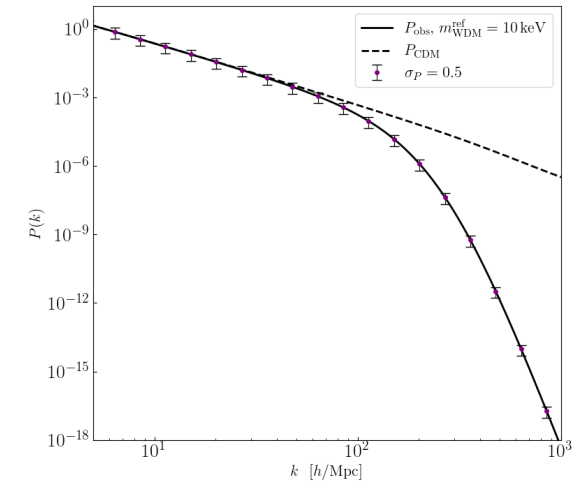
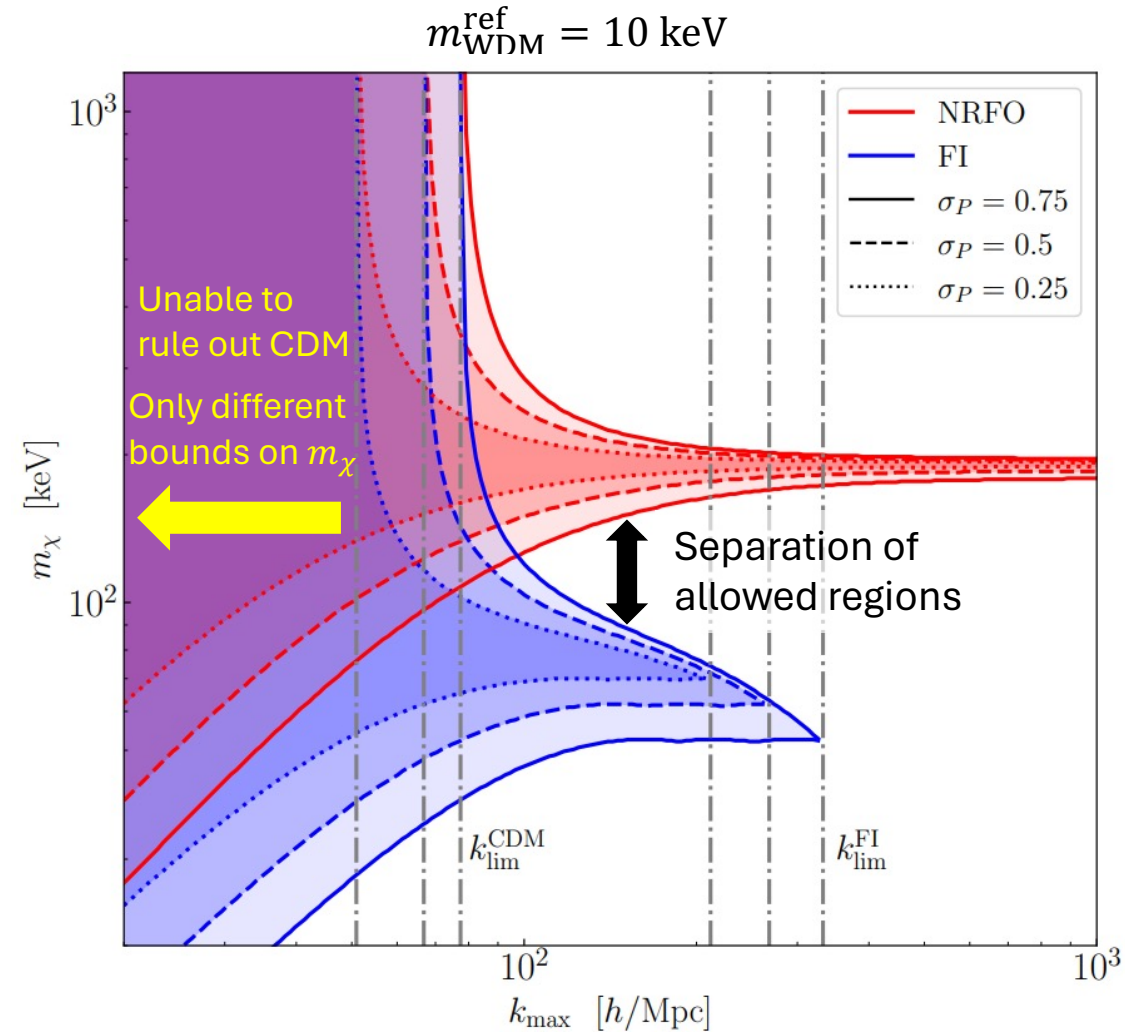
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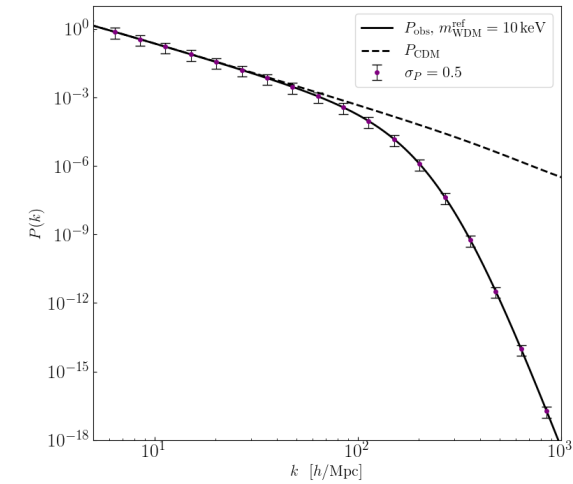
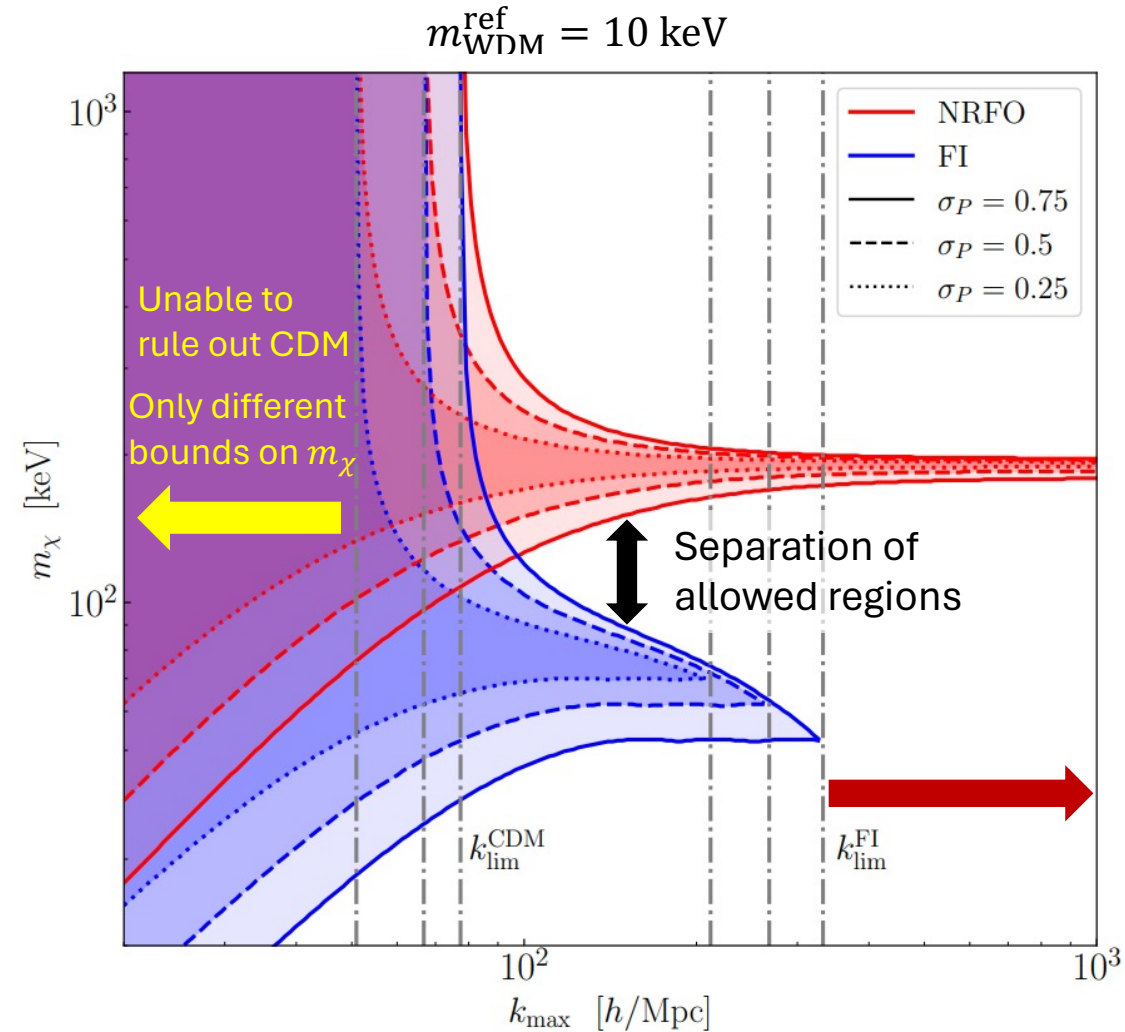
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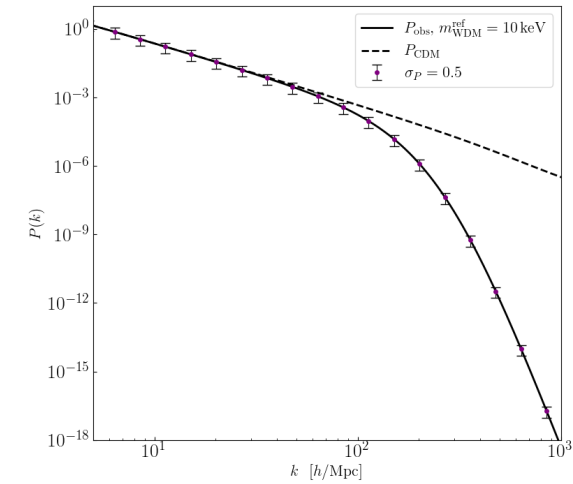
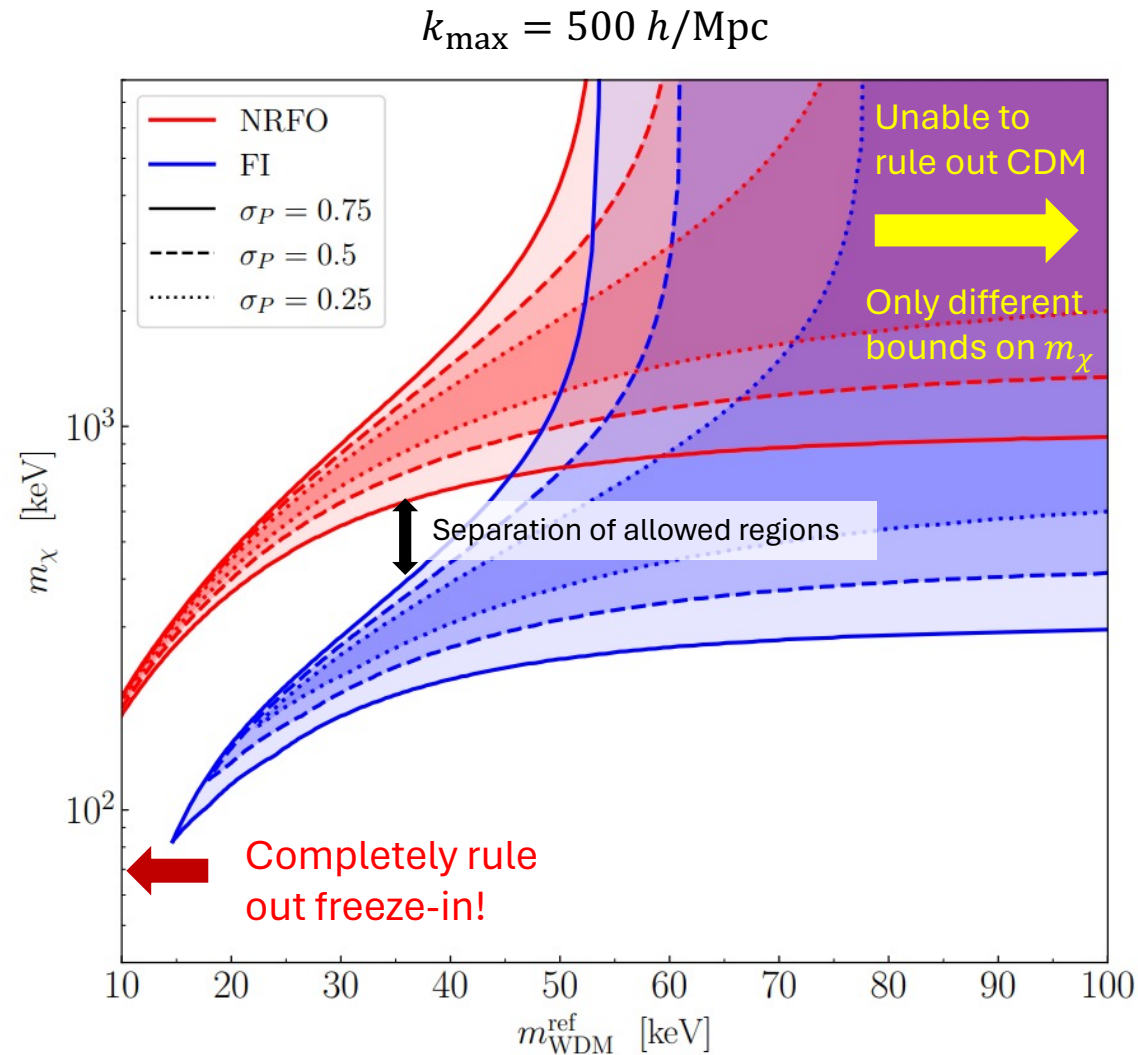
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Conclusion

- The cosmic structure contains information about the thermal history of dark matter
- Current WDM bounds can be re-interpreted to place constraints on decoupling temperatures (physical quantities relevant for the early universe) in different scenarios
- For different scenarios, these constraints imply different bounds on DM mass for different scenarios.
- If future data observes deviation from CDM predictions, it is potentially possible to discriminate different production mechanisms if there is sufficient precision

