

Oliver Buchmueller, Imperial College London

DIRECT SEARCHES FOR SUPERSYMMETRY – WHERE ARE WE TODAY?

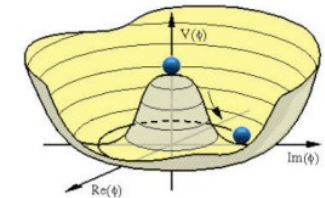
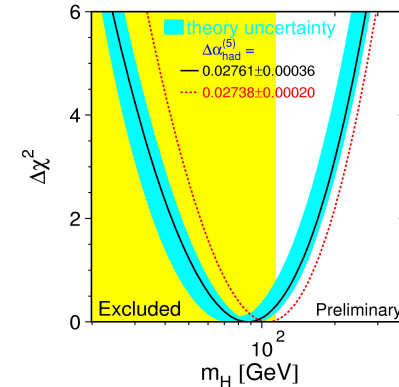


Fundamental Open Questions in Particle Physics

Direct SUSY Searches, O. Buchmüller

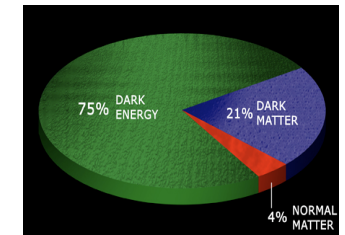
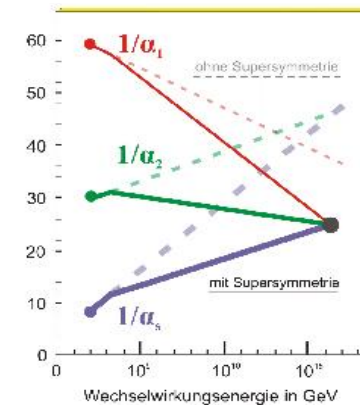
I. What is the origin of mass?

- Why are the vector bosons Z and W massive whereas the photon is massless?
- Is there a Higgs boson - or even more of them ?



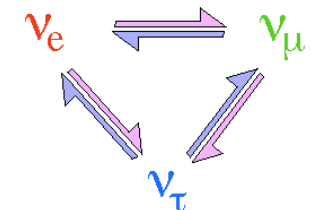
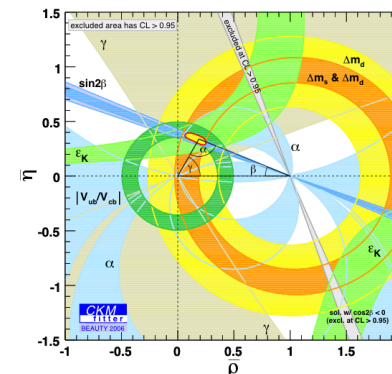
I. Is there a new symmetry - Supersymmetry ?

- Can we get experimental evidence to support the Grand Unification of all fundamental forces?
- What is the origin of Dark Matter in the Universe?
→ Is a fundamental particle responsible for it?



III. What is the origin of the matter-anti-matter asymmetry in our Universe?

- Does the answer lie in CP violation?
- Neutrino masses and mixing - how do they fit in the picture?

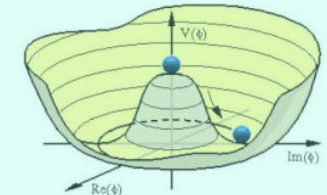
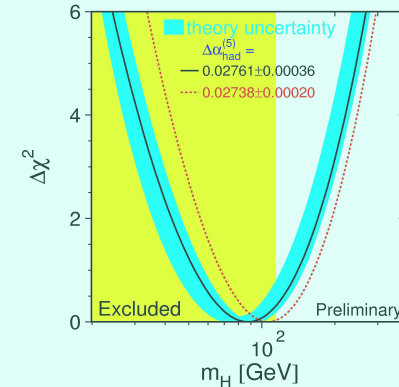


Fundamental Open Questions in Particle Physics

Direct SUSY Searches, O. Buchmüller

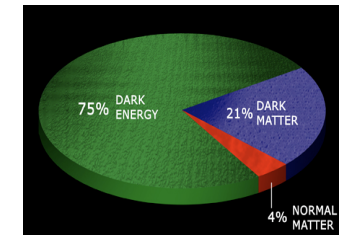
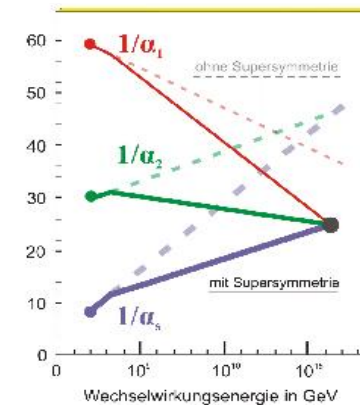
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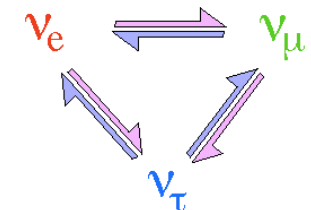
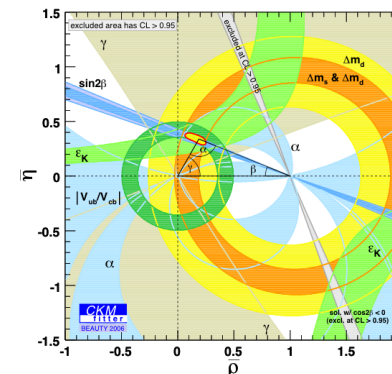
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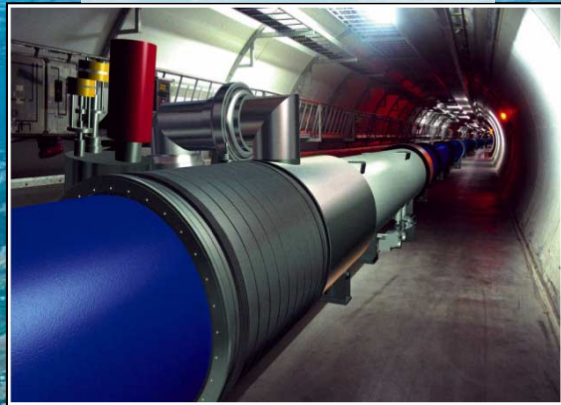
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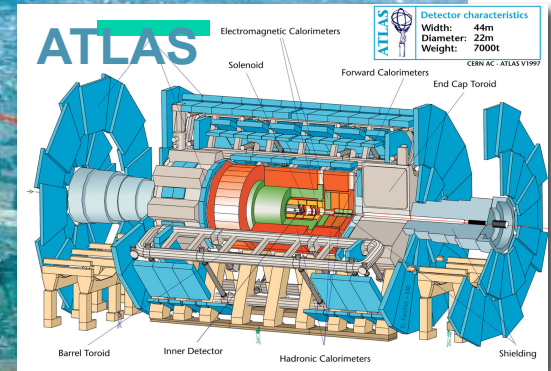
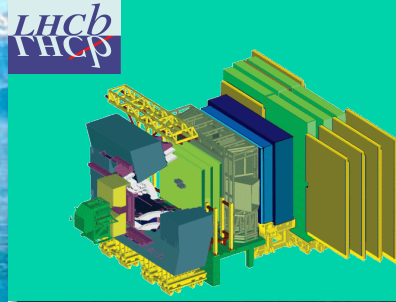
The Large Hadron Collider at CERN

Direct SUSY Searches, O. Buchmüller

LHC : 27 km long
100m underground

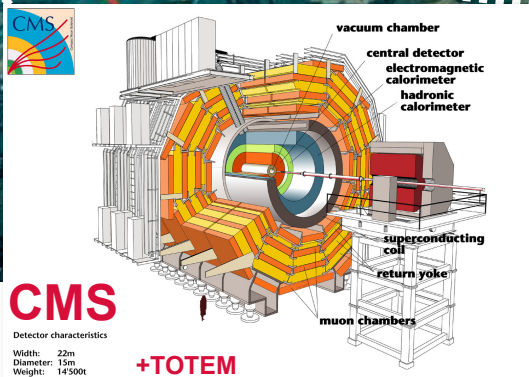


pp, B-Physics,
CP Violation

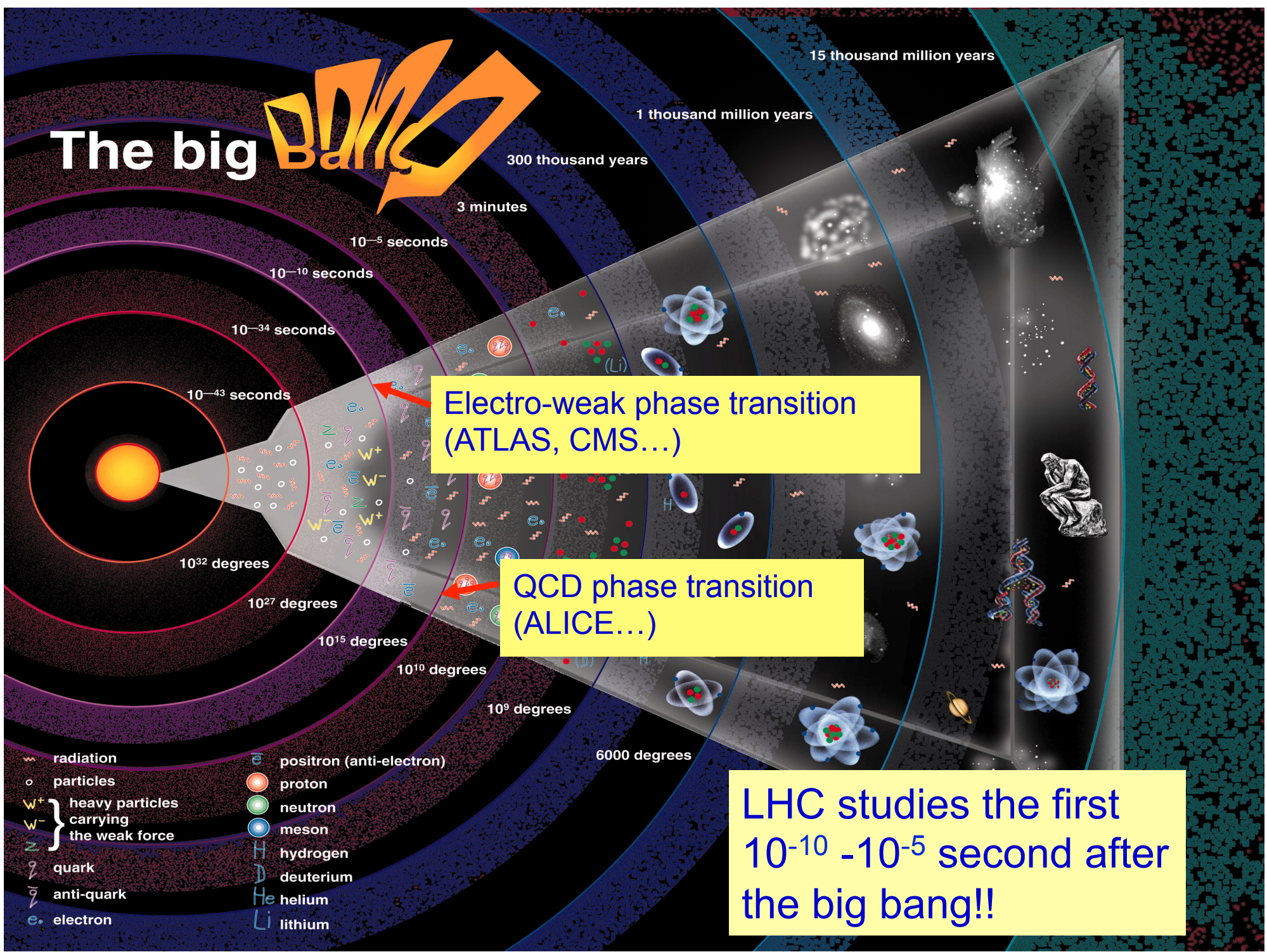


General Purpose,
pp, heavy ions

Heavy ions, pp



The big Bang

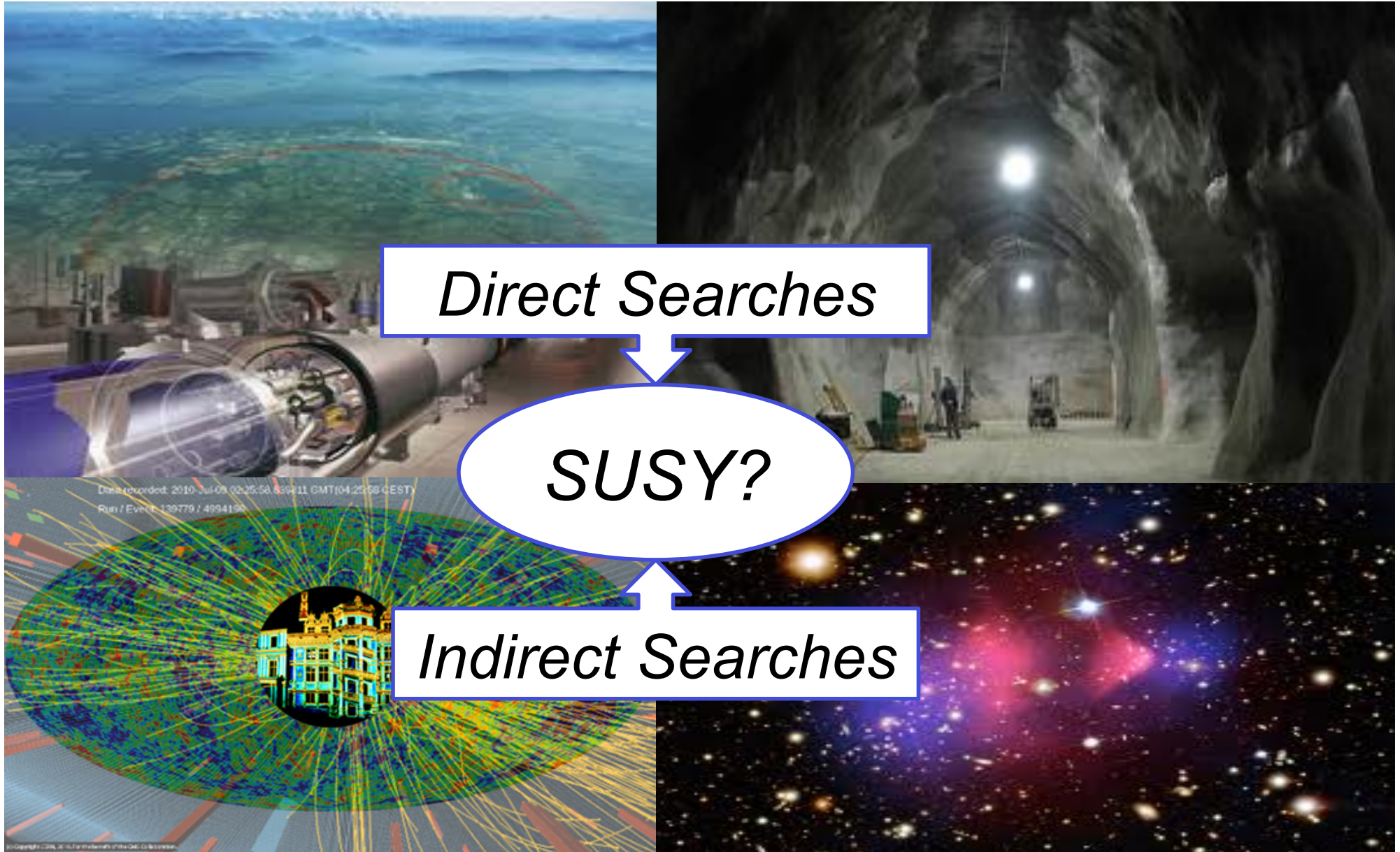


Electro-weak phase transition (ATLAS, CMS...)

QCD phase transition (ALICE...)

LHC studies the first 10^{-10} - 10^{-5} second after the big bang!!

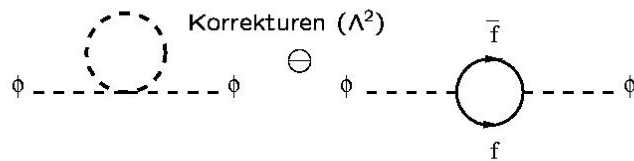
Searches for SUSY



Why is SUSY so attractive?

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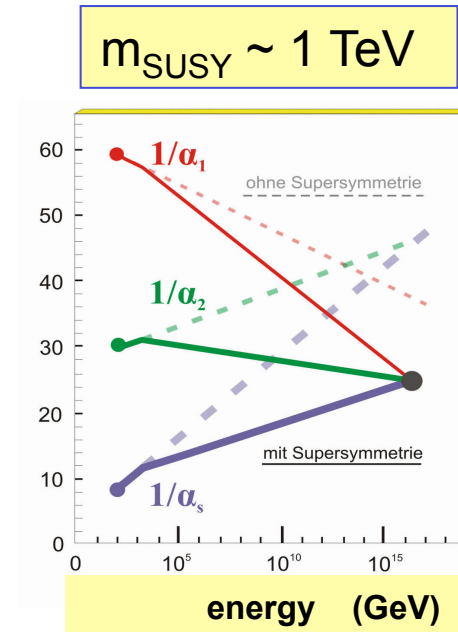
1. Quadratically divergent quantum corrections to the Higgs boson mass are avoided



$$\Delta m_H = f(m_B^2 - m_f^2)$$

(Hierarchy or naturalness problem)

2. Unification of coupling constants of the three interactions seems possible

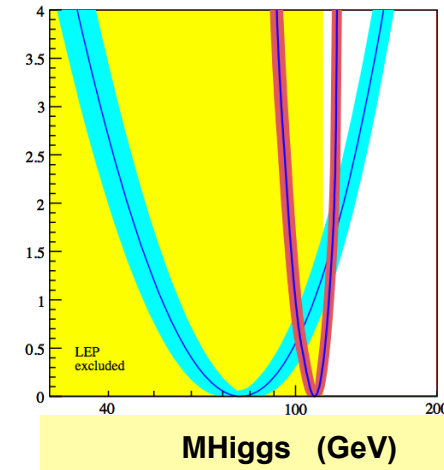


3. SUSY provides a candidate for dark matter,



**The lightest
SUSY particle
(LSP)**

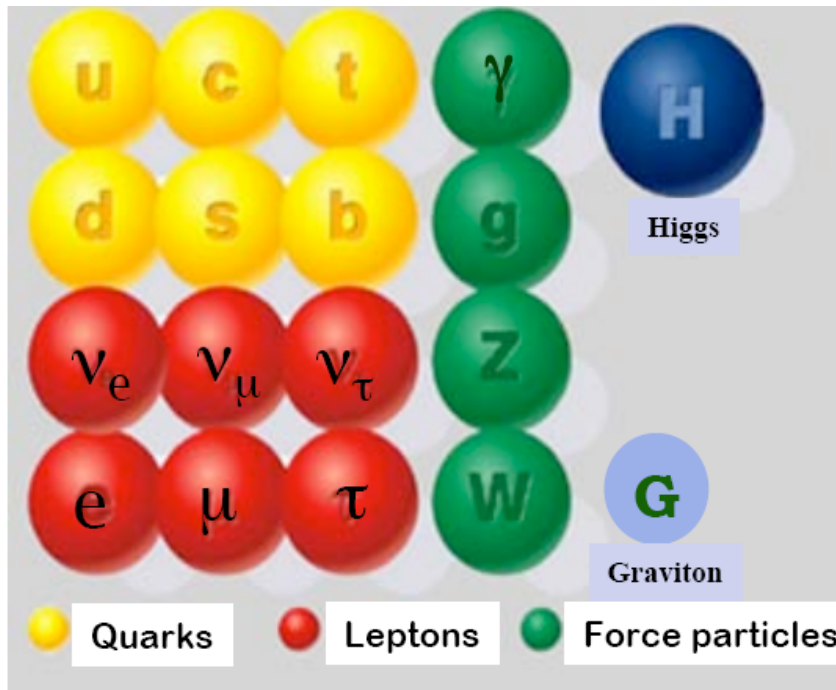
4. A SUSY extension is a small perturbation, consistent with the electroweak precision data



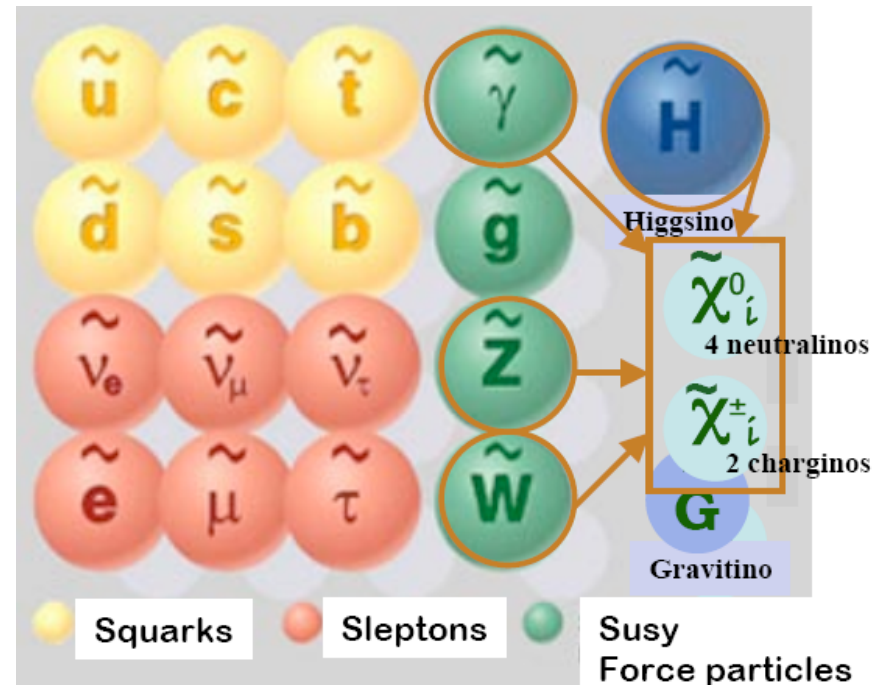
Supersymmetry

Extension of the Standard Model: Introduce a new symmetry
Spin 1/2 matter particles (fermions) \leftrightarrow Spin 1 force carriers (bosons)

Standard Model particles



SUSY particles



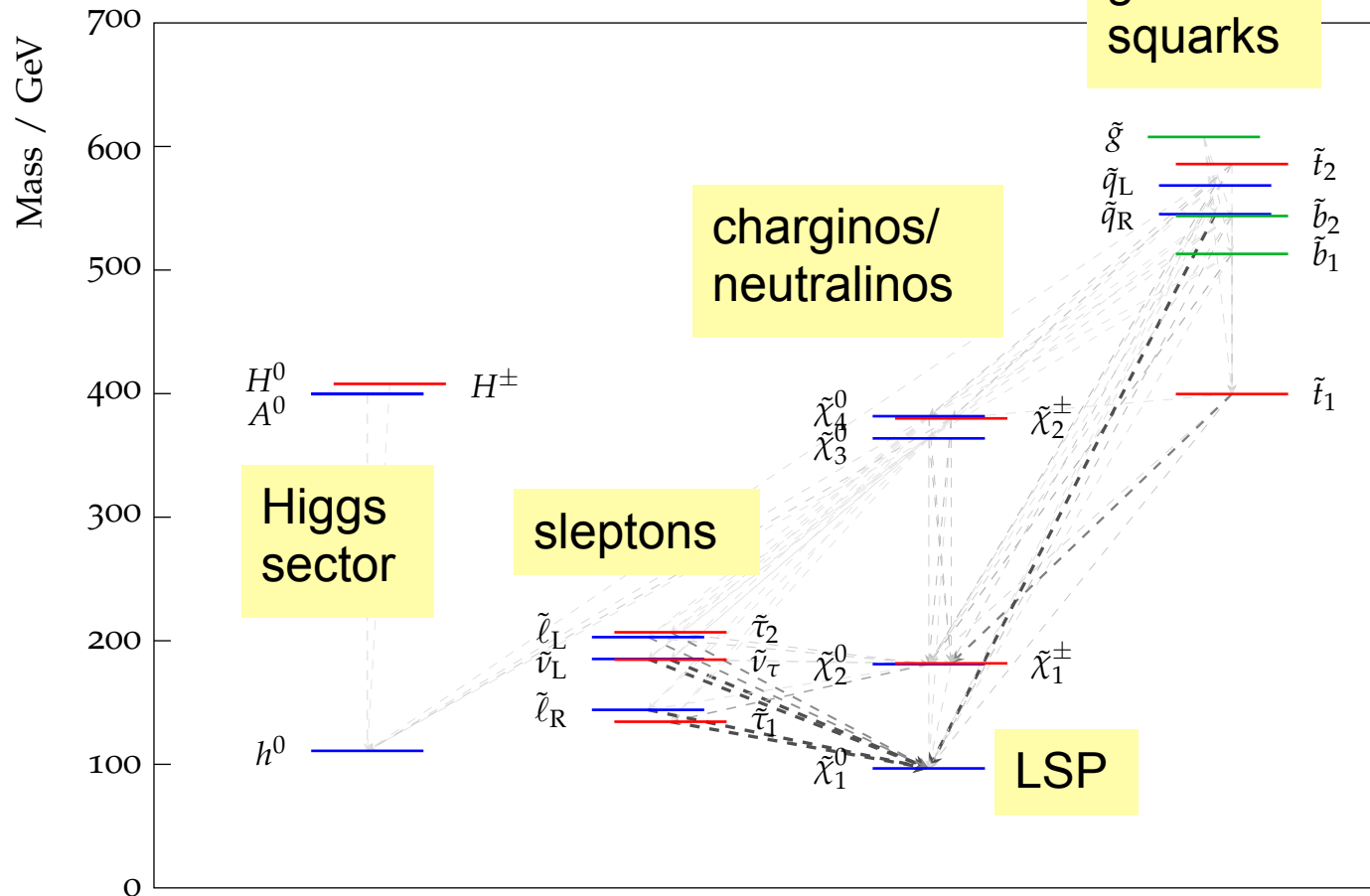
New Quantum number: R-parity: $R_p = (-1)^{B+L+2s} = +1$ SM particles
 R-parity conservation: -1 SUSY particles

- SUSY particles are produced in pairs
- The lightest SUSY particle (LSP) is stable

A “typical” SUSY Spectrum

Use the famous SPS1a benchmark point for illustration
 $[m_0=100, m_{1/2}=250, \tan\beta=10, A_0=-100, \mu>0]$

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CMSSM
 $m_0, m_{1/2}, \tan\beta, A_0, \text{sign}(\mu)$

Advantage:

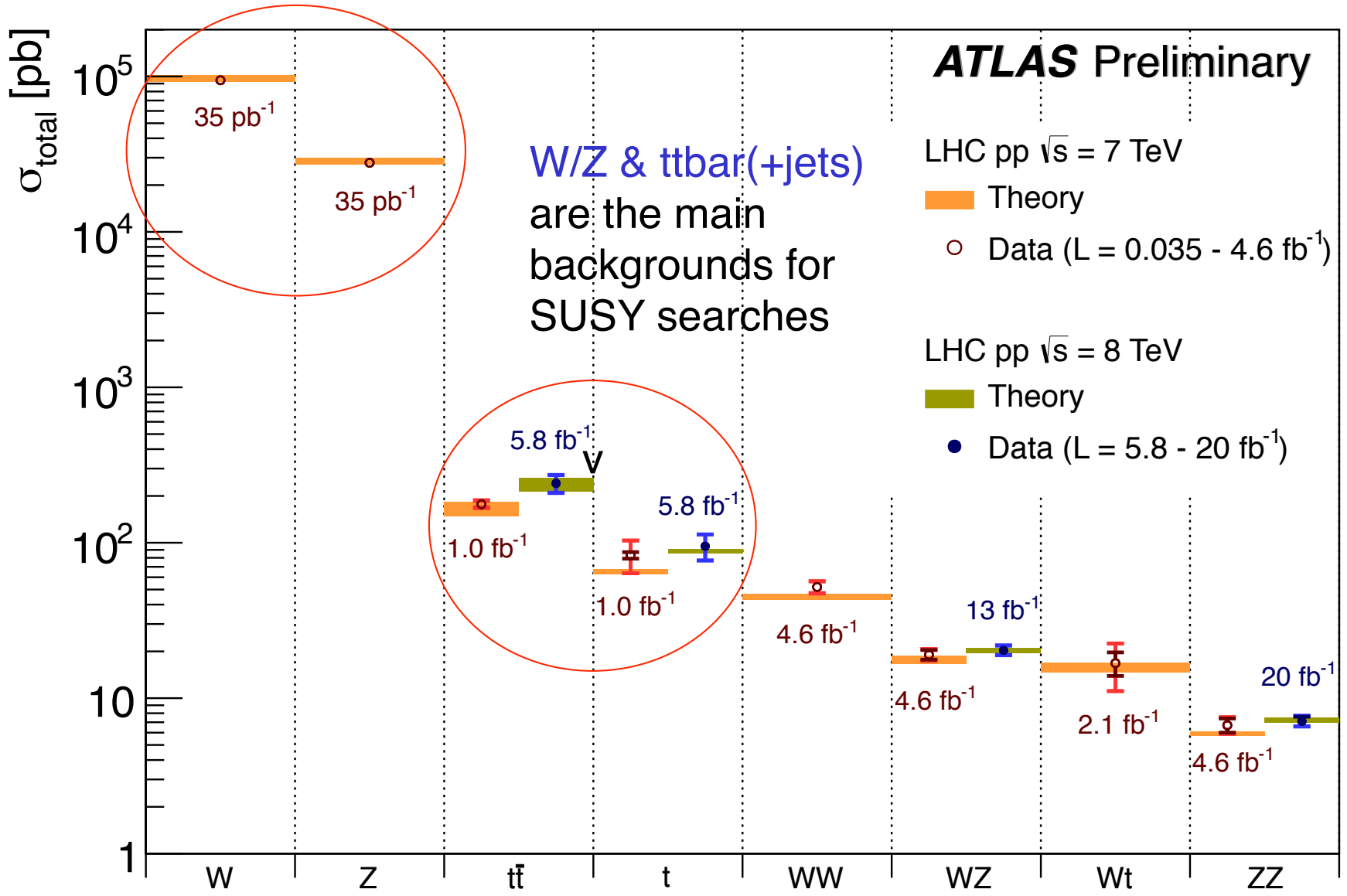
- Only four free parameters (when $\text{sign}(\mu)$ fixed)
- One of the most studied incarnations of the MSSM

Disadvantage:

- Not generally representative of SUSY (e.g. fixed mass relation between M_{gluon} and M_{LSP})

Rediscovery of the SM at a new energy frontier

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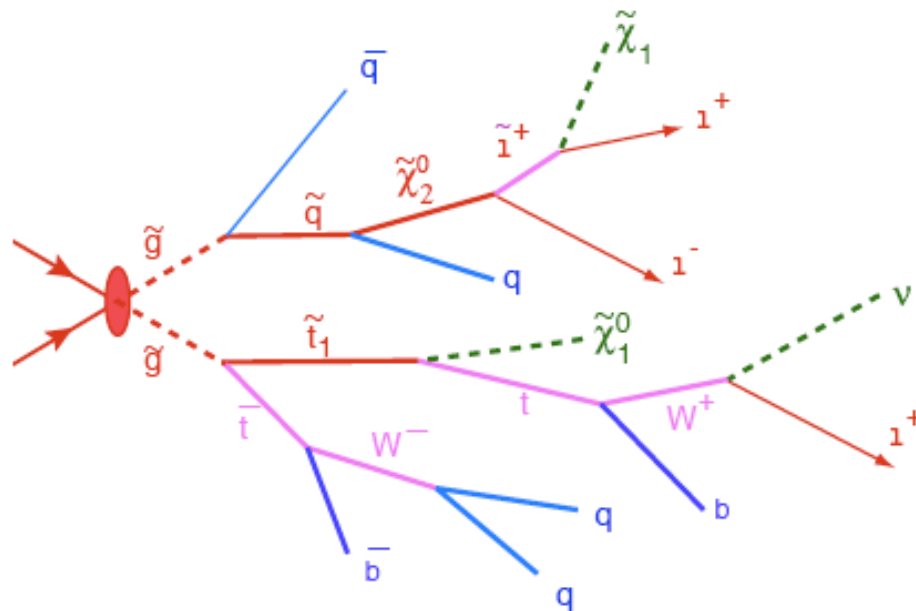
What do we call a “SUSY search”?

The definition is purely derived from the experimental signature.

Therefore, a “SUSY search signature” is characterized by

Lots of missing energy, many jets, and possibly leptons in the final state

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Missing Energy:

- from LSP

Multi-Jet:

- from cascade decay (gaugino)

Multi-Leptons:

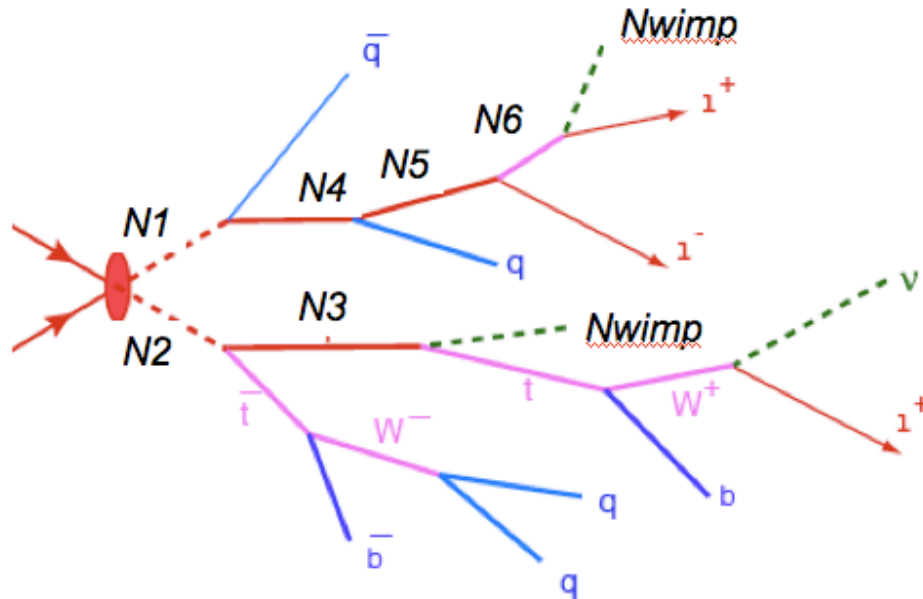
- from decay of charginos/neutralinos

RP-Conserving SUSY is a very prominent example predicting this famous signature but ...

What is its experimental signature?

... by no means is it the only New Physics model predicting this experimental pattern. Many other NP models predict this genuine signature

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Missing Energy:

- *Nwimp* - end of the cascade

Multi-Jet:

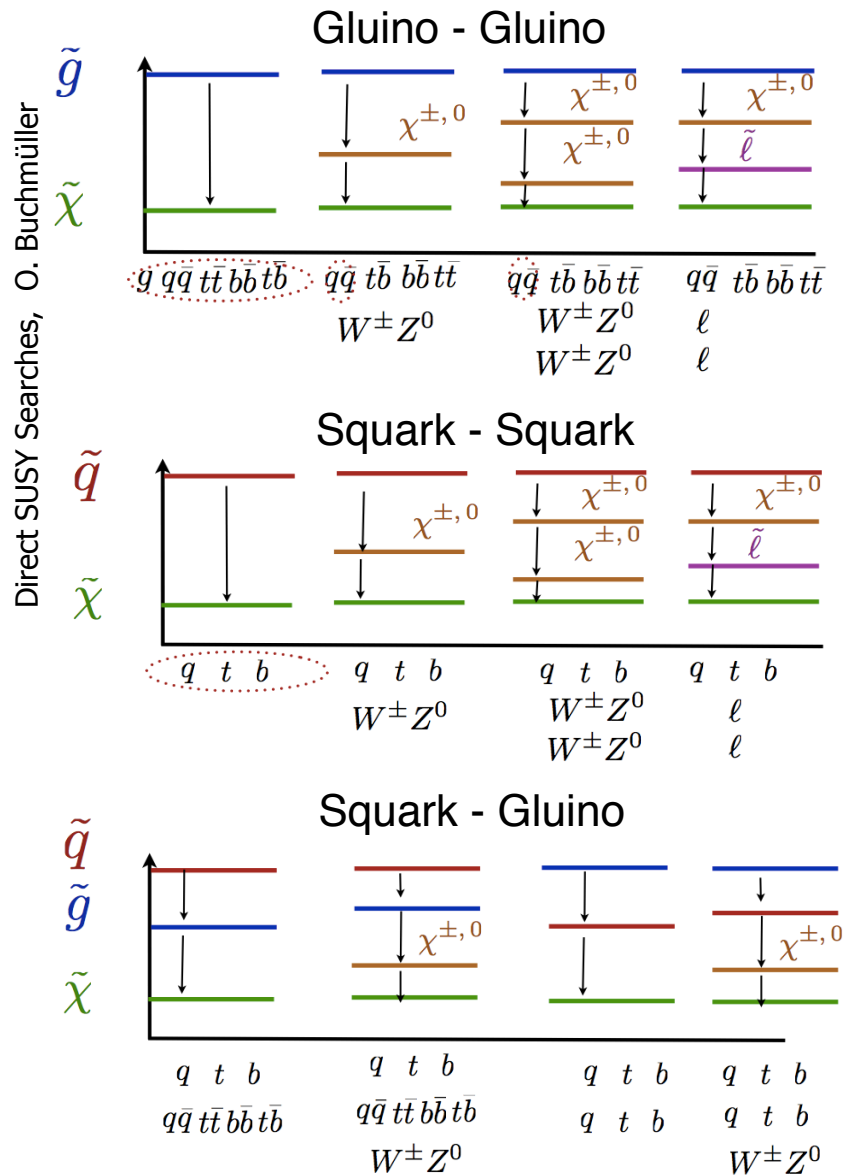
- from decay of the *N*s (possibly via heavy SM particles like top, W/Z)

Multi-Leptons:

- from decay of the *N*'s

Model examples are Extra dimensions, Little Higgs, Technicolour, etc
but a more generic definition for this signature is as follows.

Early SUSY Search Strategy at the LHC



Search Signatures

- SUSY-like decay chains range from short to long and simple to very complicated.
- All physics objects, MET, jets, leptons, photons, b's taus, tops, W, Z, etc are involved
- Comprehensive coverage of all possible signature requires a topology oriented search strategy:

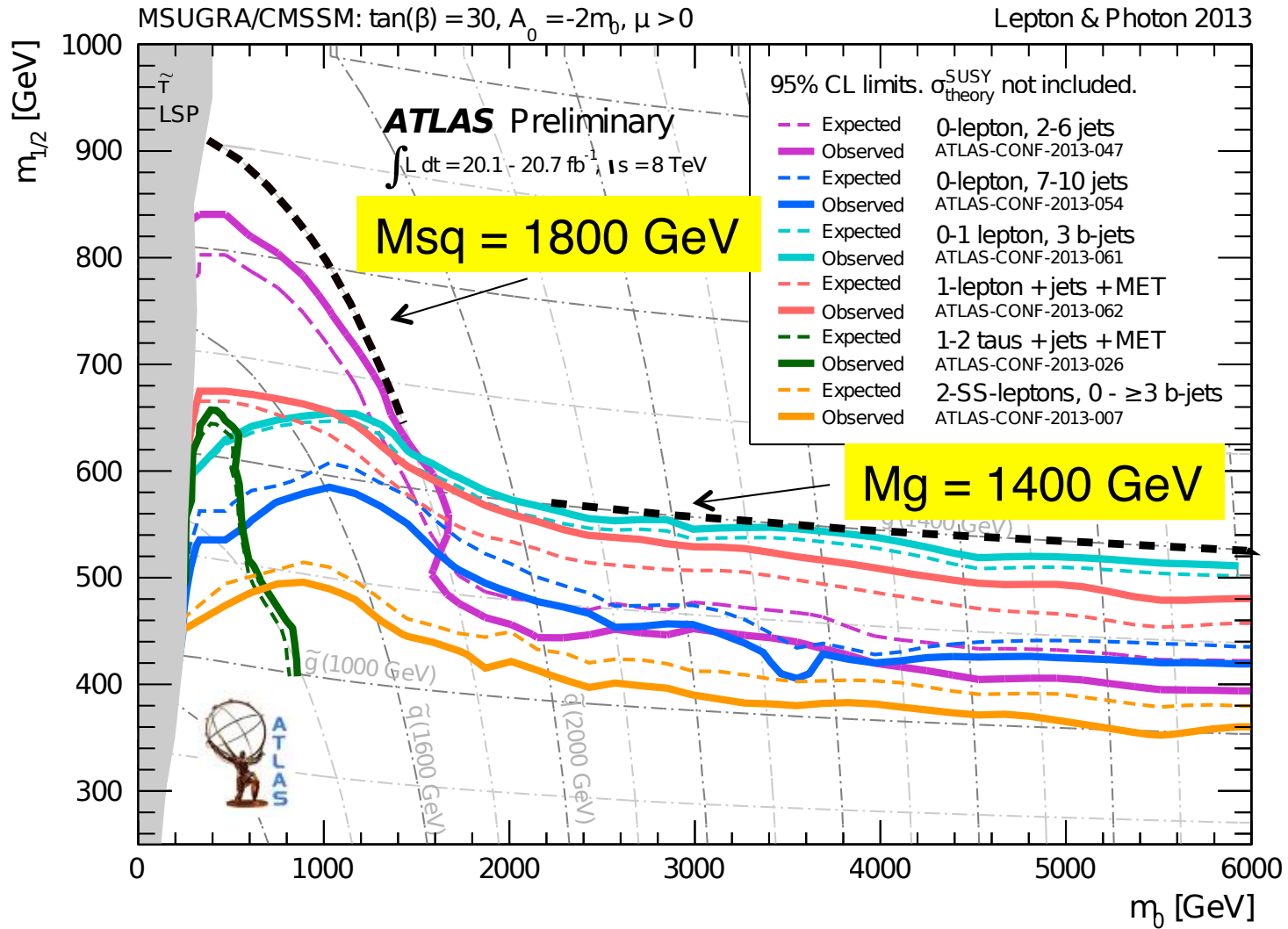
References Analyses

0-leptons	1-lepton	OSDL	SSDL	≥ 3 leptons	2-photons	γ +lepton
Jets + MET	Single lepton + Jets + MET	Opposite-sign di-lepton + jets + MET	Same-sign di-lepton + jets + MET	Multi-lepton	Di-photon + jet + MET	Photon + lepton + MET

Already in less than two years of operation ATLAS & CMS managed to carry out the full list of these core "SUSY References Analyses"!

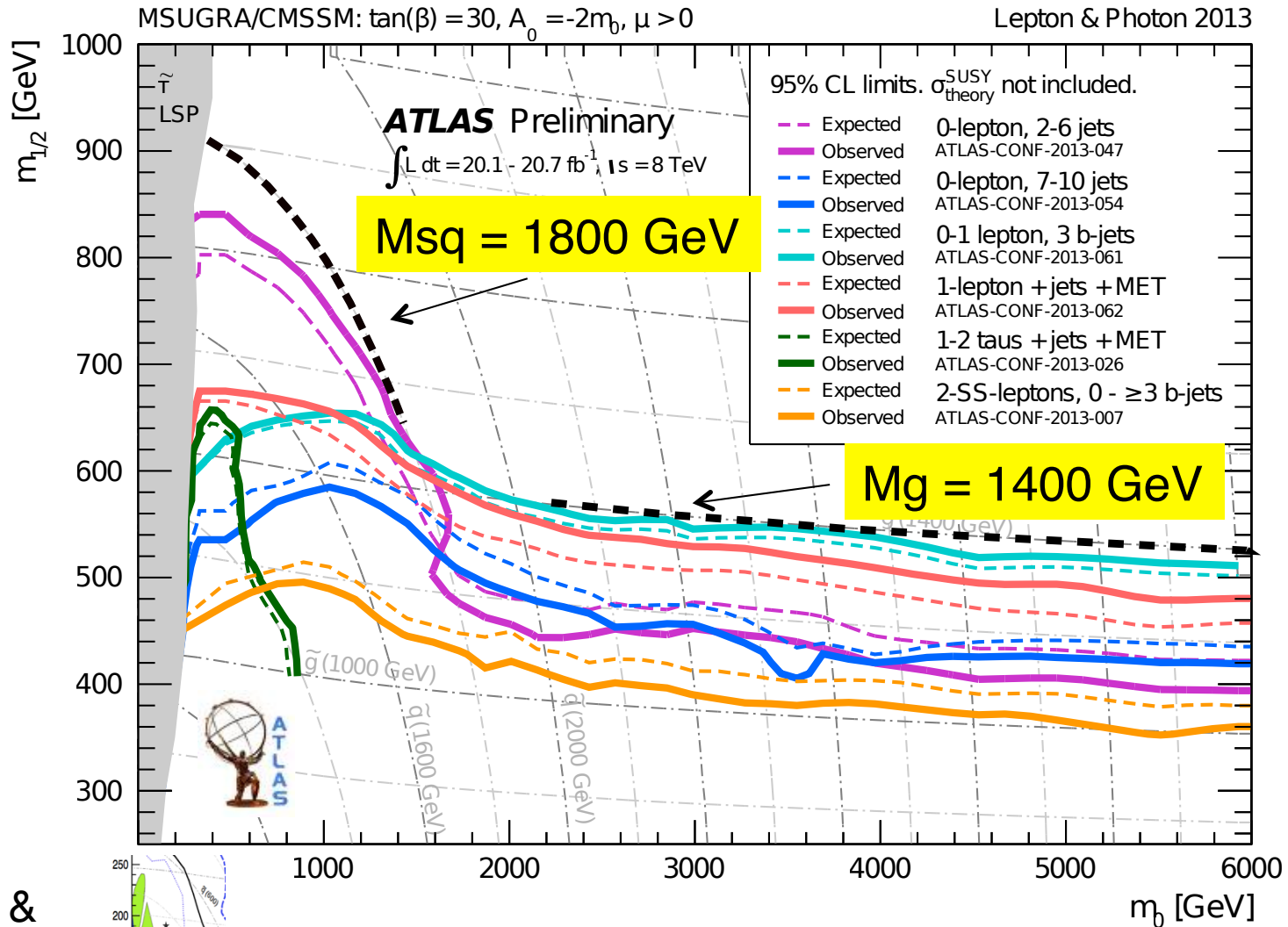
Inclusive SUSY Searches in 2013

Direct SUSY Searches, O. Buchmüller



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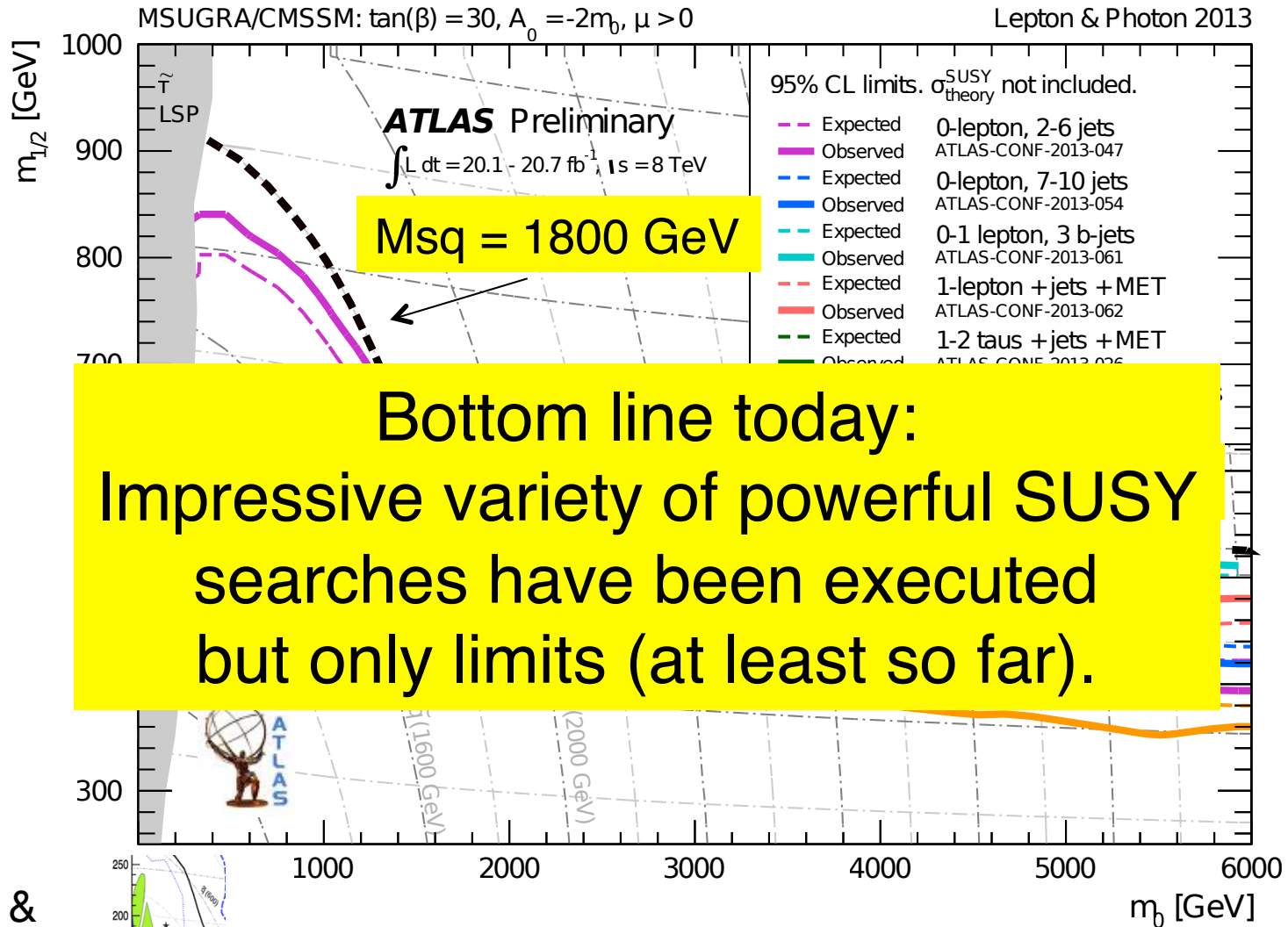
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The LHC has pushed the mass scale in constraint SUSY models to a new level!

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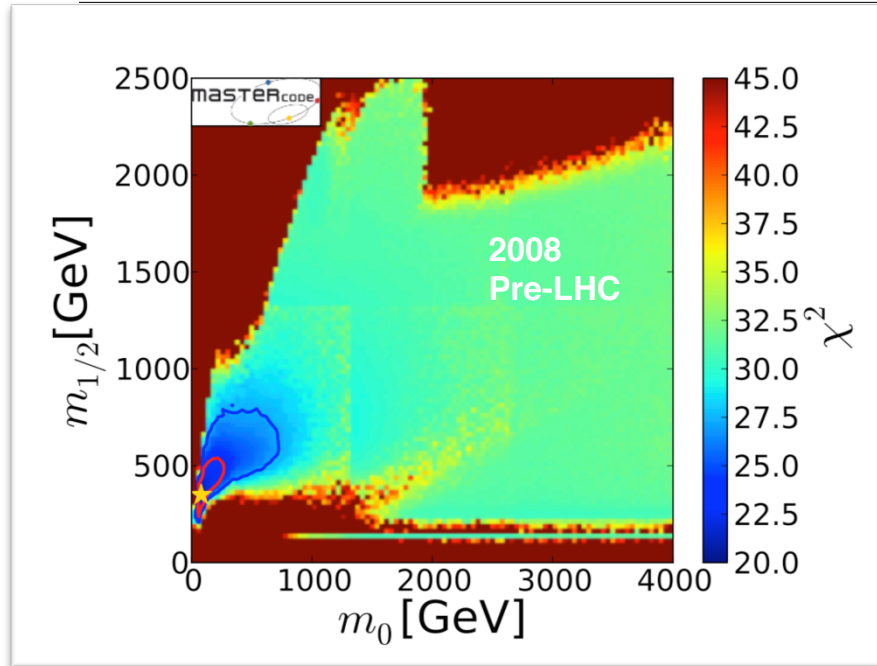
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The LHC has pushed the mass scale in constraint SUSY models to a new level!

CMSSM: Evolution with time

Direct SUSY Searches, O. Buchmüller



χ^2 increase from
bluish to reddish



Source:
<http://mastercode.web.cern.ch/mastercode/>

Observable	Source Th./Ex.	Constraint	$\Delta\chi^2$ (CMSSM)	$\Delta\chi^2$ (NUHM1)	$\Delta\chi^2$ ("SM")
m_t [GeV]	[43]	173.2 ± 0.90	0.05	0.06	-
$\Delta\alpha_{had}^{(5)}(M_Z)$	[42]	0.02749 ± 0.00010	0.009	0.004	-
M_Z [GeV]	[44]	91.1875 ± 0.0021	2.7×10^{-2}	0.26	-
Γ_Z [GeV]	[26] / [44]	$2.4952 \pm 0.0023 \pm 0.001_{SUSY}$	0.078	0.047	0.14
$\sigma_{had}^{in}(nb)$	[26] / [44]	41.540 ± 0.037	2.50	2.57	2.54
R_t	[26] / [44]	20.767 ± 0.025	1.05	1.08	1.08
$A_{fb}(\ell)$	[26] / [44]	0.01714 ± 0.00095	0.72	0.69	0.81
$A_{FB}(P_T)$	[26] / [44]	0.1465 ± 0.0032	0.11	0.13	0.07
R_b	[26] / [44]	0.21629 ± 0.00066	0.26	0.29	0.27
R_c	[26] / [44]	0.1721 ± 0.0030	0.002	0.002	0.002
$A_{FB}(b)$	[26] / [44]	0.0992 ± 0.0016	7.17	7.37	6.63
$A_{FB}(c)$	[26] / [44]	0.0707 ± 0.0035	0.86	0.88	0.80
A_b	[26] / [44]	0.923 ± 0.020	0.36	0.36	0.35
A_c	[26] / [44]	0.670 ± 0.027	0.005	0.005	0.005
$A_{FB}(SLD)$	[26] / [44]	0.1513 ± 0.0021	3.16	3.03	3.51
$\sin^2 \theta_{eff}^l(Q_{FB})$	[26] / [44]	0.2324 ± 0.0012	0.63	0.64	0.59
M_W [GeV]	[26] / [44]	$80.399 \pm 0.023 \pm 0.010_{SUSY}$	1.77	1.39	2.08
$a_{\mu}^{exp} - a_{\mu}^{SM}$	[53] / [42,54]	$(30.2 \pm 8.8 \pm 2.0_{SUSY}) \times 10^{-10}$	4.35	1.82	11.19 (N/A)
M_h [GeV]	[28] / [53,56]	$> 114.4[\pm 1.5_{SUSY}]$	0.0	0.0	0.0
$BR_{b \rightarrow sy}^{EXP/SM}$	[45] / [46]	$1.117 \pm 0.076_{EXP} \pm 0.082_{SM} \pm 0.050_{SUSY}$	1.83	1.09	0.94
$BR(B_s \rightarrow \mu^+ \mu^-)$	[29] / [41]	CMS & LHCb	0.04	0.44	0.01
$BR_{B \rightarrow \tau \nu}^{EXP/SM}$	[29] / [46]	$1.43 \pm 0.43_{EXP+TH}$	1.43	1.59	1.00
$BR(B_d \rightarrow \mu^+ \mu^-)$	[29] / [46]	$< 4.6[\pm 0.1_{SUSY}] \times 10^{-9}$	0.0	0.0	0.0
$BR_{\mu \rightarrow X \gamma}^{EXP/SM}$	[47] / [46]	0.99 ± 0.32	0.02	$\ll 0.01$	$\ll 0.01$
$BR_{K \rightarrow \mu \nu}^{EXP/SM}$	[29] / [48]	$1.008 \pm 0.014_{EXP+TH}$	0.39	0.42	0.33
$BR_{K \rightarrow \pi \nu \nu}^{EXP/SM}$	[49] / [50]	< 4.5	0.0	0.0	0.0
$\Delta M_{B_s}^{EXP/SM}$	[49] / [51,52]	$0.97 \pm 0.01_{EXP} \pm 0.27_{SM}$	0.02	0.02	0.01
$\Delta M_{B_d}^{EXP/SM}$	[29] / [46,51,52]	$1.00 \pm 0.01_{EXP} \pm 0.13_{SM}$	$\ll 0.01$	0.33	$\ll 0.01$
$\Delta c_K^{EXP/SM}$	[49] / [51,52]	$1.08 \pm 0.14_{EXP+TH}$	0.27	0.37	0.33
$\Omega_{CDM} h^2$	[31] / [13]	$0.1120 \pm 0.0056 \pm 0.012_{SUSY}$	8.4×10^{-4}	0.1	N/A
σ_p^{21}	[25]	$(m_{\tilde{g}}, \sigma_p^{21})$ plane	0.13	0.13	N/A
jets + B_T	[18,20]	$(m_0, m_{1/2})$ plane	1.55	2.20	N/A
$H/A, H^\pm$	[21]	$(M_A, \tan \beta)$ plane	0.0	0.0	N/A
Total $\chi^2/d.o.f.$	All	All	28.8/22	27.3/21	32.7/23 (21.5/22)
p-values			15%	16%	9% (49%)

Global Fit to indirect and direct constraints on SUSY!

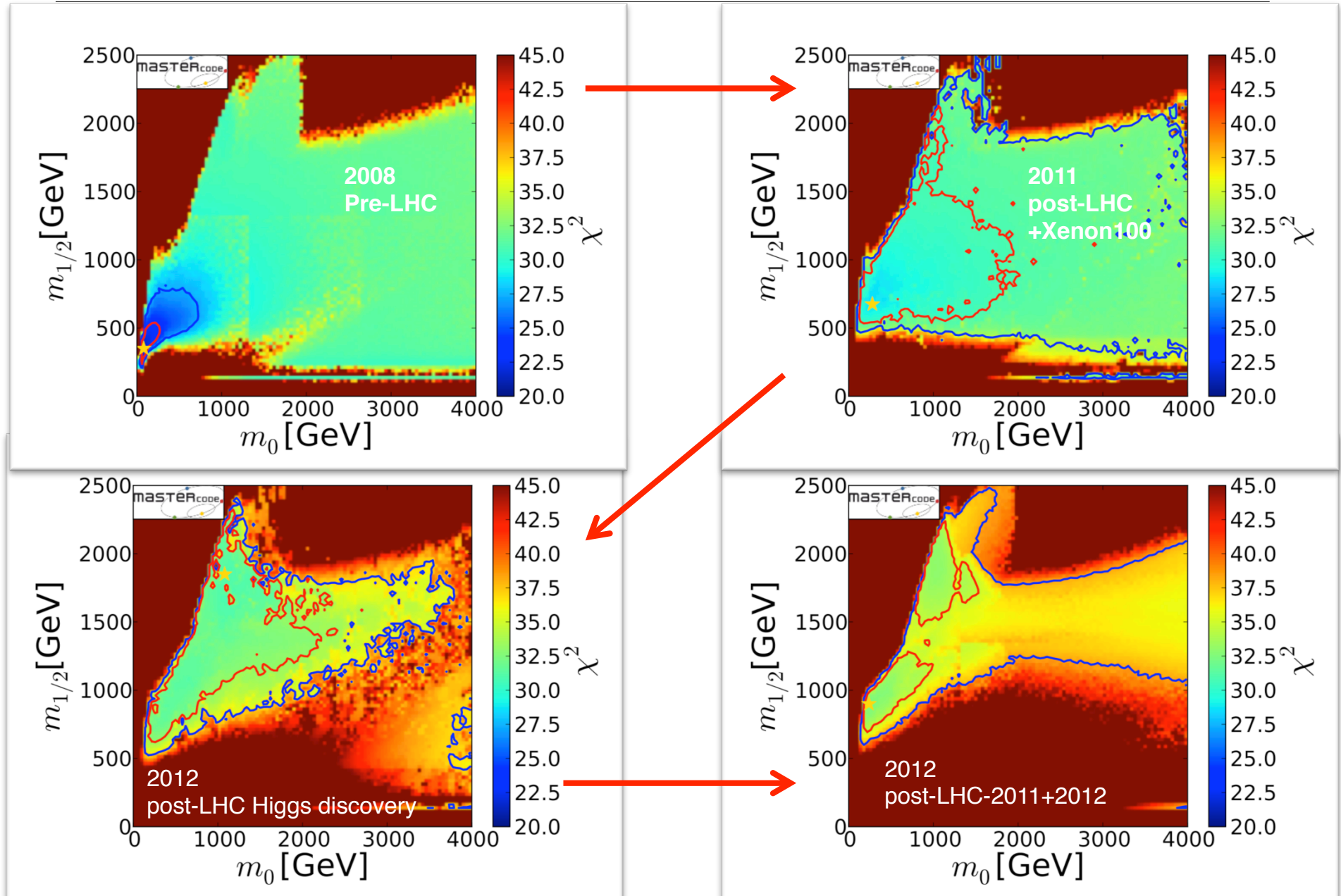
Other "fitter" groups find very similar results: e.g.

SuperBayeS: [arXiv:1212.2636](https://arxiv.org/abs/1212.2636)

Fittino group: [arXiv:1204.4199](https://arxiv.org/abs/1204.4199)

CMSSM: Evolution with time

Direct SUSY Searches, O. Buchmüller

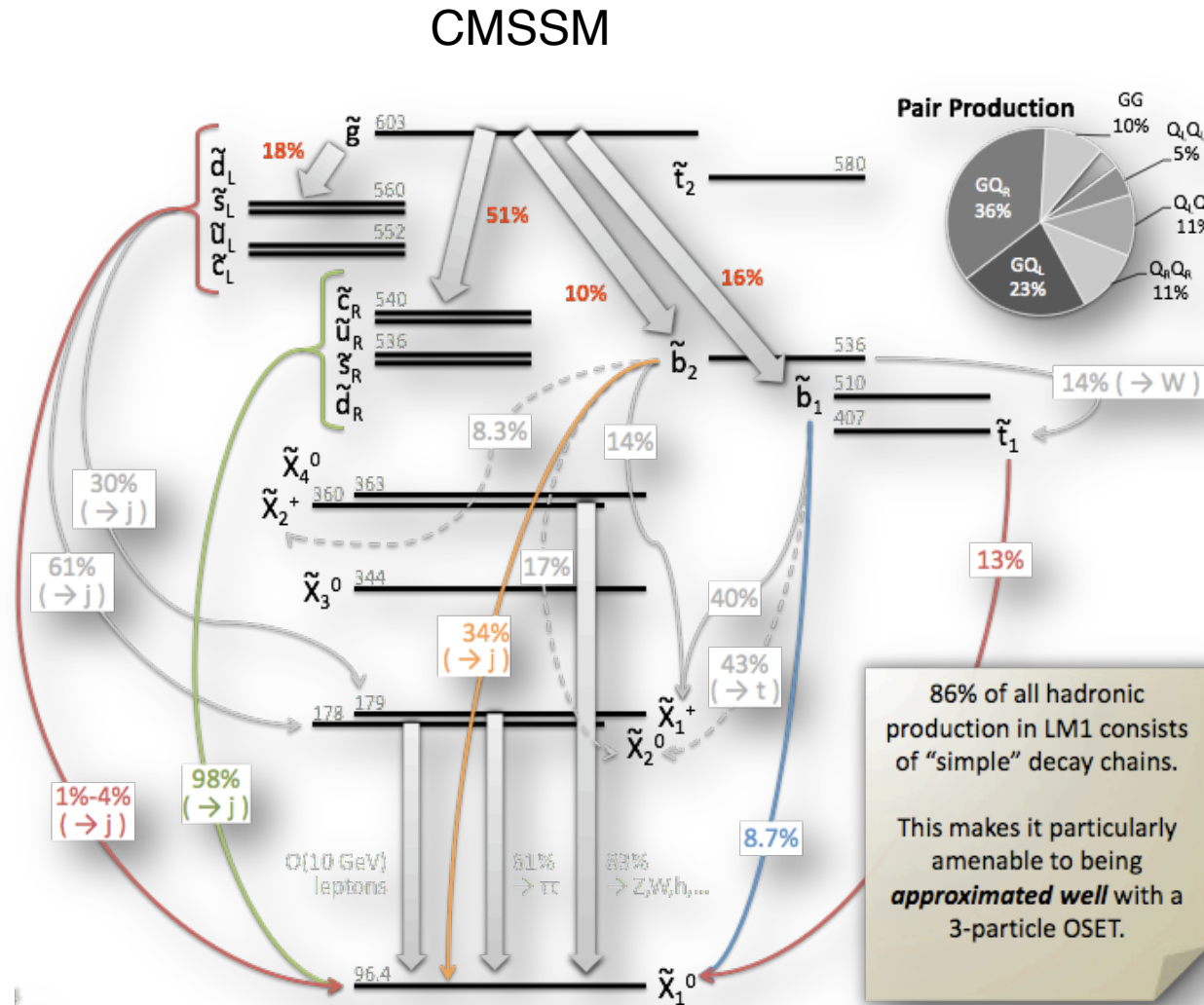


SUSY Status – post 7 TeV LHC data

- Constrained SUSY models like the CMSSM are severely put under pressure by the LHC limits!
- Experiments define new benchmarks and less complex SUSY models in order to present the interpretation of their searches.
- Aided by the discovery of a Higgs boson, the focus of the experimental search strategy and corresponding interpretation shifts towards other scenarios like “Natural SUSY” (i.e. 3rd generation squark searches).

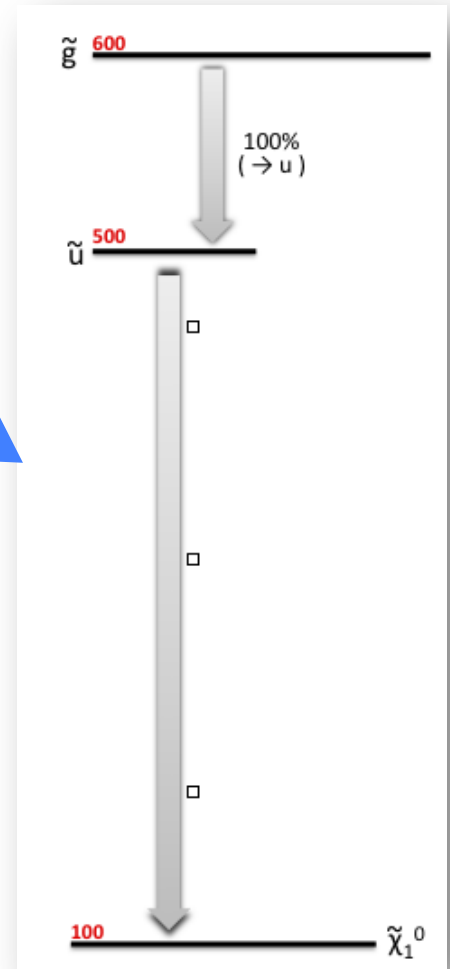
Interpretation in Simplified Models

Direct SUSY Searches, O. Buchmüller



What the individual searches are sensitive to is much more simple...

86% of all hadronic production in LM1 consists of "simple" decay chains. This makes it particularly amenable to being approximated well with a 3-particle OSET.

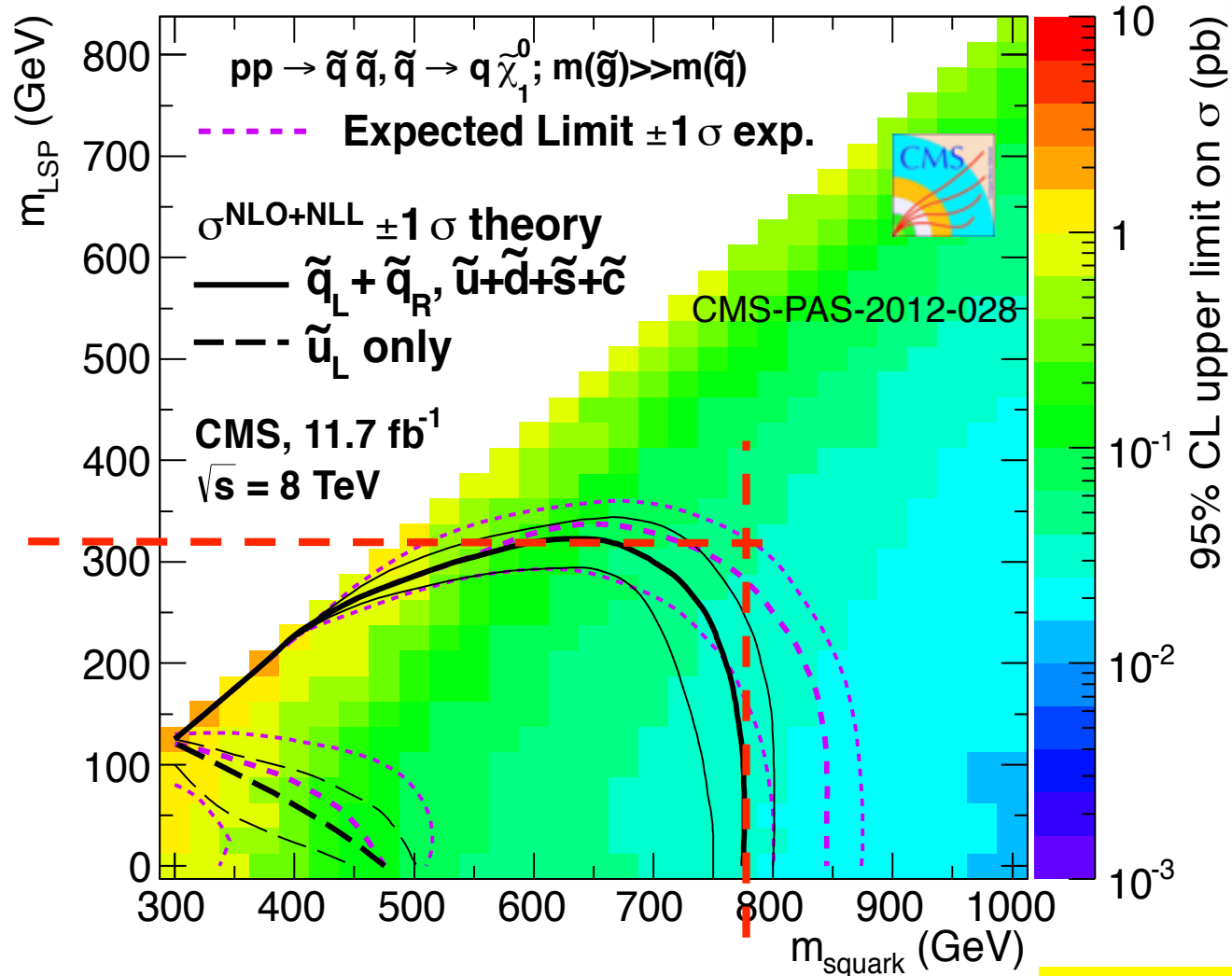


Simplified model spectrum (SMS)
with 3 particles, 2 decay modes

SMS: a few interesting features

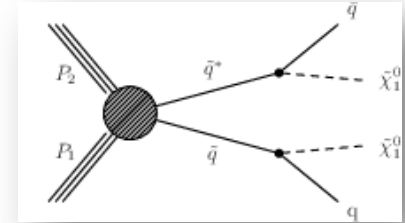
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$m_{\text{LSP}}^{\text{max}} \approx 0.3 \text{ TeV}$: LSP mass above
which there is NO limit anymore



$m_G^{\text{max}} \approx 0.8 \text{ TeV}$: Best limit in plane

Assumes 100%
BR for decay chain
considered.



$$\tilde{q}\tilde{q} \rightarrow q\tilde{\chi}_1^0\bar{q}\tilde{\chi}_1^0$$

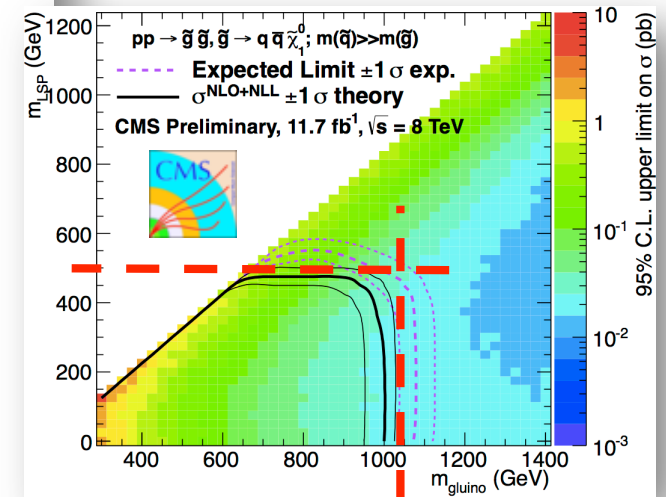
How to summarize SMS limits?

Approach taken in the 2012 Experimental SUSY PDG review
[OB & Paul De Jong]:

<http://pdg.lbl.gov/2012/reviews/rpp2012-rev-susy-2-experiment.pdf>

Direct SUSY Searches, O. Buchmüller

Model	Assumption	$m_{\tilde{q}}$	$m_{\tilde{g}}$
CMSSM	$m_{\tilde{q}} \approx m_{\tilde{g}}$	1400	1400
	all $m_{\tilde{q}}$	-	800
	all $m_{\tilde{g}}$	1300	-
Simplified model $\tilde{g}\tilde{g}$	$m_{\tilde{\chi}_1^0} = 0$	-	900
	$m_{\tilde{\chi}_1^0} > 300$	-	no limit
Simplified model $\tilde{q}\tilde{q}$	$m_{\tilde{\chi}_1^0} = 0$	750	-
	$m_{\tilde{\chi}_1^0} > 250$	no limit	-
Simplified model $\tilde{g}\tilde{q}, \tilde{g}\tilde{\tilde{q}}$	$m_{\tilde{\chi}_1^0} = 0, m_{\tilde{q}} \approx m_{\tilde{g}}$	1500	1500
	$m_{\tilde{\chi}_1^0} = 0, \text{all } m_{\tilde{g}}$	1400	-
	$m_{\tilde{\chi}_1^0} = 0, \text{all } m_{\tilde{q}}$	-	900



This was an appropriate approach for the rather limited amount of inclusive searches and corresponding SMS interpretations available in 2011 (7 TeV).

How to summarize SMS limits?

Approach taken in the 2012 Experimental SUSY PDG review
[OB & Paul De Jong]:

<http://pdg.lbl.gov/2012/reviews/rpp2012-rev-susy-2-experiment.pdf>

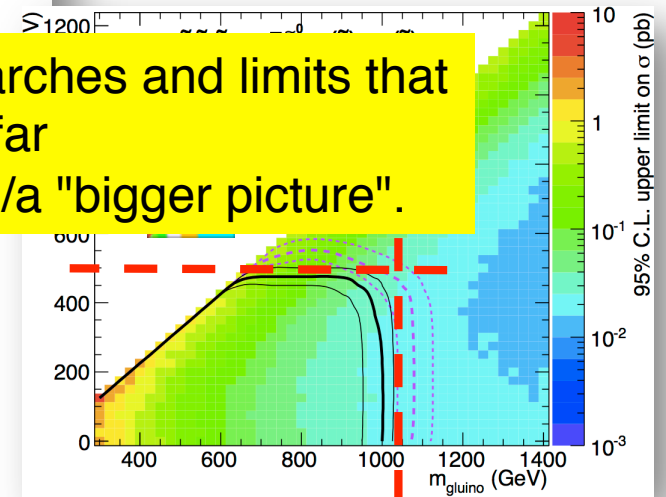
Direct SUSY Searches, O. Buchmüller

Model	Assumption	$m_{\tilde{q}}$	$m_{\tilde{g}}$
CMSSM	$m_{\tilde{q}} \approx m_{\tilde{g}}$	1400	1400
	all $m_{\tilde{a}}$	-	800

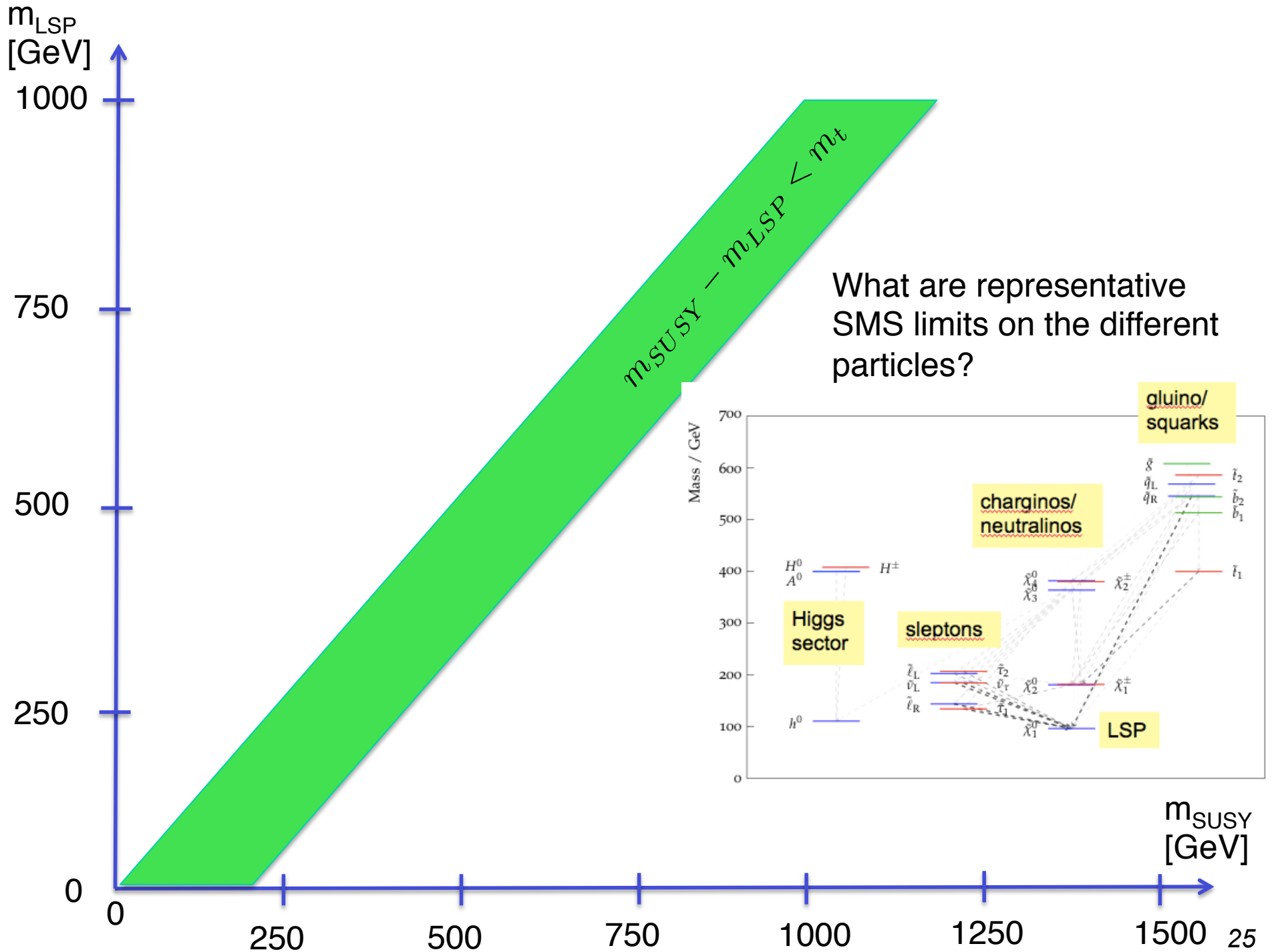
Simplified

Simplified model $\tilde{q}\tilde{q}$	$m_{\tilde{\chi}_1^0} = 0$	750	-
	$m_{\tilde{\chi}_1^0} > 250$	no limit	-
Simplified model $\tilde{g}\tilde{q}, \tilde{g}\tilde{\tilde{q}}$	$m_{\tilde{\chi}_1^0} = 0, m_{\tilde{q}} \approx m_{\tilde{g}}$	1500	1500
	$m_{\tilde{\chi}_1^0} = 0, \text{ all } m_{\tilde{g}}$	1400	-
	$m_{\tilde{\chi}_1^0} = 0, \text{ all } m_{\tilde{q}}$	-	900

It is a challenge to do justice to the many searches and limits that have been established so far
- even more so to put it all together into the/a "bigger picture".

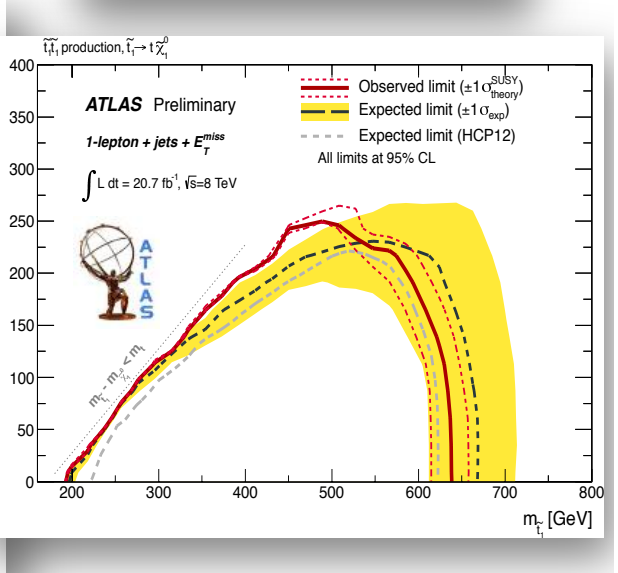
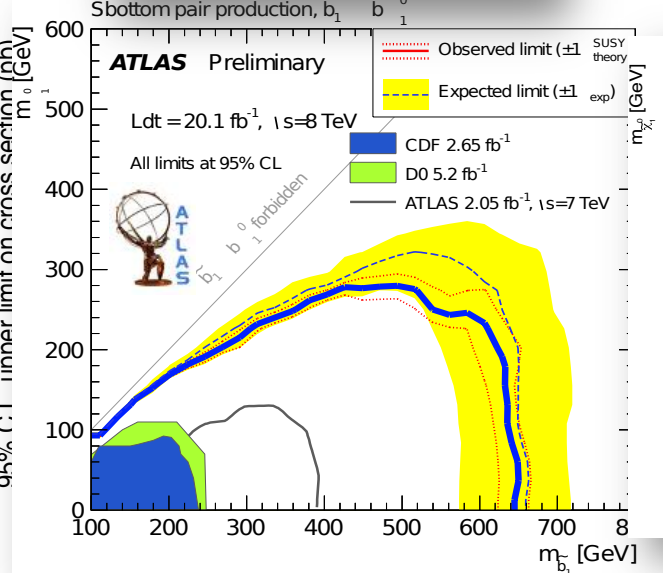
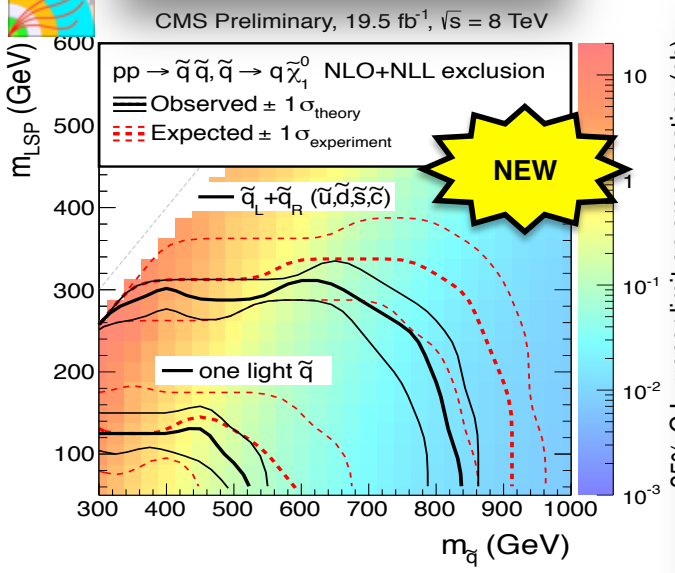
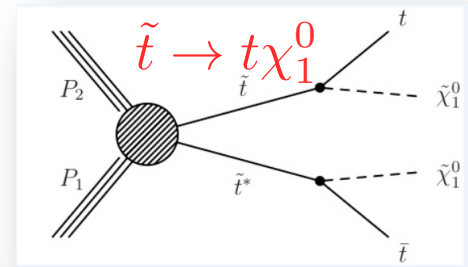
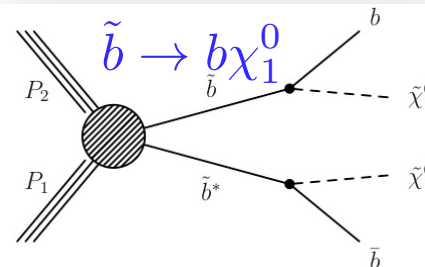
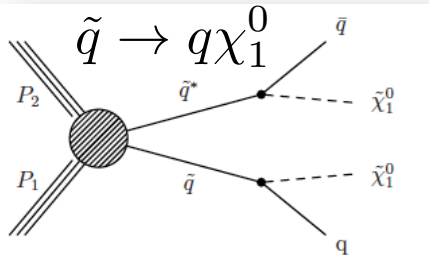


This was an appropriate approach for the rather limited amount of inclusive searches and corresponding SMS interpretations available in 2011 (7 TeV).



Direct squark production – chosen limits

Direct SUSY Searches, O. Buchmüller



CMS-SUS-PAS-13-012

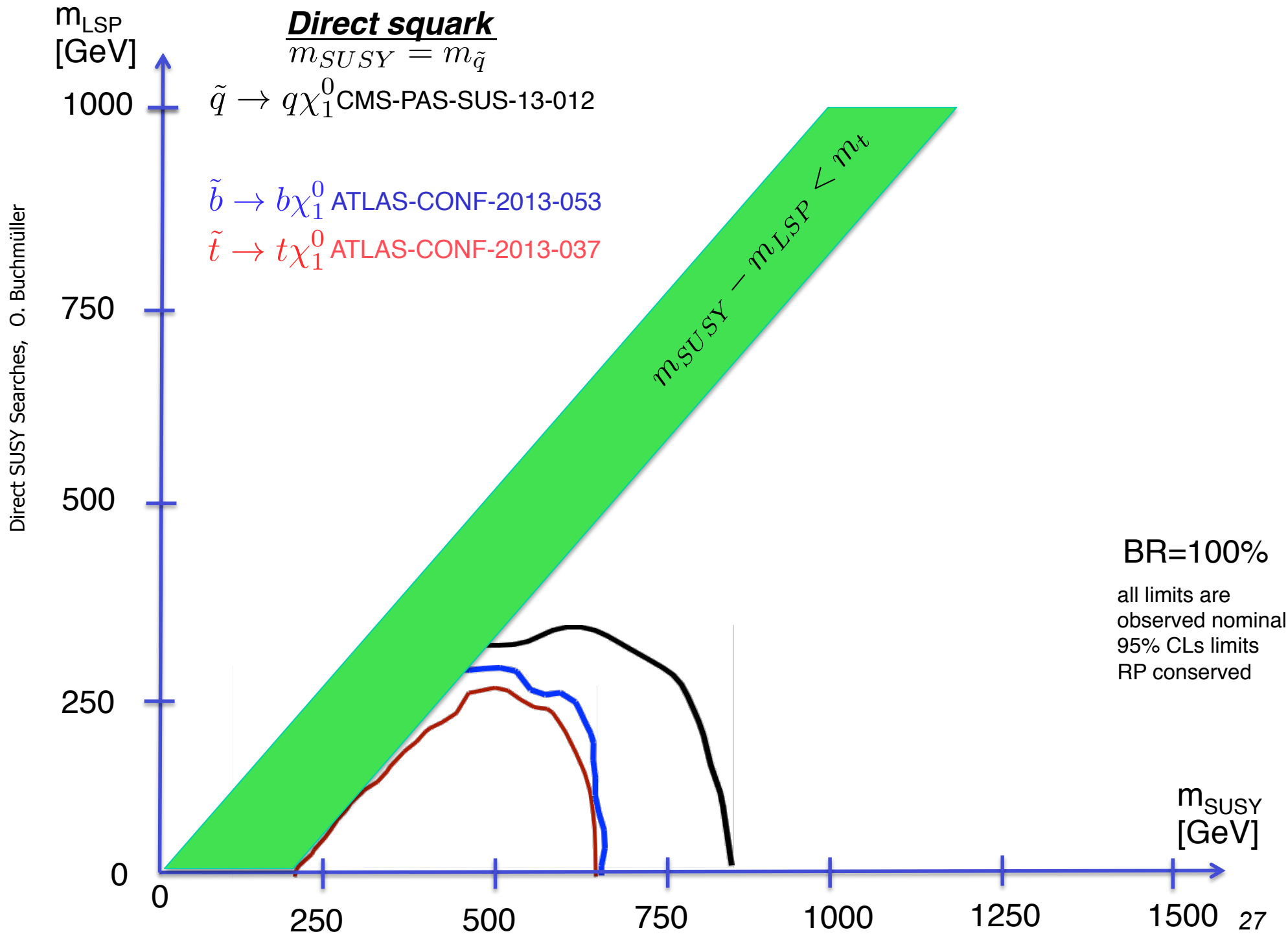
Signature: Jets + E_T^{miss} + H_T
Limit assumes all 1st & 2nd gen squarks to be mass degenerate or only one light squark!

ATLAS-CONF-2013-053

Signature: 2 b-jets + E_T^{mis}

ATLAS-CONF-2013-037

Signature: 1 Lepton + jets + E_T^{mis}



m_{LSP}
[GeV]

Direct squark

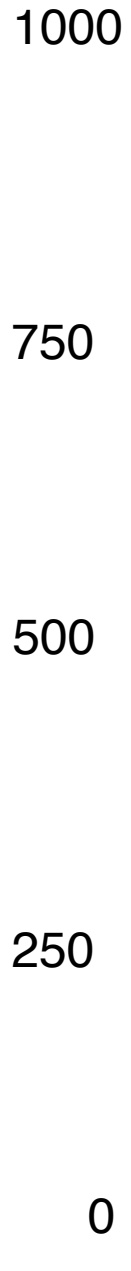
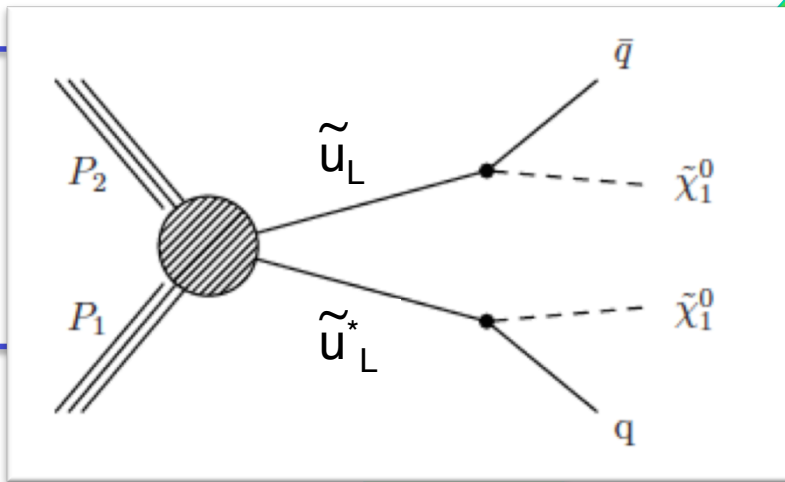
$m_{SUSY} = m_{\tilde{q}}$

$\tilde{q} \rightarrow q\chi_1^0$ CMS-PAS-SUS-13-012 (8x mass)

$\tilde{u}_L \rightarrow q\chi_1^0$ CMS-PAS-SUS-13-012

$\tilde{b} \rightarrow b\chi_1^0$ ATLAS-CONF-2013-053

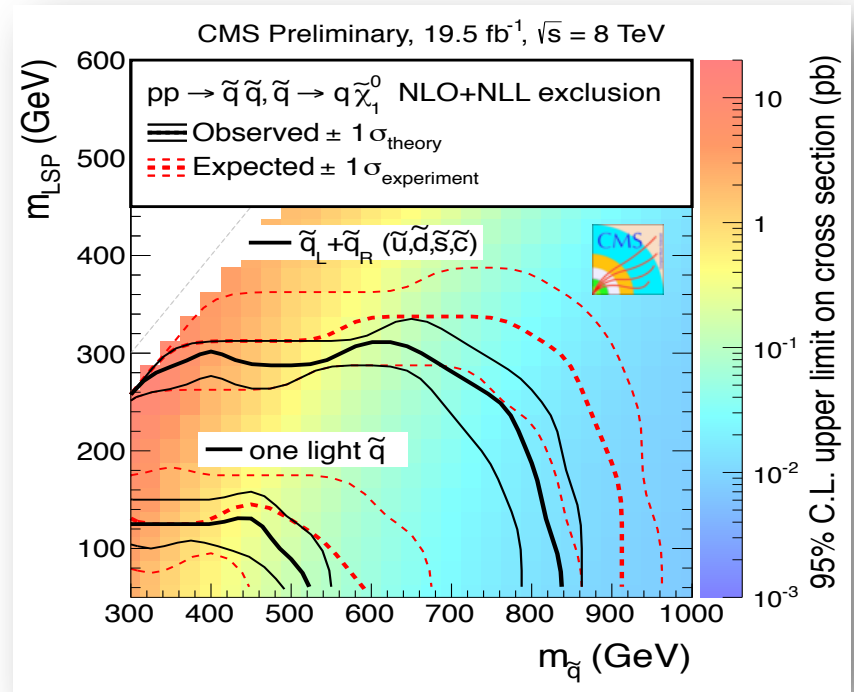
$\tilde{t} \rightarrow t\chi_1^0$ ATLAS-CONF-2013-037



CMS-SUS-PAS-13-012

Signature: Jets + E_t^{miss}

Limit assumes only one light squark (e.g. u_L) and decoupled gluino (as before).



m_{SUSY}
[GeV]

m_{LSP}
[GeV]

Direct squark
 $m_{SUSY} = m_{\tilde{q}}$

Direct squark	$\tilde{q} \rightarrow q\chi_1^0$	$\tilde{u}_L \rightarrow q\chi_1^0$	$\tilde{b} \rightarrow b\chi_1^0$	$\tilde{t} \rightarrow t\chi_1^0$
Best limit: [GeV]	~850	~500	~650	~650
No limit for M_{LSP} [GeV]	~300	~120	~270	~260

$\tilde{q} \rightarrow q\chi_1^0$ CMS-PAS-SUS-13-015

$\tilde{u}_L \rightarrow q\chi_1^0$ CMS-PAS-SUS-13-015

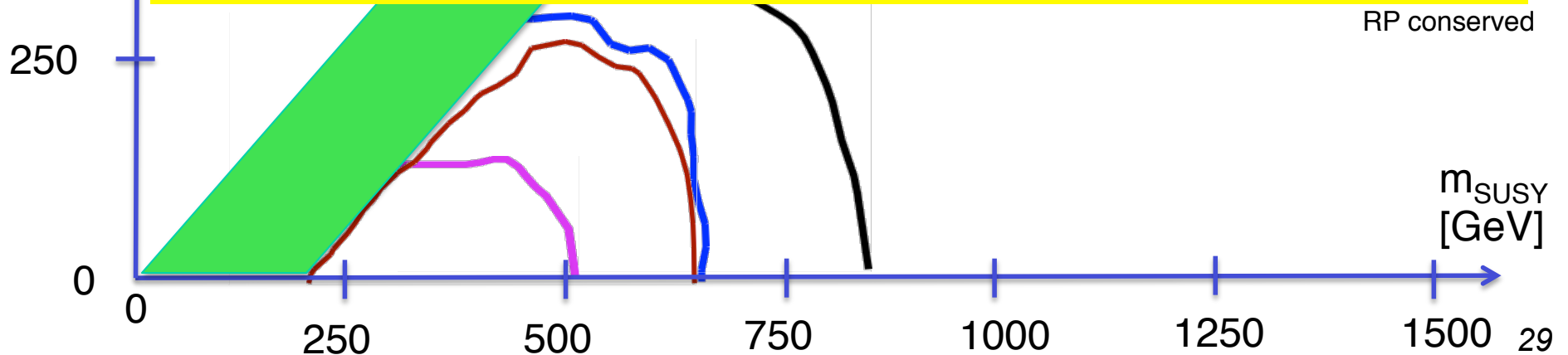
$\tilde{b} \rightarrow b\chi_1^0$ ATLAS-CONF-2013-037

$\tilde{t} \rightarrow t\chi_1^0$ ATLAS-CONF-2013-037

Direct squarks:
Better control of background estimates via b-tagging and special topology signatures like many jets per event. This allows for a higher sensitivity on direct stop and sbottom production.

1st & 2nd generation squark limits are only better than the 3rd generation when assuming **eight-fold mass degeneracy!**

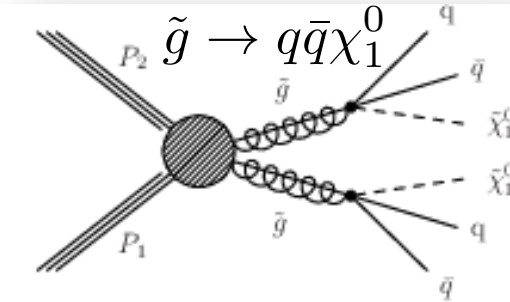
Attention: limits on single squarks are rather weak!



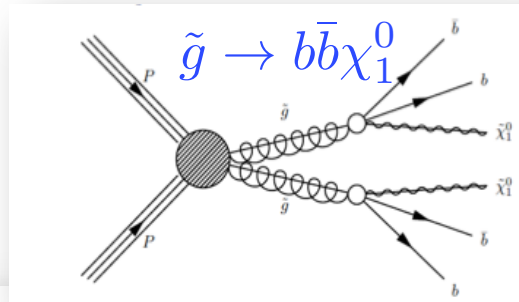
%
final
s
RP conserved

Glino mediated squark production – limits chosen

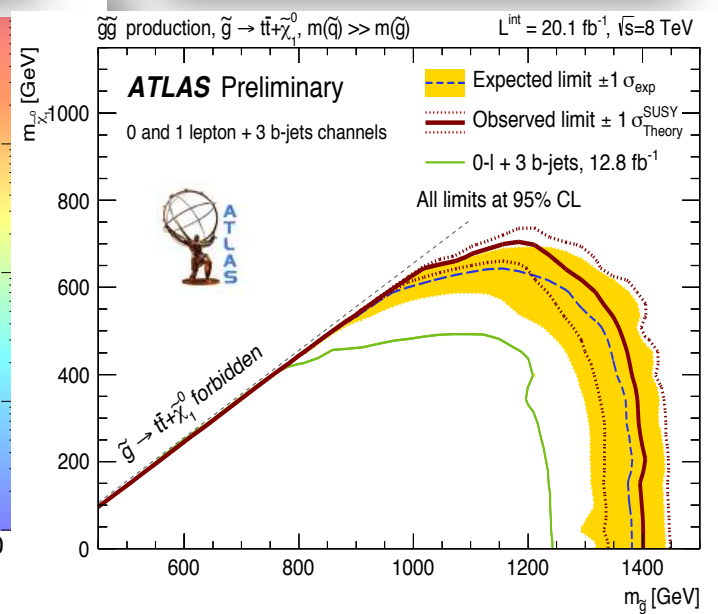
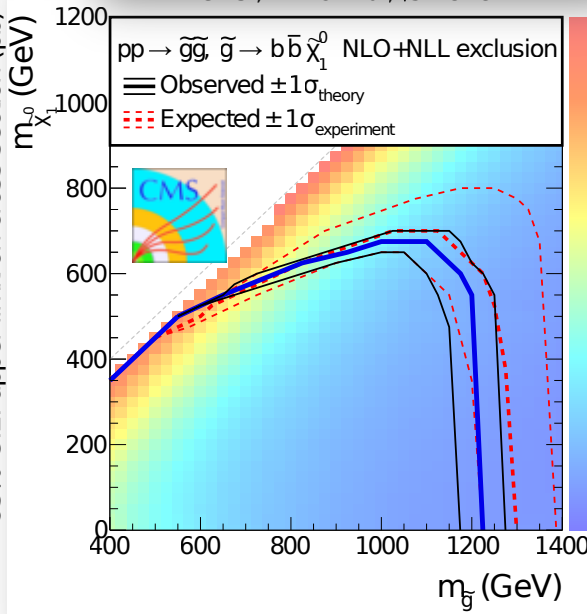
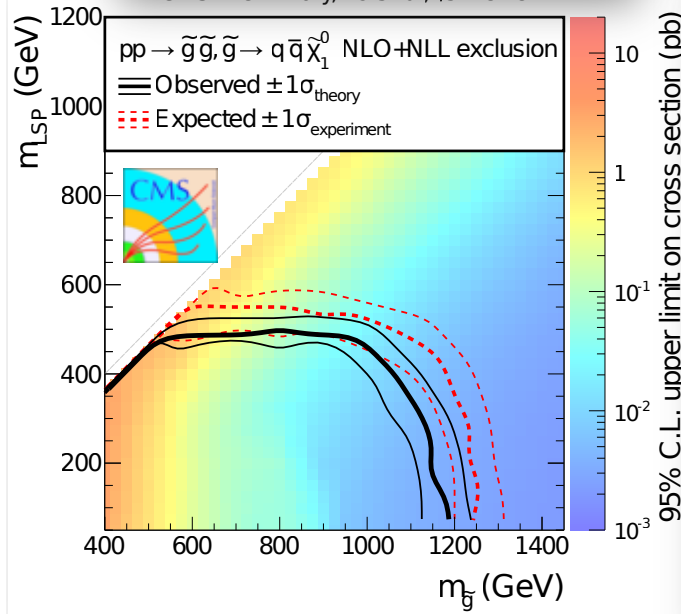
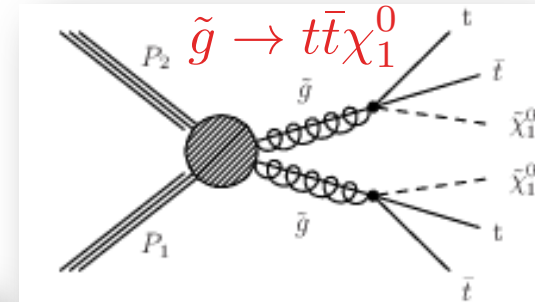
D. Buchmüller



CMS Preliminary, 19.5 fb⁻¹, $\sqrt{s} = 8$ TeV



CMS, L = 19.4 fb⁻¹, $\sqrt{s} = 8$ TeV



CMS-SUS-PAS-13-012

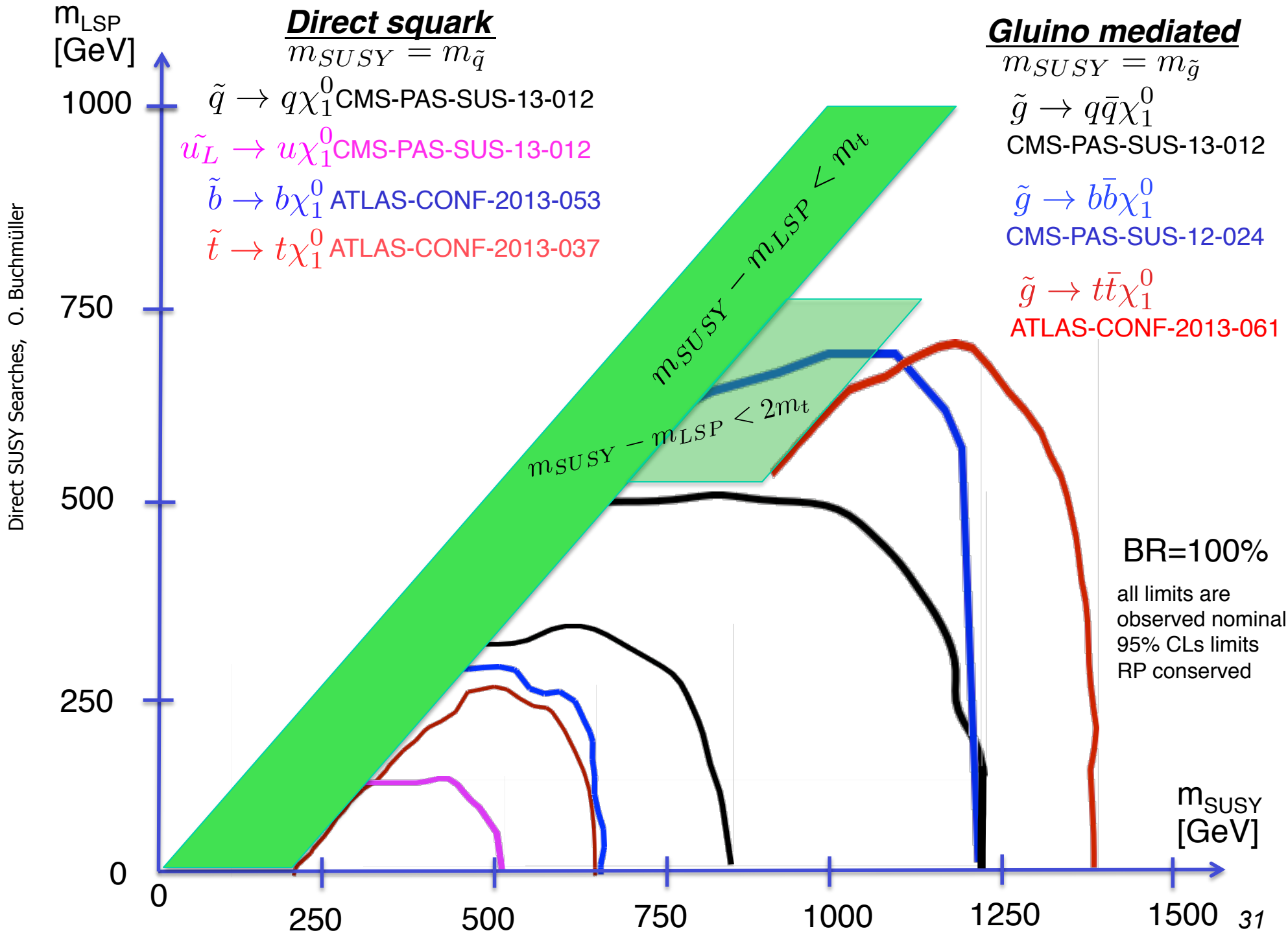
Signature: Jets + H_T + E_t^{miss}

CMS-SUS-PAS-12-024

Signature: : Jets + b-tag + E_t^{miss}

ATLAS-CONF-2013-061

Signature: 0/1 Leptons +
3 b-tag + E_t^{mis}



Gluino mediated:

Similar to direct squark production; better control of background via b-tagging and special topology signatures (e.g. many jets) provide higher sensitivity on gluino decay chains involving stops and sbottoms.

Gluino mediated

$$m_{SUSY} = m_{\tilde{g}}$$

$$\tilde{g} \rightarrow q\bar{q}\chi_1^0$$

CMS-PAS-SUS-13-012

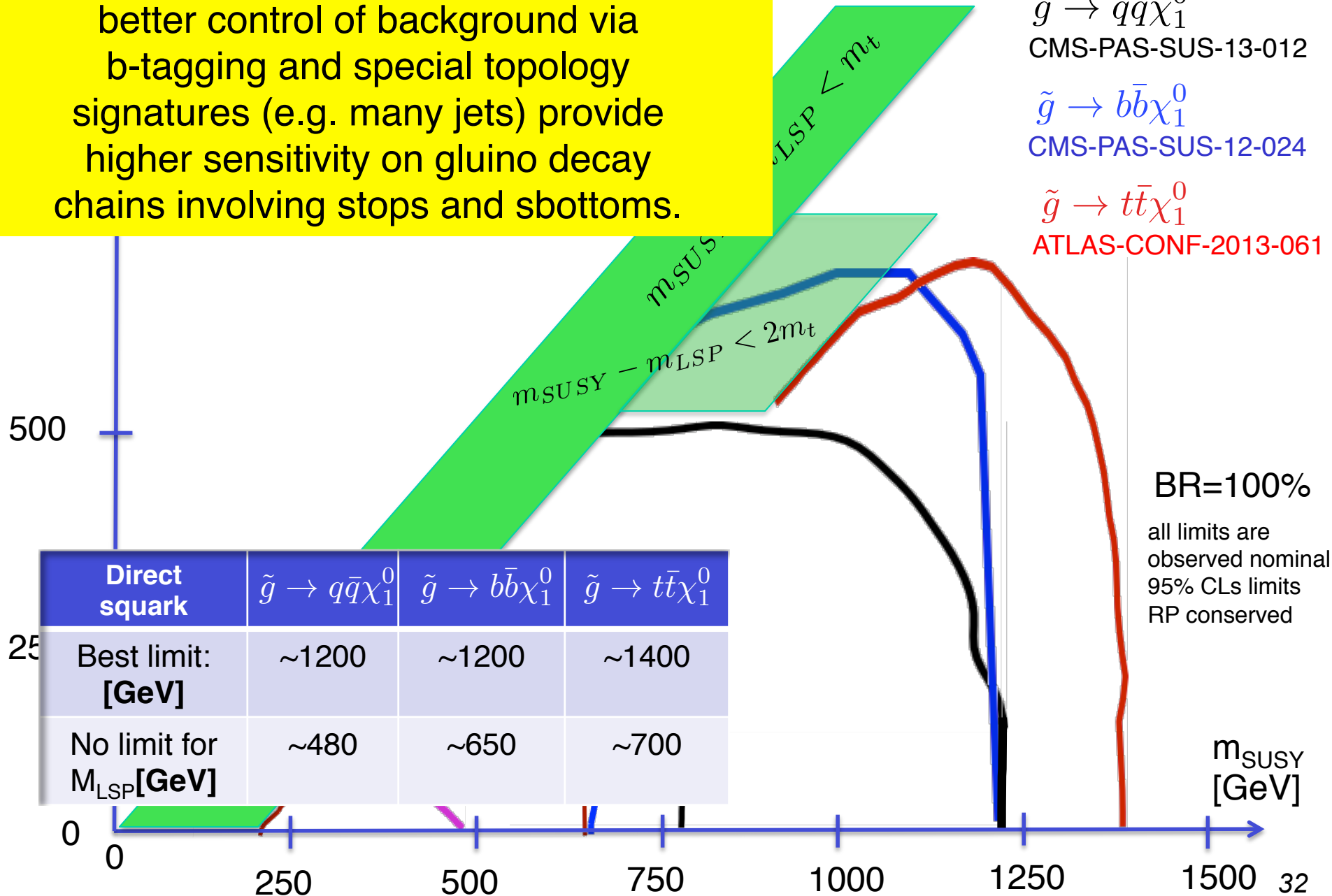
$$\tilde{g} \rightarrow b\bar{b}\chi_1^0$$

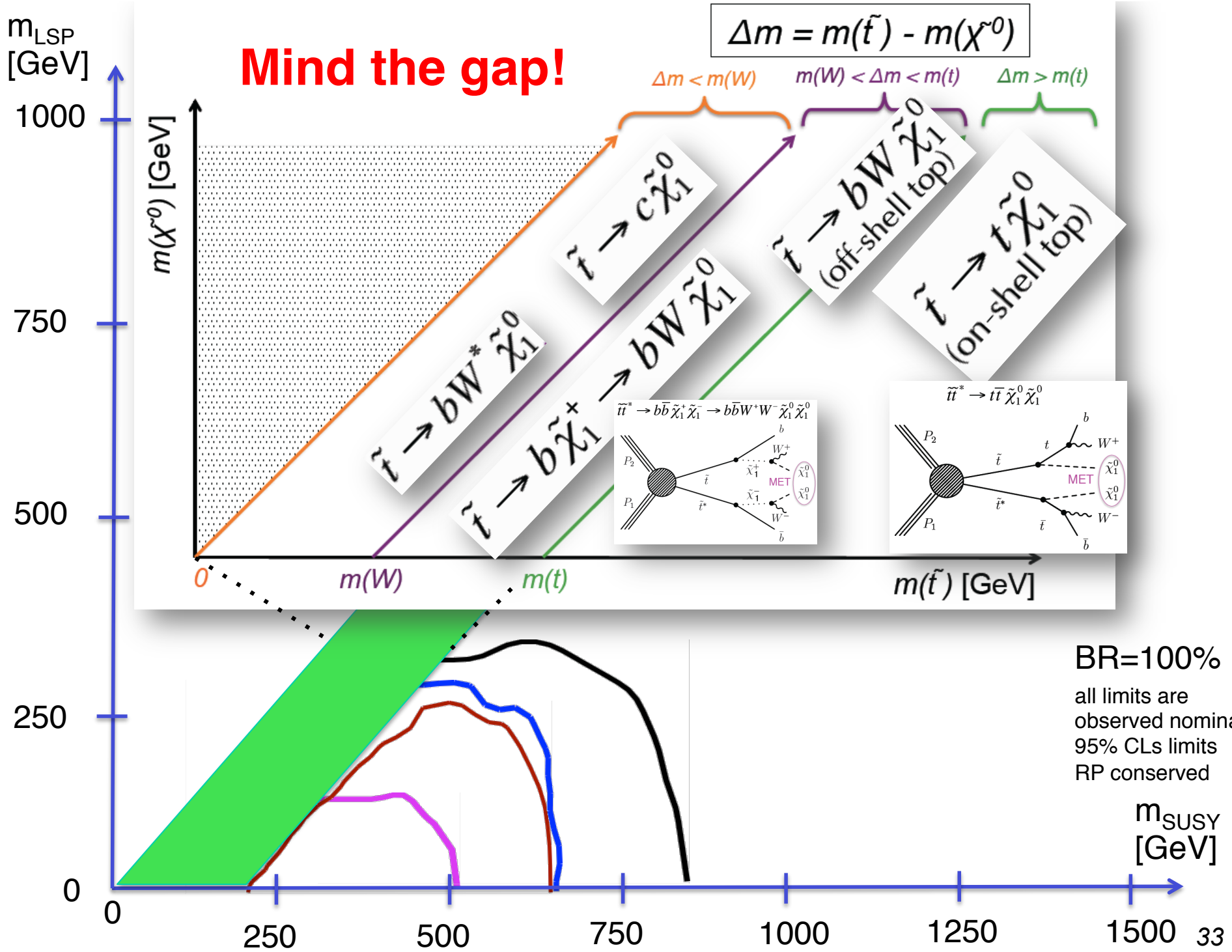
CMS-PAS-SUS-12-024

$$\tilde{g} \rightarrow t\bar{t}\chi_1^0$$

ATLAS-CONF-2013-061

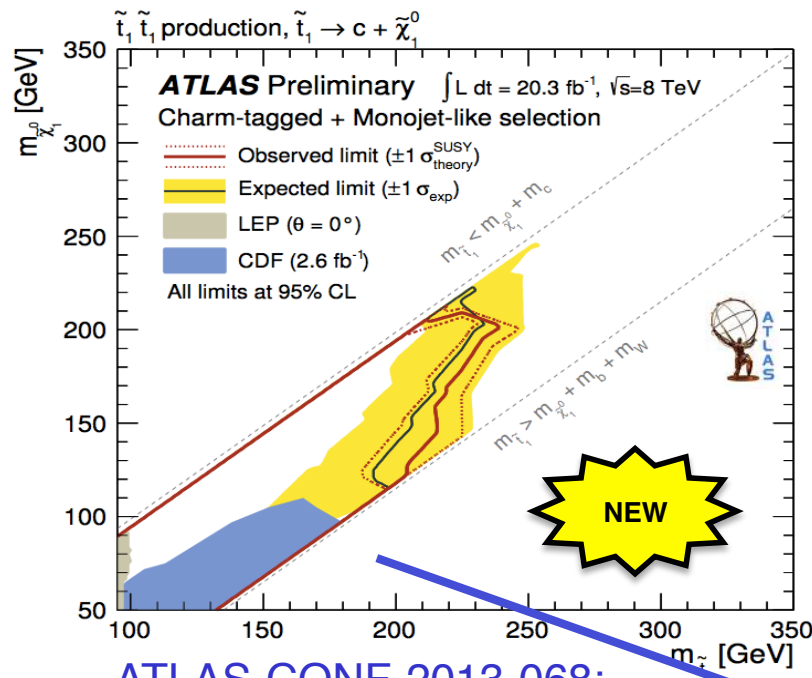
Direct SUSY Searches, C. Buchmüller



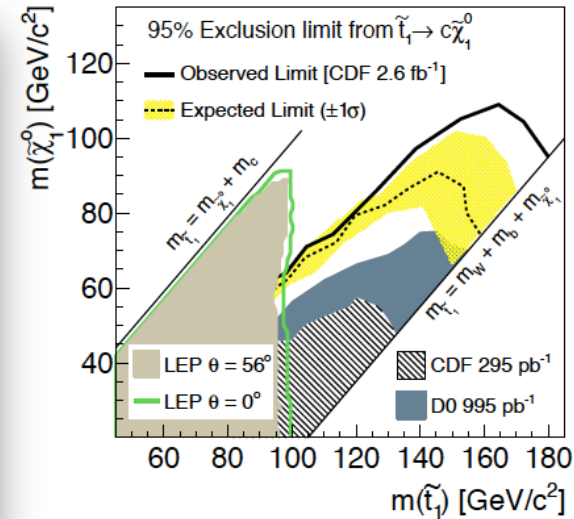


Dedicated searches for direct stop-pair production

Direct SUSY Searches, O. Buchmüller



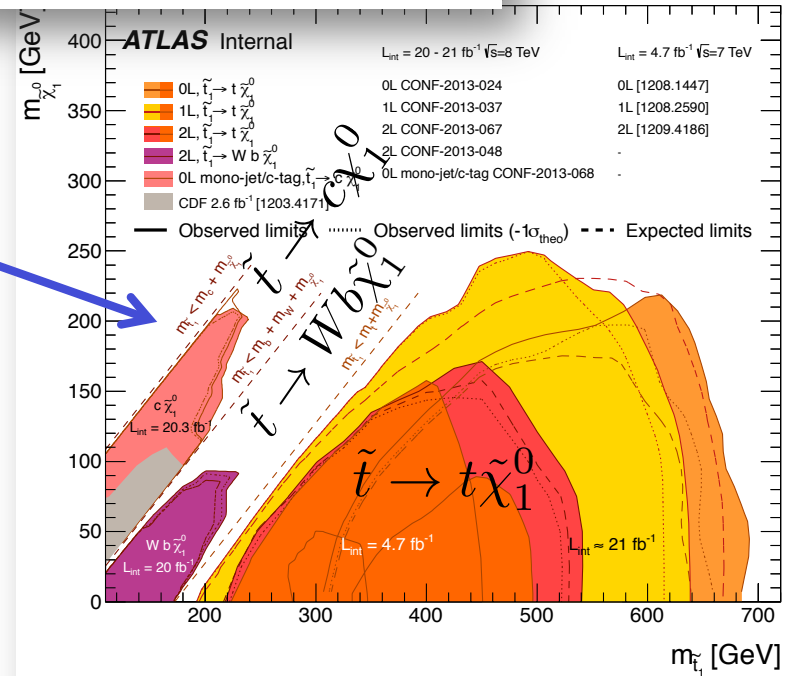
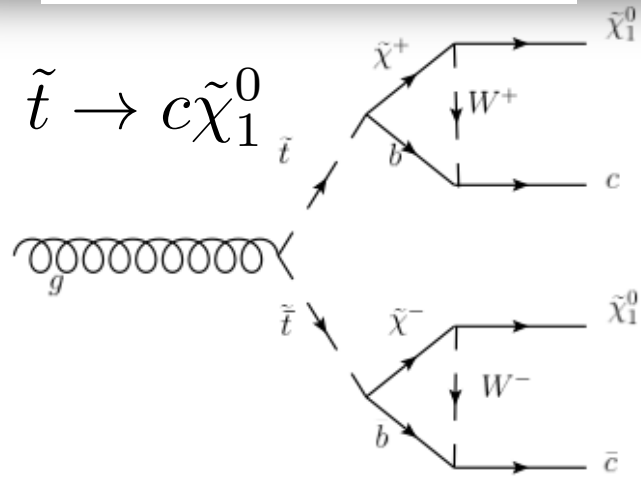
ATLAS-CONF-2013-068:



LEP:
LEPSUSY
WG/04-01.1

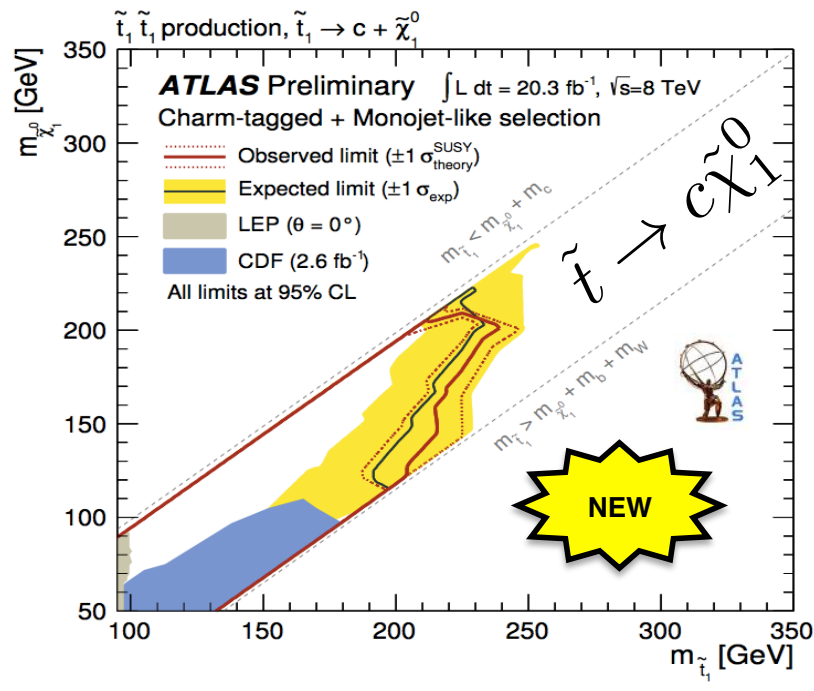
CDF:
1203.4171
D0:
0803.2263

Status: July 2013



Stop decay to charm and neutralino

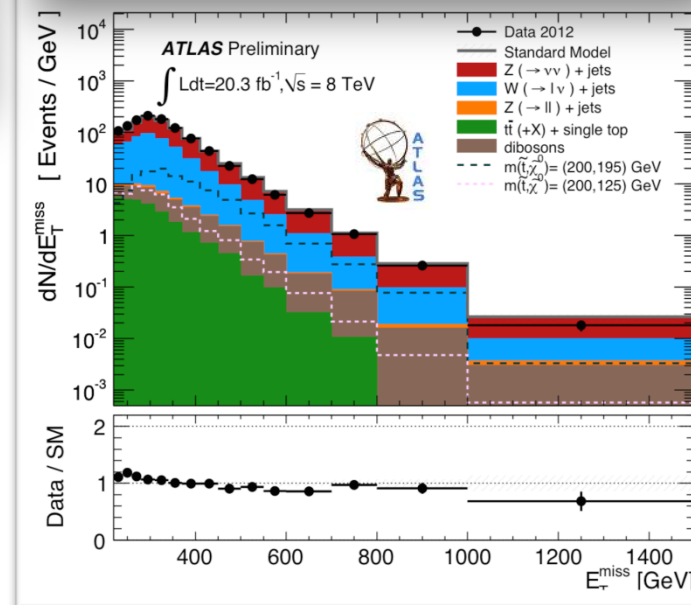
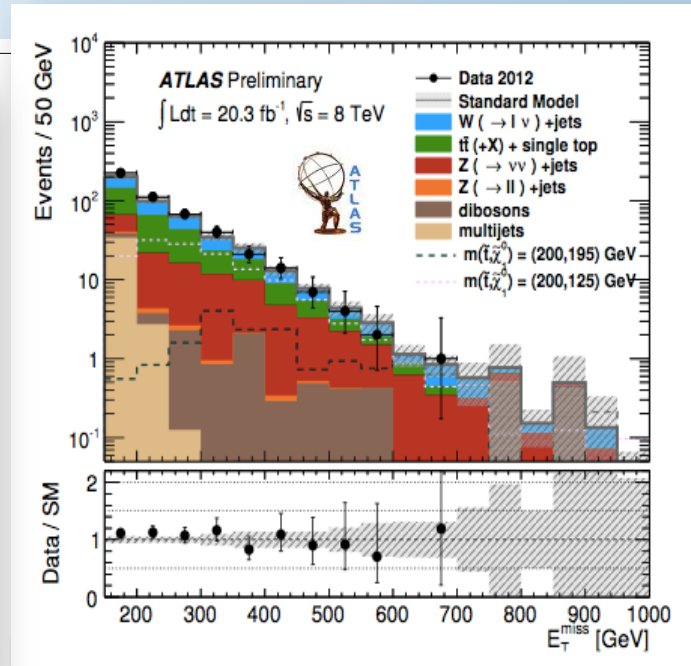
Direct SUSY Searches, O. Buchmüller



ATLAS-CONF-2013-068:

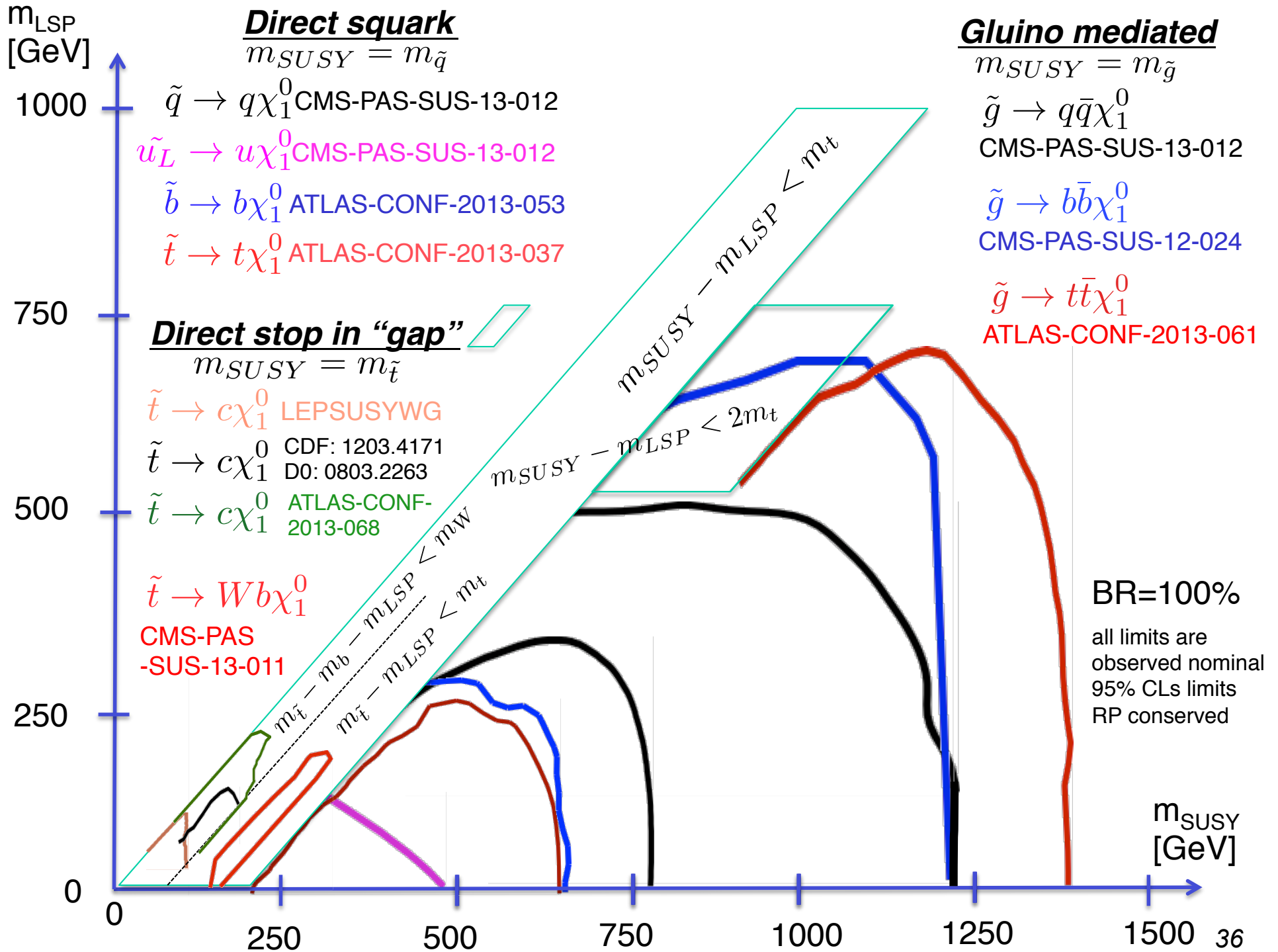
Two different selections:

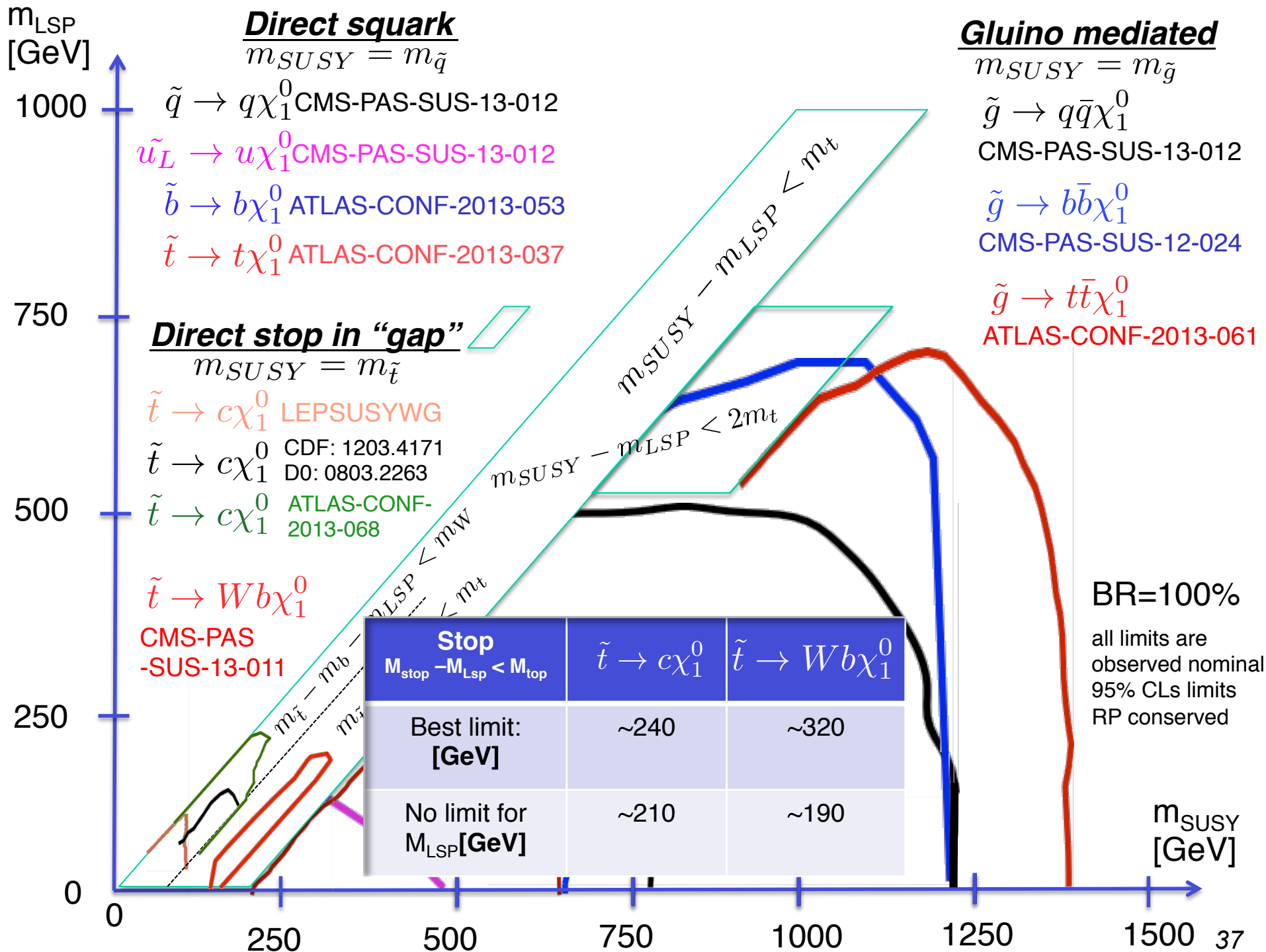
- Monojet-like selection to cover region close to ‘diagonal’
- MVA based c-tag selection for remaining region

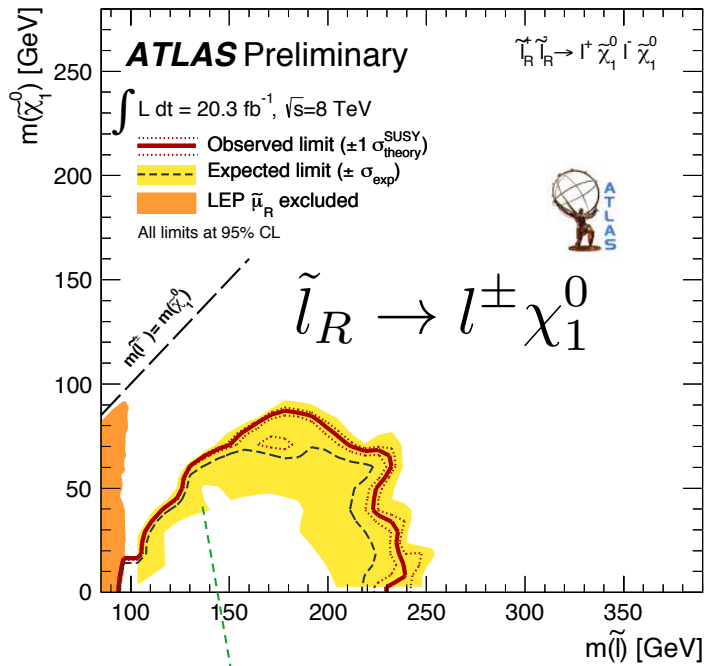


c-Tag:
 95% excl.
 visible XS:
 $\epsilon \sigma = 0.7 \text{ fb}$
 signal obs:
 13 events
 signal exp:
 14^{+5}_{-4} events
 CL_B :
 0.45

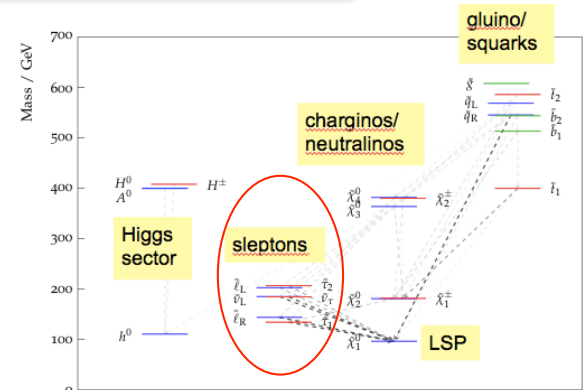
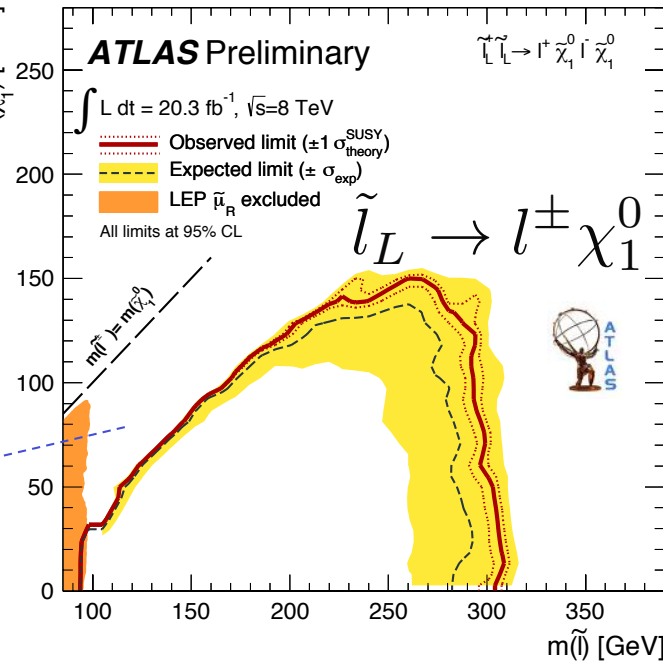
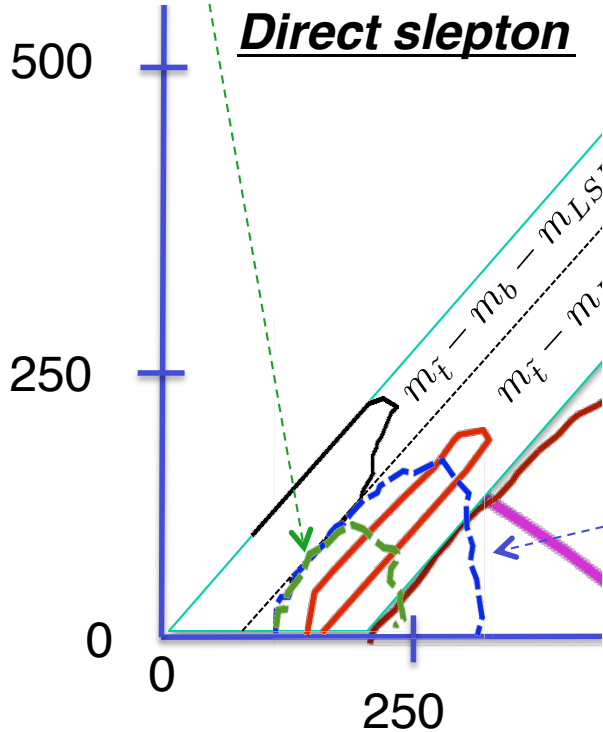
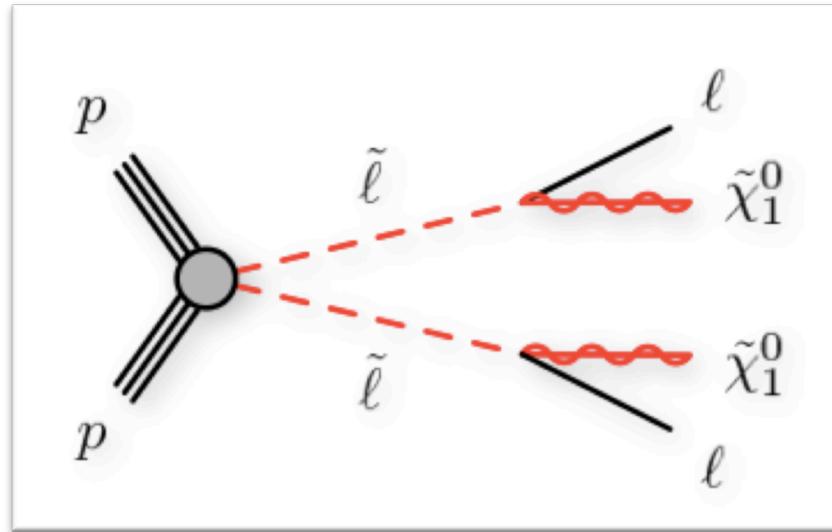
Monojet
 95% excl.
 visible XS:
 $\epsilon \sigma = 136 \text{ fb}$
 signal obs:
 2770 events
 signal exp:
 2060^{+780}_{-560} events
 CL_B :
 0.86



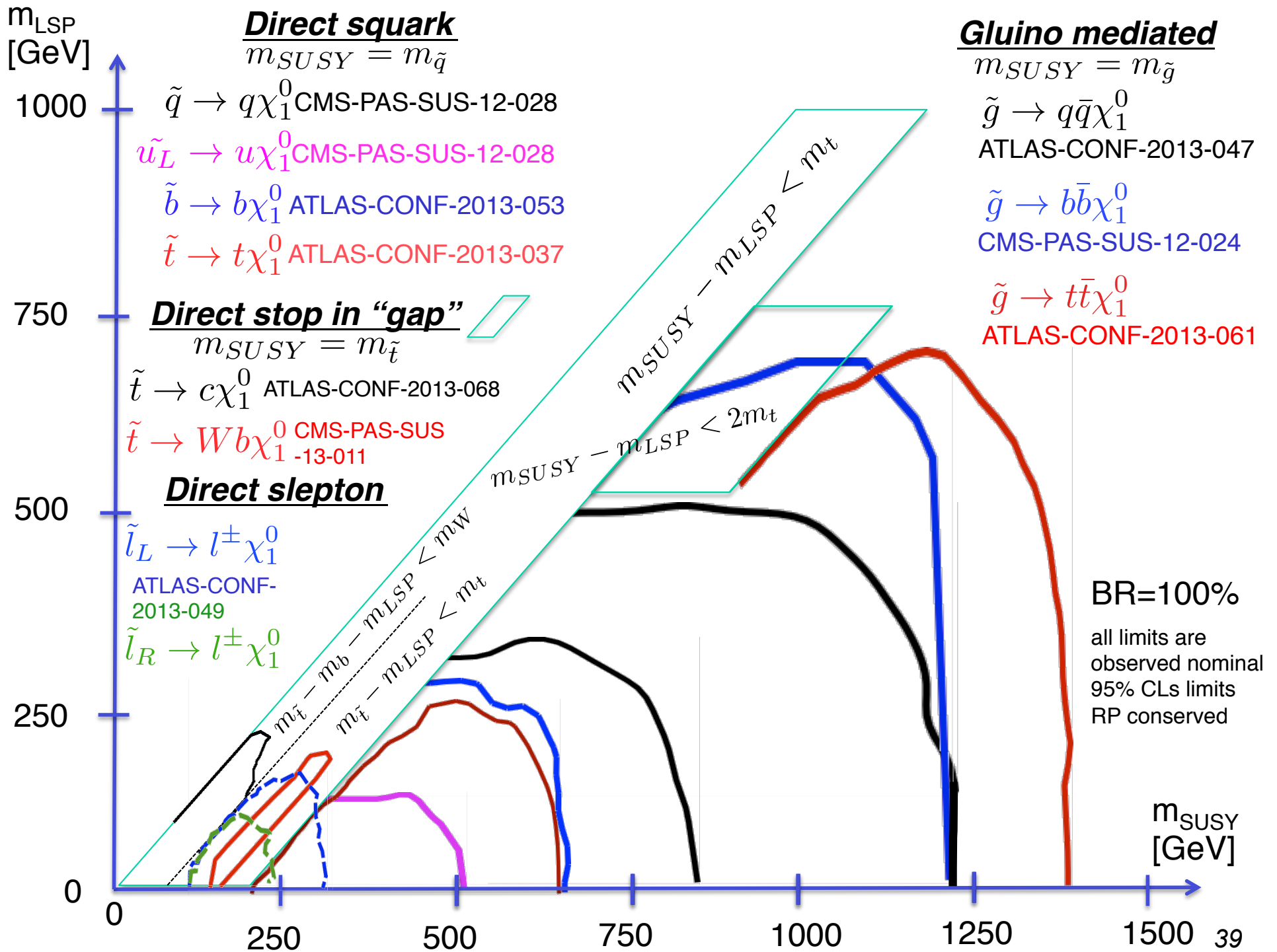


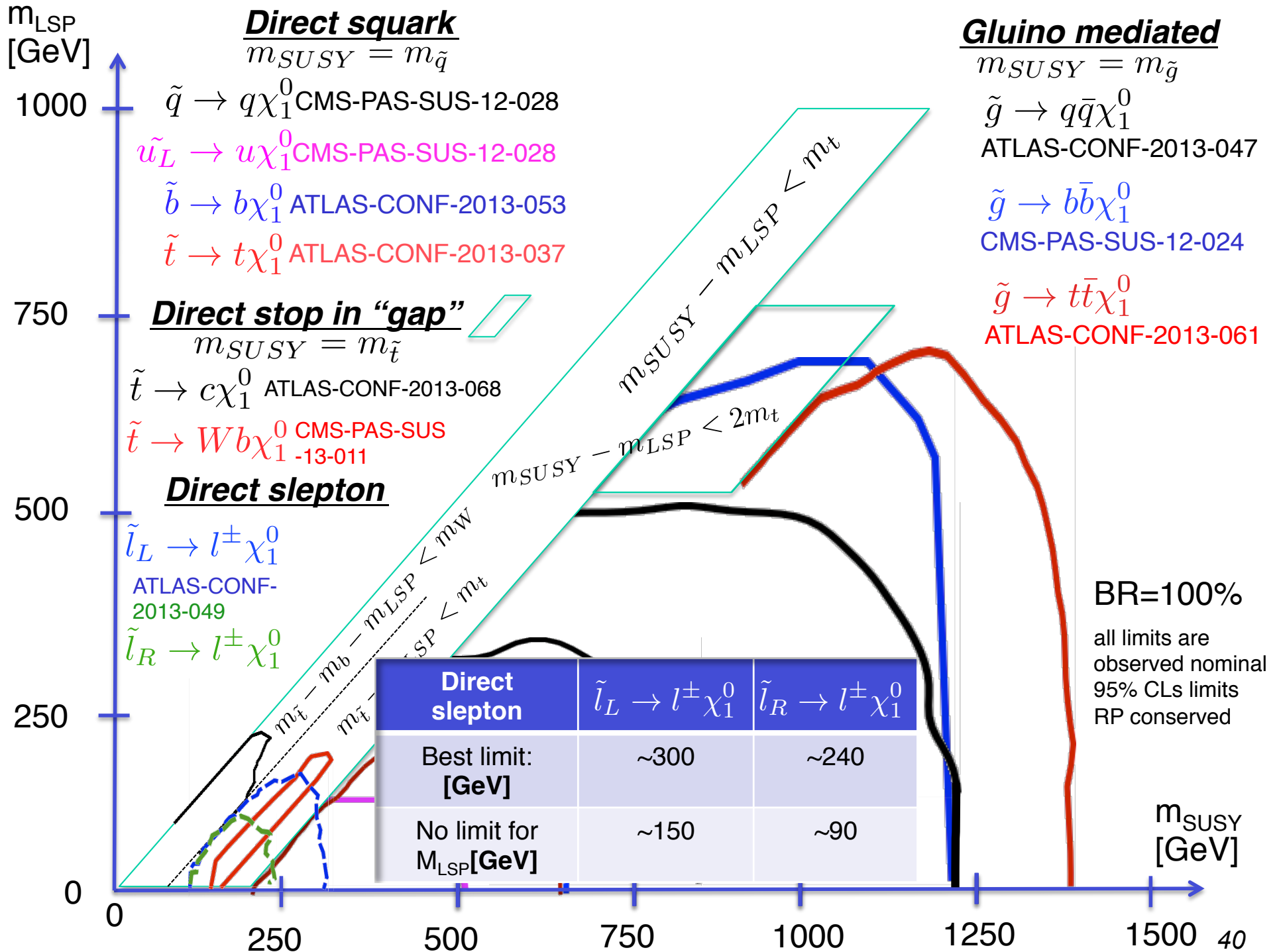


Direct slepton production

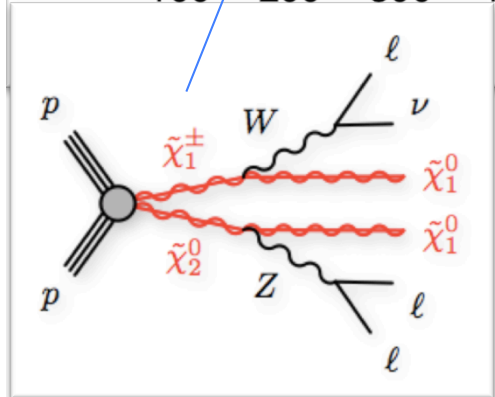
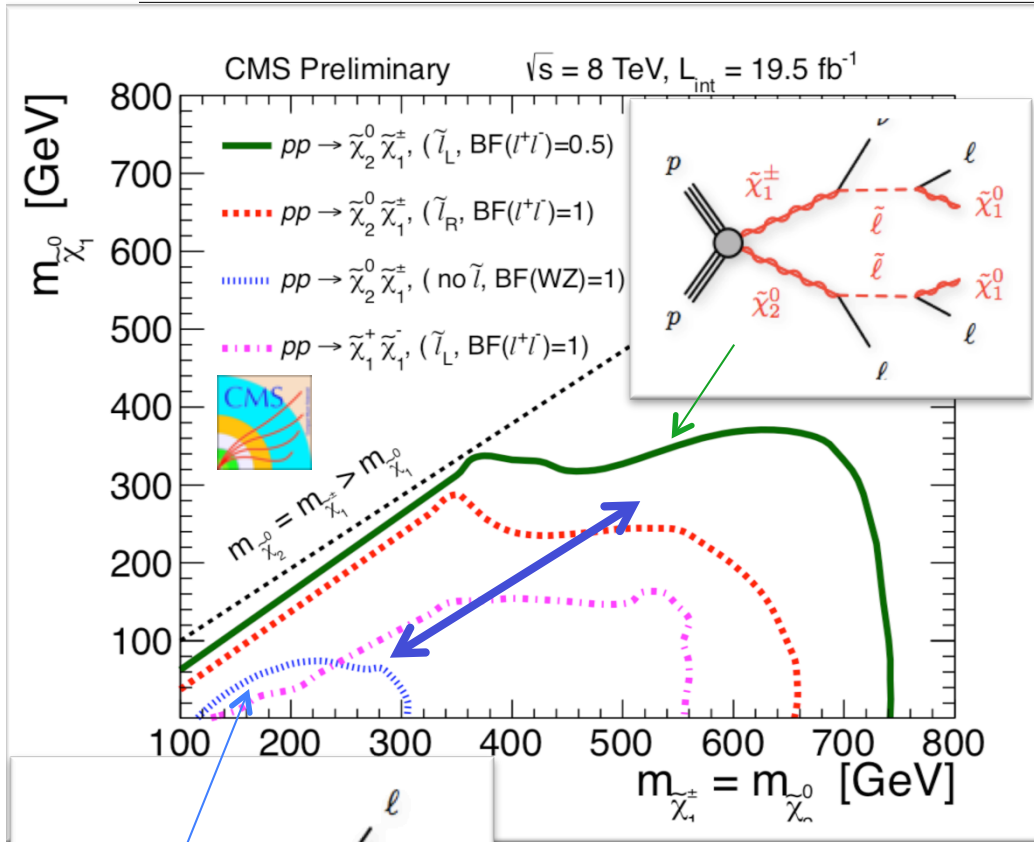


ATLAS-CONF-2013-049
Signature
2 lepton + E_T^{miss}





Direct chargino/neutralino production



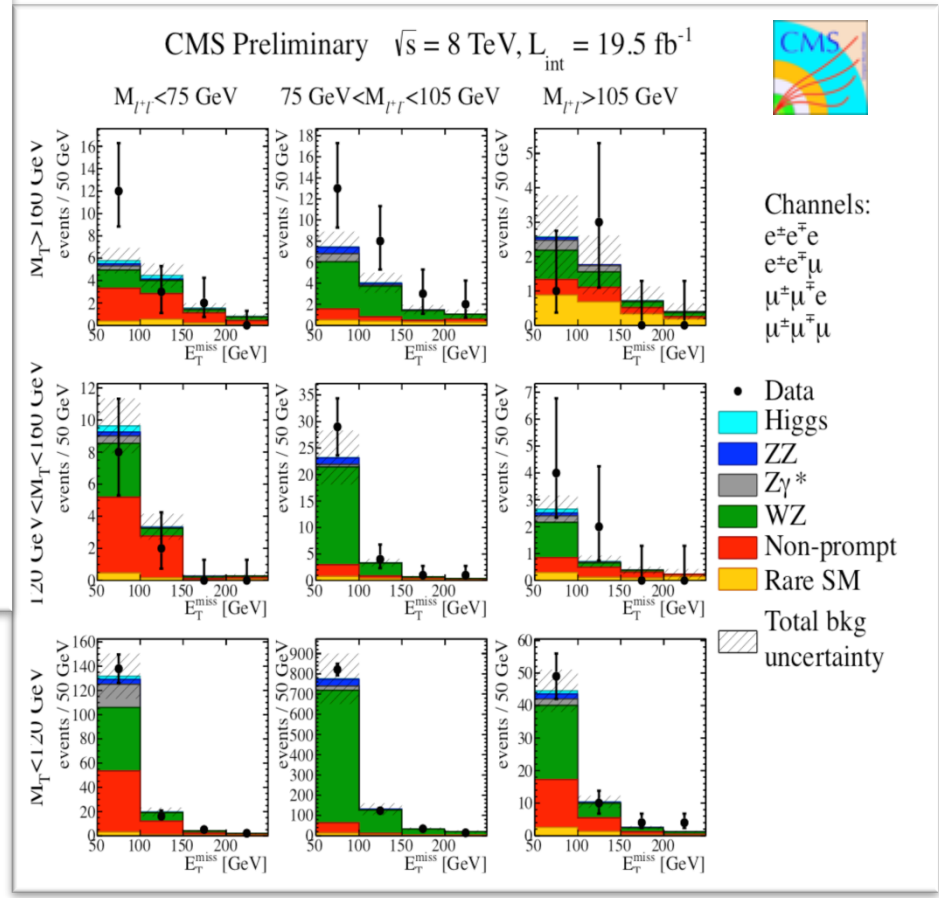
heavy slepton “hard(er)”

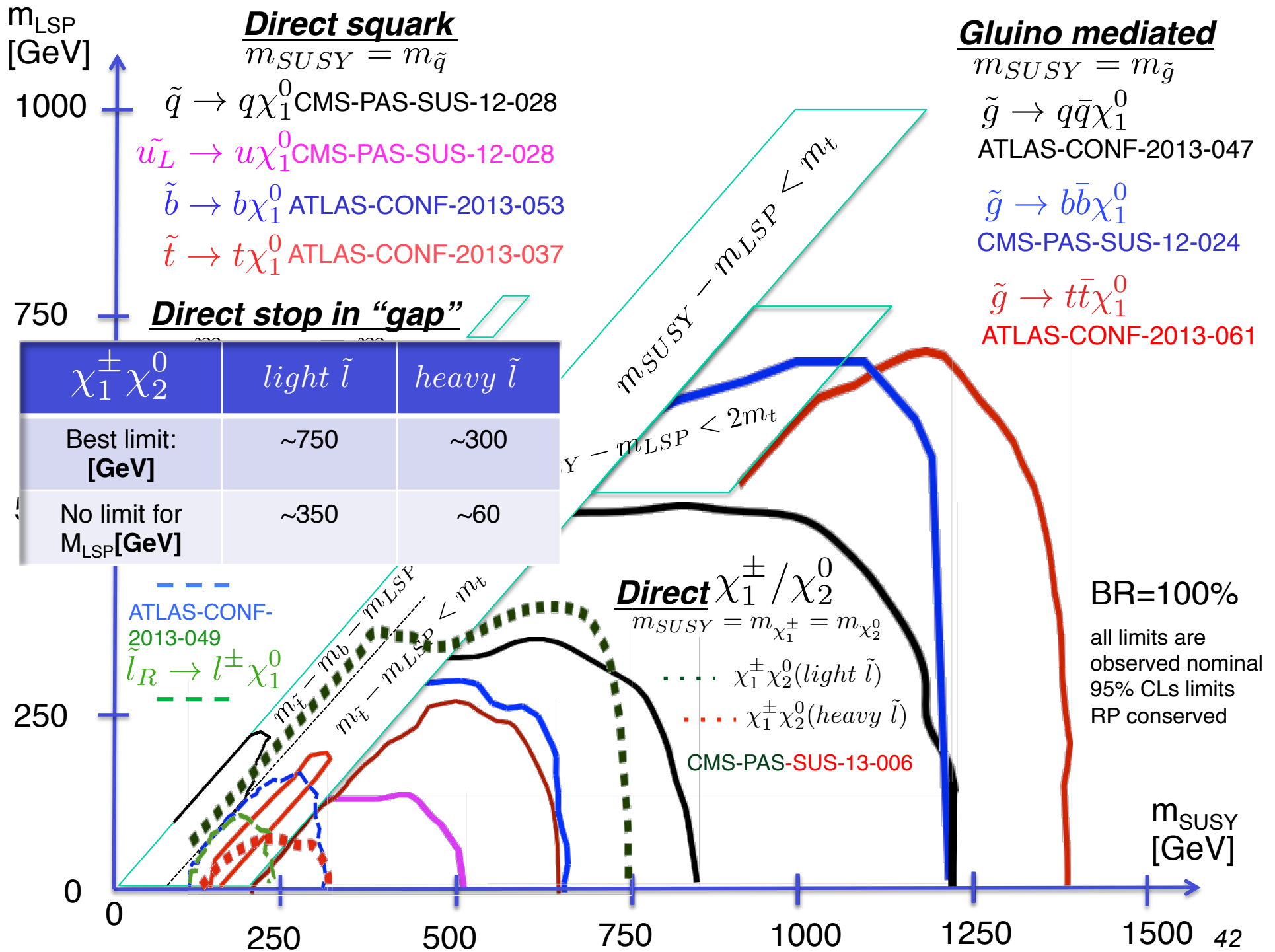
Add $Z(l^+l^-)+2\text{jets}$ topology in bins of E_t^{miss} to increase sensitivity for “heavy” slepton case

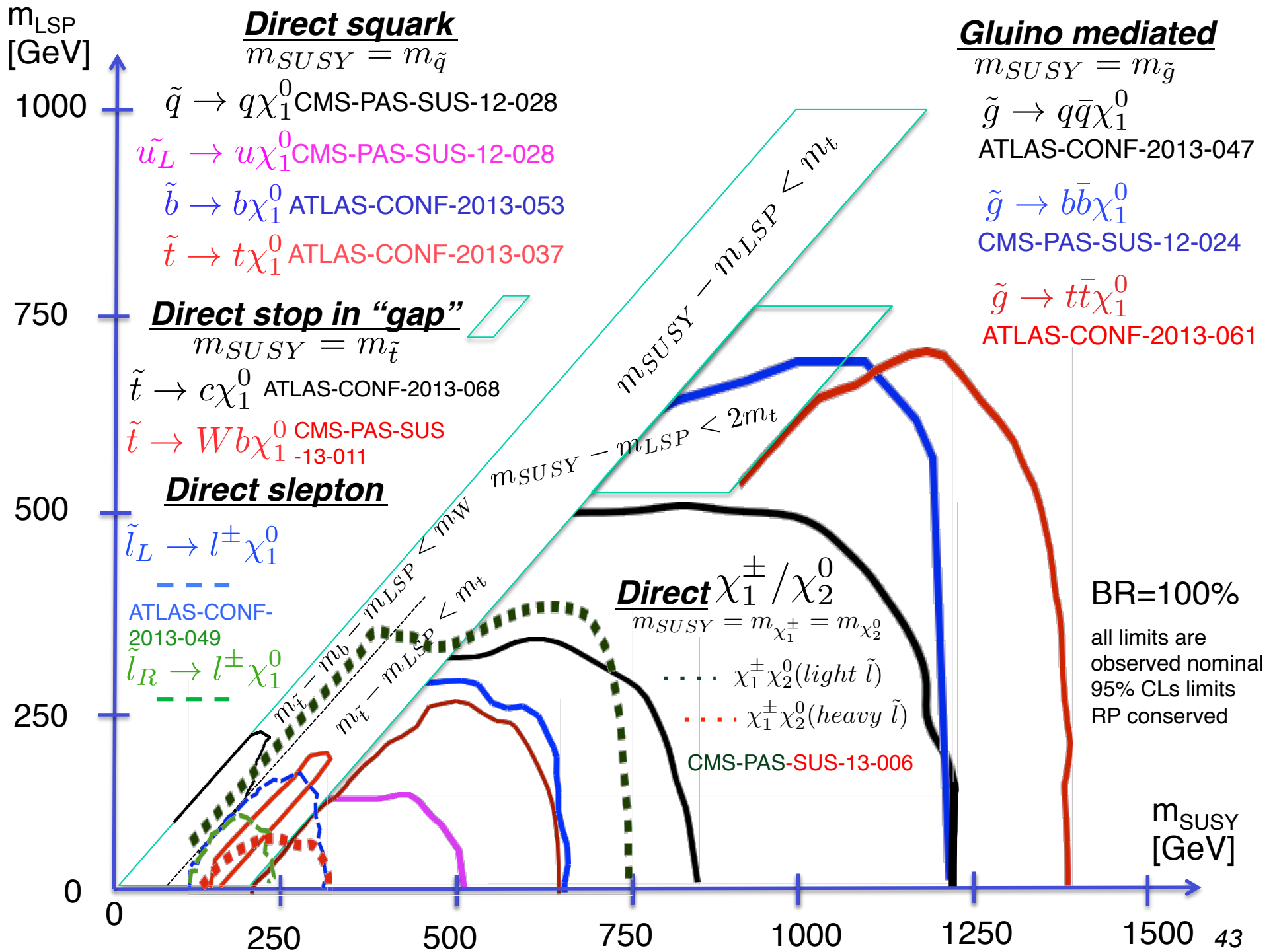
light slepton “easy”

CMS-PAS-SUS-13-006

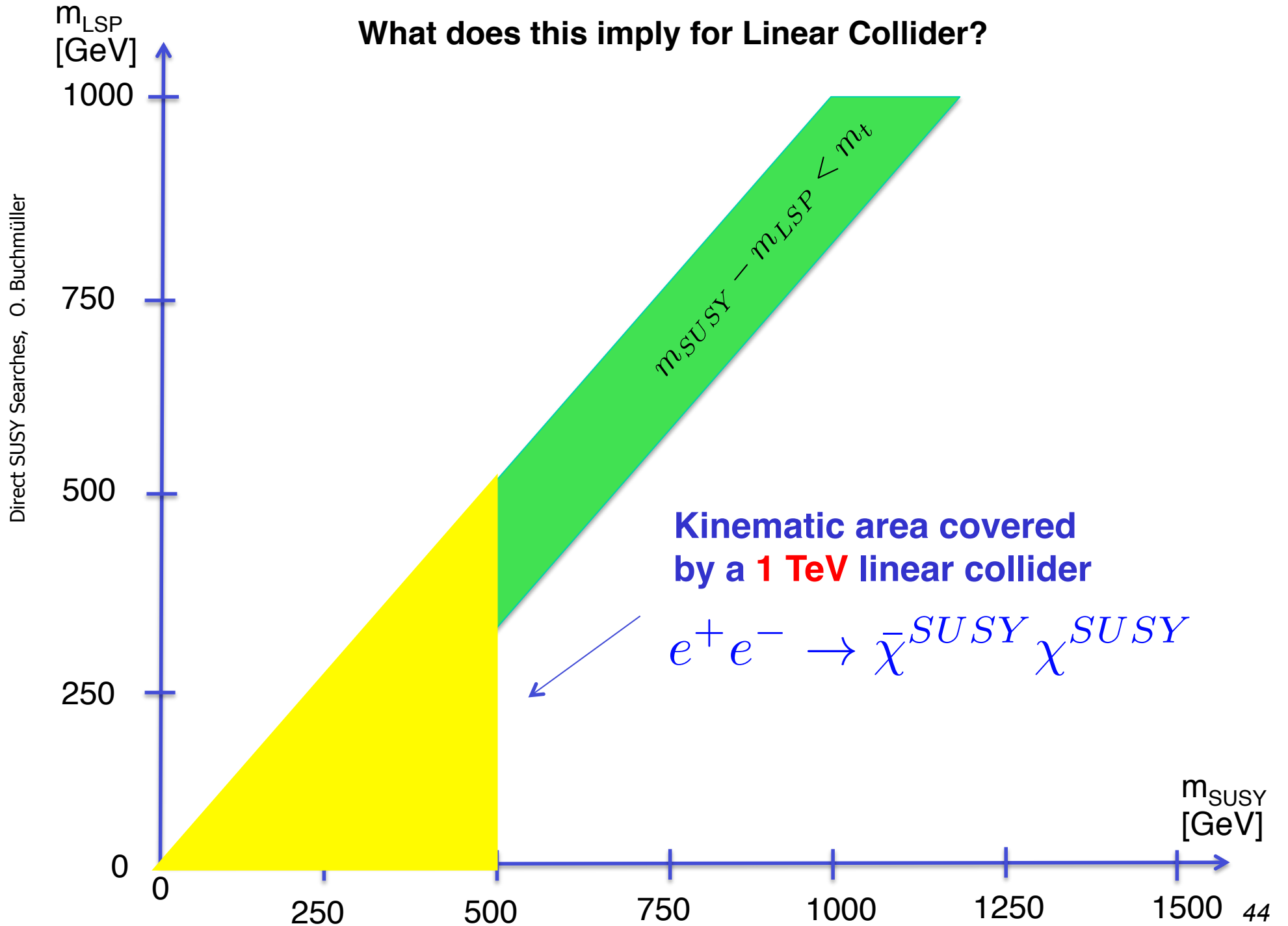
3l OSSF categorised in bins of $M_{ll}, MT,$ and ET^{mis}

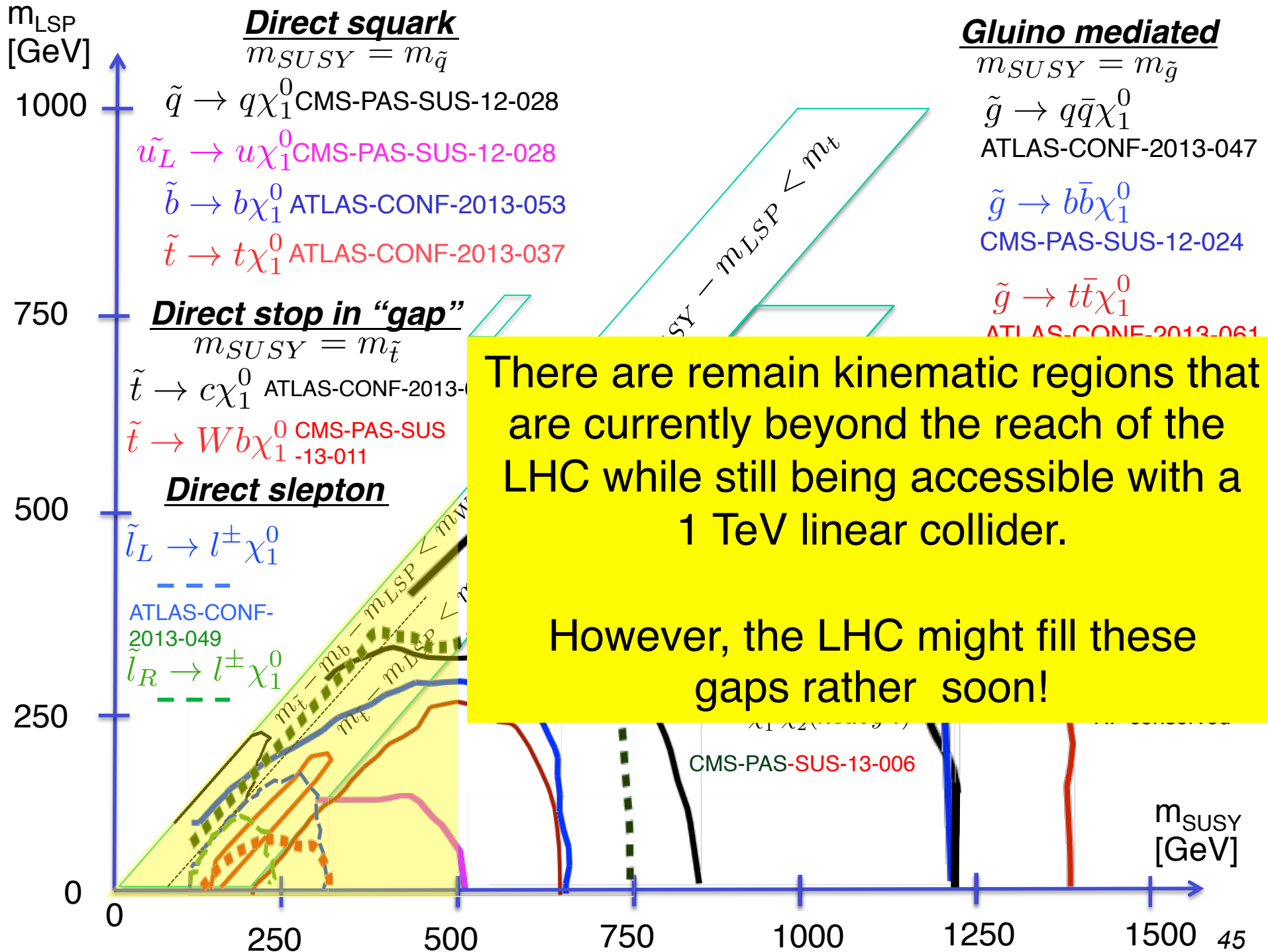






What does this imply for Linear Collider?

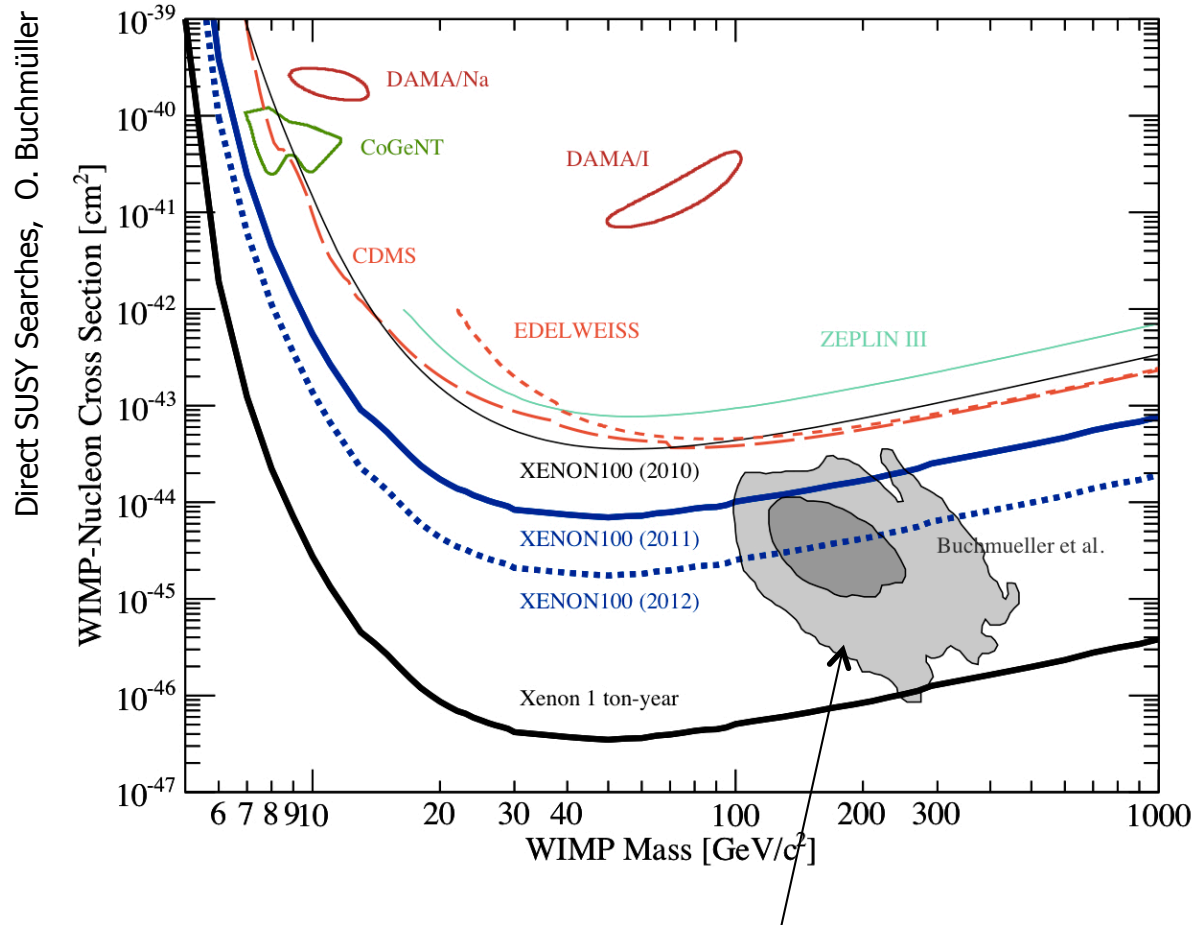




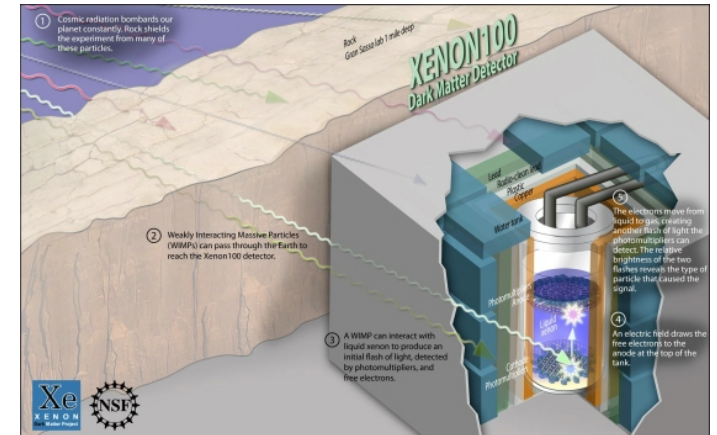
There are remain kinematic regions that are currently beyond the reach of the LHC while still being accessible with a 1 TeV linear collider.

However, the LHC might fill these gaps rather soon!

Direct Dark Matter Detection Experiments



Expectation from global
SUSY fits in the CMSSM



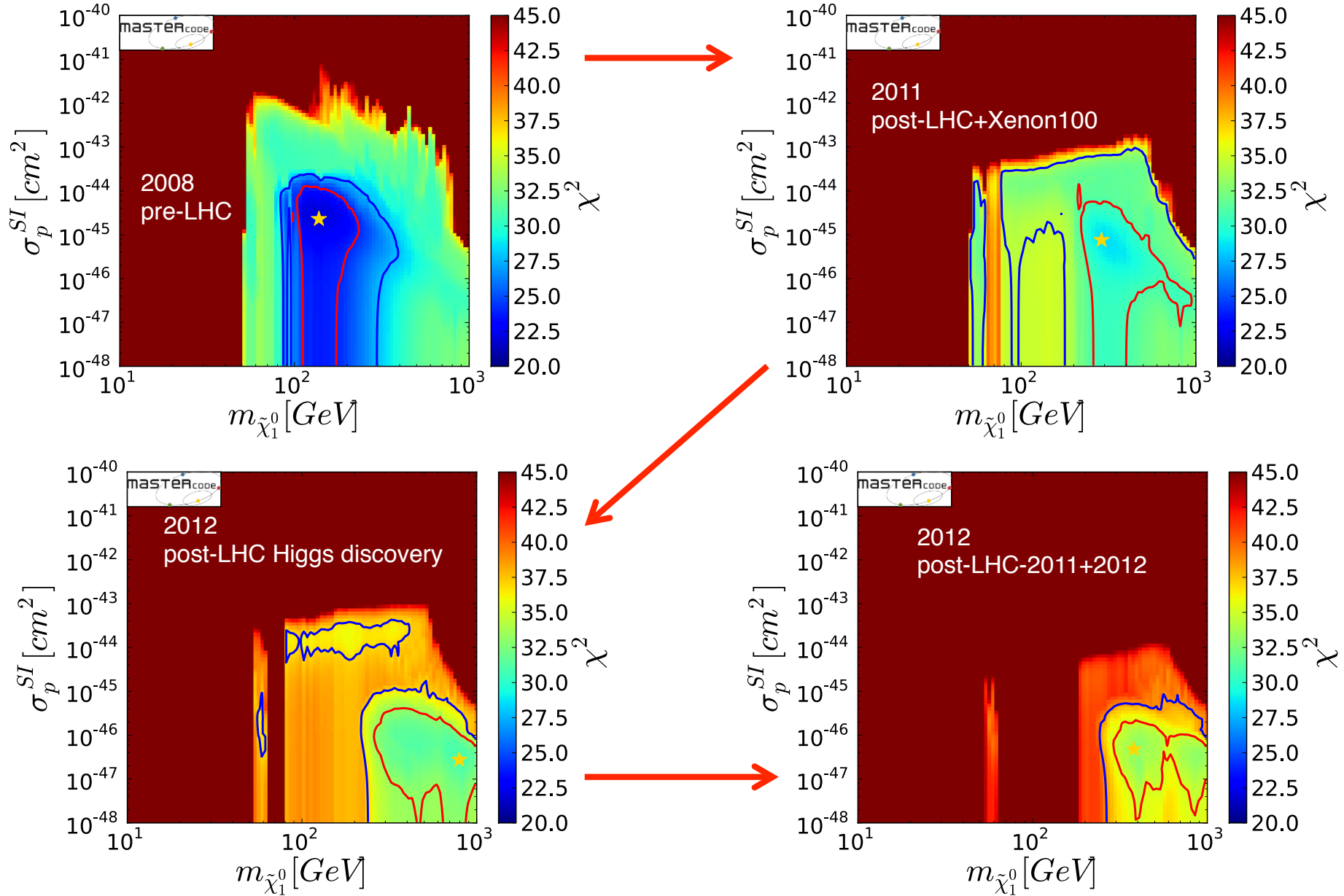
Example:
The XENON100 experiment is
located deep underground at the
Gran Sasso National Laboratory in Italy.

62 kg liquid Xenon target
100 days of data taking
1.8±0.6 events expected
3 events observed

⇒ **Exclude $7.0 \times 10^{-45} \text{ cm}^2$**
for a $M_{\text{WIMP}} = 50 \text{ GeV}$ at 90% CL.

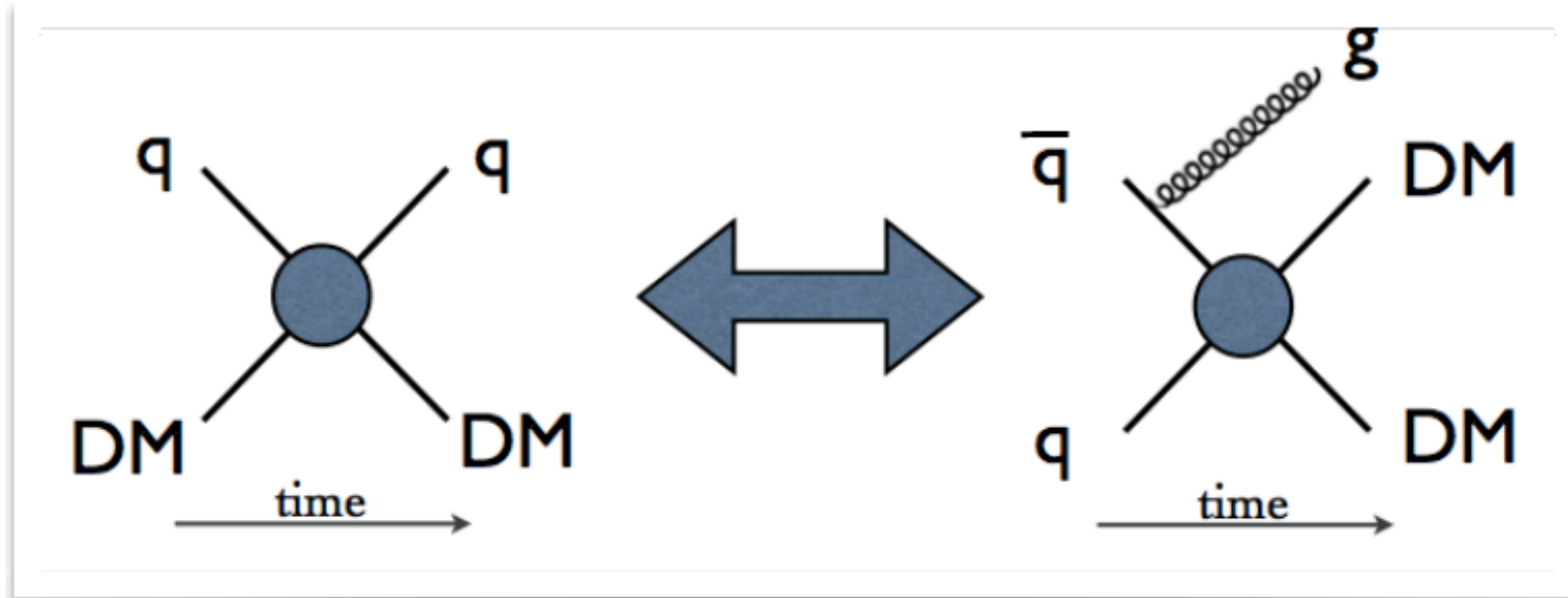
SUSY & Dark Matter: Evolution with time

Direct SUSY Searches, O. Buchmüller



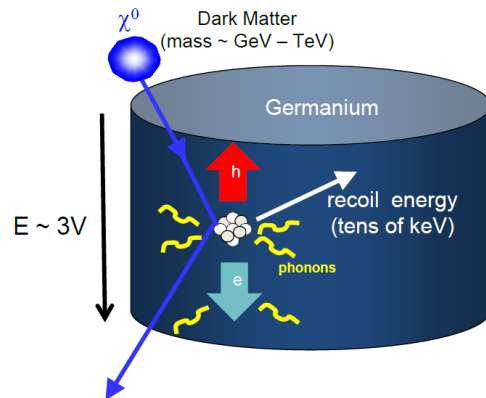
Dark Matter Searches: Direct Detection vs Colliders

Direct SUSY Searches, O. Buchmüller



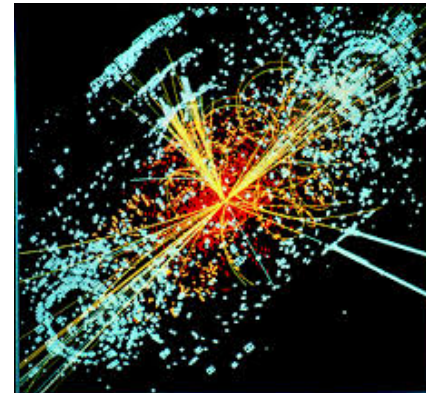
Direct Detection Experiments

- DM-nucleus scattering



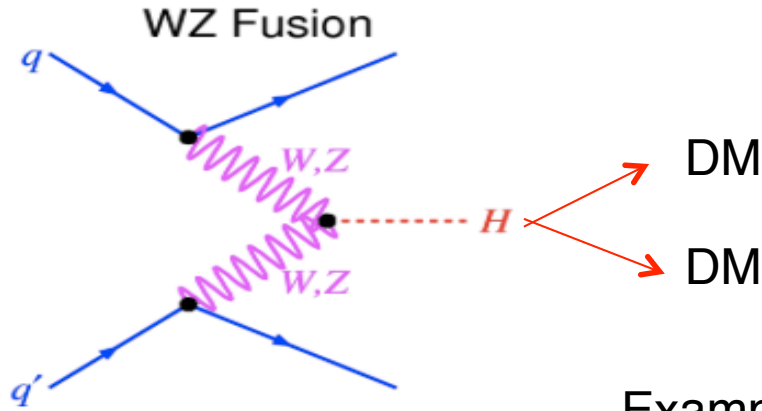
Collider Experiments

- Pair-production of DM
- missing energy signature

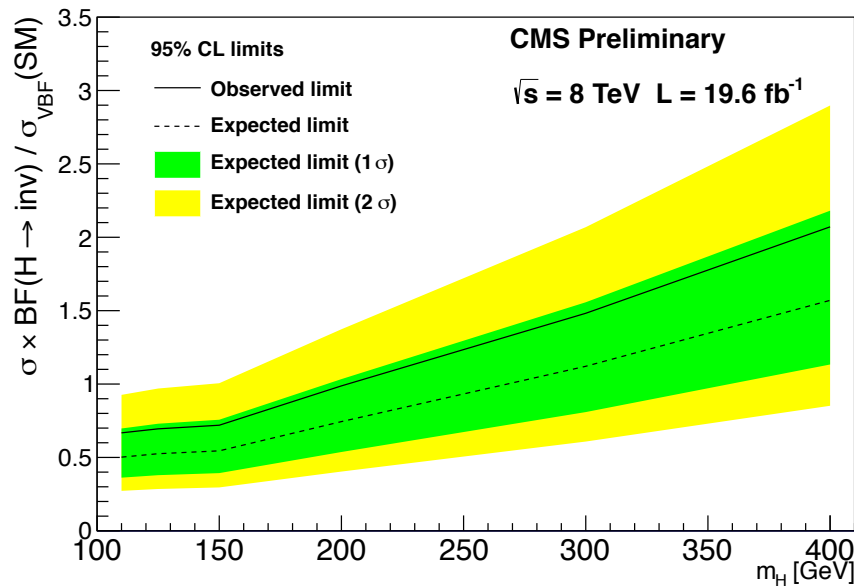
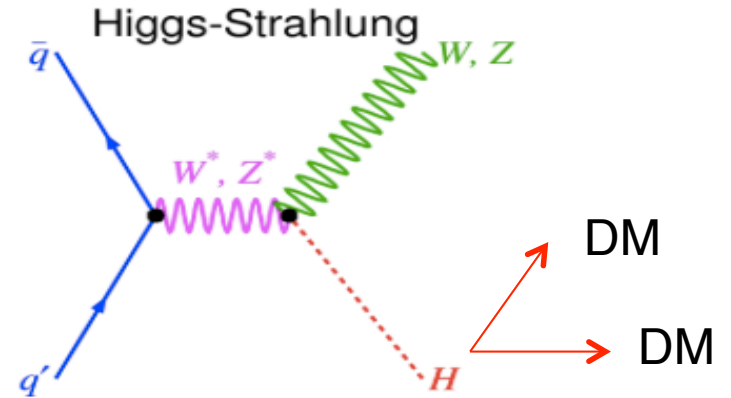


Dark Matter from invisible Higgs searches

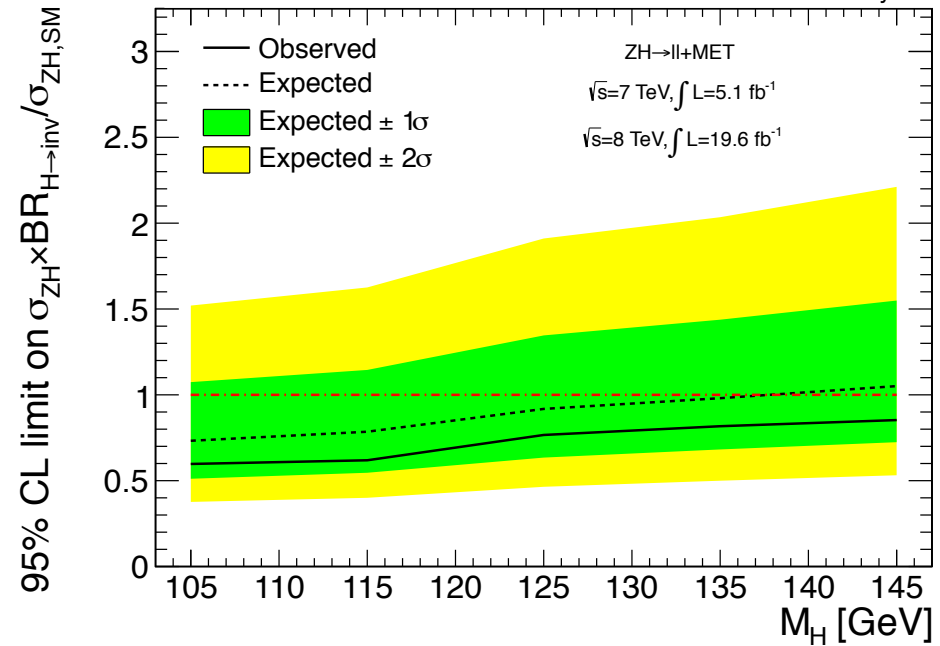
Direct SUSY Searches, O. Buchmüller



Example CMS



$\text{BR}(H \rightarrow \text{invisible}) < 68\% \text{ @ } 95\% \text{CL}$



$\text{BR}(H \rightarrow \text{invisible}) < 75\% \text{ @ } 95\% \text{CL}$

Dark Matter from invisible Higgs searches

Status 2012 CMS only:

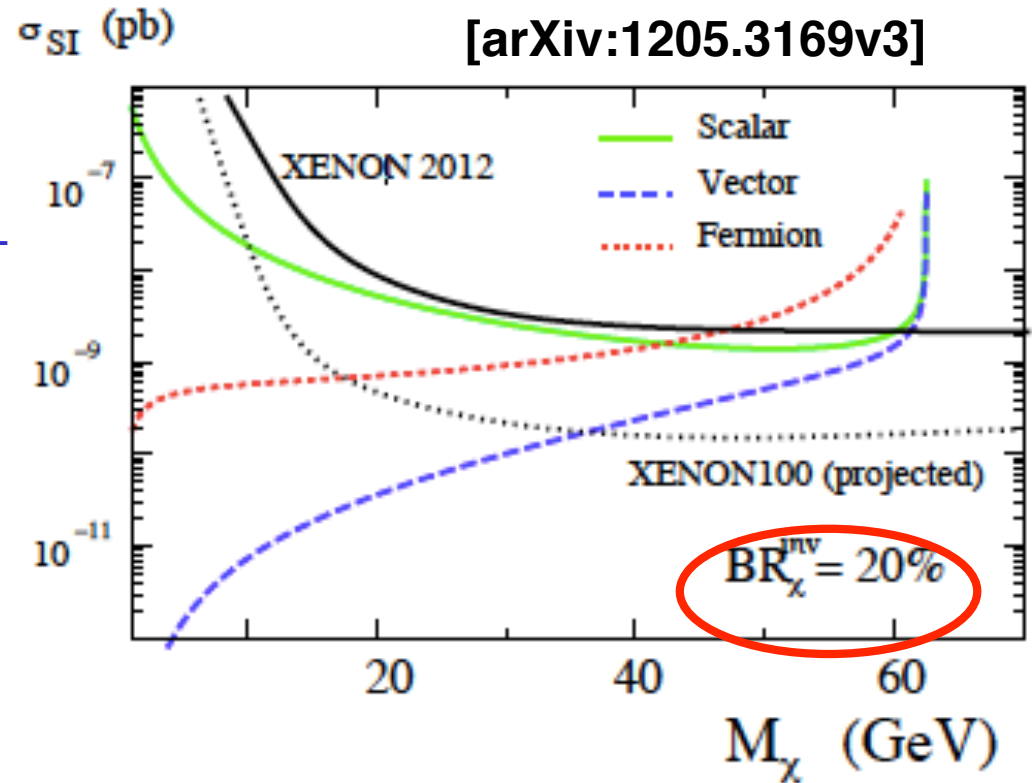
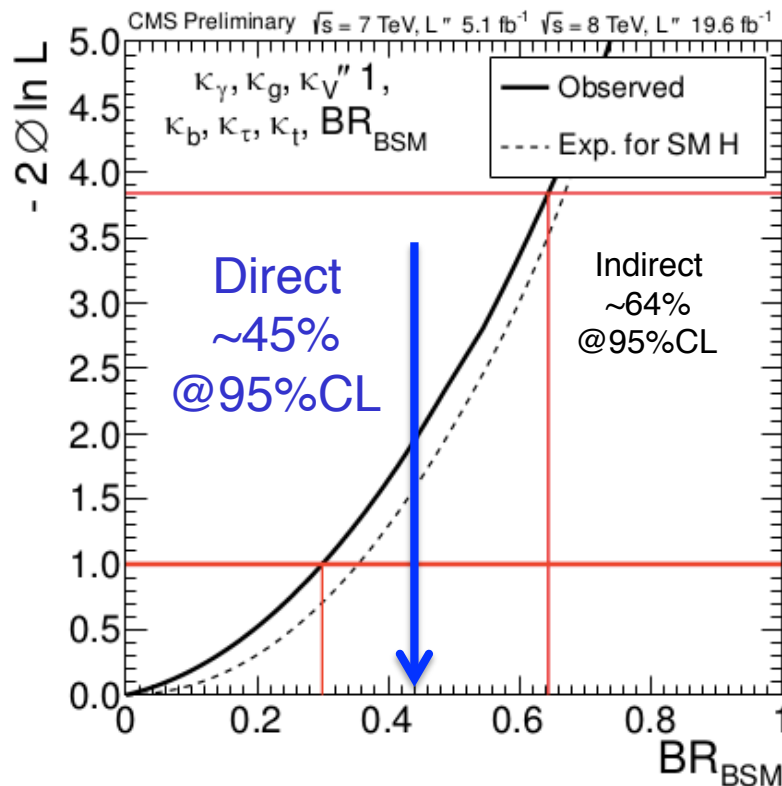
VBF: $BR_{H \rightarrow \text{invisible}} < 68\% @ 95\%CL$

VH: $BR_{H \rightarrow \text{invisible}} < 75\% @ 95\%CL$

Naïve combination: $\sim <45\% @ 95\% CL$

Direct SUSY Searches, O. Buchmüller

$BR_{H \rightarrow \text{invisible}}$ Direct vs Indirect

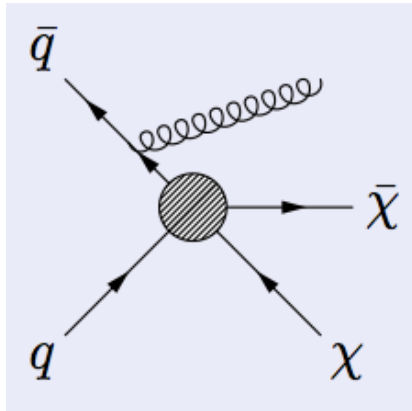


Assuming the experiments are able to maintain trigger and analysis acceptances, the LHC will provide a VERY powerful comparison of indirect & direct measurement of $\Gamma_{H \rightarrow \text{invisible}}$. In the (near) future this might provide a stringent constraint for $M_{DM} < M_H/2$

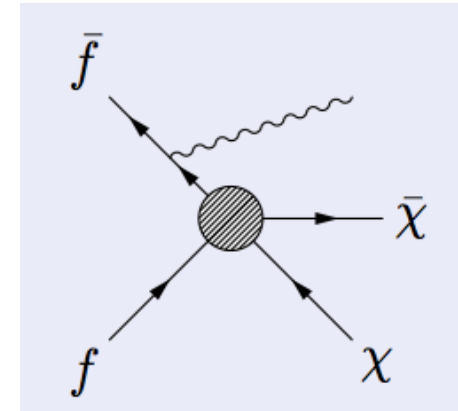
Monojet and Monophoton (plus E_T^{miss})

Direct SUSY Searches, O. Buchmüller

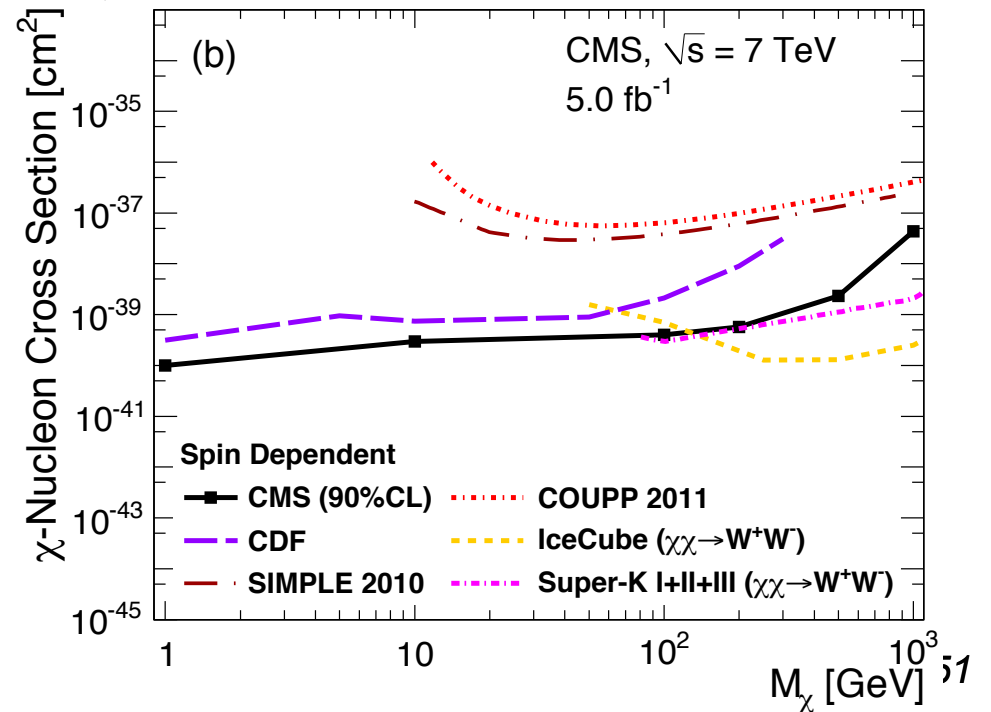
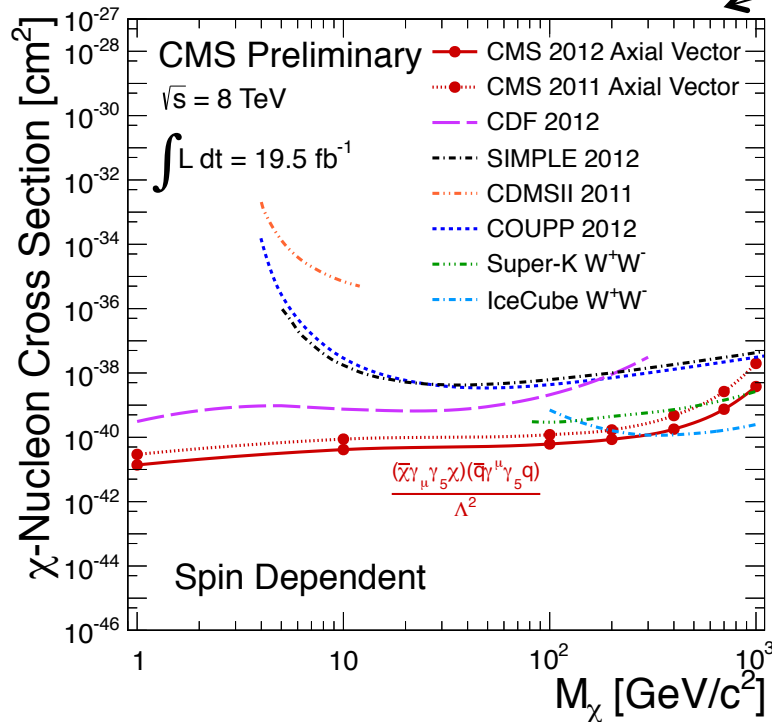
Monojet: hard jet + E_T^{miss}



Monophoton: hard photon + E_T^{miss}



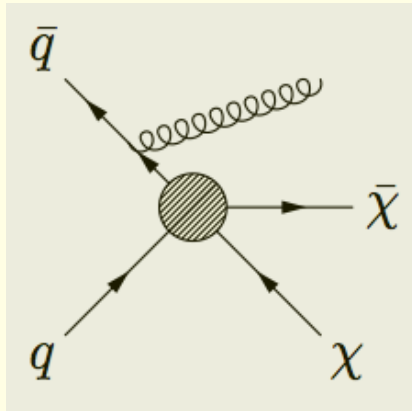
Note different y scale



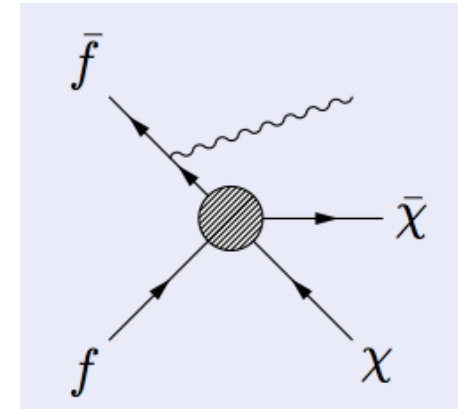
Monojet and Monophoton (plus E_T^{miss})

Direct SUSY Searches, O. Buchmüller

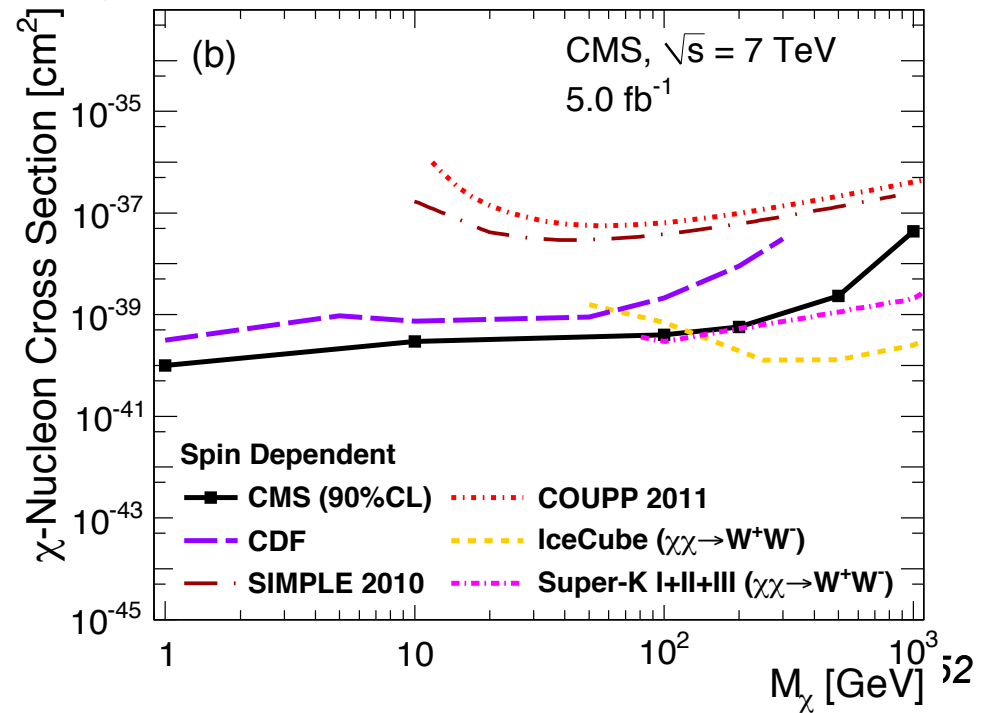
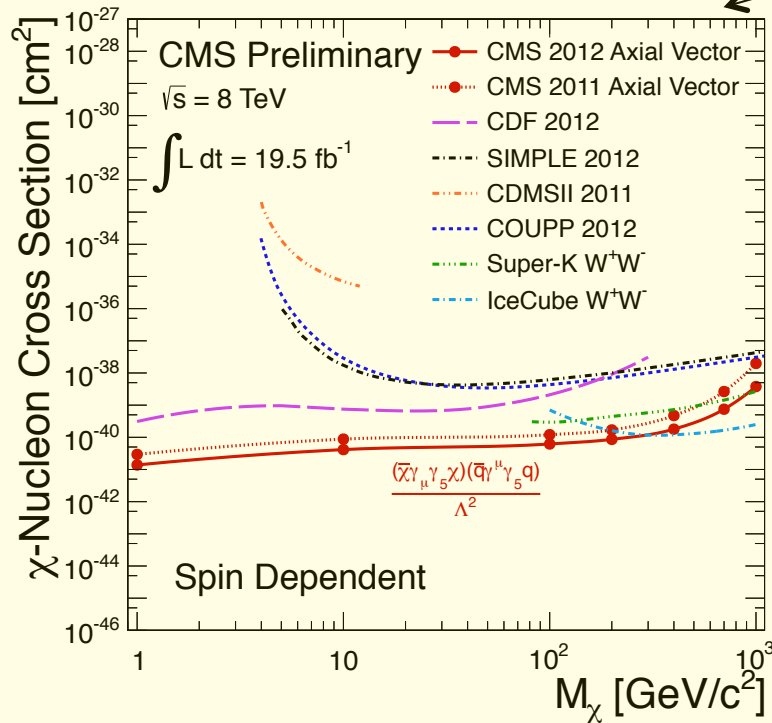
Monojet: hard jet + E_T^{miss}



Monophoton: hard photon + E_T^{miss}

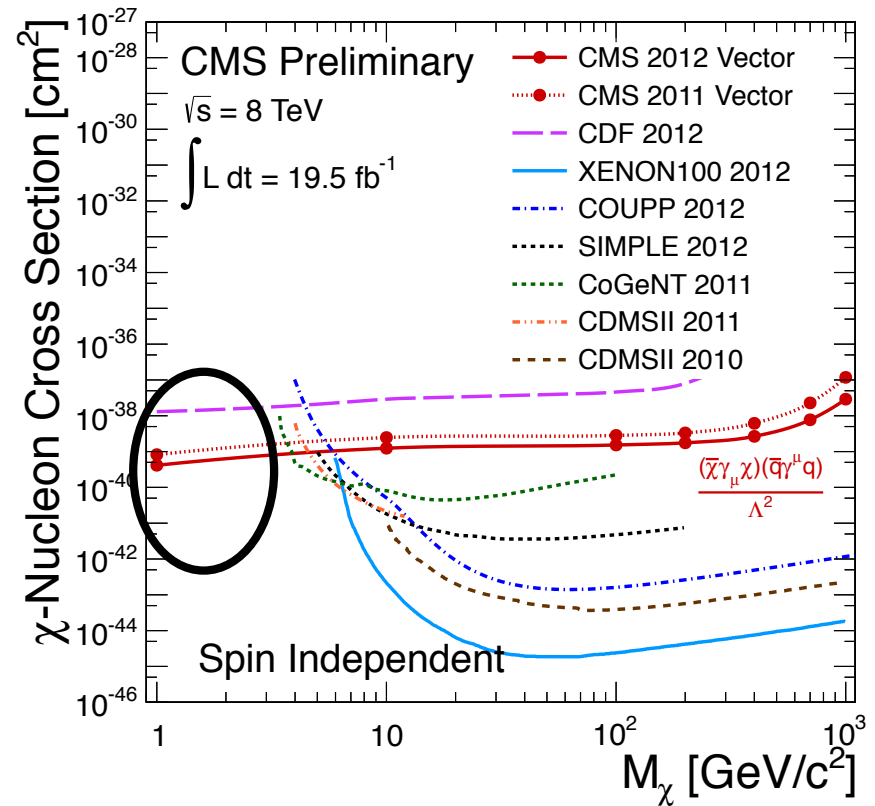
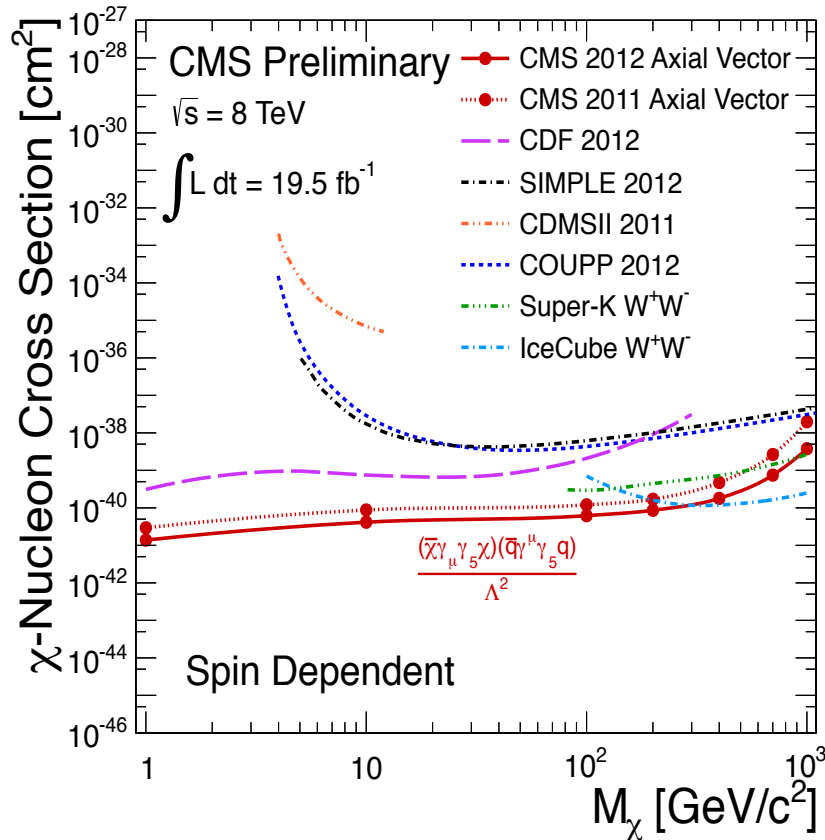


Note different y scale



Monojet analyses better than direct detection?!

Direct SUSY Searches, O. Buchmüller



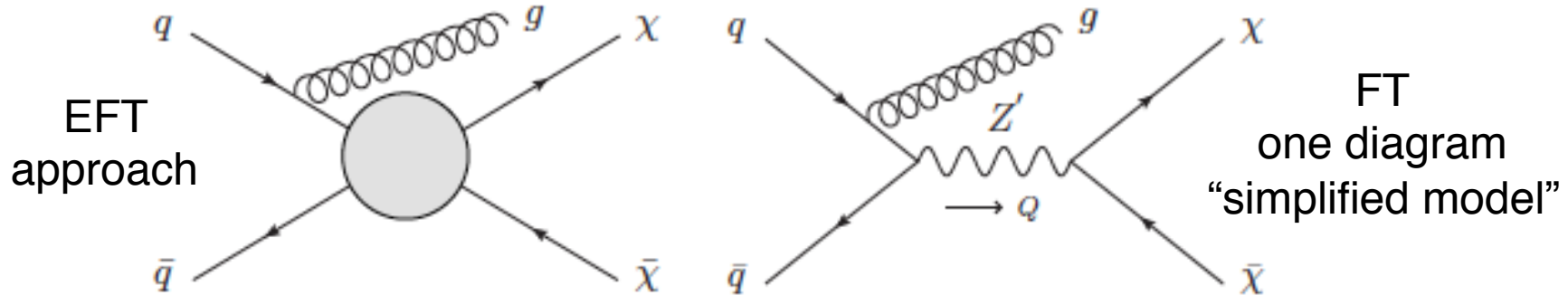
Claim [often made]:

For **low mass** and the entire **spin-dependent** case monojet limits are stronger than direct detection limits!

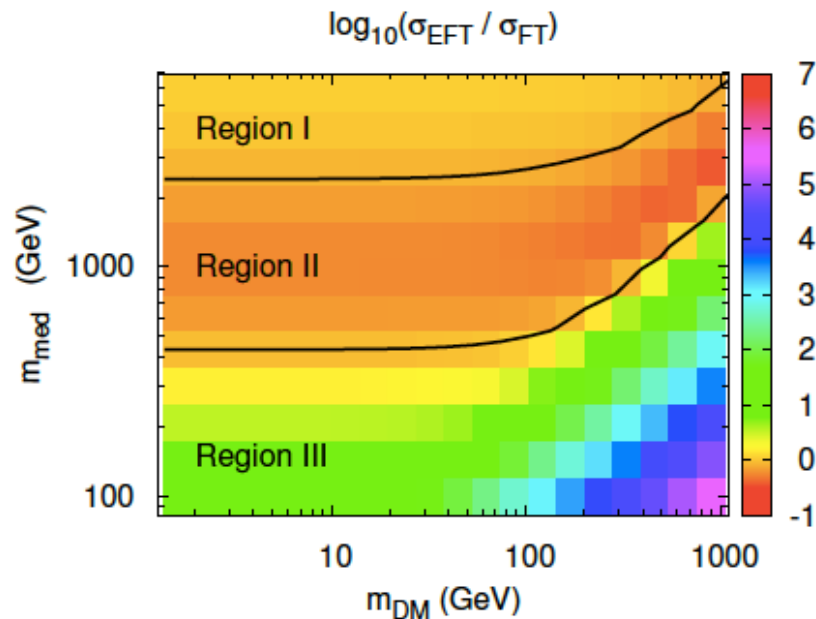
Validity of Effective Field Theory Limits

Recent work from OB, M.Dolan, C.McCabe: arXiv:1308.6799

➤ Compare Effective Field Theory (EFT) with Full Theory (FT)



Use vector and axial-vector mediators (e.g. Z') as example - scalar are similar in conclusion!



Compare prediction of FT with EFT in $m_{\text{med}} - m_{\text{DM}}$ plane. Three regions become visible:

Region I: EFT and FT agree better than 20%

➤ EFT is valid!

Region II: EFT yields significant weaker limits than FT

➤ EFT limits are too conservative!

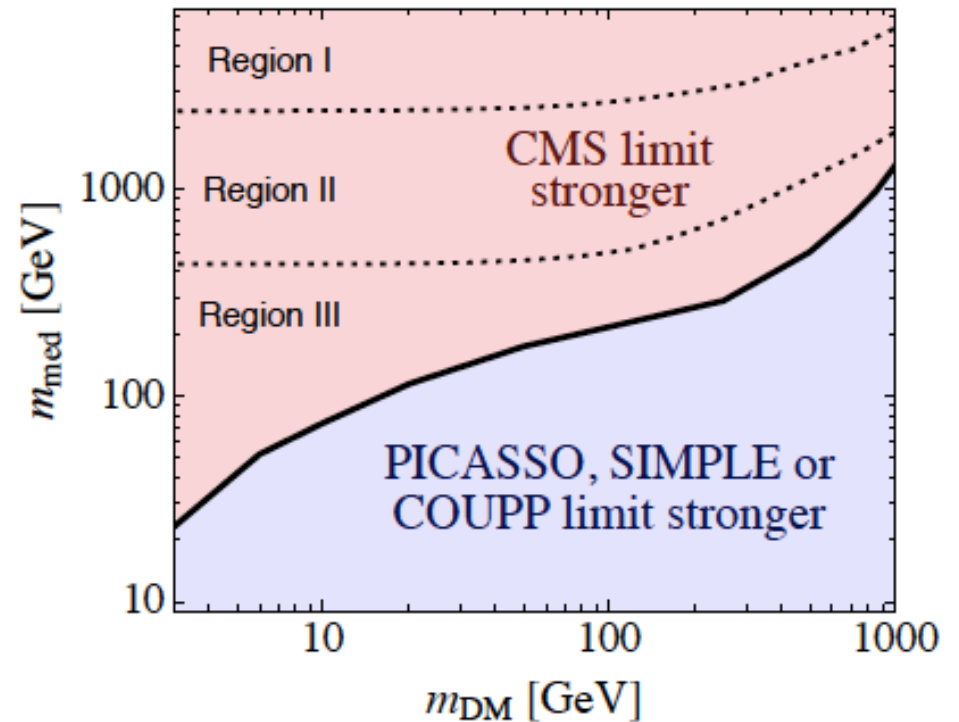
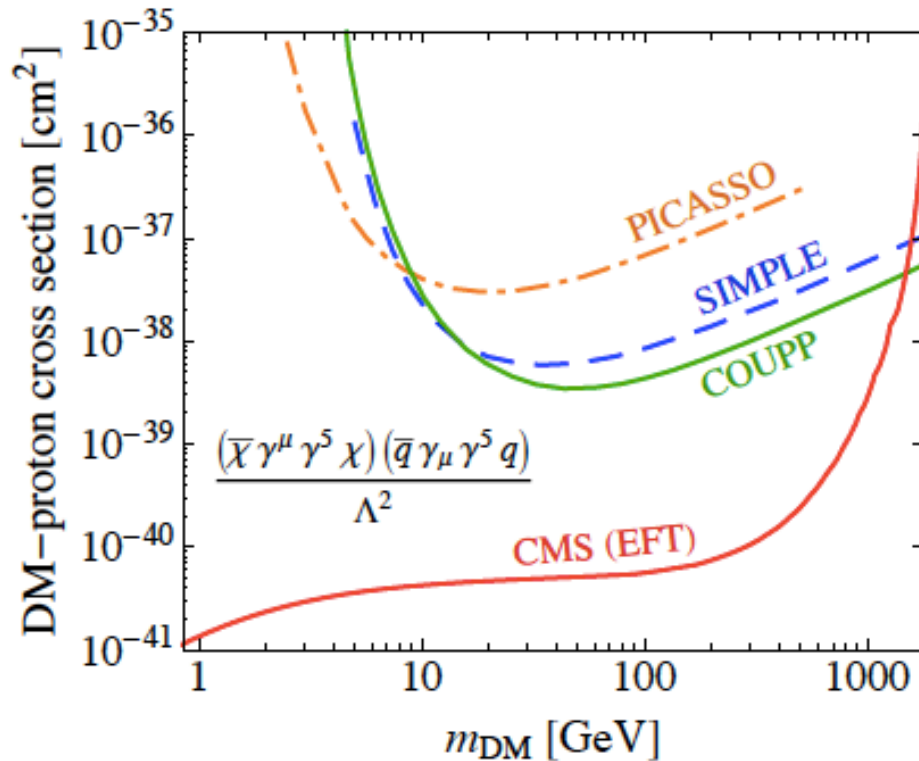
Region III: EFT yields significant stronger limits than FT

➤ EFT limits are too aggressive!

Beyond EFT limits: Simplified models

Working out the complementarity between direct DM detection experiments and collider based DM searches!

üller



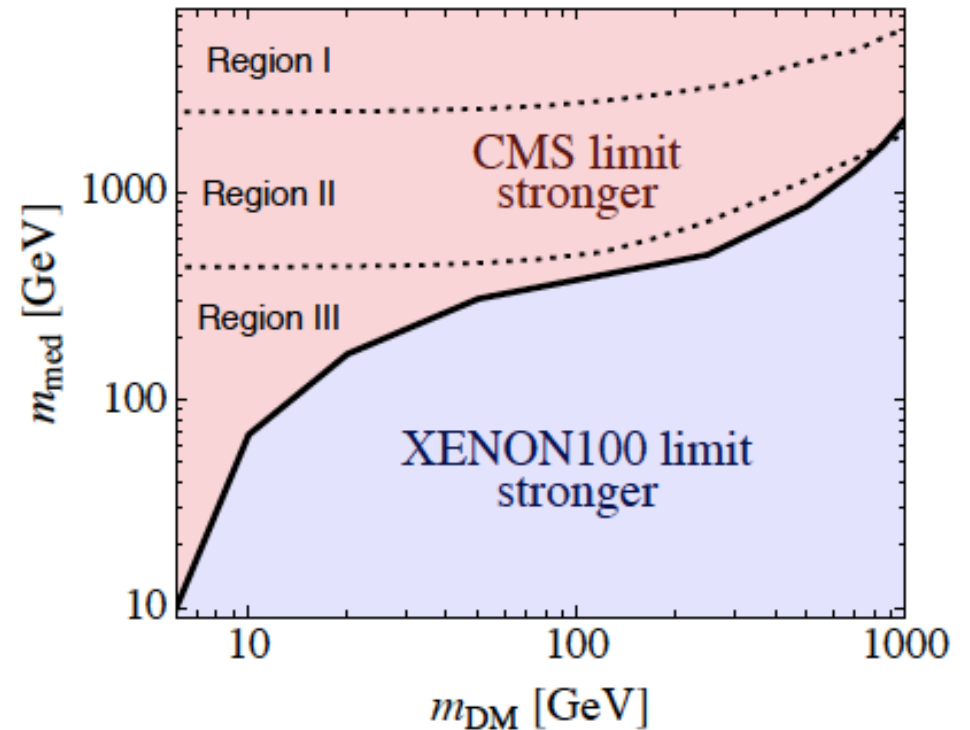
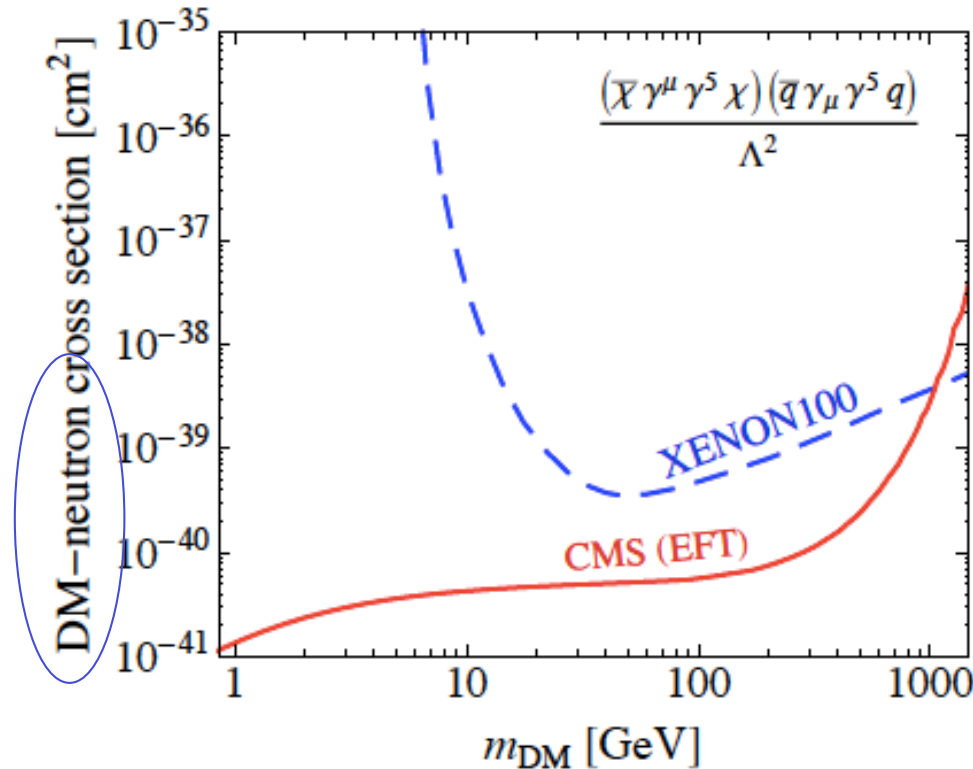
EFT limits give the impression that monojets searches outperform direct detection BUT EFT only applies a VERY small class of DM models.

Simplified model limits give a much better Account of the REAL complementarity and thus seem superior for a comparison.

Beyond EFT limits: Simplified models

Working out the complementarity between direct DM detection experiments and collider based DM searches!

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EFT limits give the impression that monojets searches outperform direct detection BUT EFT only applies a VERY small class of DM models.

Simplified model limits give a much better Account of the REAL complementarity and thus seem superior for a comparison.

Outlook: 8 TeV vs 14 TeV

Use parton luminosities to illustrate the gain of 14 vs 8 TeV

Direct SUSY Searches, O. Buchmüller

Higgs:

$pp \rightarrow H$, $H \rightarrow WW, ZZ$ and $\gamma\gamma$
mainly gg : factor ~ 2

SUSY – 3rd Generation:

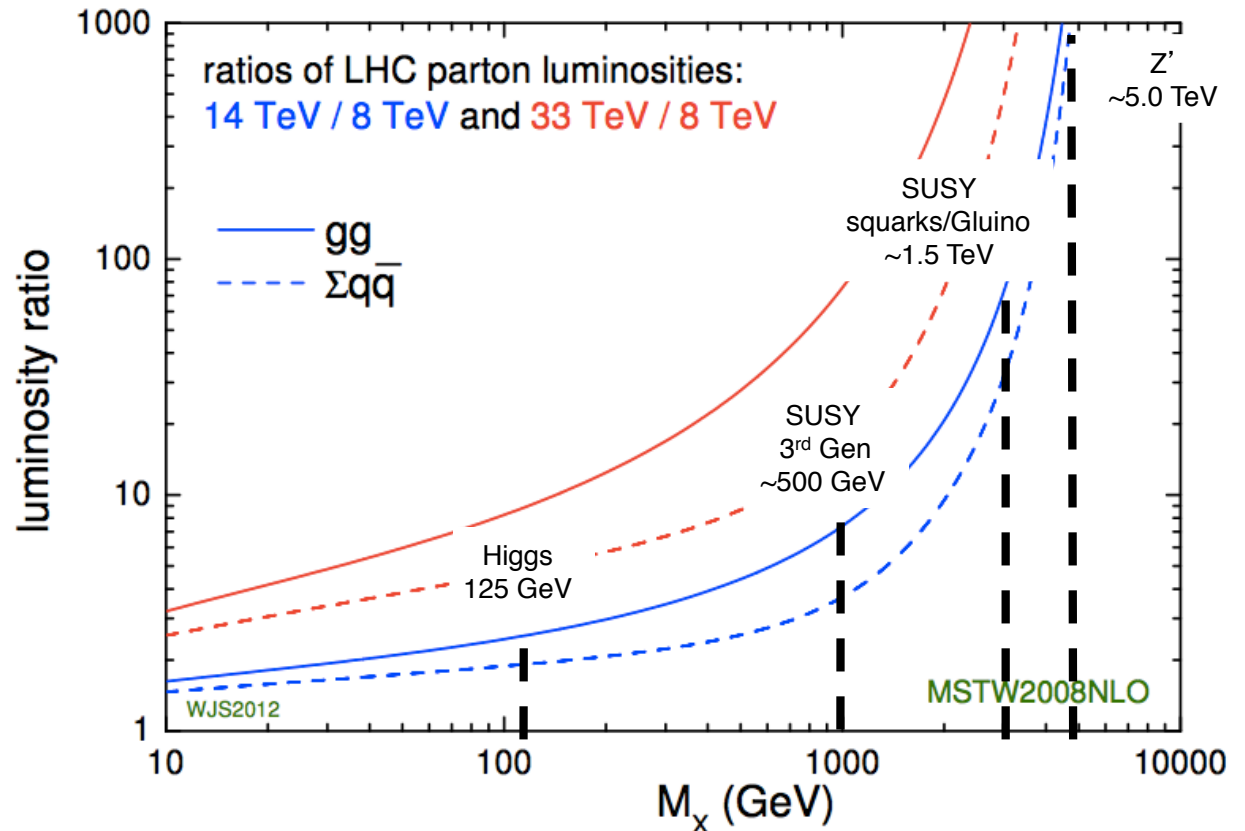
Mass scale ~ 500 GeV
 qq and gg : factor ~ 3 to 6

SUSY – Squarks/Gluino:

Mass scale ~ 1.5 TeV
 qq, gg, qg : factor ~ 40 to 80

Z' :

Mass scale ~ 5 TeV
 qq : factor ~ 1000



Increase in energy will help a lot!
Not just for SUSY...

Summary

- **So far SUSY has not revealed itself!**
 - Even by 2010 the LHC has entered new territory for SUSY searches and since pushed the (coloured) SUSY mass scale to the ~ 1 TeV scale
 - We were well prepared for an early discovery but we also knew that it could take more time and ingenuity before we can claim a discovery (if SUSY exist)
- **The LHC experiments have established an impressive variety of very powerful direct searches for many different final states (very hard work)!**
 - Based on these results we need to establish the “big picture” in order to understand find out if/where our search strategy might have weak spots or even holes!
- **The high energy running of the LHC starting 2015 will be our next very (as in very) real chance for discovery!**

The SUSY (&DM) story continues ... stay tuned!

BACKUP

ATLAS & CMS public results

All results presented in this talk (and many more) can be accessed via the public page of the ATLAS and CMS experiments:

ATLAS SUSY: <https://twiki.cern.ch/twiki/bin/view/AtlasPublic/SupersymmetryPublicResults>

CMS SUSY : <https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsSUS>

ATLAS Summary

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: EPS 2013

ATLAS Preliminary

$$\int \mathcal{L} dt = (4.4 - 22.9) \text{ fb}^{-1} \quad \sqrt{s} = 7, 8 \text{ TeV}$$

Direct SUSY Searches, O. Buchmüller

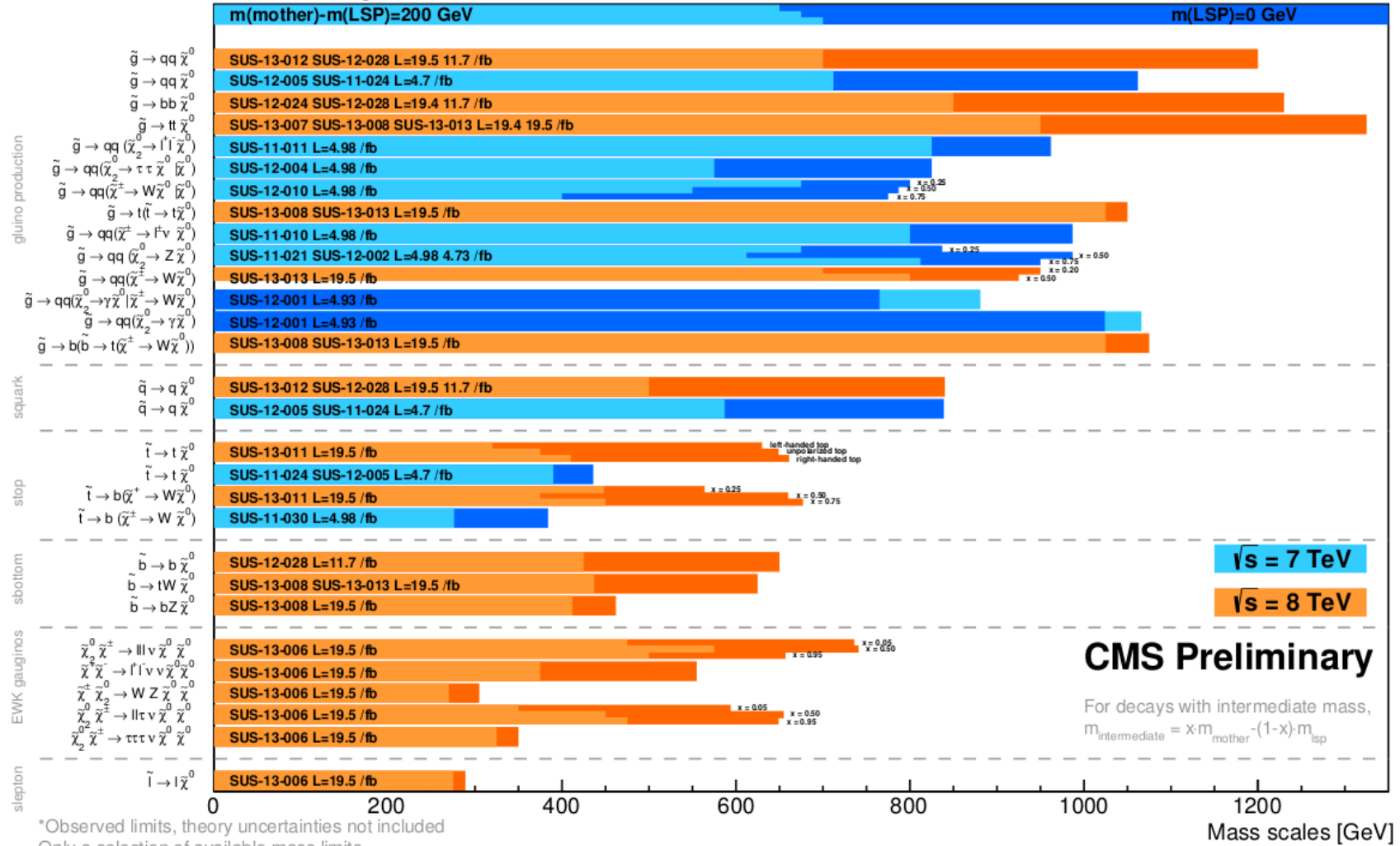
Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference	
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	\tilde{q}, \tilde{g} 1.7 TeV	$m(\tilde{q})=m(\tilde{g})$
	MSUGRA/CMSSM	1 e, μ	3-6 jets	Yes	20.3	\tilde{g} 1.2 TeV	any $m(\tilde{q})$
	MSUGRA/CMSSM	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	any $m(\tilde{q})$
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q} 740 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g} 1.3 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^{\pm} \rightarrow qqW^{\pm}\tilde{\chi}_1^0$	1 e, μ	3-6 jets	Yes	20.3	\tilde{g} 1.18 TeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{\chi}^{\pm})=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$
	$\tilde{g}\tilde{g} \rightarrow qqql(\ell\ell)\tilde{\chi}_1^0\tilde{\chi}_1^0$	2 e, μ (SS)	3 jets	Yes	20.7	\tilde{g} 1.1 TeV	$m(\tilde{\chi}_1^0) < 650 \text{ GeV}$
	GMSB ($\tilde{\ell}$ NLSP)	2 e, μ	2-4 jets	Yes	4.7	\tilde{g} 1.24 TeV	$\tan\beta < 15$
	GMSB ($\tilde{\tau}$ NLSP)	1-2 τ	0-2 jets	Yes	20.7	\tilde{g} 1.4 TeV	$\tan\beta > 18$
	GGM (bino NLSP)	2 γ	0	Yes	4.8	\tilde{g} 1.07 TeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$
	GGM (wino NLSP)	1 $e, \mu + \gamma$	0	Yes	4.8	\tilde{g} 619 GeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$
	GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	\tilde{g} 900 GeV	$m(\tilde{\chi}_1^0) > 220 \text{ GeV}$
	GGM (higgsino NLSP)	2 e, μ (Z)	0-3 jets	Yes	5.8	\tilde{g} 690 GeV	$m(H) > 200 \text{ GeV}$
Gravitino LSP	0	mono-jet	Yes	10.5	$E_T^{1/2}$ scale 645 GeV	$m(\tilde{g}) > 10^{-4} \text{ eV}$	
3 rd gen. \tilde{g} med.	$\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0	3 b	Yes	20.1	\tilde{g} 1.2 TeV	$m(\tilde{\chi}_1^0) < 600 \text{ GeV}$
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.3	\tilde{g} 1.14 TeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^{\pm}$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.34 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$
	$\tilde{g} \rightarrow b\tilde{t}\tilde{\chi}_1^{\pm}$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.3 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}$
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1 100-630 GeV	$m(\tilde{\chi}_1^0) < 100 \text{ GeV}$
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^{\pm}$	2 e, μ (SS)	0-3 b	Yes	20.7	\tilde{b}_1 430 GeV	$m(\tilde{\chi}_1^0)=2 m(\tilde{\chi}_1^{\pm})$
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^{\pm}$	1-2 e, μ	1-2 b	Yes	4.7	\tilde{t}_1 167 GeV	$m(\tilde{\chi}_1^0)=55 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$	2 e, μ	0-2 jets	Yes	20.3	\tilde{t}_1 220 GeV	$m(\tilde{\chi}_1^0)=m(\tilde{t}_1)-m(W)-50 \text{ GeV}, m(\tilde{t}_1) < m(\tilde{\chi}_1^{\pm})$
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	2 e, μ	2 jets	Yes	20.3	\tilde{t}_1 225-525 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^{\pm}$	0	2 b	Yes	20.1	\tilde{t}_1 150-580 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{\chi}_1^{\pm})-m(\tilde{\chi}_1^0)=5 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	1 e, μ	1 b	Yes	20.7	\tilde{t}_1 200-610 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^{\pm}$	0	2 b	Yes	20.5	\tilde{t}_1 320-660 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	mono-jet/c-tag	Yes	20.3	\tilde{t}_1 200 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0) < 85 \text{ GeV}$
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	2 e, μ (Z)	1 b	Yes	20.7	\tilde{t}_1 500 GeV	$m(\tilde{\chi}_1^0) > 150 \text{ GeV}$
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ (Z)	1 b	Yes	20.7	\tilde{t}_2 520 GeV	$m(\tilde{t}_1)=m(\tilde{\chi}_1^0)+180 \text{ GeV}$	
EIW direct	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$ 85-315 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$
	$\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^{\pm} \rightarrow \tilde{\ell}\nu(\tilde{\ell}\bar{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^{\pm}$ 125-450 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^{\pm})+m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^{\pm} \rightarrow \tilde{\tau}\nu(\tilde{\tau}\bar{\nu})$	2 τ	0	Yes	20.7	$\tilde{\chi}_1^{\pm}$ 180-330 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\tau}, \bar{\nu})=0.5(m(\tilde{\chi}_1^{\pm})+m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^{\pm}\tilde{\chi}_1^0 \rightarrow \tilde{\ell}\nu\tilde{\ell}\ell(\tilde{\nu}\bar{\nu}), \tilde{\nu}\tilde{\ell}\ell(\tilde{\nu}\bar{\nu})$	3 e, μ	0	Yes	20.7	$\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^0$ 600 GeV	$m(\tilde{\chi}_1^0)=m(\tilde{\chi}_1^{\pm}), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^{\pm})+m(\tilde{\chi}_1^0))$
	$\tilde{\chi}_1^{\pm}\tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0 Z^* \tilde{\chi}_1^0$	3 e, μ	0	Yes	20.7	$\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^0$ 315 GeV	$m(\tilde{\chi}_1^0)=m(\tilde{\chi}_1^{\pm}), m(\tilde{\chi}_1^0)=0, \text{ sleptons decoupled}$
Long-lived particles	Direct $\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}$ prod., long-lived $\tilde{\chi}_1^{\pm}$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^{\pm}$ 270 GeV	$m(\tilde{\chi}_1^0)-m(\tilde{\chi}_1^{\pm})=160 \text{ MeV}, \tau(\tilde{\chi}_1^{\pm})=0.2 \text{ ns}$
	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	22.9	\tilde{g} 857 GeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$	1-2 μ	0	-	15.9	$\tilde{\chi}_1^0$ 475 GeV	$10 < \tan\beta < 50$
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma G$, long-lived $\tilde{\chi}_1^0$	2 γ	0	Yes	4.7	$\tilde{\chi}_1^0$ 230 GeV	$0.4 < \tau(\tilde{\chi}_1^0) < 2 \text{ ns}$
$\tilde{\chi}_1^0 \rightarrow q\tilde{q}$ (RPV)	1 μ	0	Yes	4.4	\tilde{q} 700 GeV	1 mm $< c\tau < 1 \text{ m}, \tilde{g}$ decoupled	
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e + \mu$	2 e, μ	0	-	4.6	$\tilde{\nu}_\tau$ 1.61 TeV	$\lambda'_{311}=0.10, \lambda_{132}=0.05$
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e(\mu) + \tau$	1 $e, \mu + \tau$	0	-	4.6	$\tilde{\nu}_\tau$ 1.1 TeV	$\lambda'_{311}=0.10, \lambda_{1(2)33}=0.05$
	Bilinear RPV CMSSM	1 e, μ	7 jets	Yes	4.7	\tilde{q}, \tilde{g} 1.2 TeV	$m(\tilde{q})=m(\tilde{g}), c\tau_{LSP} < 1 \text{ mm}$
	$\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^{\pm} \rightarrow W\tilde{\chi}_1^0\tilde{\chi}_1^0 \rightarrow ee\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$	4 e, μ	0	Yes	20.7	$\tilde{\chi}_1^{\pm}$ 760 GeV	$m(\tilde{\chi}_1^0) > 300 \text{ GeV}, \lambda_{121} > 0$
	$\tilde{\chi}_1^{\pm}\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^{\pm} \rightarrow W\tilde{\chi}_1^0\tilde{\chi}_1^0 \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 $e, \mu + \tau$	0	Yes	20.7	$\tilde{\chi}_1^{\pm}$ 350 GeV	$m(\tilde{\chi}_1^0) > 80 \text{ GeV}, \lambda_{133} > 0$
	$\tilde{g} \rightarrow qq\tilde{q}$	0	6 jets	-	4.6	\tilde{g} 666 GeV	
$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$	2 e, μ (SS)	0-3 b	Yes	20.7	\tilde{g} 880 GeV		
Other	Scalar gluon	0	4 jets	-	4.6	sgluon 100-287 GeV	incl. limit from 1110.2693
	WIMP interaction (D5, Dirac χ)	0	mono-jet	Yes	10.5	M^* scale 704 GeV	$m(\chi) < 80 \text{ GeV}$, limit of $\sim 687 \text{ GeV}$ for D8

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

CMS Summary

Direct SUSY Searches, O. Buchmüller

Summary of CMS SUSY Results* in SMS framework EPSHEP 2013



Summary

Slide from 2007
EPS Plenary talk in Manchester
LHC Detectors
Commissioning & Physics
pre-accident & 14 TeV assumption



- LHC&Experiments are on track for first collisions in 2008
 - Challenge: commissioning of machine and detectors of unprecedented complexity, technology and performance
- The LHC will discover (or exclude) the Higgs by ~2010
 - Electro Weak Symmetry Breaking
 - Large phase space can already be excluded with only $\sim 1\text{fb}^{-1}$
- The LHC will discover low energy SUSY (if it exists)
 - Could be easy; could also take more time and ingenuity before we can claim a discovery
 - First signals might emerge already in the first data
 - 1-3 TeV can be covered already with $< 10\text{fb}^{-1}$
- The LHC will cover a new physics scale of 1-3 TeV
 - Many new physics models; Black hole, Extra Dimensions, Little Higgs, Split Susy, New Bosons, Technicolour, etc ...

In other words; the next five years will be an exciting time for particle physics ...

(Best) mass limits in a nutshell (RP conserving)

Direct SUSY Searches, O. Buchmüller

Direct squark	$\tilde{q} \rightarrow q\chi_1^0$	$\tilde{u}_L \rightarrow q\chi_1^0$	$\tilde{b} \rightarrow b\chi_1^0$	$\tilde{t} \rightarrow t\chi_1^0$
Best limit: [GeV]	~850	~500	~650	~650
No limit for M_{LSP} [GeV]	~300	~120	~270	~260

coloured sparticle production

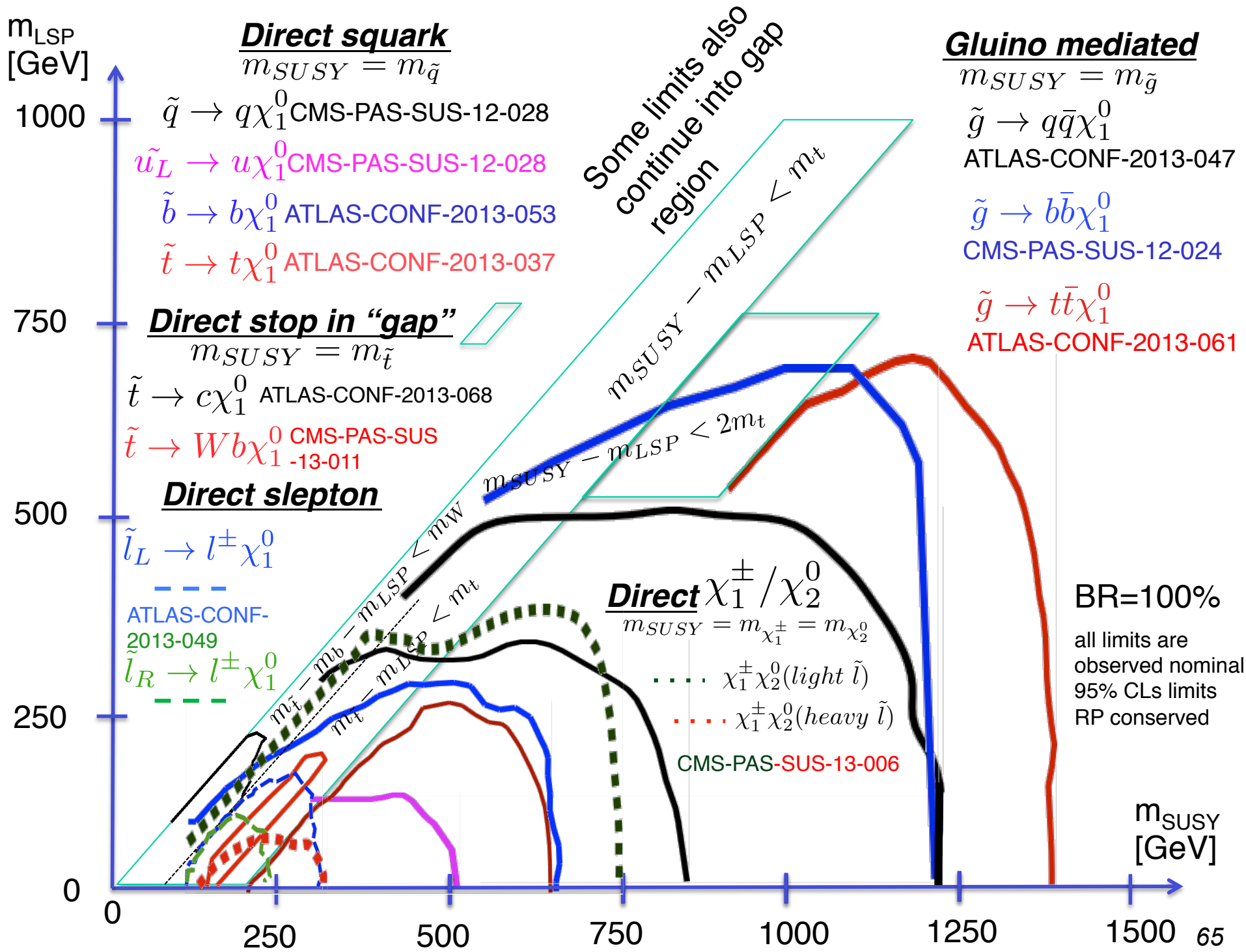
Direct squark	$\tilde{g} \rightarrow q\bar{q}\chi_1^0$	$\tilde{g} \rightarrow b\bar{b}\chi_1^0$	$\tilde{g} \rightarrow t\bar{t}\chi_1^0$
Best limit: [GeV]	~1200	~1200	~1400
No limit for M_{LSP} [GeV]	~480	~650	~700

Stop $M_{stop} - M_{LSP} < M_{top}$	$\tilde{t} \rightarrow c\chi_1^0$	$\tilde{t} \rightarrow Wb\chi_1^0$
Best limit: [GeV]	~240	~320
No limit for M_{LSP} [GeV]	~210	~190

EWK sparticle production

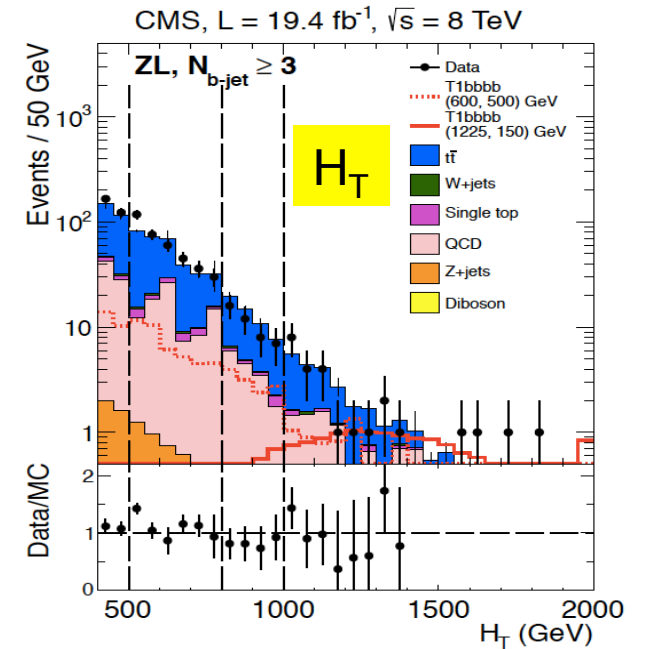
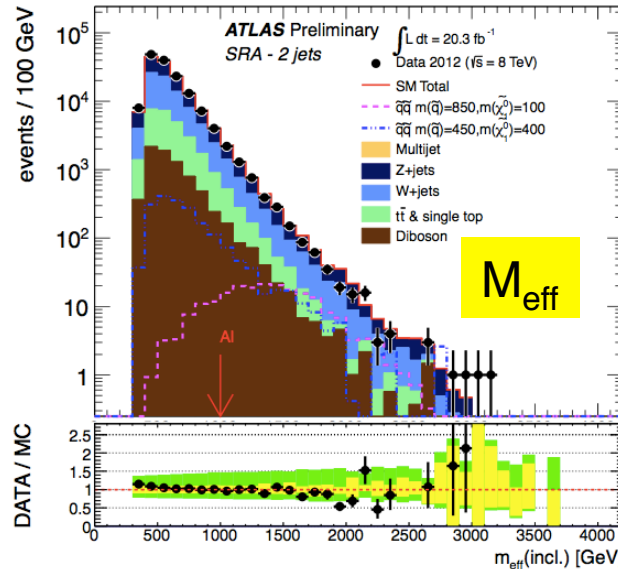
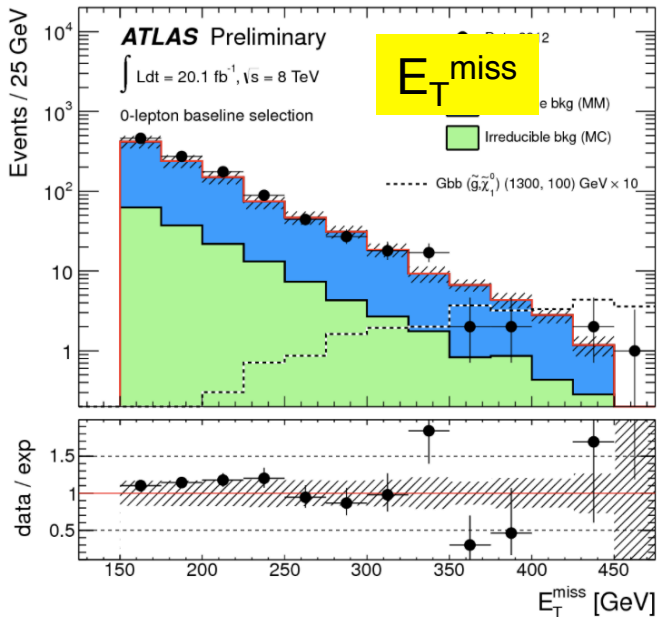
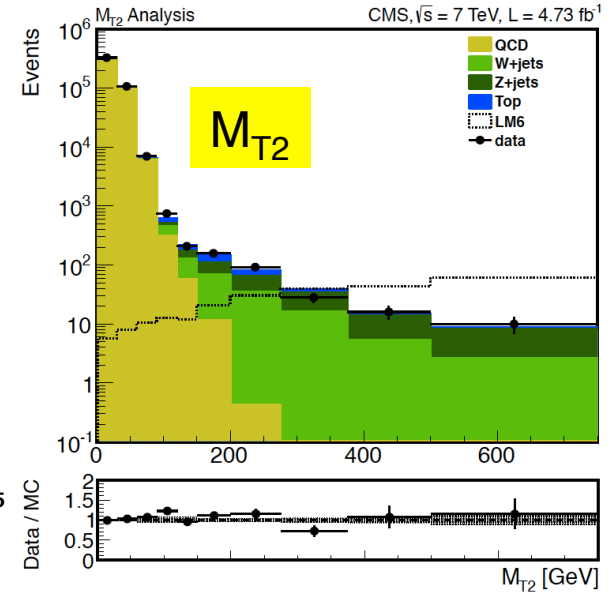
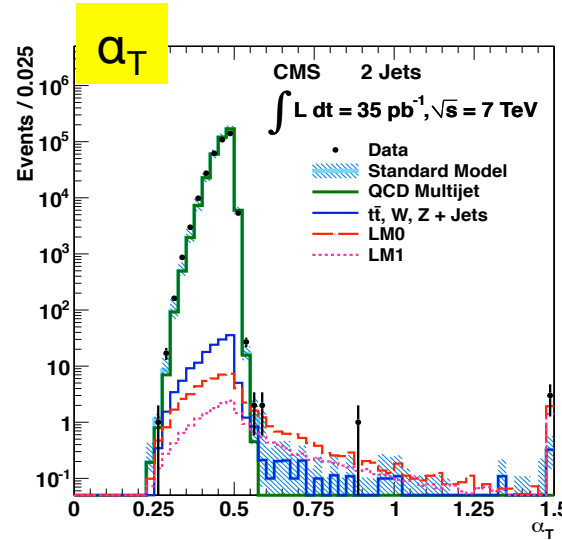
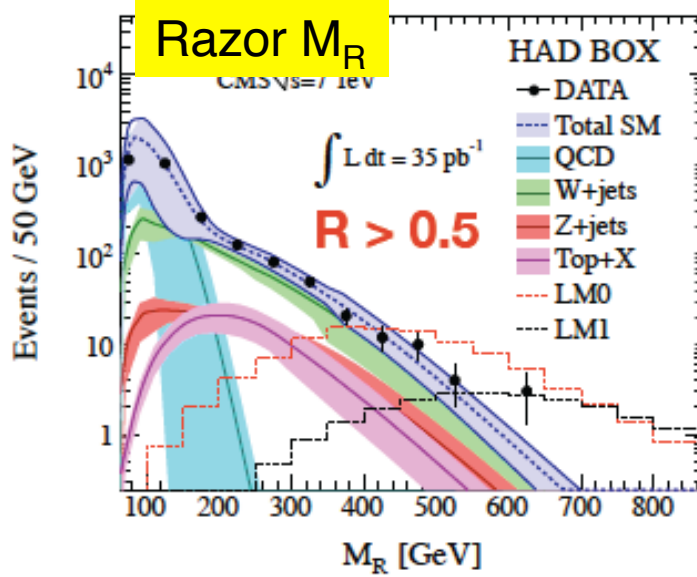
Direct slepton	$\tilde{l}_L \rightarrow l^\pm\chi_1^0$	$\tilde{l}_R \rightarrow l^\pm\chi_1^0$
Best limit: [GeV]	~300	~240
No limit for M_{LSP} [GeV]	~150	~90

$\chi_1^\pm\chi_2^0$	light \tilde{l}	heavy \tilde{l}
Best limit: [GeV]	~750	~300
No limit for M_{LSP} [GeV]	~350	~60

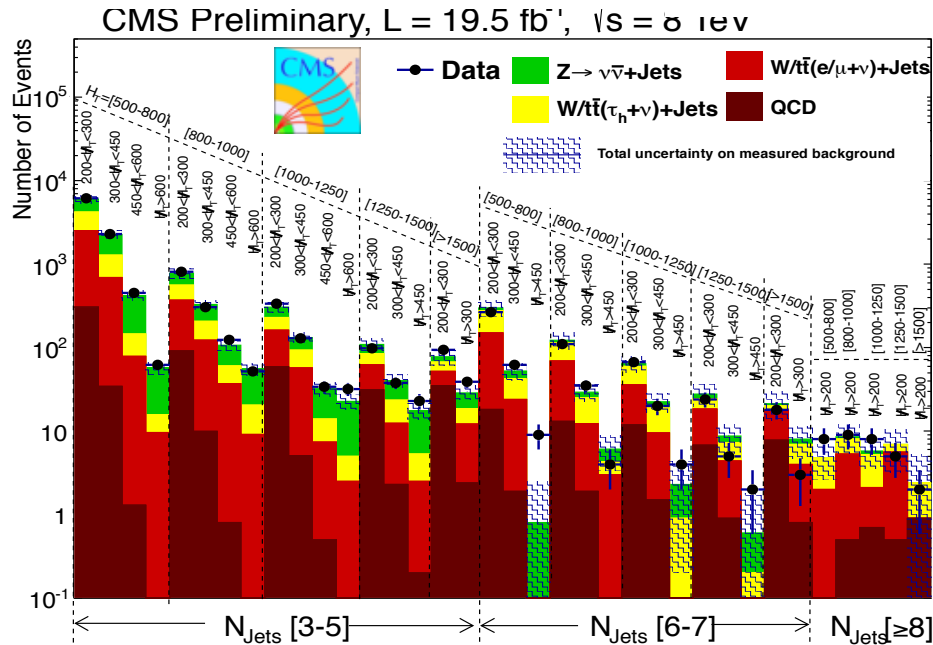
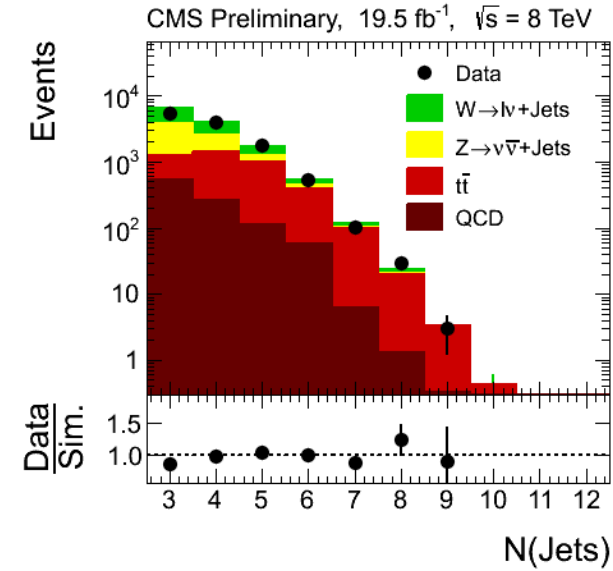
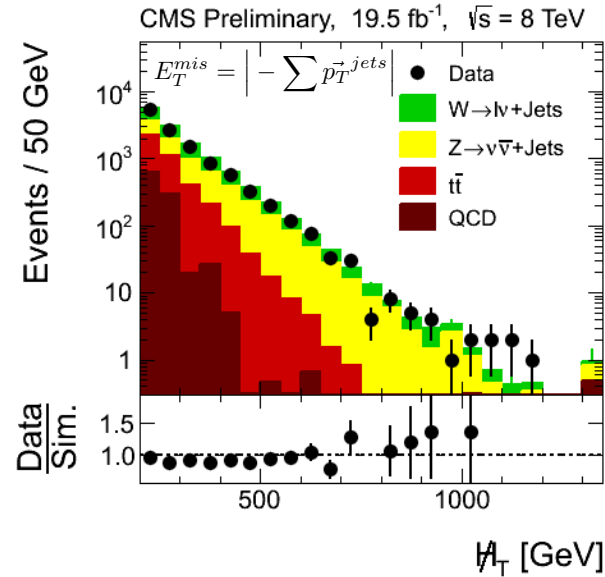
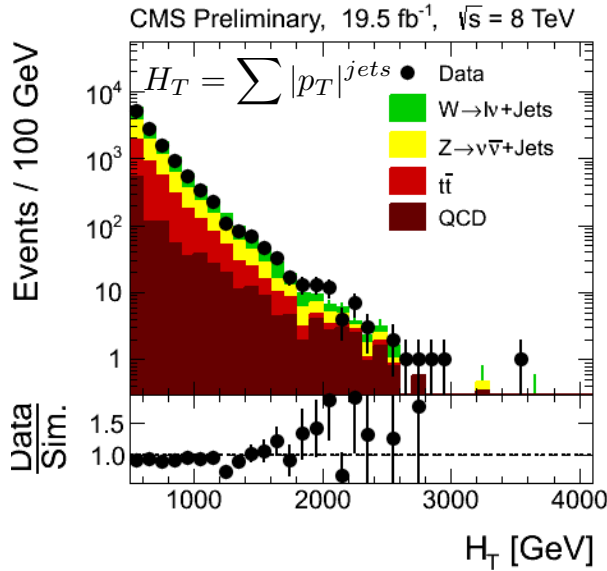


Many Different Kinematic Variables

Direct SUSY Searches, O. Buchmüller



Multijets & missing energy search



Traditional inclusive Jets + E_T^{mis} search, which uses simple kinematic variables to categorize the events.

Main backgrounds QCD, W/Z+jet and t \bar{t} are estimated using data-driven techniques.

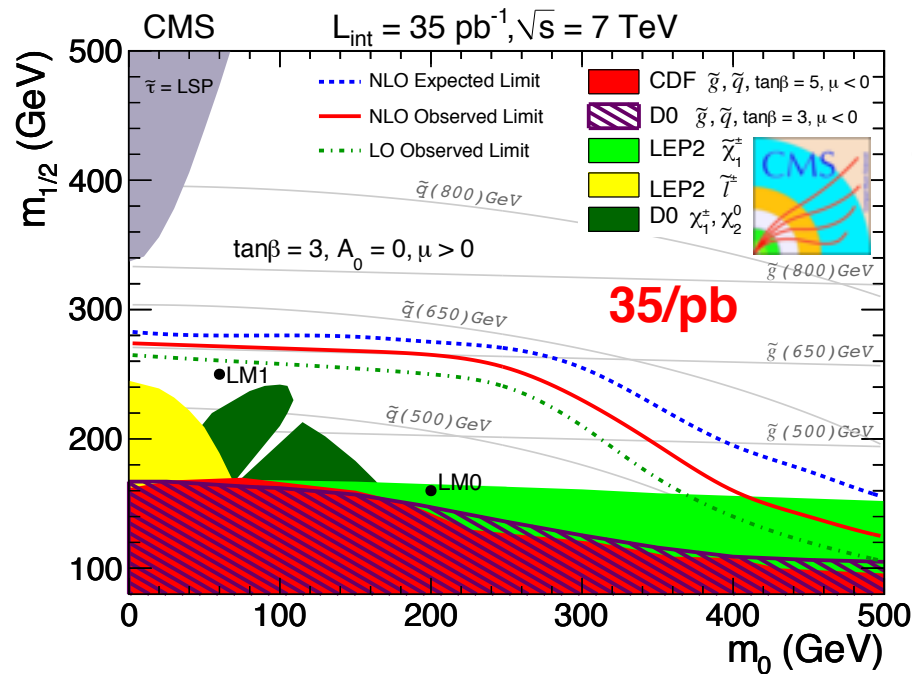
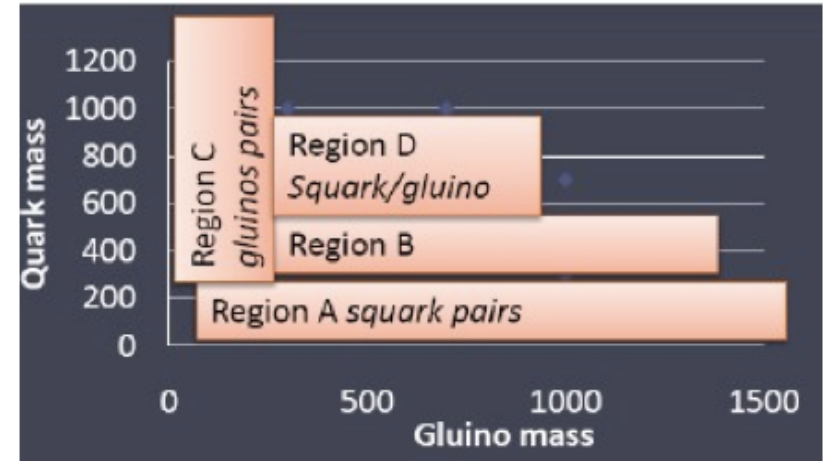
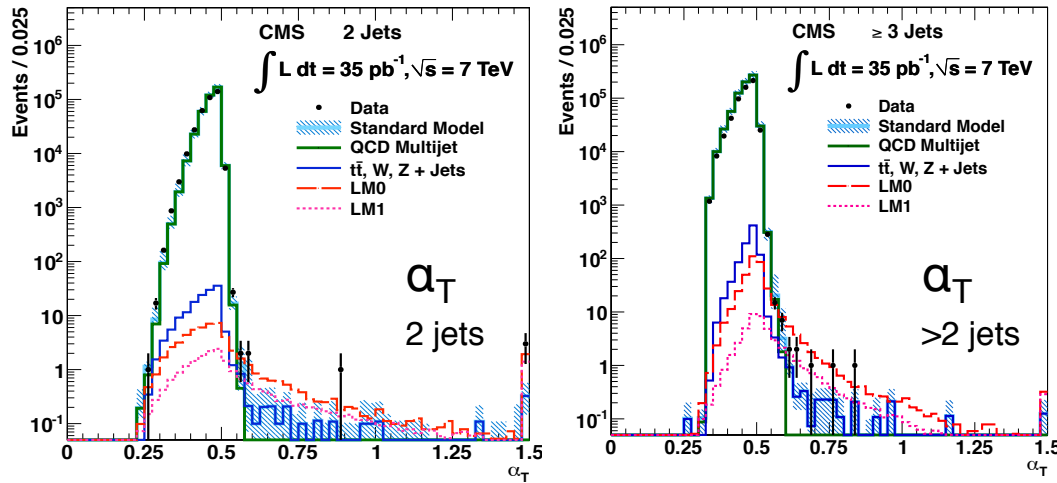


CMS-SUS-PAS-13-012

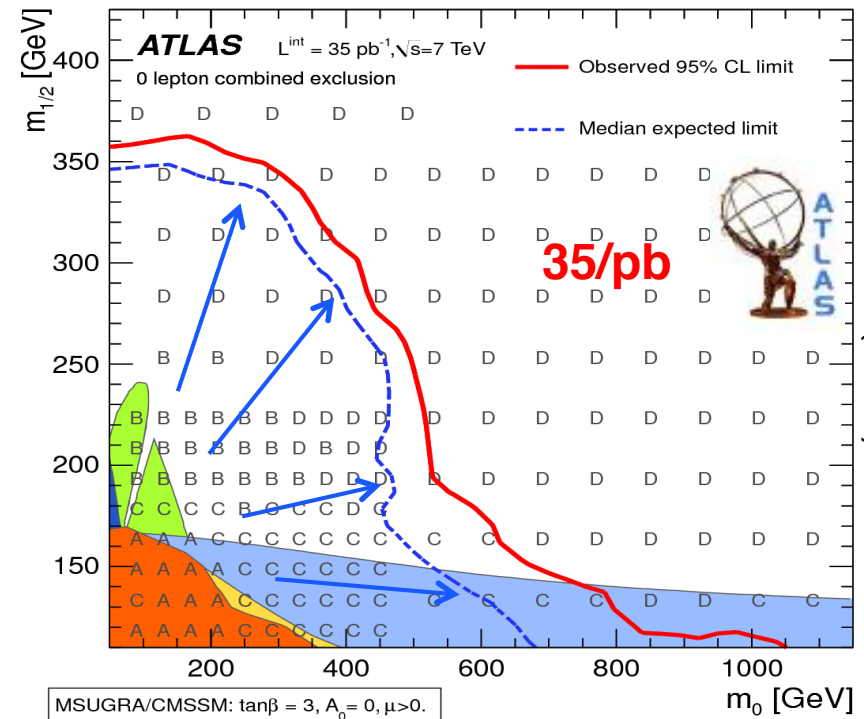
See parallel talk for details:

C. Autermann

2010: Entering New Territory at the LHC!



Phys. Lett. B698:196-218

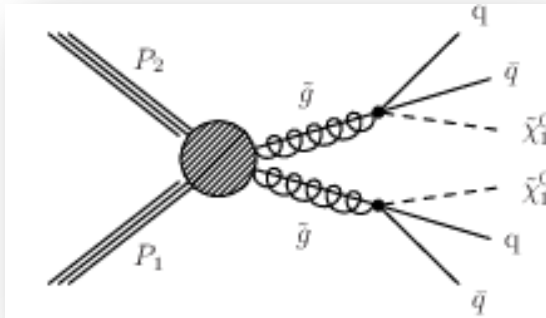


PLB 701 (2011) 186

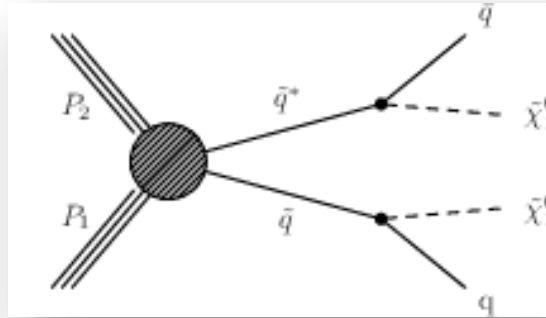
MSUGRA/CMSSM: $\tan\beta = 3, A_0 = 0, \mu > 0$.

Simplified Model Spectra (SMS)

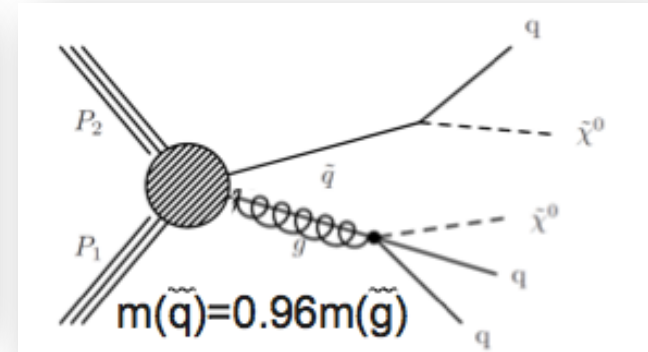
earches, O. Buchmüller



$$\tilde{g}\tilde{g} \rightarrow q\bar{q}\tilde{\chi}^0 q\bar{q}\tilde{\chi}^0$$



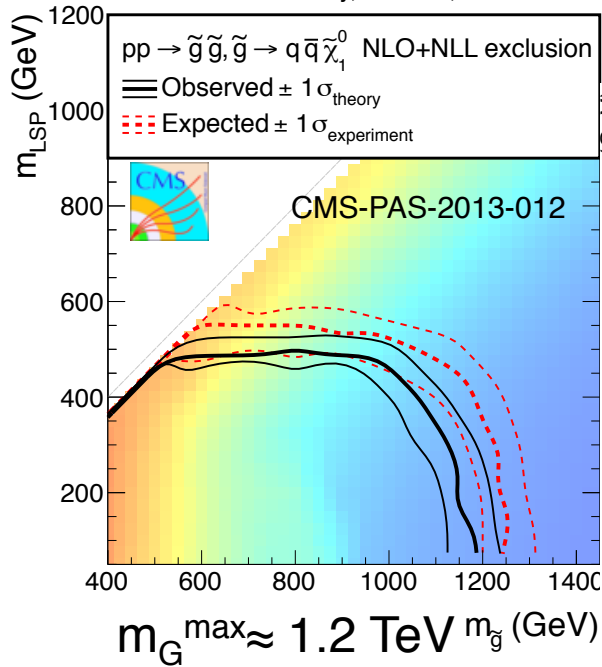
$$\tilde{q}\tilde{q} \rightarrow q\tilde{\chi}^0 \bar{q}\tilde{\chi}^0$$



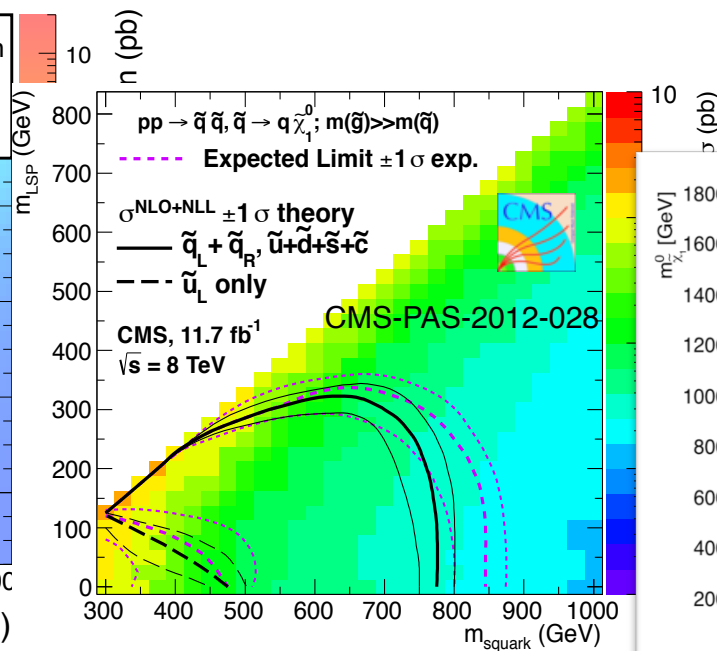
$$\tilde{q} \rightarrow q\tilde{\chi}_1^0, \tilde{g} \rightarrow qq\tilde{\chi}_1^0$$

$$m(\tilde{q}) = 0.96m(\tilde{g})$$

CMS Preliminary, 19.5 fb⁻¹, √s = 8 TeV

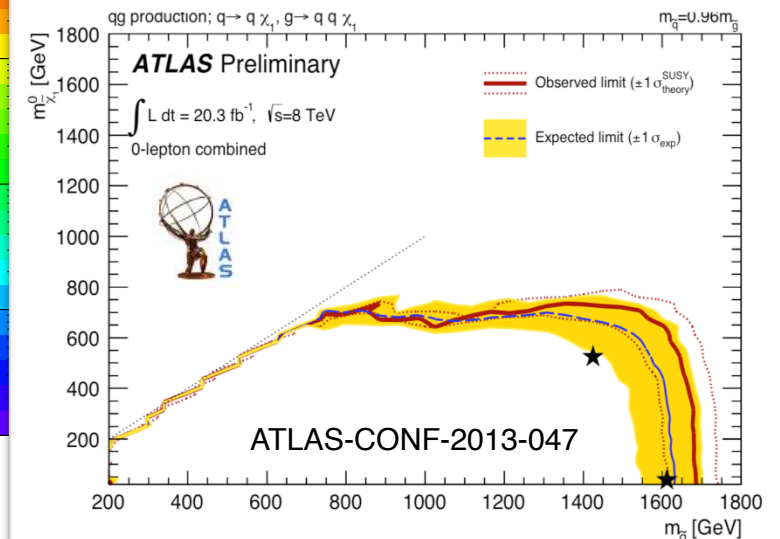


$$m_G^{\max} \approx 1.2 \text{ TeV}$$

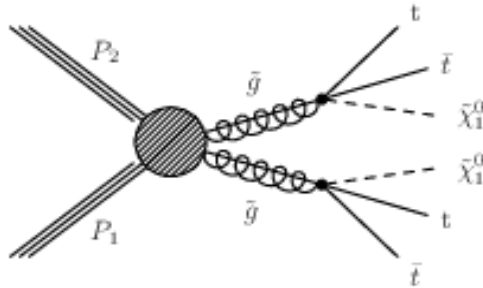


$$m_{sq}^{\max} \approx 0.8 \text{ TeV}$$

$$m_G^{\max} \approx 1.7 \text{ TeV}$$

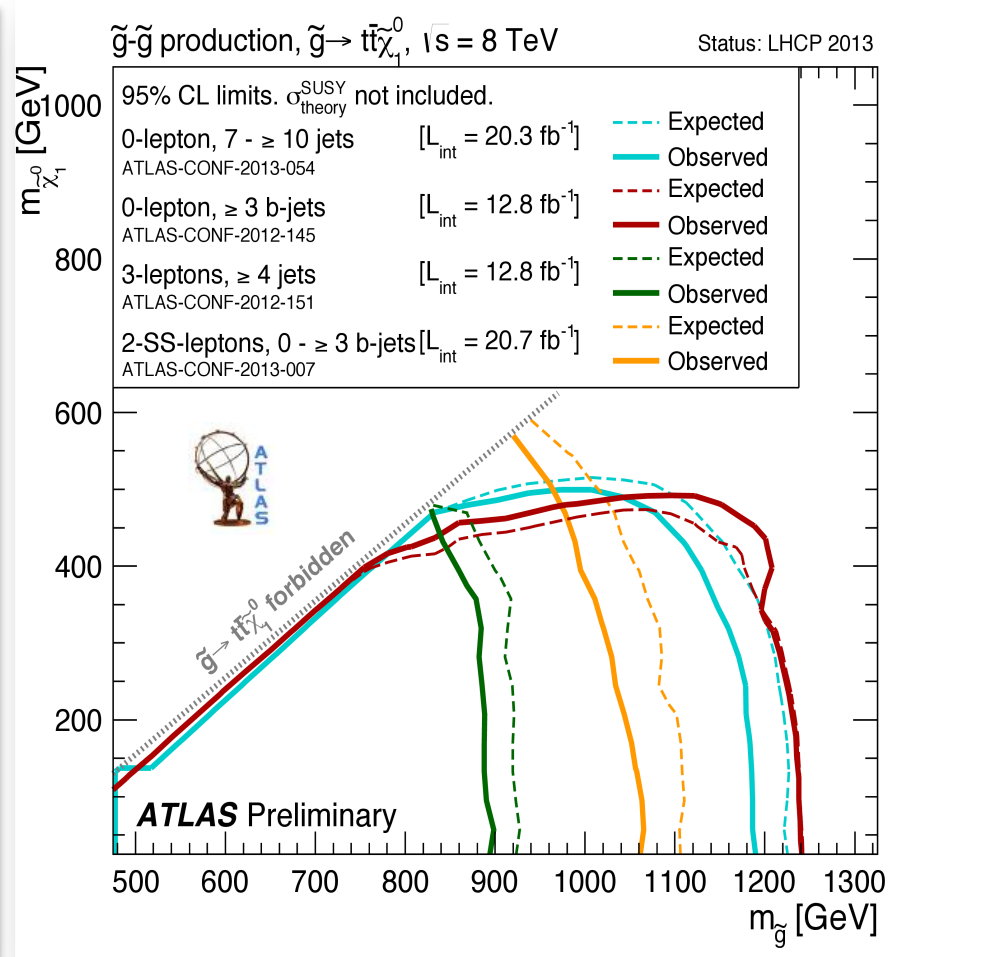
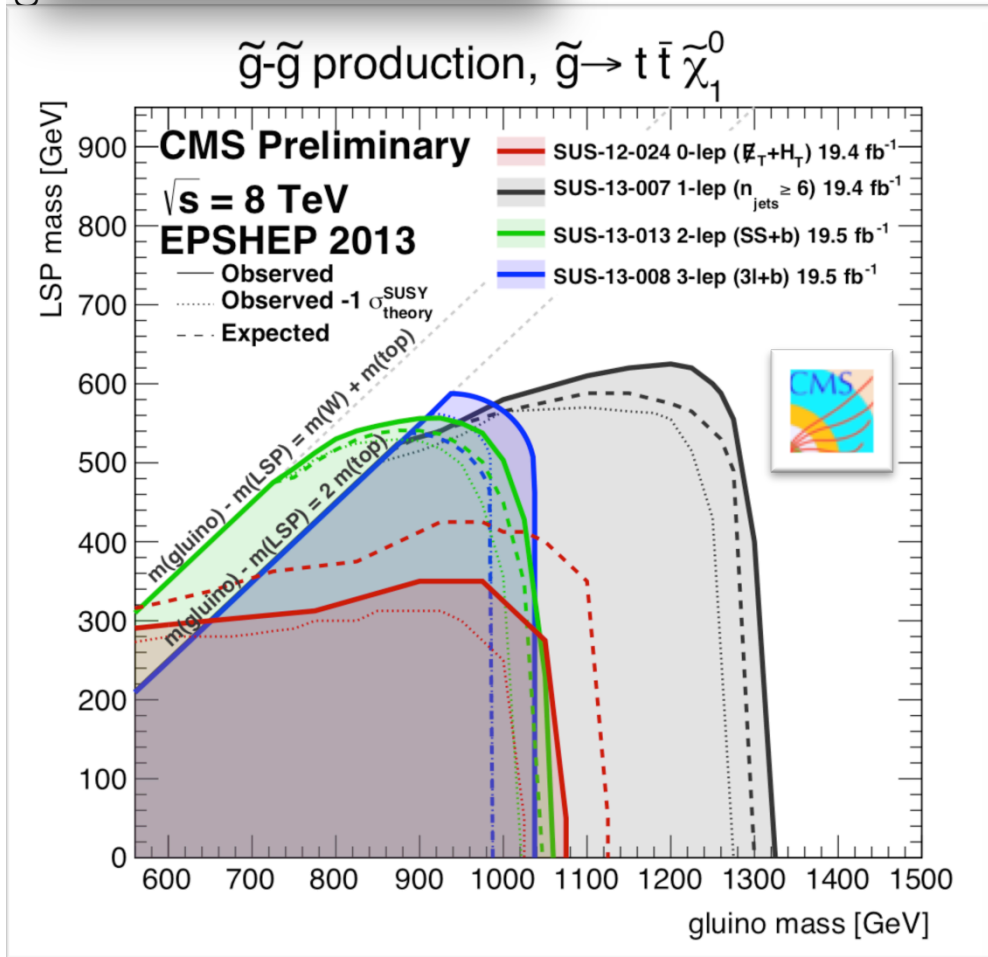


Today: many more SMS and many more searches

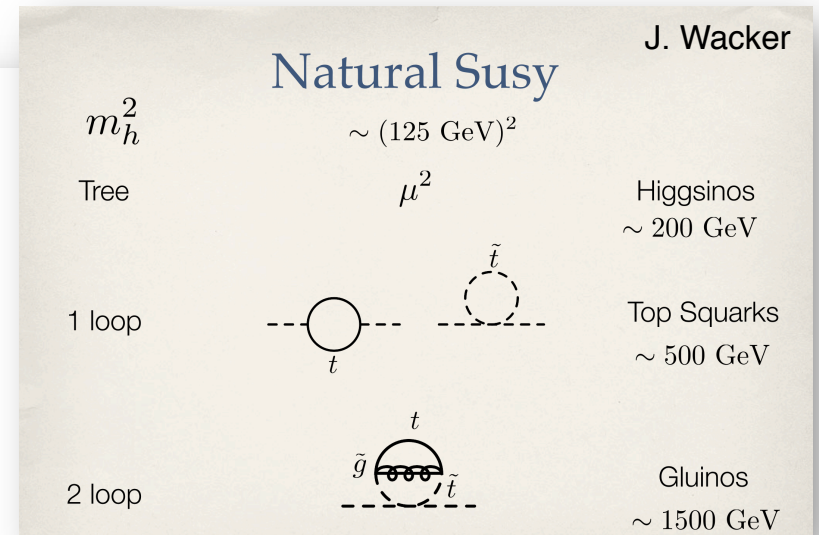
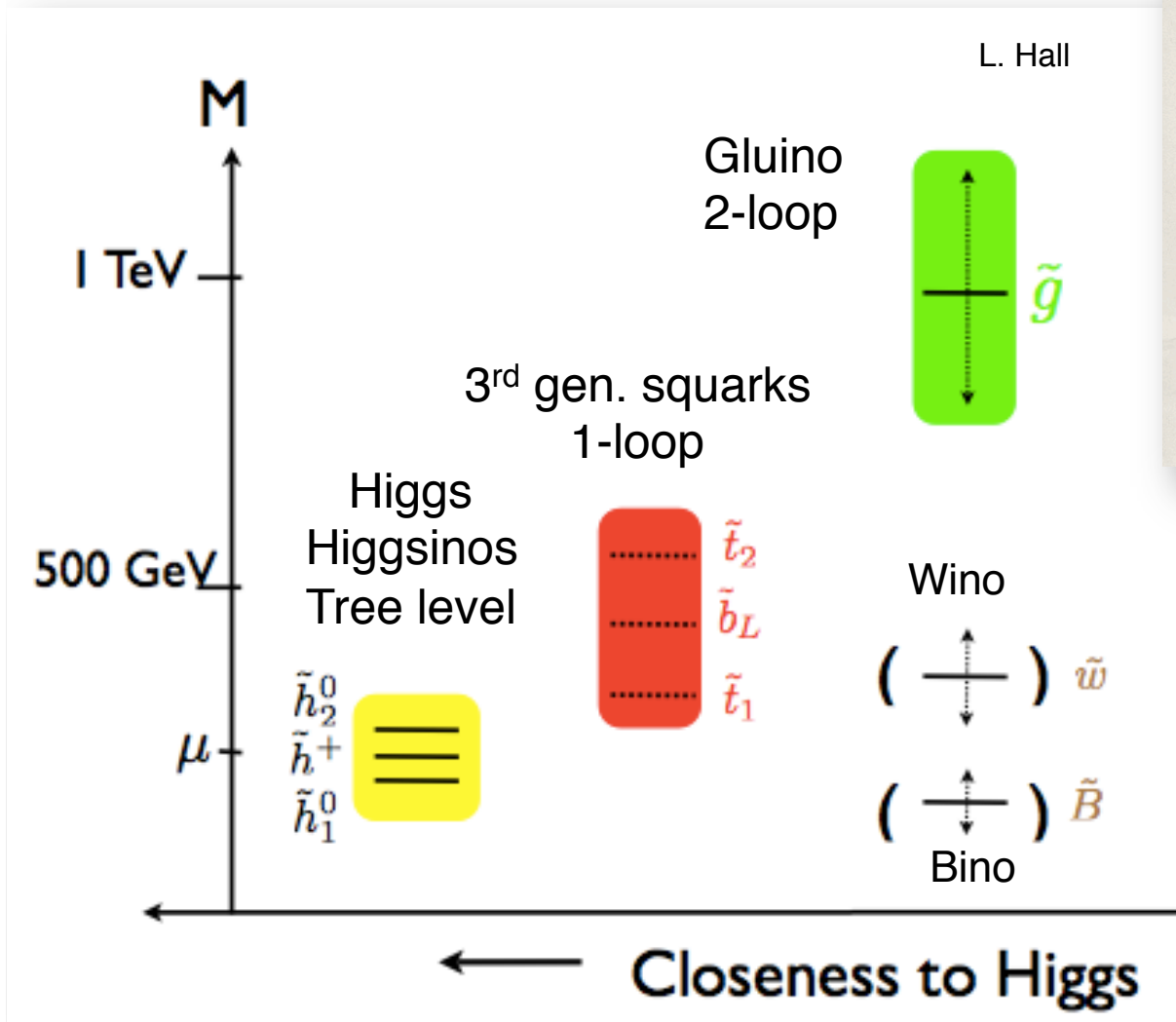


Example: $\tilde{g}\tilde{g} \rightarrow t\bar{t}\bar{t}\bar{t}\chi_1^0\chi_1^0$

Several searches are interpreted in this particular SMS!



(Minimal) Natural SUSY Spectrum



Use the argument of “naturalness” (i.e. fine-tuning) to motivate light **3rd generation squarks** (especially stop) and a rather light **gluino!**

More in Gian Giudice talk!

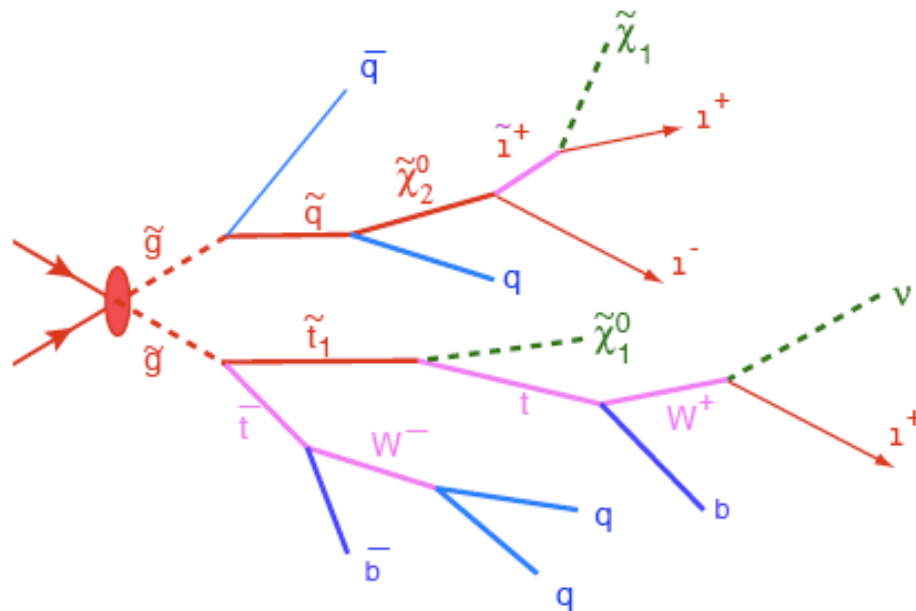
What do we call a “SUSY search”?

The definition is purely derived from the experimental signature.

Therefore, a “SUSY search signature” is characterized by

Lots of missing energy, many jets, and possibly leptons in the final state

Direct SUSY Searches, O. Buchmüller



Missing Energy:

- from LSP

Multi-Jet:

- from cascade decay (gaugino)

Multi-Leptons:

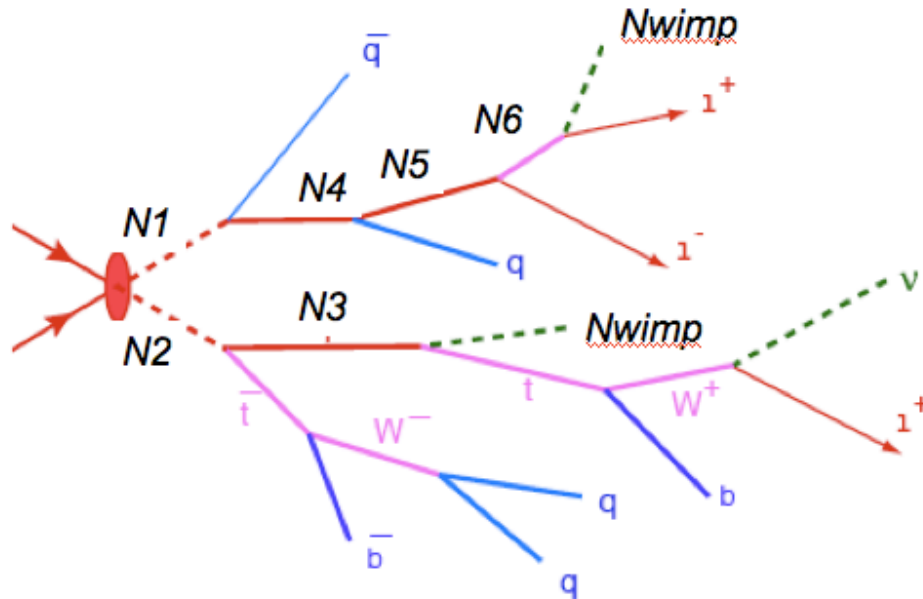
- from decay of charginos/neutralios

RP-Conserving SUSY is a very prominent example predicting this famous signature but ...

What is its experimental signature?

... by no means is it the only New Physics model predicting this experimental pattern. Many other NP models predict this genuine signature

Direct SUSY Searches, O. Buchmüller



Missing Energy:

- *Nwimp* - end of the cascade

Multi-Jet:

- from decay of the *N*s (possibly via heavy SM particles like top, *W/Z*)

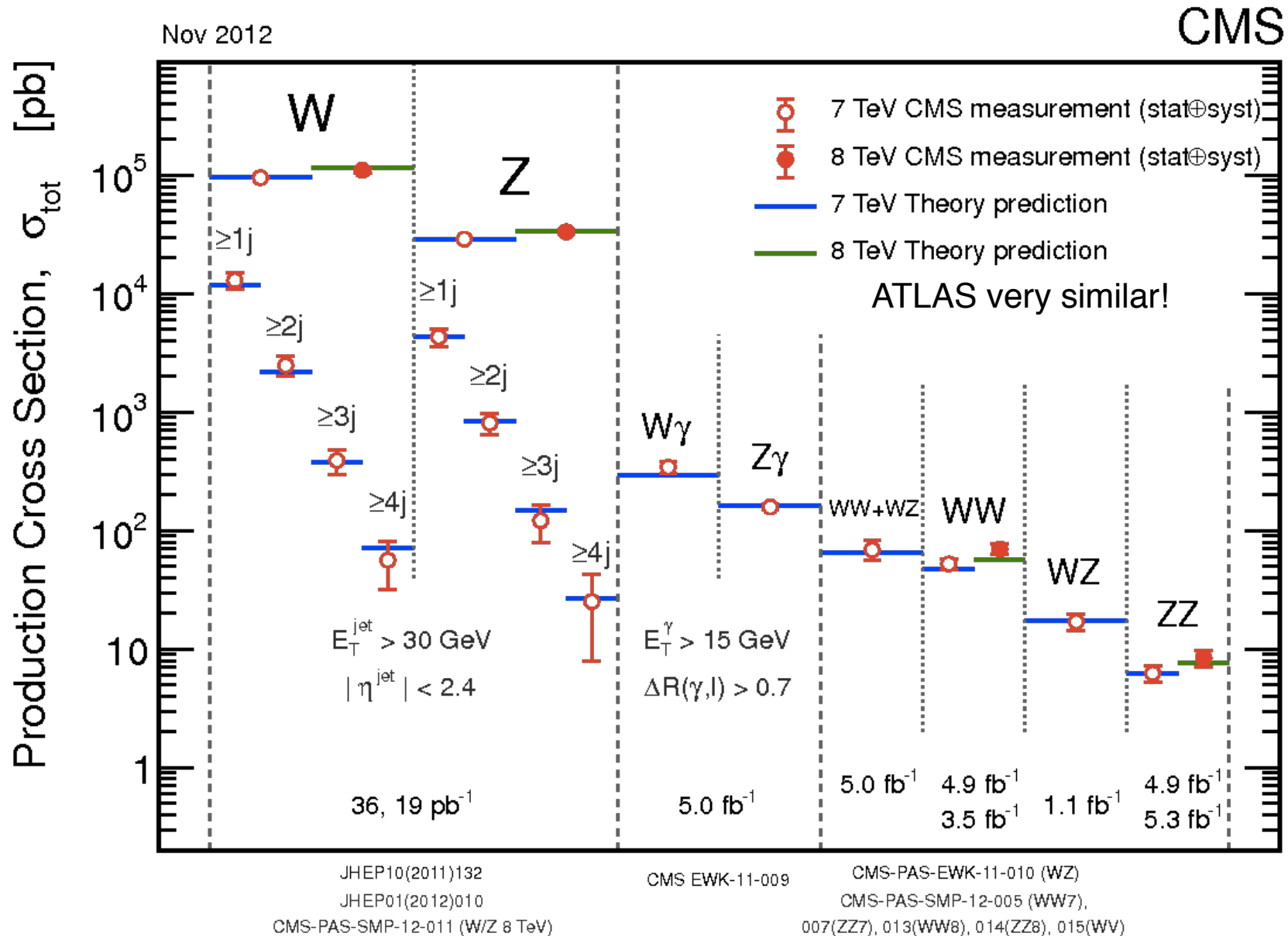
Multi-Leptons:

- from decay of the *N*'s

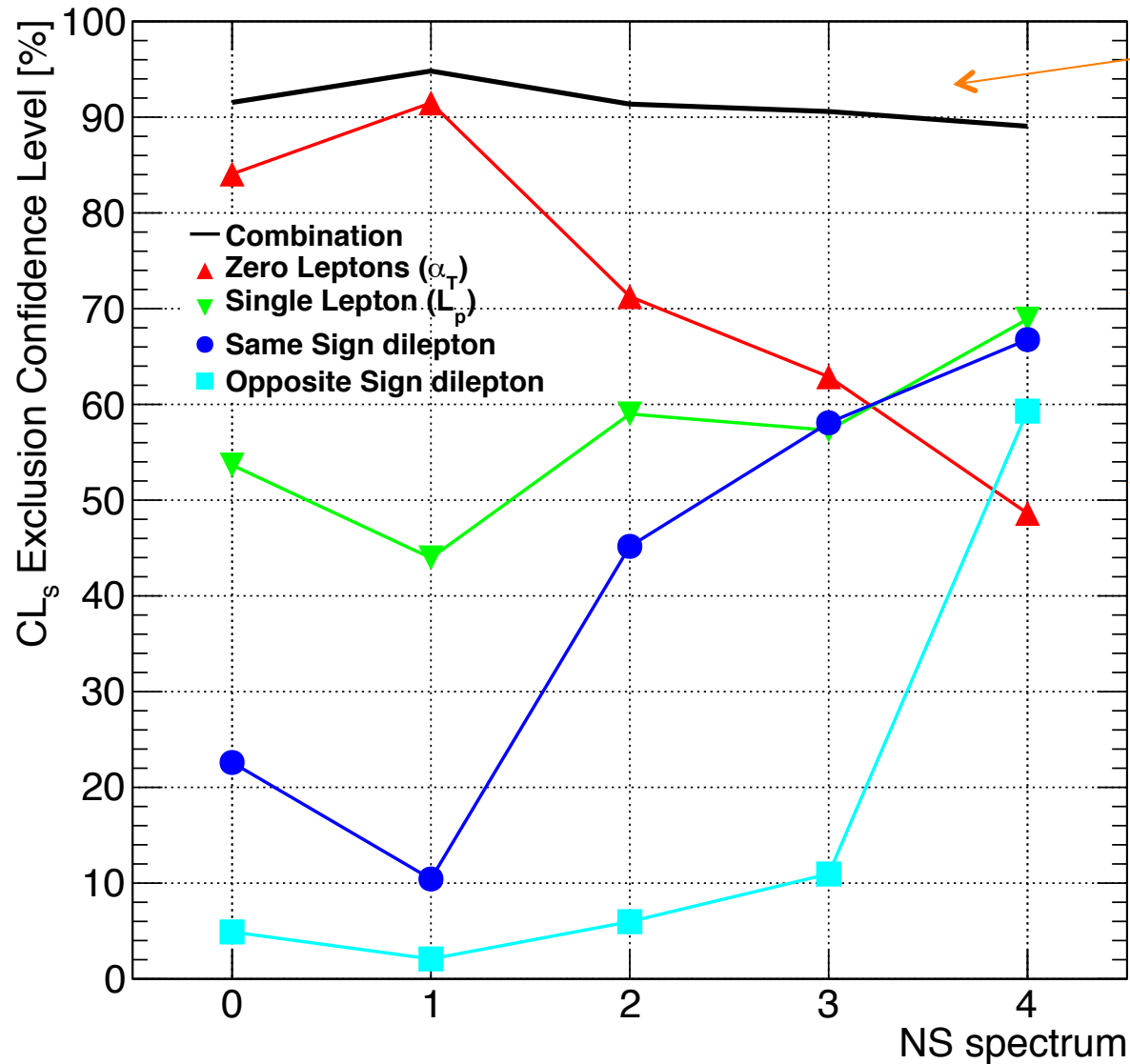
Model examples are Extra dimensions, Little Higgs, Technicolour, etc
but a more generic definition for this signature is as follows.

Rediscovery of the SM at a new energy frontier

Direct SUSY Searches, O. Buchmüller



Combination vs individual search



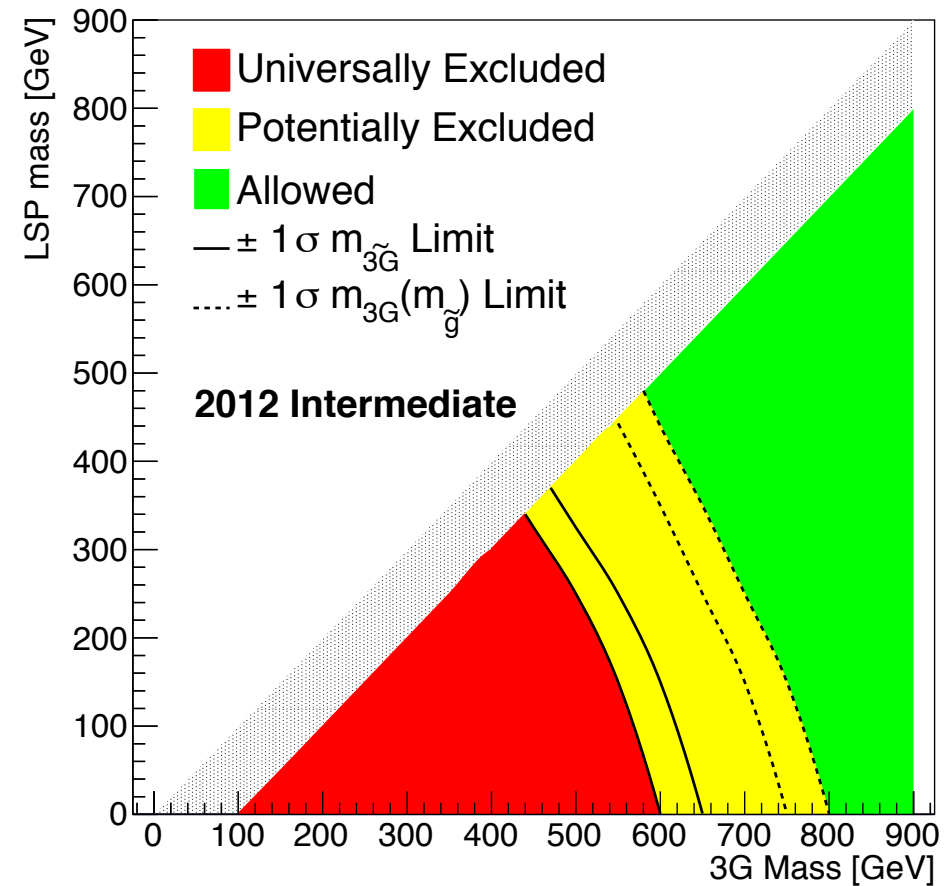
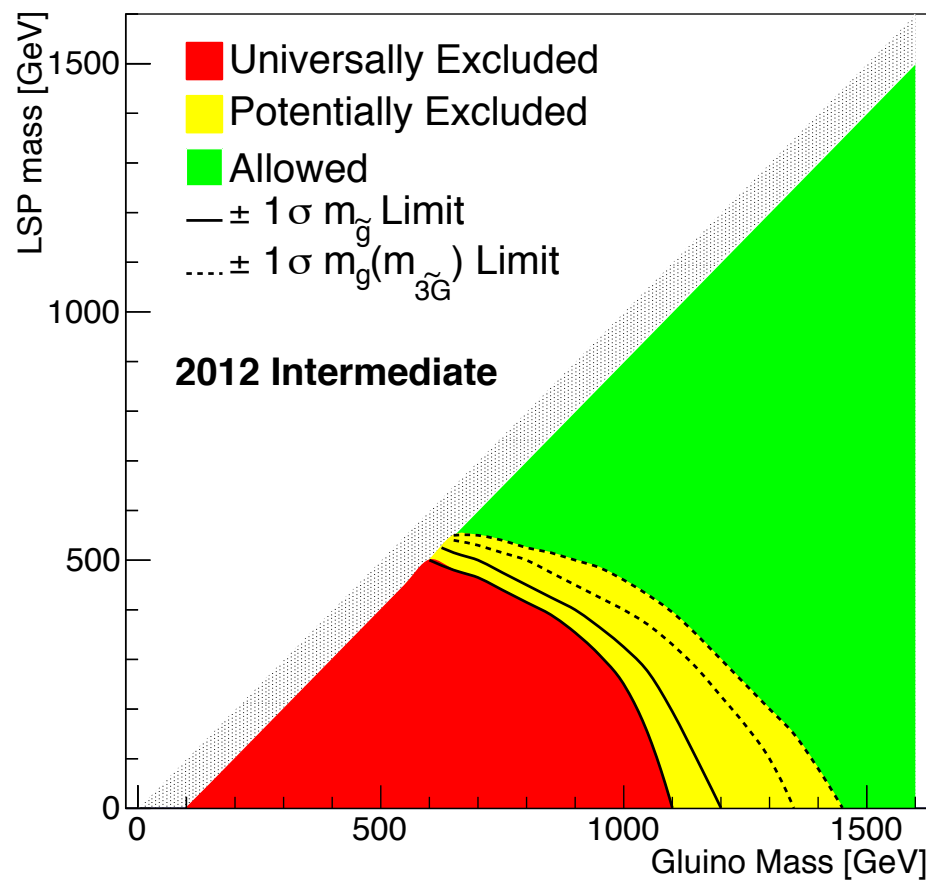
Combination of searches stable

Individual searches exhibit large variations

Combinations is stable vs. complexity while individual searches are NOT!

Natural SUSY: universal limits

If the gluino mass OR 3G mass lies in the red band, the point is excluded.
If the gluino mass AND 3G mass lie in the yellow band the point may or may not be excluded. Otherwise the point is not excluded.



Combining with the latest published 8 TeV results:

Outlook: 8 TeV vs 14 TeV

Use 30/fb for 2011/2012 for comparison

Direct SUSY Searches, O. Buchmüller

Higgs:

pp \rightarrow H, H \rightarrow WW, ZZ and $\gamma\gamma$
mainly gg: factor ~ 2

SUSY – 3rd Generation:

Mass scale \sim 500 GeV
qq and gg: factor ~ 3 to 6

SUSY – Squarks/Gluino:

Mass scale \sim 1.5 TeV
qq,gg,qq: factor ~ 40 to 80

Z' :

Mass scale \sim 5 TeV
qq: factor ~ 1000

Higgs:

15/fb@14 TeV to match 2011/2012
mainly gg: factor ~ 2

SUSY – 3rd Generation:

5/fb to 10/fb@14 TeV to match 2011/2012
qq and gg: factor ~ 3 to 6

SUSY – Squarks/Gluino:

0.4/fb to 0.8/fb@14 TeV to match 2011/2012
qq,gg,qq: factor ~ 40 to 80

Z' :

0(1/pb) @14 TeV to match 2011/2012
qq: factor ~ 1000

RP-violating searches/interpretation

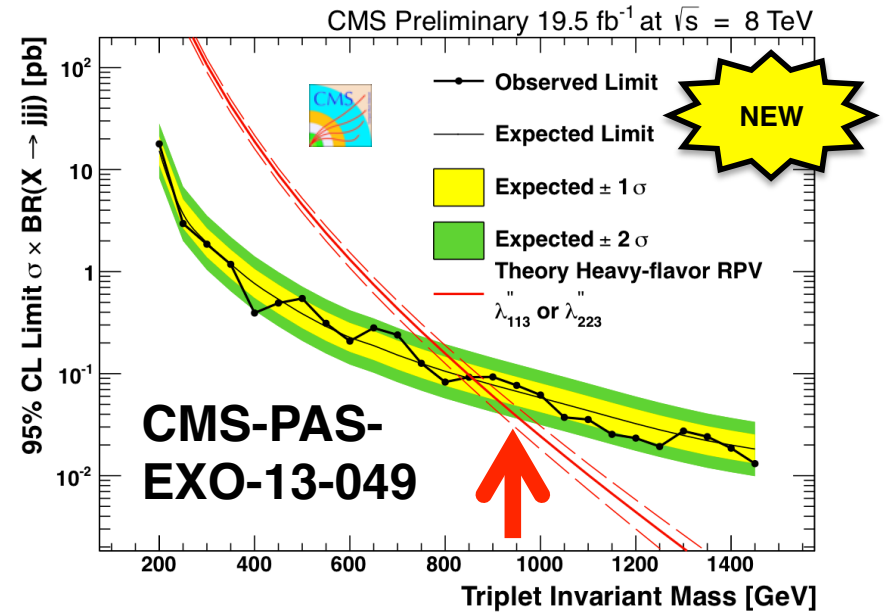
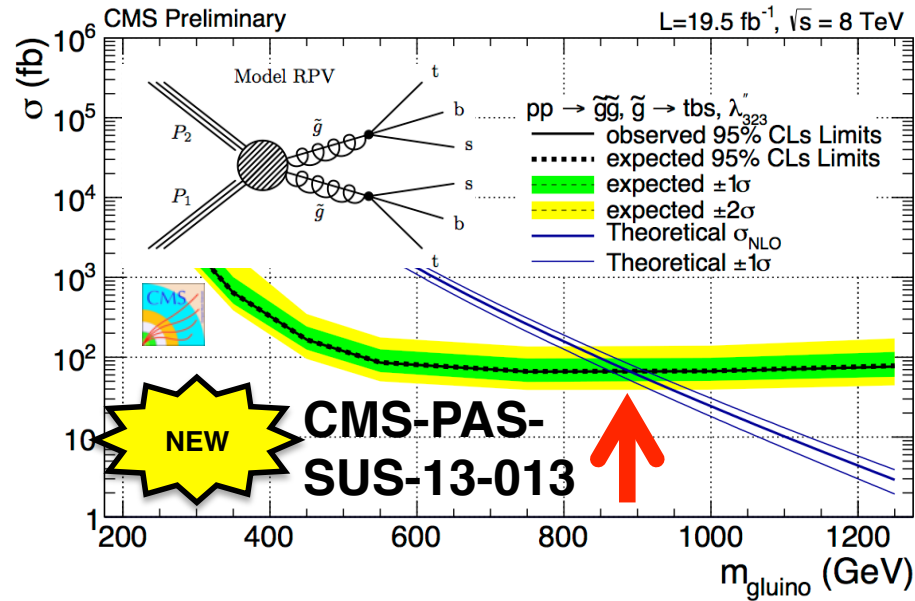
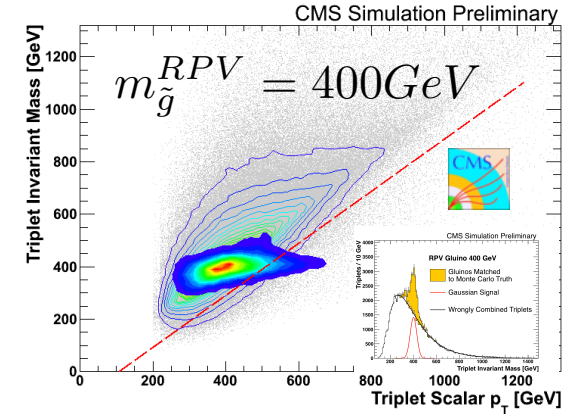
Searches, O. Buchmüller

Generic same-sign di-lepton search with different signal regions

SR	Expected	Observed
RPV0	38 ± 14	35
RPV2	5.3 ± 2.1	5
SStop1	160 ± 59	152
SStop1++	90 ± 32	92
SStop2	40 ± 13	52
SStop2++	22 ± 8	25

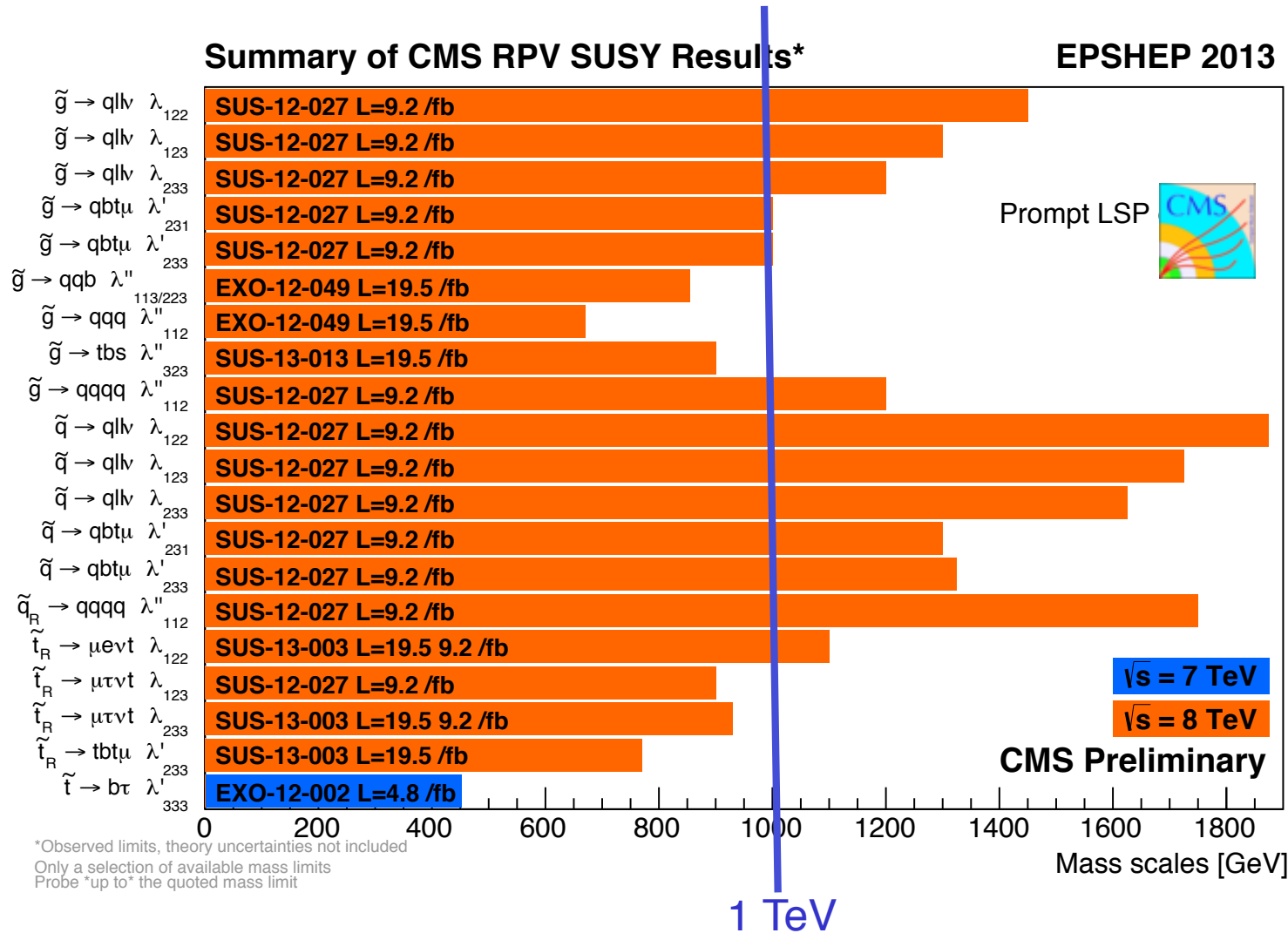
$$\tilde{g} \rightarrow 3 \text{ jets}$$

Take all triplets,
QCD: $M_{3j} \sim \Sigma P_T^j$;
SUSY: $M_{3j} \sim M_g$
 $M_{3j} < \Sigma P_T^j - 160 \text{ GeV}$



RP violation searches: Summary

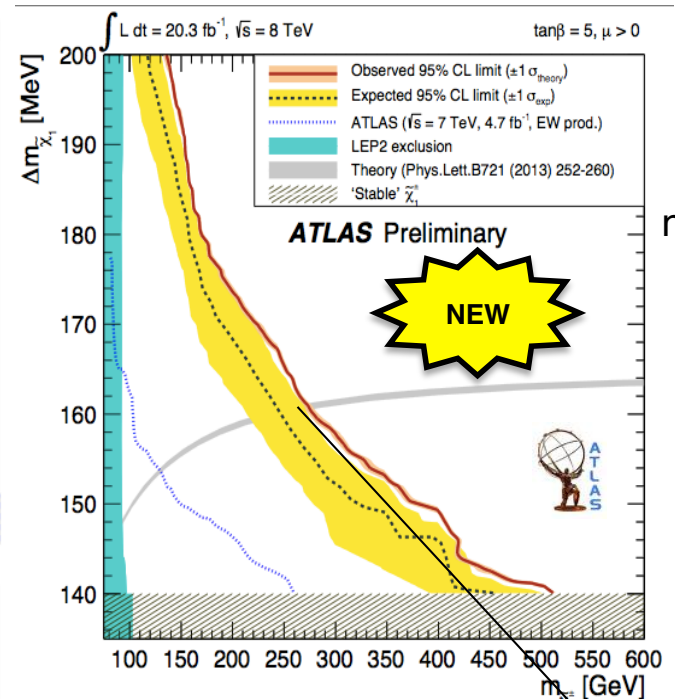
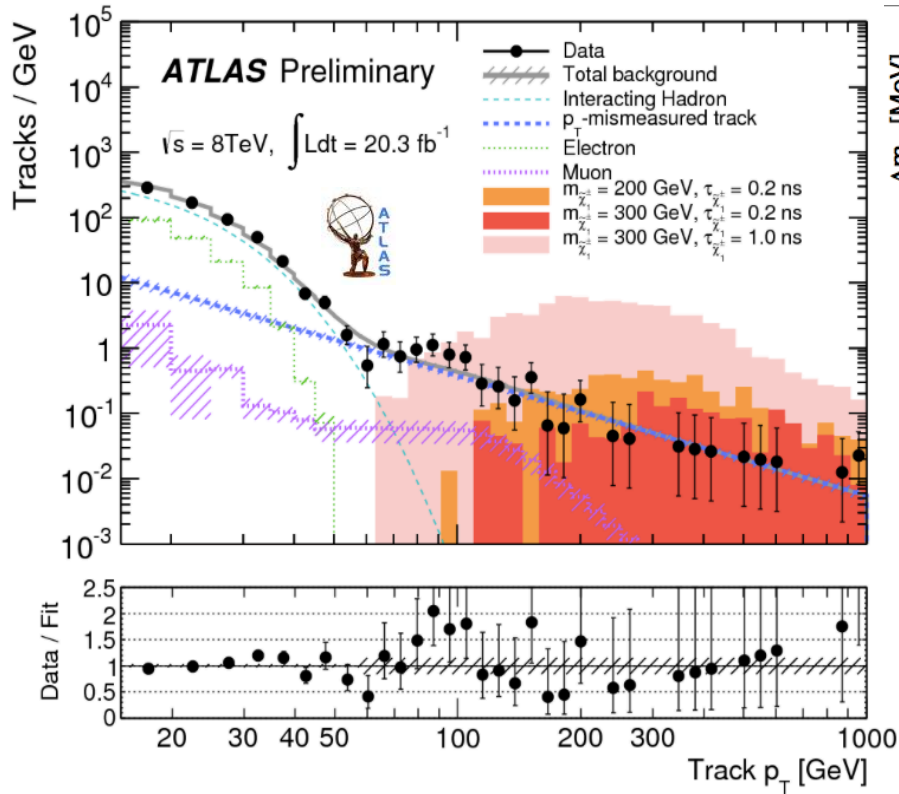
Direct SUSY Searches, O. Buchmüller



ATLAS
similar

Like RP conserving searches, these searches are also probing the 1 TeV scale and even beyond!

Long-lived particle (SUSY) searches



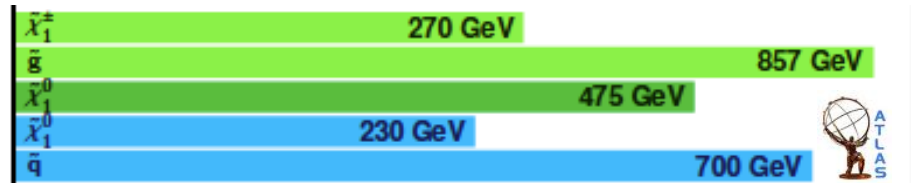
**ATLAS-CONF
-2013-069**
nearly mass-degenerate

$$\chi_1^\pm \chi_1^0$$

search based on
disappearing-track
signature

CMS similar
See parallel talk from
L. Quertenmont

Long-lived particles	Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Stable, stopped \tilde{g} R-hadron	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1^0$	$\tilde{\chi}_1^0 \rightarrow qq\mu$ (RPV)	Disapp. trk	1 jet	Yes	20.3
					0	1-5 jets	Yes	22.9	
					1-2 μ	0	-	15.9	
					2 γ	0	Yes	4.7	
					1 μ	0	Yes	4.4	



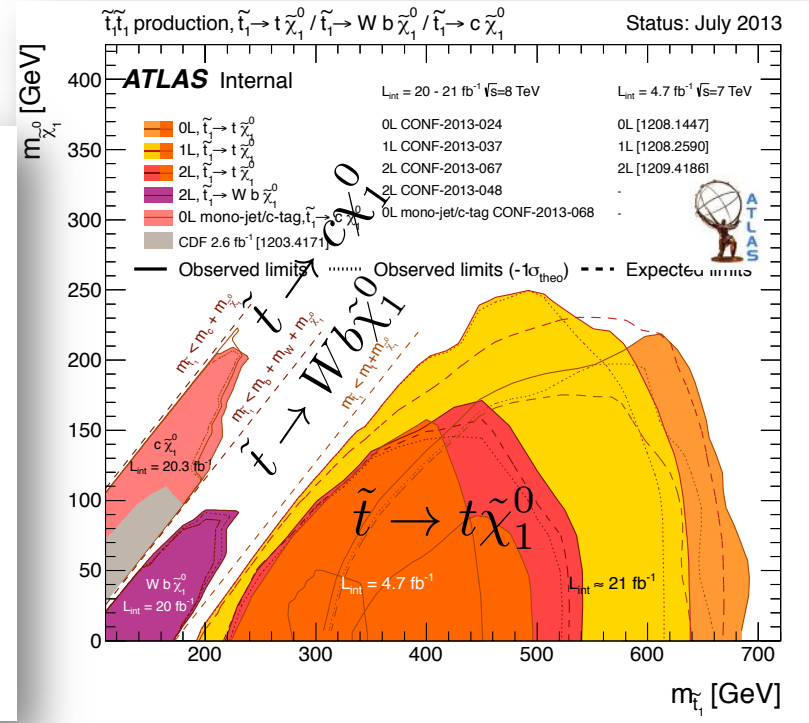
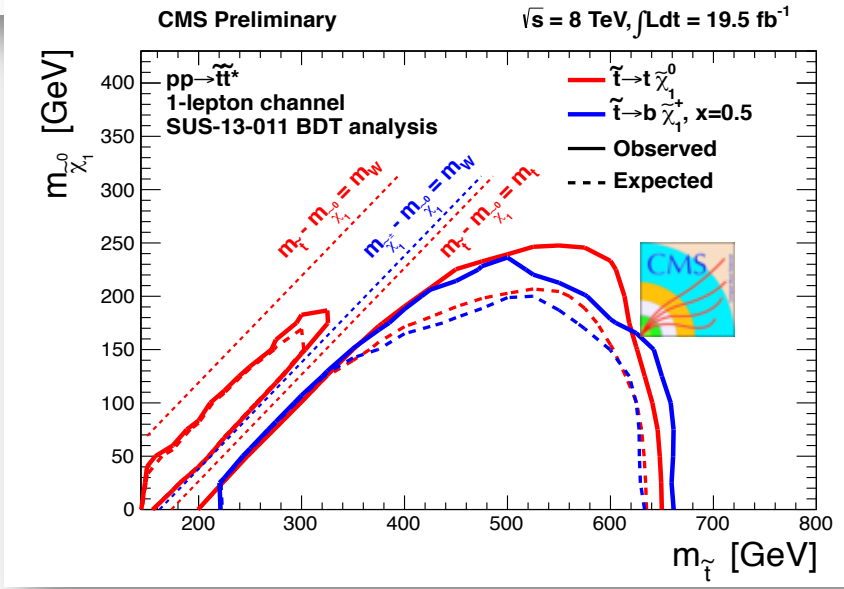
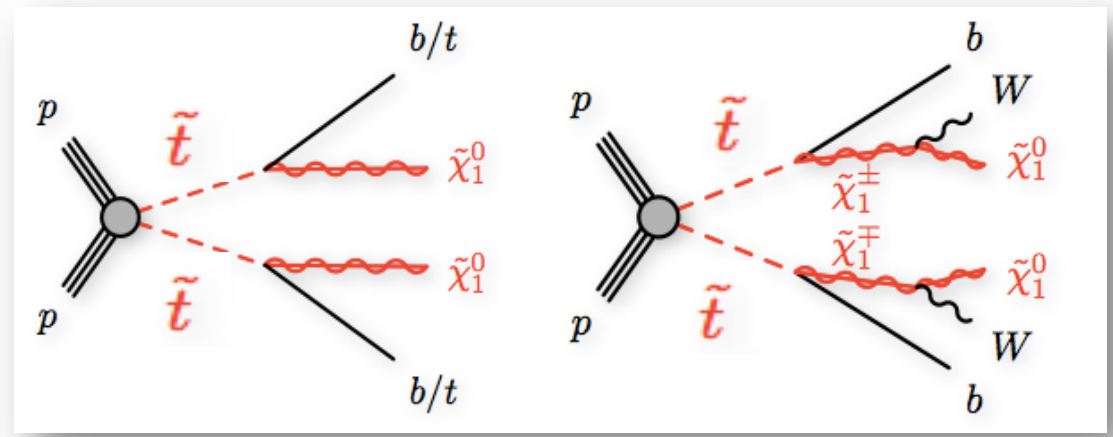
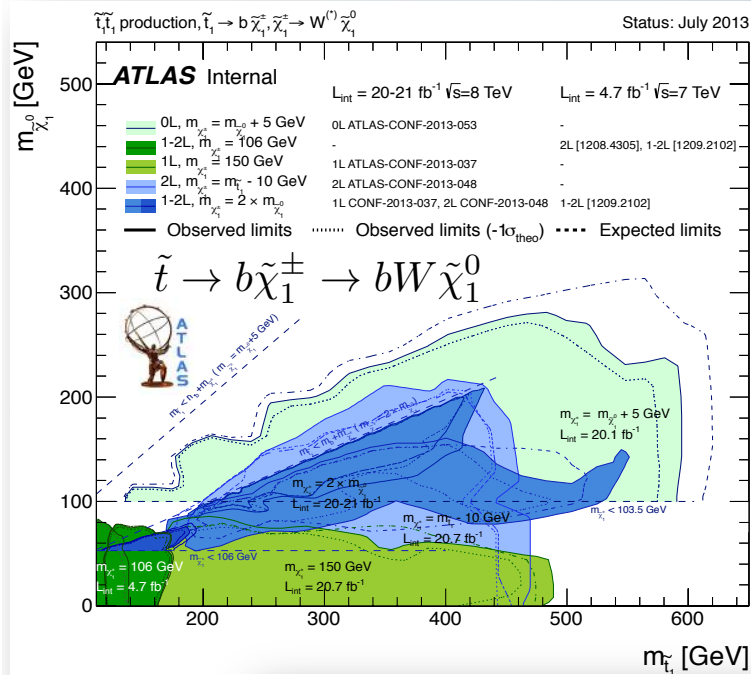
$m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm) = 0.2 \text{ ns}$
 $m(\tilde{\chi}_1^0) = 100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$
 $10 < \tan\beta < 50$
 $0.4 < \tau(\tilde{\chi}_1^0) < 2 \text{ ns}$
 $1 \text{ mm} < c\tau < 1 \text{ m}, \tilde{g} \text{ decoupled}$

ATLAS-CONF-2013-069
 ATLAS-CONF-2013-057
 ATLAS-CONF-2013-058
 1304.6310
 1210.7451

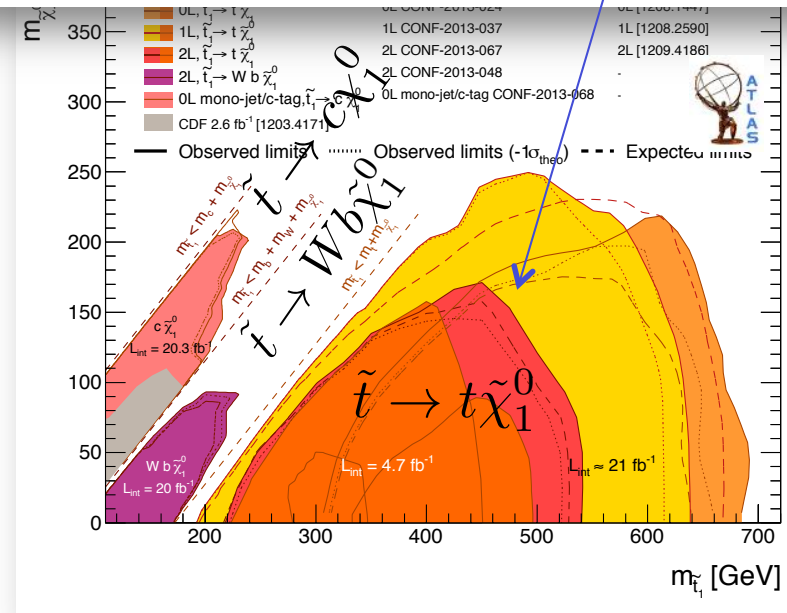
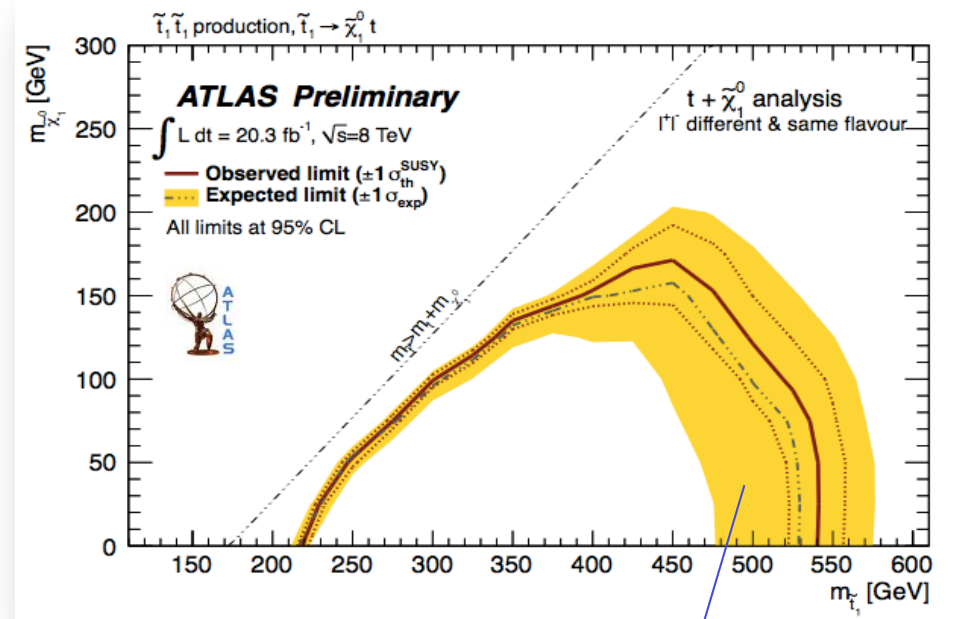
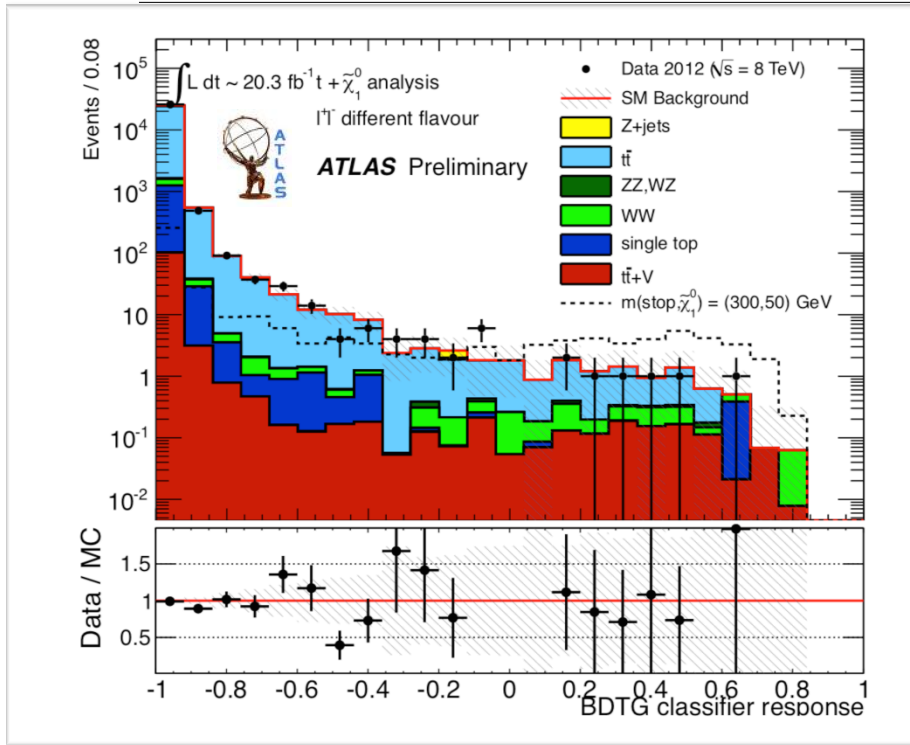
**About to probe
the 1 TeV scale**

1 TeV

Dedicated searches for direct stop-pair production



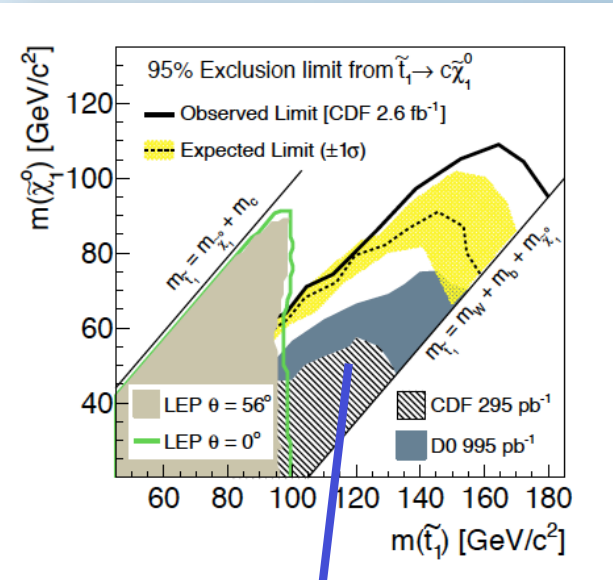
Dedicated searches for direct stop-pair production



ATLAS-CONF-2013-065:
 Scalar stop analysis with two leptons in the
 final state using a MVA technique.

Dedicated searches for direct stop-pair production

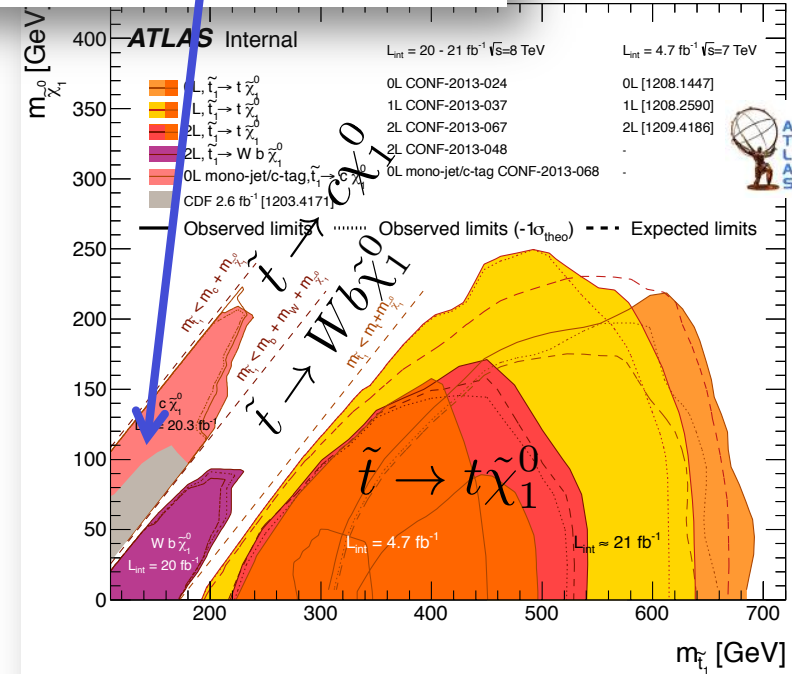
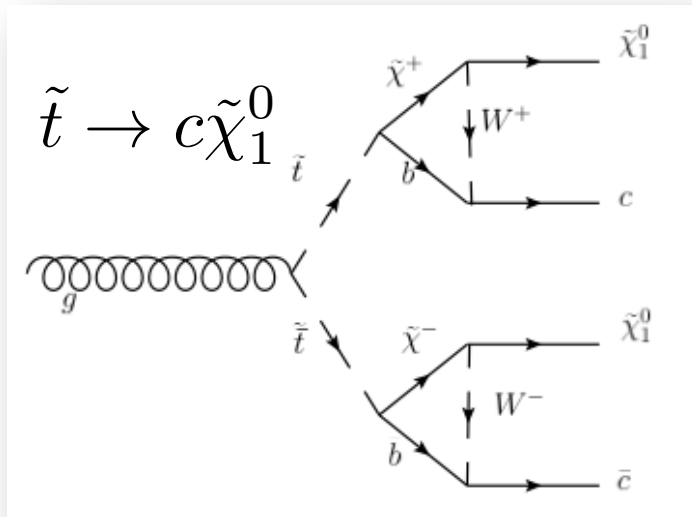
Direct SUSY Searches, O. Buchmüller



LEP:
LEPSUSY
WG/04-01.1

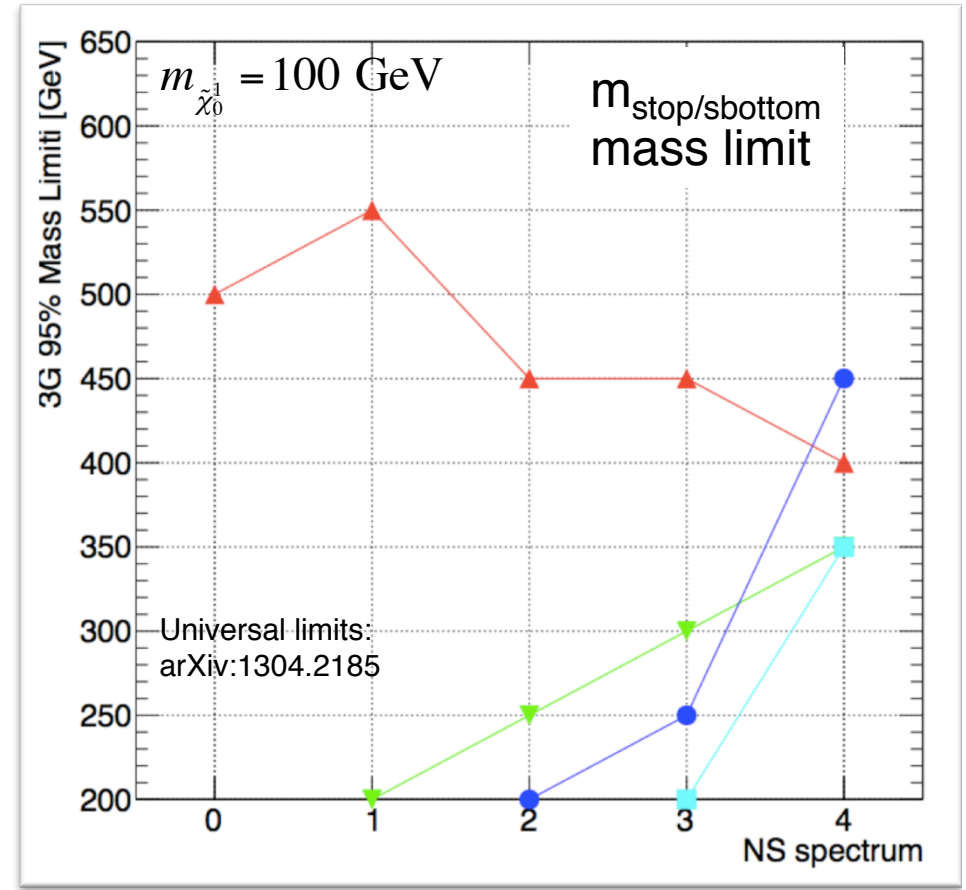
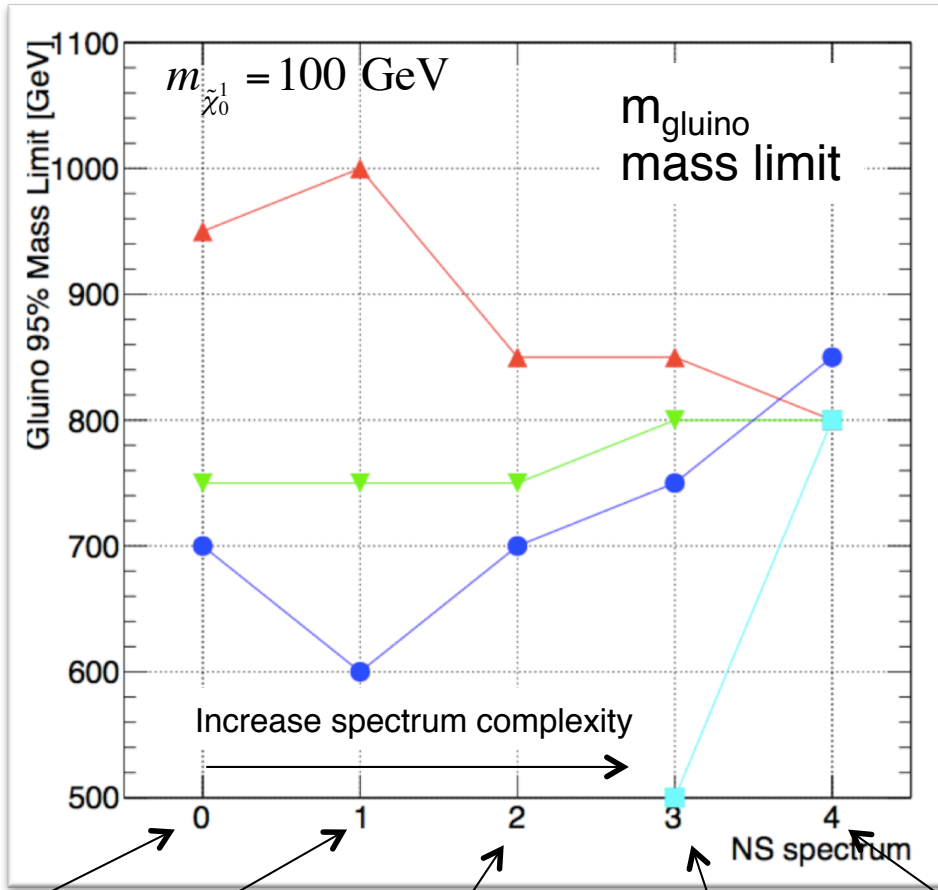
CDF:
1203.4171
D0:
0803.2263

Status: July 2013



SMS limits: A word of caution!

Direct SUSY Searches, O. Buchmüller



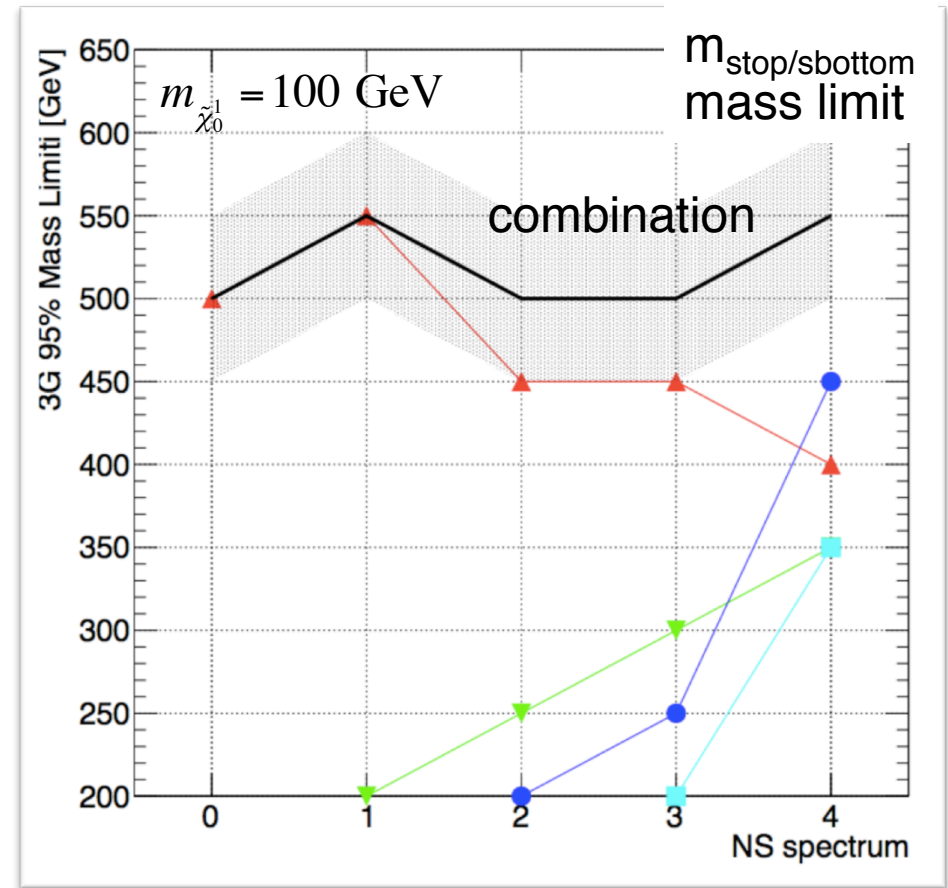
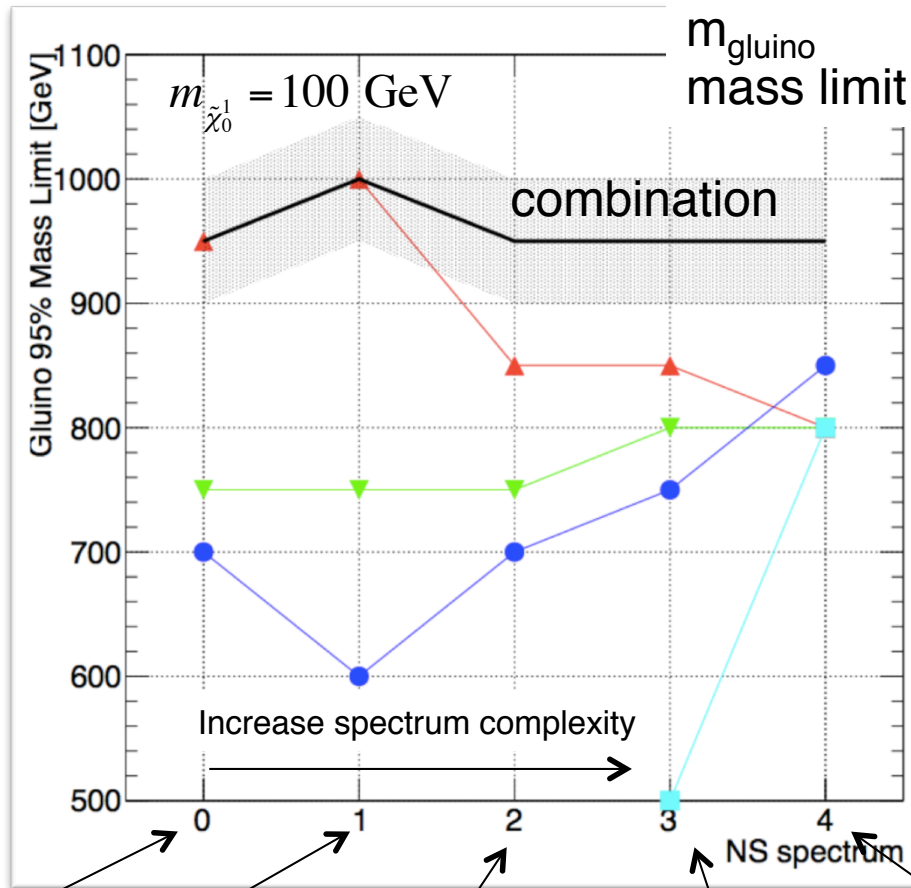
NS0	NS1	NS2	NS3	NS4
\tilde{g}	\tilde{g}	\tilde{g}	\tilde{g}	\tilde{g}
\tilde{t}_1, \tilde{t}_2	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1, \tilde{b}_2$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1, \tilde{b}_2$
		$\tilde{\chi}_0^2$	$\tilde{\chi}_0^2$	$\tilde{\chi}_0^2$
		$\tilde{\chi}^\pm$	$\tilde{\chi}^\pm$	$\tilde{\chi}^\pm, \tilde{\ell}_{L,R}$
$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$

Used inclusive searches from 2011:

- 0-Lepton CMS-SUS-11-022
- 1-Lepton CMS-SUS-12-010
- 2-Lepton SS CMS-SUS-11-010
- 2-Lepton OS CMS-SUS-11-011

Combining Searches = less model dependence

Direct SUSY Searches, O. Buchmüller



NS0	NS1	NS2	NS3	NS4
\tilde{g}	\tilde{g}	\tilde{g}	\tilde{g}	\tilde{g}
\tilde{t}_1, \tilde{t}_2	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1, \tilde{b}_2$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1, \tilde{b}_2$
		$\tilde{\chi}_0^2$	$\tilde{\chi}_0^2$	$\tilde{\chi}_0^2$
		$\tilde{\chi}^\pm$	$\tilde{\chi}^\pm$	$\tilde{\chi}^\pm, \tilde{\ell}_{L,R}$
$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$

Used inclusive searches from 2011:

0-Lepton CMS-SUS-11-022

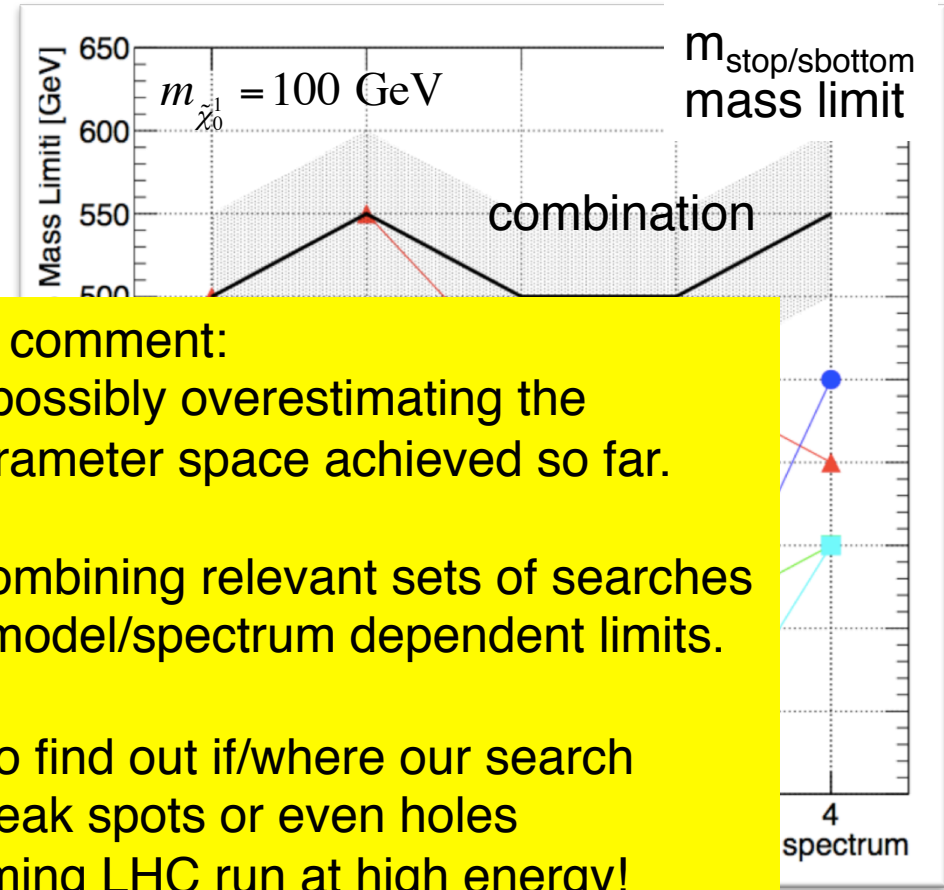
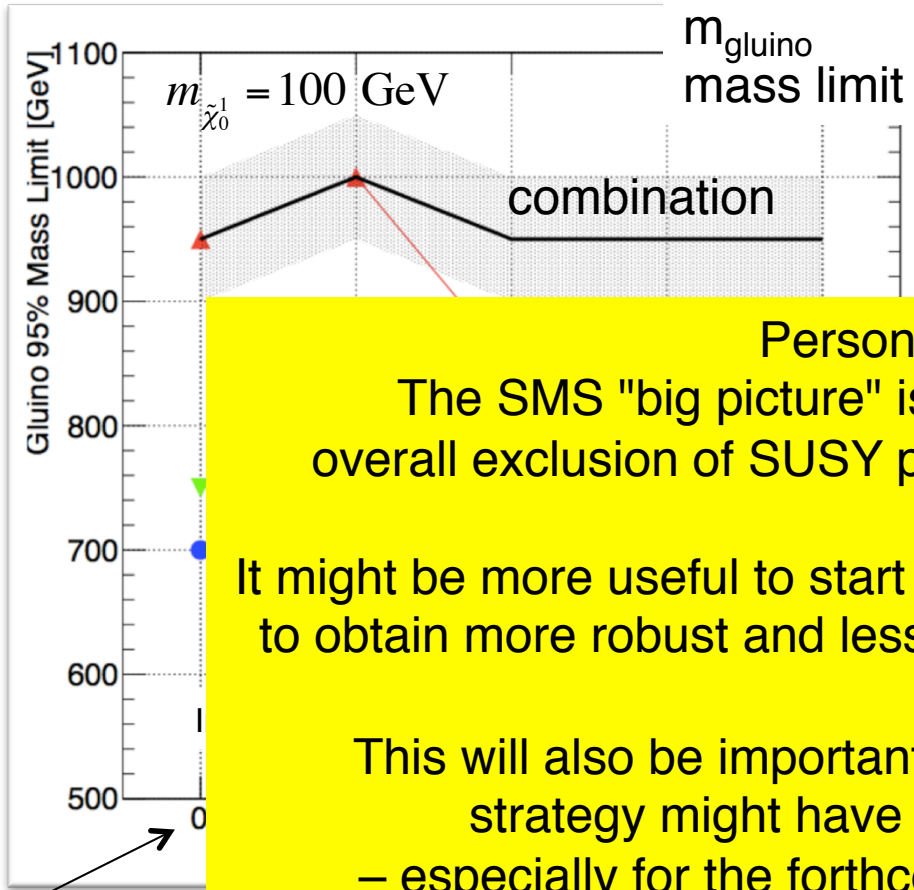
1-Lepton CMS-SUS-12-010

2-Lepton SS CMS-SUS-11-010

2-Lepton OS CMS-SUS-11-011

Combining Searches = less model dependence

Direct SUSY Searches, O. Buchmüller



Personal comment:
 The SMS "big picture" is possibly overestimating the overall exclusion of SUSY parameter space achieved so far.
 It might be more useful to start combining relevant sets of searches to obtain more robust and less model/spectrum dependent limits.
 This will also be important to find out if/where our search strategy might have weak spots or even holes – especially for the forthcoming LHC run at high energy!

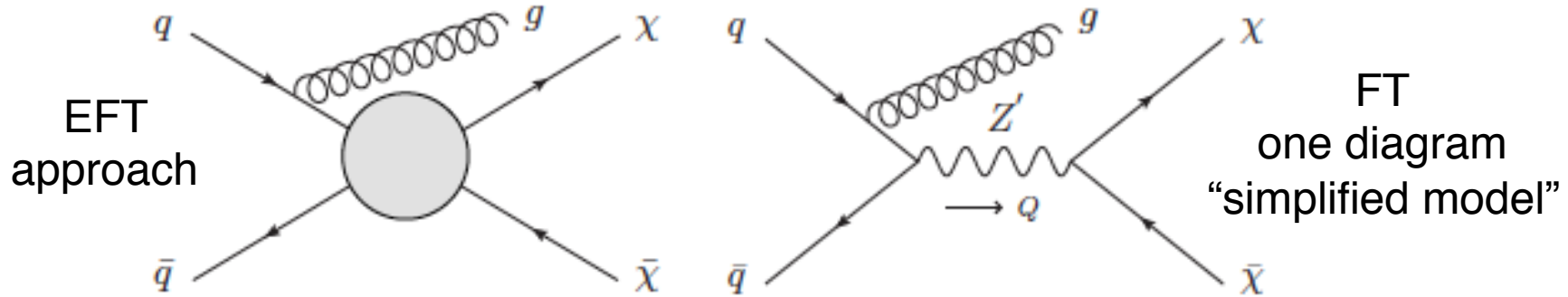
NS0				
\tilde{g}	\tilde{g}	\tilde{g}	\tilde{g}	\tilde{g}
\tilde{t}_1, \tilde{t}_2	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1, \tilde{b}_2$	$\tilde{t}_1, \tilde{t}_2, \tilde{b}_1, \tilde{b}_2$
		$\tilde{\chi}_0^2$	$\tilde{\chi}_0^2$	$\tilde{\chi}_0^2$
		$\tilde{\chi}^\pm$	$\tilde{\chi}^\pm$	$\tilde{\chi}^\pm, \tilde{\ell}_{L,R}$
$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$	$\tilde{\chi}_0^1$

- 2011:
- 0-Lepton CMS-SUS-11-022
 - 1-Lepton CMS-SUS-12-010
 - 2-Lepton SS CMS-SUS-11-010
 - 2-Lepton OS CMS-SUS-11-011

Validity of Effective Field Theory Limits

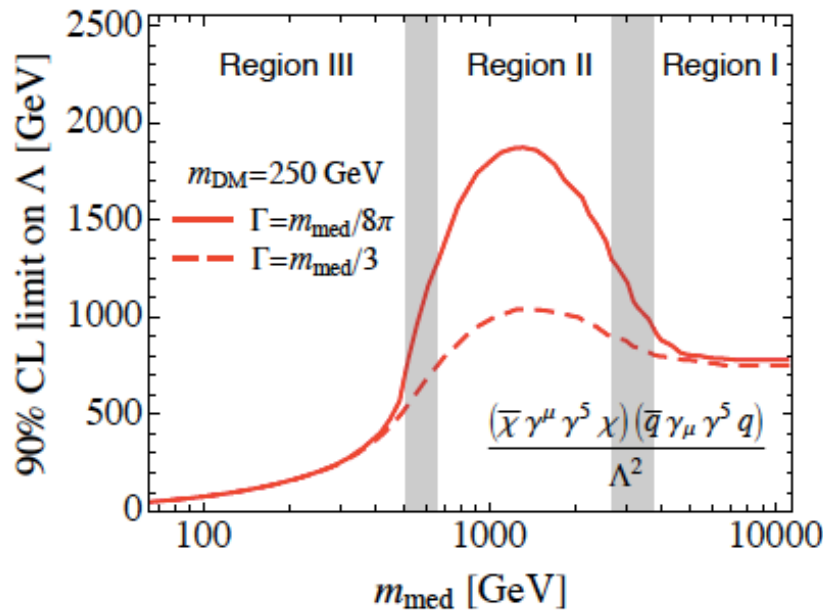
Recent work from OB, M.Dolan, C.McCabe: arXiv:1308.6799

➤ Compare Effective Field Theory (EFT) with Full Theory (FT)



Use vector and axial-vector mediators (e.g. Z') as example - scalar are similar in conclusion!

Direct SUSY Searches, O. Buchmüller



Three Regions as function of mediator mass:

Region I: Heavy m_{med}

➤ EFT is valid!

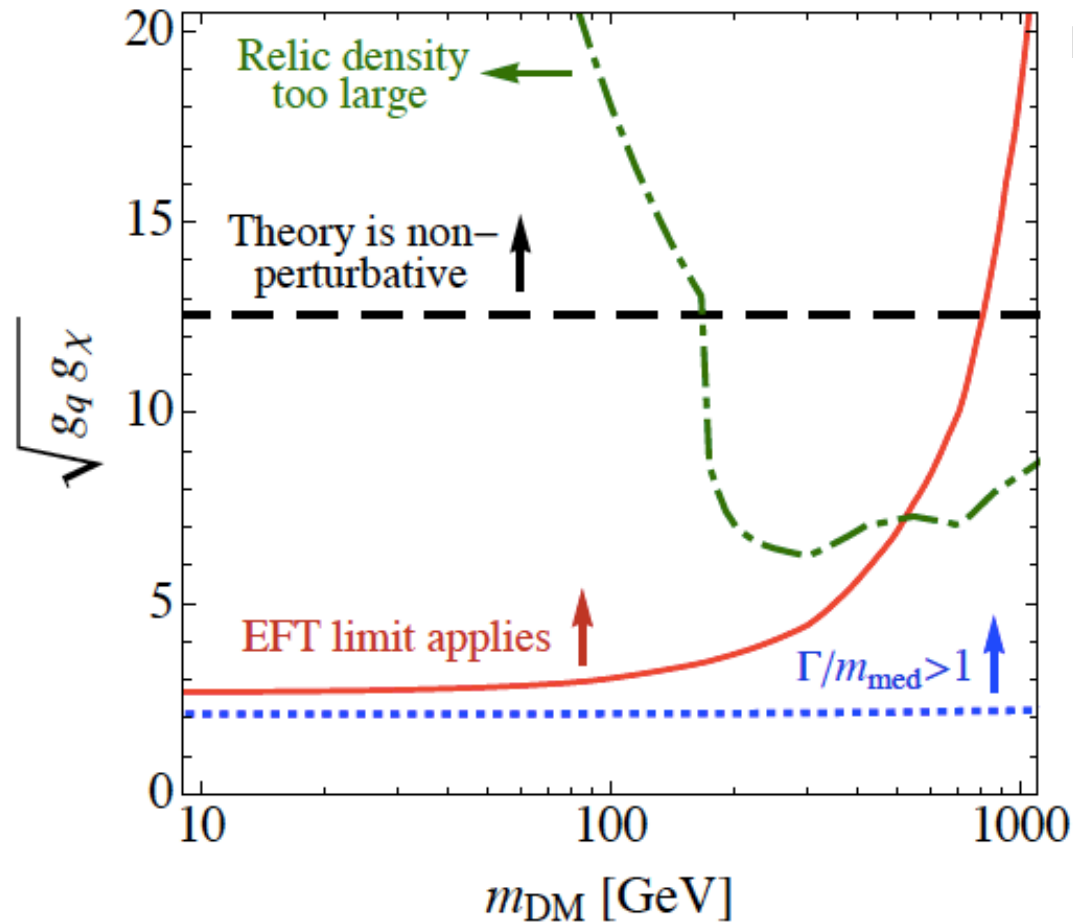
Region II: Medium m_{med} – Resonant enhancement

➤ EFT limits are too conservative!

Region III: Low m_{med}

➤ EFT limits are too aggressive!

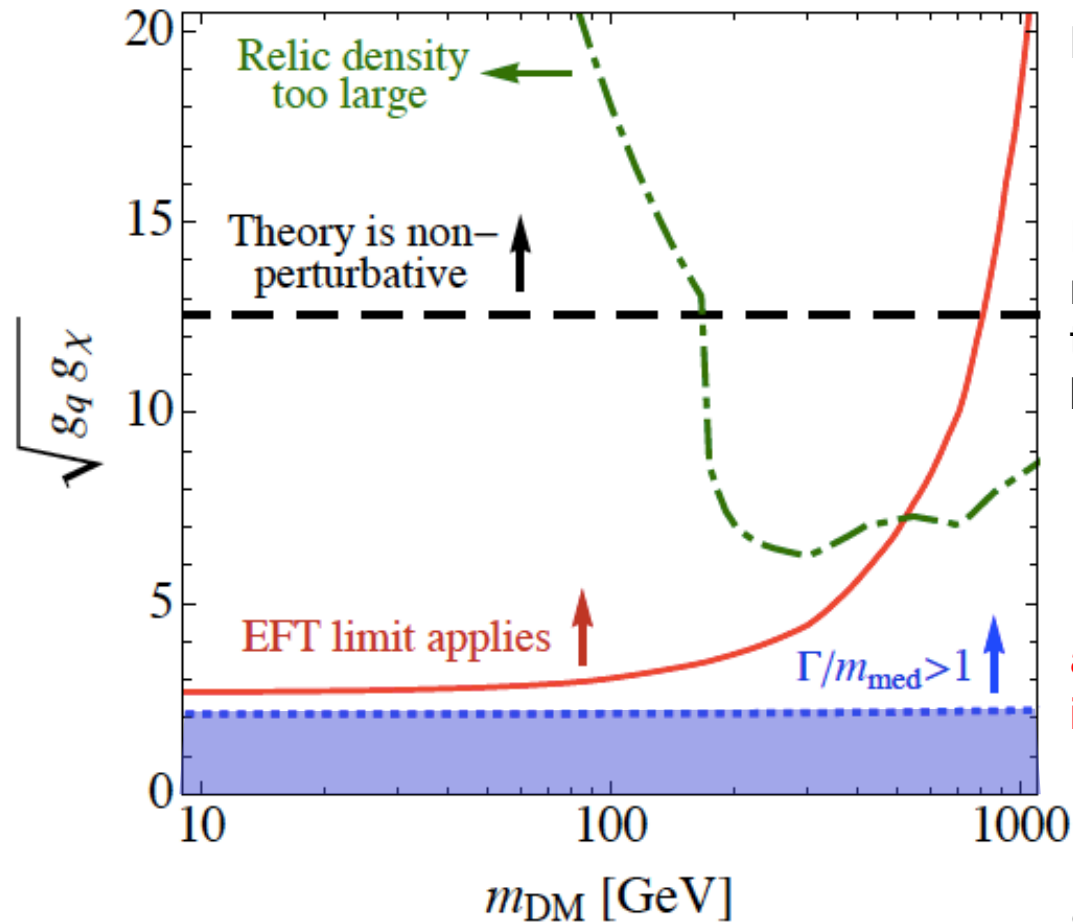
What those this imply on model-dependences of EFT limits?



Look at EFT validity in $m_{\text{DM}} - \text{coupling}^*$ plane!

* Coupling chose such that CMS EFT limit on Λ applies to FT

What those this imply on model-dependences of EFT limits?



Look at EFT validity in $m_{\text{DM}} - \text{coupling}^*$ plane!

1. Must require $m_{\text{med}} < \Gamma_{\text{med}}$

First we define the region in which mediator mass is larger than its width. Therefore, only this region a well-defined theory of PF can be straightforwardly established.

In the remaining part of the plot:

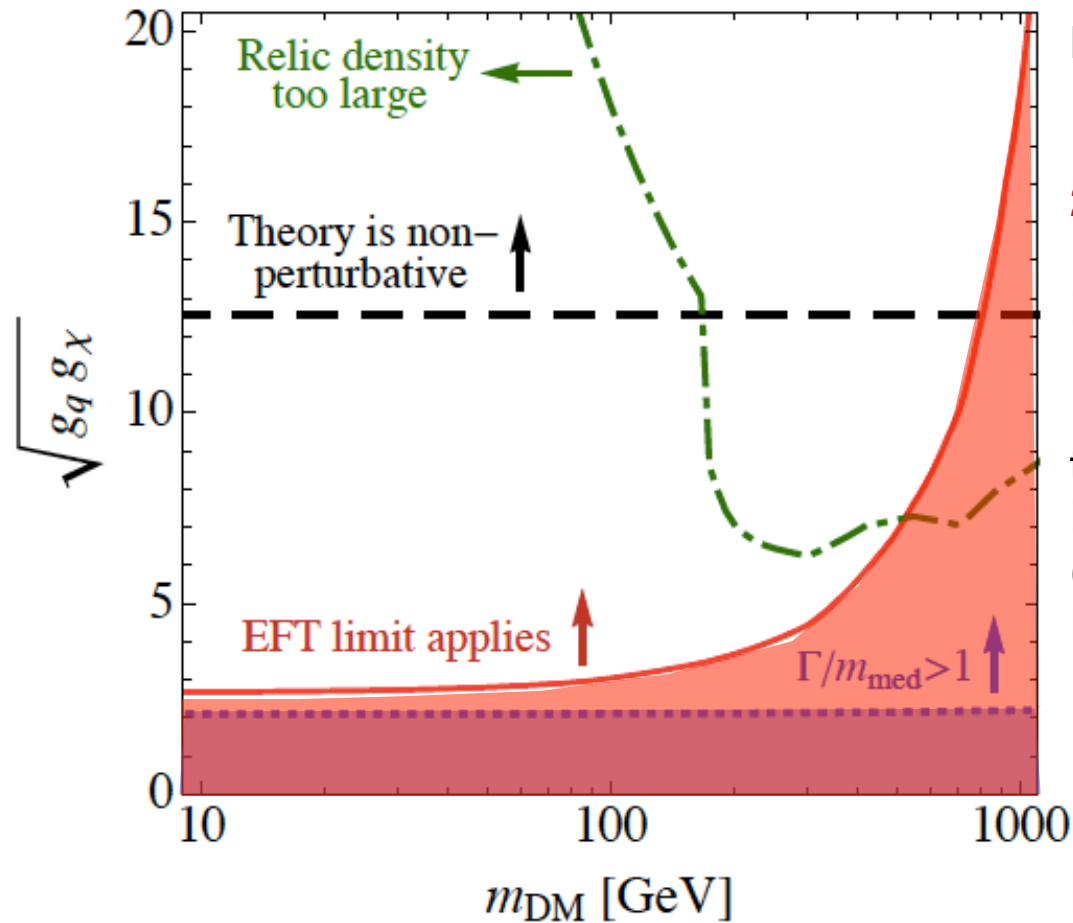
$$\sqrt{g_q g_\chi} > 2$$

a particle-like interpretation of the mediator is doubtful!

See discussion about equation 3.5 in arXiv:1308.6799 for further details.

* Coupling chose such that CMS EFT limit on Λ applies to FT

What those this imply on model-dependences of EFT limits?



Look at EFT validity in $m_{DM} - \text{coupling}^*$ plane!

1. Must require $m_{med} < \Gamma_{med}$
2. Region in which EFT is valid

For this we calculate the minimum coupling

$$\sqrt{g_q g_\chi} = m_{med} / \Lambda_{CMS}$$

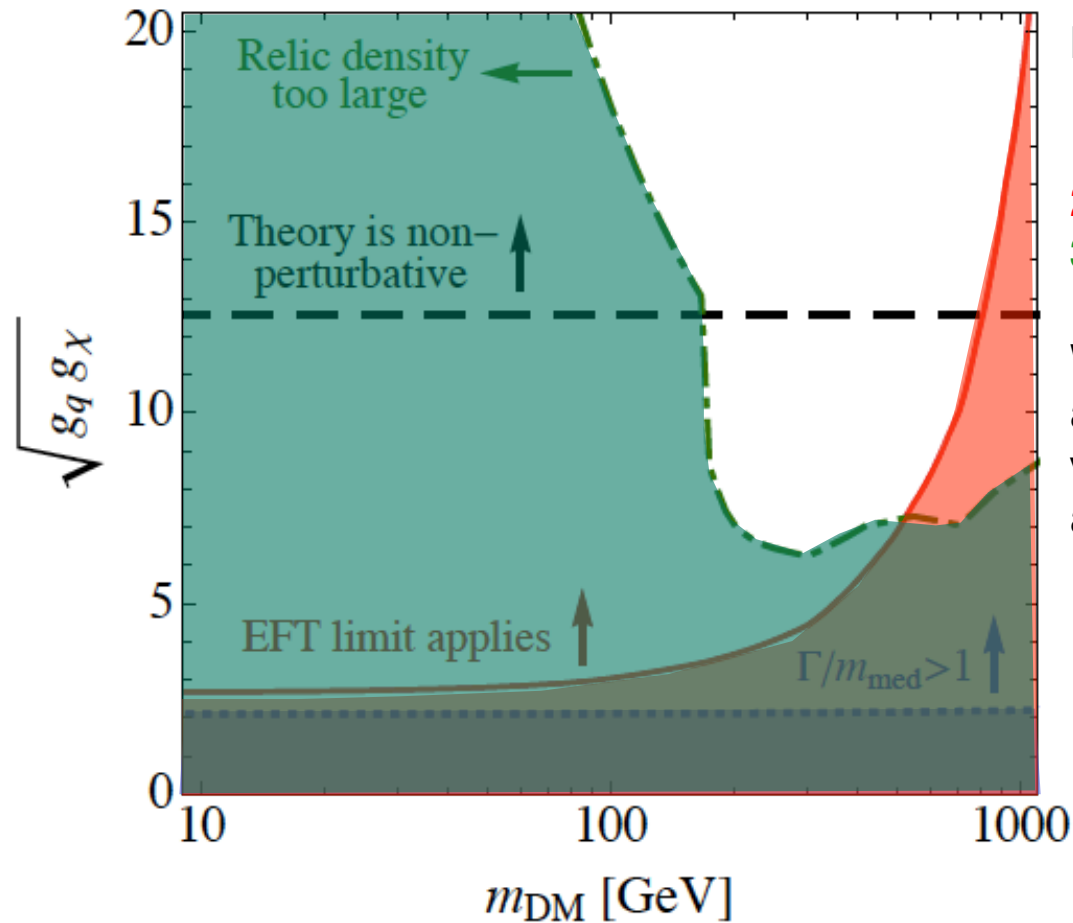
that the simplified model must have for the EFT limits to apply. This is defined by region I (i.e. better than 20% agreement of FT and EFT).

First important observation:

The EFT is only valid in the region in which mediator width is larger than mediator mass!

* Coupling chose such that CMS EFT limit on Λ applies to FT

What those this imply on model-dependences of EFT limits?



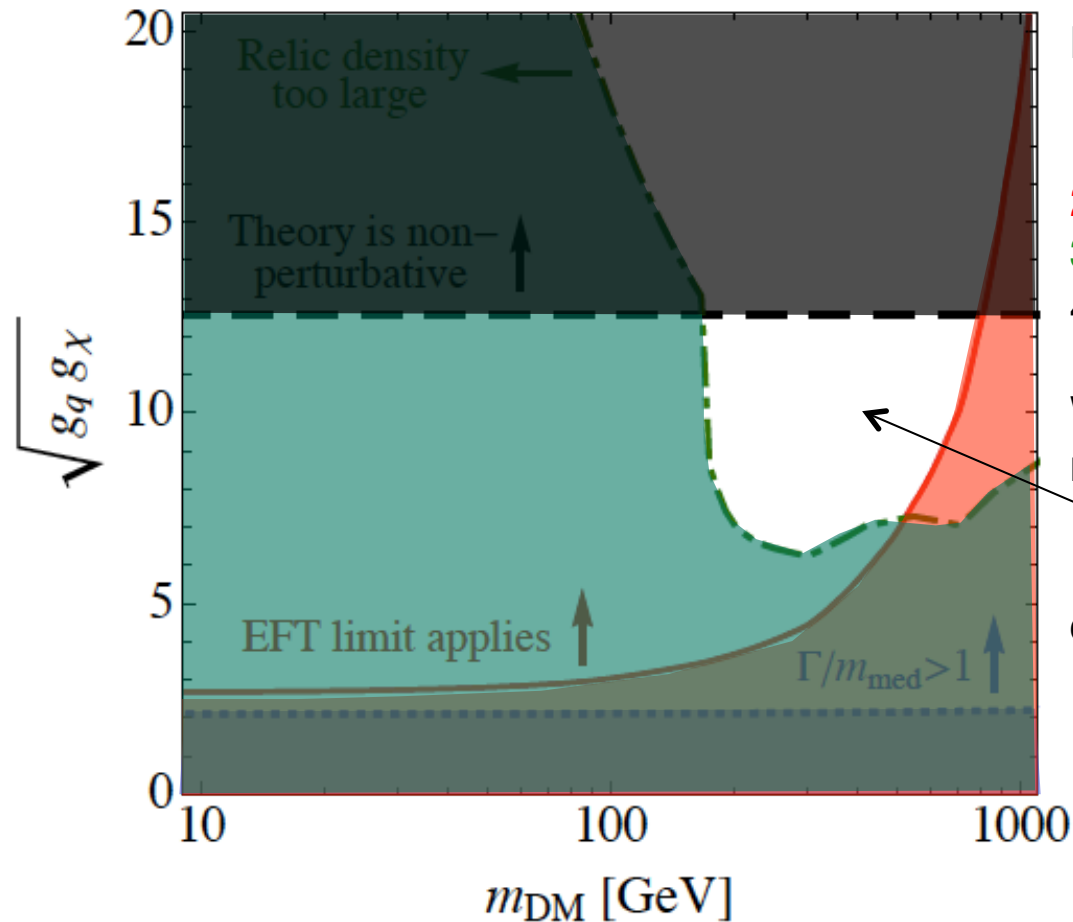
Look at EFT validity in $m_{\text{DM}} - \text{coupling}^*$ plane!

1. Must require $m_{\text{med}} < \Gamma_{\text{med}}$
2. Region in which EFT is valid (20%)
3. Require compatibility with relic density

When exclude the region in which relic abundance is larger then the observed value of $\Omega_{\text{xx}} h^2 = 0.119$ only mediator masses above a few hundred GeV fulfill this.

* Coupling chose such that CMS EFT limit on Λ applies to FT

What those this imply on model-dependences of EFT limits?



Look at EFT validity in $m_{\text{DM}} - \text{coupling}^*$ plane!

1. Must require $m_{\text{med}} < \Gamma_{\text{med}}$
2. Region in which EFT is valid (20%)
3. Require compatibility with relic density
4. Require theory to be perturbative ($< 4\pi$)

When we also require that the region/theory must be perturbative:

$$\sqrt{g_q g_\chi} < 4\pi$$

only a very small region is left!

This together with the observation that all DM theories for which the EFT is valid must have $m_{\text{med}} < \Gamma_{\text{med}}$ leads to the conclusion the the EFT only applies to a very (as in VERY) small class of DM models.

EFT limits of monojet searches are therefore highly model-dependent!